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# The wave equation with oscillating density : observability at low frequency

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## 1 Introduction and results

Let  $\Omega$  be a smooth bounded domain in  $\mathbb{R}^d$ , and  $\rho(x, y)$  a smooth function on  $\mathbb{R}^d \times \mathbb{R}^d$ , such that

$$(1.1) \quad 0 < \rho_{\min} \leq \rho(x, y) \leq \rho_{\max} \quad \forall (x, y)$$

$$(1.2) \quad \rho \text{ is } 2\pi\text{-periodic with respect to the second variable, i.e.}$$

$$\rho(x, y) = \rho(x, y + 2\pi\ell) \quad \forall \ell \in \mathbb{Z}^d.$$

For  $\varepsilon > 0$ , let  $(\omega_n^\varepsilon, e_n^\varepsilon(x))$  be the spectrum of the Dirichlet problem for the operator  $-\rho^{-1}(x, x/\varepsilon)\Delta_g$  on  $L^2(\Omega; \rho(x, x/\varepsilon)d_g x)$  normalized in the form

$$(1.3) \quad \begin{cases} \rho(x, x/\varepsilon)(\omega_n^\varepsilon)^2 e_n^\varepsilon(x) = -\Delta_g e_n^\varepsilon(x) & \text{in } \Omega \\ e_n^\varepsilon(x) = 0 & \text{on } \partial\Omega \\ \int_\Omega e_n^\varepsilon(x) \overline{e_m^\varepsilon(x)} \rho(x, x/\varepsilon) d_g x = \delta_{n,m} & ; \quad 0 < \omega_1^\varepsilon \leq \omega_2^\varepsilon \leq \dots \end{cases}$$

Here,  $\Delta_g$  denotes the Laplace operator for some fixed smooth metric  $g$  on  $\overline{\Omega}$ , and  $d_g x$  is the volume form associated to  $g$ .

For any given  $\gamma_0 > 0$ , we shall denote by  $J_{\gamma_0}^\varepsilon$  the space of solutions  $u^\varepsilon(t, x)$  of the wave equation with oscillating density  $\rho$

$$(1.4) \quad \begin{cases} (\rho(x, x/\varepsilon)\partial_t^2 - \Delta_g) u^\varepsilon(t, x) = 0 & \text{in } \mathbb{R}_+ \times \Omega \\ u^\varepsilon(t, x)|_{x \in \partial\Omega} = 0 \end{cases}$$

with maximum frequency less than  $\gamma_0/\varepsilon$ .

In other words,  $J_{\gamma_0}^\varepsilon$  is the set

$$(1.5) \quad J_{\gamma_0}^\varepsilon = \left\{ u^\varepsilon(t, x) = \sum_{\varepsilon\omega_n^\varepsilon \leq \gamma_0} (u_{+,n} e^{it\omega_n^\varepsilon} + u_{-,n} e^{-it\omega_n^\varepsilon}) e_n^\varepsilon(x) \right\}.$$

Let  $\{u_k^{\varepsilon_k}\}$  be a bounded sequence (in  $L_{\text{loc}}^2(\mathbb{R}_t, L^2(\Omega))$ ), of solutions of (1.4), with  $\lim \varepsilon_k = 0$ . It is well known that any weak limit of this sequence will satisfy the homogenized wave equation in  $\Omega$

$$(1.6) \quad \begin{cases} (\underline{\rho}(x)\partial_t^2 - \Delta_g) u(t, x) = 0 & \text{in } \mathbb{R} \times \Omega \\ u(t, x)|_{x \in \partial\Omega} = 0 \end{cases}$$

where  $\underline{\rho}(x) = \oint \rho(x, y) dy$  is the mean value of  $\rho$ .

Let  $V$  be an open subset of  $\Omega$ , and  $T_0 > 0$ .

One says that waves solution of (1.6) are observable from  $V$  in time  $T_0$  if there exists a constant  $C_0$  s.t for any  $L^2$ -solution of (1.6) one has

$$(1.7) \quad \int_0^{T_0} \int_\Omega |u(t, x)|^2 \underline{\rho}(x) dt d_g x \leq C_0 \int_0^{T_0} \int_V |u(t, x)|^2 \underline{\rho}(x) dt d_g x.$$

If  $u = \sum_{\pm, n} u_{\pm, n} e^{\pm it\omega_n} e_n(x)$  is the Fourier series of  $u$  in the spectral decomposition of  $(-\underline{\rho})^{-1}(x)\Delta_g$ , this condition is equivalent to the following

$$(1.8) \quad \begin{cases} \exists C_0 \text{ s.t. } \forall (u_{+,n}, u_{-,n})_n \in \ell^2 \times \ell^2 \\ \sum_n |u_{+,n}|^2 + |u_{-,n}|^2 \leq C_0 \int_0^{T_0} \int_V |u(t, x)|^2 \underline{\rho}(x) dt d_g x \end{cases}.$$

It is proved in [B.L.R.] that (1.7) holds true under the geometric-control hypothesis

$$(1.9) \quad \begin{cases} 1) & \text{there is no infinite order of contact between the boundary} \\ & \partial\Omega \text{ and the bicharacteristics of } \underline{\rho}(x)\partial_t^2 - \Delta_g \\ 2) & \text{any generalized bicharacteristic of } \underline{\rho}(x)\partial_t^2 - \Delta_g \\ & \text{parameterized by } t \in ]0, T_0[ \text{ meets } \bar{V} \end{cases}$$

Here the generalized bicharacteristic flow is the one defined by Melrose and Sjöstrand in [M-S].

Our main result is the following theorem, which asserts that the estimate (1.7) remains true under the hypothesis (1.9) for  $\underline{\rho}(x)$ , for solutions of (1.4) in  $J_{\gamma_0}^\varepsilon$ , if  $\gamma_0$  is small enough.

**Theorem 1.1** *Let the hypothesis (1.9) be satisfied. There exist small positive constants  $\gamma_0, \varepsilon_0$  and a constant  $C_0$ , such that for any  $\varepsilon \in ]0, \varepsilon_0[$  and any  $u^\varepsilon \in J_{\gamma_0}^\varepsilon$*

$$(1.10) \quad \int_0^{T_0} \int_{\Omega} |u^\varepsilon(t, x)|^2 \rho(x, x/\varepsilon) dt d_g x \leq C_0 \int_0^{T_0} \int_V |u^\varepsilon(t, x)|^2 \rho(x, x/\varepsilon) dt d_g x.$$

This is clearly a stability result of the observability estimate (1.7) under the singular perturbation  $\underline{\rho}(x) \rightarrow \rho(x, x/\varepsilon)$ . Let us recall that Theorem 1.1 has been proved in the 1-d case by C. Castro and E. Zuazua [C-Z], and that in the 1-d case, the counter-example of Avallaneda-Bardos-Rauch shows that (1.10) failed for  $\gamma_0$  large. Indeed, in the 1-d case, when  $\rho = \rho(x/\varepsilon)$ , C. Castro [C] has shown that the greatest value of  $\gamma_0$  such that (1.10) holds true for some  $T_0$  (when  $V \Subset [a, b] = \Omega$ ) is related with the first instability interval of the Hillé equation on the line  $(\frac{d}{dy})^2 + \omega^2 \rho(y)$ . In the multi-d case, the understanding of the best value of  $\gamma_0$  such that (1.10) holds true will clearly involved the understanding of the localization and propagation of Bloch waves for the boundary value problem (1.4) : this highly difficult problem is out of the scope of the present work.

The conserved energy for solutions of (1.4) is

$$(1.11) \quad E(u^\varepsilon) = \frac{1}{2} \int_{\Omega} \{ |\partial_t u^\varepsilon|^2 \rho(x, x/\varepsilon) + |\nabla_g u^\varepsilon|^2 \} d_g x.$$

Applying the estimate (1.10) to  $\partial_t u^\varepsilon$ , one easily gets the energy observability estimate

**Corollary 1.1** *Under the hypothesis and with the notations of Theorem 1.1, there exists a constant  $C_0$  s.t. for any  $\varepsilon \in ]0, \varepsilon_0[$  and any  $u^\varepsilon \in J_{\gamma_0}^\varepsilon$  one has*

$$(1.12) \quad E(u^\varepsilon) \leq C_0 \int_0^{T_0} \int_V |\partial_t u^\varepsilon|^2 \rho(x, x/\varepsilon) dt d_g x.$$

## 2 Sketch of proof

1. Reduction to a semi-classical estimate
2. The Bloch-wave
3. Lopatinski estimate
4. Propagation estimate

1. In the first part, using a Littlewood-Paley decomposition, we reduce the proof of the inequality (1.10) to the assertion

$$(2.1) \quad \left\{ \begin{array}{l} \text{There exist } \gamma_0, \varepsilon_0, h_0, C_0 \text{ such that for any } \varepsilon \in ]0, \varepsilon_0[, \text{ and} \\ h \in [\varepsilon/\gamma_0, h_0] \text{ the inequality (1.10) holds true for any } u^\varepsilon \in I_h^\varepsilon, \\ \text{where } I_h^\varepsilon = \left\{ u^\varepsilon = \sum_{0.9 \leq \omega_n^\varepsilon h \leq 2.1} (u_{+,n} e^{it\omega_n^\varepsilon} + u_{-,n} e^{-it\omega_n^\varepsilon}) e_n^\varepsilon(x) \right\}. \end{array} \right.$$

2. In the second part, we choose a coordinate system near the boundary

$$(2.2) \quad \left\{ \begin{array}{l} \partial\Omega \times [0, r_0] \xrightarrow{\Theta} \mathbb{R}^d \\ (x', x_d) \mapsto \Theta(x', x_d) \end{array} \right.$$

which satisfies

$$(2.3) \quad \left\{ \begin{array}{l} i) \Theta(\partial\Omega \times [0, r_0]) \subset \bar{\Omega} \\ ii) \text{ for } x_d \text{ small, } x_d \mapsto \Theta(x', x_d) \text{ is the geodesic normal to the} \\ \text{boundary at } x' \in \partial\Omega, \text{ for the metric } g \text{ on } \bar{\Omega}. \end{array} \right.$$

In these coordinates, the Laplace operator takes the form

$$(2.4) \quad \left\{ \begin{array}{l} \Delta_g = \frac{\partial}{\partial x_d} \left( A_0(x) \frac{\partial}{\partial x_d} + A_1(x, \partial_{x'}) \right) + A_2(x, \partial_{x'}) ; \\ x = (x', x_d), x' \in \partial\Omega \end{array} \right.$$

where  $A_j(x, \partial_{x'})$  are differential operators of order  $j$  on  $\partial\Omega$ , with  $x_d$  as parameter. Let  $a_j(x, \xi')$  be the principal symbol of  $A_j$ . The dual metric  $g^{-1}(x, \xi) \stackrel{\text{def}}{=} \|\xi\|_x^2$  on the cotangent bundle  $T^*\Omega$  is

$$(2.5) \quad \|\xi\|_x^2 = a_0(x) \xi_d^2 + a_1(x, \xi') \xi_d + a_2(x, \xi').$$

Let  $\mathbb{T}^d = (\mathbb{R}/2\pi\mathbb{Z})^d$  be the  $d$ -dimensional torus and for  $\varepsilon > 0$ ,  $S_\varepsilon \subset \partial\Omega \times [0, r_0] \times \mathbb{T}_y^d$  the subvariety

$$(2.6) \quad S_\varepsilon = \{(x, y); y = \Theta(x)/\varepsilon \bmod (2\pi\mathbb{Z})^d\}$$

Let  $f(x)$  be a function on  $\partial\Omega \times [0, r_0]$ . We define a distribution  $T(f)$  on  $\partial\Omega \times [0, r_0] \times \mathbb{T}_y^d$  by the formula

$$(2.7) \quad T(f) = \sum_{\ell \in \mathbb{Z}^d} e^{i\ell(y - \Theta(x)/\varepsilon)} f(x) = (2\pi)^d \delta_{y = \Theta(x)/\varepsilon} \otimes f(x).$$

If  $X$  is a vector field on  $\partial\Omega \times [0, r_0]$ , we shall denote by  $X_\varepsilon^*$  the lift of  $X$  on  $S_\varepsilon$ . If  $x' = (x_1, \dots, x_{d-1})$  is a local coordinate system on  $\partial\Omega$ , and  $(\Theta_1(x), \dots, \Theta_d(x)) = \Theta(x)$  are the Cartesian coordinates of  $\Theta$ , one has

$$(2.8) \quad \left(\frac{\partial}{\partial x_k}\right)_\varepsilon^* = \frac{\partial}{\partial x_k} + \frac{1}{\varepsilon} \sum_{j=1}^d \frac{\partial \Theta_j}{\partial x_k}(x) \frac{\partial}{\partial y_j} \quad \text{for } 1 \leq k \leq d.$$

and

$$(2.9) \quad \left(\frac{\partial}{\partial x_k}\right)_\varepsilon^* T(f) = T\left(\frac{\partial}{\partial x_k} f\right) \quad \text{for } 1 \leq k \leq d.$$

The Bloch-operator on  $\partial\Omega \times [0, r_0] \times \mathbb{T}^d$  is defined by

$$(2.10) \quad \begin{cases} \mathbb{B}_\varepsilon(x, \varepsilon \partial_x, \varepsilon \partial_t; y, \partial_y) = \hat{\rho}(x, y)(\varepsilon \partial_t)^2 - \varepsilon^2 (\Delta_g)_\varepsilon^*; & \hat{\rho}(x, y) = \rho(\Theta(x), y) \\ (\Delta_g)_\varepsilon^* = \left(\frac{\partial}{\partial x_d}\right)_\varepsilon^* \left(A_0(x) \left(\frac{\partial}{\partial x_d}\right)_\varepsilon^* + A_1(x, (\partial_{x'})_\varepsilon^*)\right) + A_2(x, (\partial_{x'})_\varepsilon^*) \end{cases}$$

It satisfies the identity

$$(2.11) \quad \mathbb{B}_\varepsilon(T(u(x, t))) = \varepsilon^2 T(\rho(\Theta(x), \Theta(x)/\varepsilon) \partial_t^2 - \Delta_g)(u(x, t))$$

Let  $\tilde{A}_j$  be the operators

$$(2.12) \quad \tilde{A}_j = A_j(x, (\partial_{x'})_\varepsilon^*)$$

and let  $e_k(x)$   $1 \leq k \leq d$  be the vectors of  $\mathbb{R}^d$

$$(2.13) \quad e_k(x) = \frac{\partial \Theta}{\partial x_k}(x)$$

If  $v(t, x, y)$  is a distribution on  $\overset{\circ}{X} \times \mathbb{T}^d$ , with  $X = \mathbb{R}_t \times (\partial\Omega \times [0, r_0])$ , we shall write the equation  $\mathbb{B}_\varepsilon(v) = 0$  as a  $2 \times 2$  system for the vector  $w = \mathcal{A}(v)$ .

$$(2.14) \quad \mathcal{A}(v) = \begin{bmatrix} w_0 \\ w_1 \end{bmatrix} \stackrel{\text{def}}{=} \begin{bmatrix} v \\ (A_0(x)(\varepsilon \frac{\partial}{\partial x_d})_\varepsilon^* + \varepsilon \tilde{A}_1)v \end{bmatrix}$$

This system takes the form

$$(2.15) \quad \begin{cases} \varepsilon \frac{\partial}{\partial x_d} w + \mathbb{M}w = 0 \\ \mathbb{M} = \begin{bmatrix} e_d(x) \cdot \partial_y + \varepsilon A_0^{-1}(x) \tilde{A}_1 & -A_0^{-1}(x) \\ \varepsilon^2 \tilde{A}_2 - \hat{\rho}(x, y)(\varepsilon \partial_t)^2 & e_d(x) \cdot \partial_y \end{bmatrix} \end{cases}$$

The operator  $\mathbb{M}$  will be seen as a semi-classical operator in  $t, x, \frac{\varepsilon}{i} \partial_{x'} = \xi', \frac{\varepsilon}{i} \partial_t = \tau$  with operator values in the fiber  $\mathbb{T}^d$

$$(2.16) \quad \mathbb{M} = \sum_{j=0}^2 \left(\frac{\varepsilon}{i}\right)^j \mathbb{M}^j(x, \xi', \tau; y, \partial_y).$$

The differential degree in  $y$  of  $\mathbb{M}^j$  is at most  $2-j$  and the principal symbol  $\mathbb{M}^0$  is the matrix

$$(2.17) \quad \mathbb{M}^0(x, \xi', \tau; y, \partial_y) = \begin{bmatrix} e_d(x) \cdot \partial_y + a_0^{-1}(x) a_1(x, i\xi' + e'(x) \cdot \partial_y) & -a_0^{-1}(x) \\ a_2(x, i\xi' + e'(x) \cdot \partial_y) + \hat{\rho}(x, y) \tau^2 & e_d(x) \cdot \partial_y \end{bmatrix}$$

Let  $E^\bullet = \{E^s, s \in \mathbb{R}\}$  be the scale of Hilbert spaces on the torus

$$(2.18) \quad E^s = H^s(\mathbb{T}^d) \oplus H^{s-1}(\mathbb{T}^d).$$

For any  $\rho = (x, \xi', \tau)$ ,  $\mathbb{M}^j(\rho, y, \partial_y)$  maps  $E^s$  into  $E^{s-1+j}$  and  $\mathbb{M}^0$  is an elliptic operator. Let  $\mathbb{M}_0^0$  be the restriction of  $\mathbb{M}^0$  to the zero section  $\xi' = \tau = 0$ .

$$(2.19) \quad \mathbb{M}_0^0(x, \partial_y) = \mathbb{M}^0(x, 0, 0, y, \partial_y) = \begin{bmatrix} e_d(x) \cdot \partial_y + a_0^{-1} a_1(x, e'(x) \cdot \partial_y) & -a_0^{-1}(x) \\ a_2(x, e'(x) \cdot \partial_y) & e_d(x) \cdot \partial_y \end{bmatrix}$$

The eigenvalues  $\lambda_{\pm, \ell}^0(x)$  of  $\frac{1}{i} \mathbb{M}_0^0(x, \partial_y)$  on the space  $e^{i\ell y} \mathbb{C}^2$ , for  $\ell \in \mathbb{Z}^d$  are the complex roots of the equation

$$(2.20) \quad a_0(x)(-\lambda + e_d \cdot \ell)^2 + (-\lambda + e_d \cdot \ell) a_1(x, e' \cdot \ell) + a_2(x, e' \cdot \ell) = 0$$

which is equivalent to

$$(2.21) \quad \|\mathit{t}d\Theta(x)(\ell) - \lambda(0, \dots, 0, 1)\|_x^2 = 0.$$

In particular we have

$$(2.22) \quad \inf_x \min_{\ell \neq 0} |\lambda_{\pm, \ell}^0(x)| > 0$$

so the double eigenvalue  $\lambda_{\pm, 0}^0(x) = 0$  is isolated in the spectrum of  $\mathbb{M}_0^0(x, \partial_y)$ .

In the sequel, we shall restrict the values of the Sobolev index of regularity  $s$  on the torus to some fixed large interval,  $s \in [-\sigma_0, \sigma_0]$ ,  $\sigma_0 \gg \frac{d}{2}$ .

Let  $X = \partial\Omega \times \mathbb{R}_t \times [0, r_0]$ . We denote by  $\mathit{t}T^*X$  the tangential cotangent bundle

$$(2.23) \quad \mathit{t}T^*X = T^*(\partial\Omega \times \mathbb{R}_t) \times [0, r_0]$$

Let  $W_1 \Subset W_0$  be two small neighborhoods of the set  $\{\xi' = \tau = 0\} \times \{t \in [-T_0, 2T_0]\}$  in  $\mathit{t}T^*X$ .

We choose a non-negative function  $\chi_0 \in C_0^\infty(W_0)$ , such that  $\chi_0 \equiv 1$  on  $W_1$ .

If  $W_0$  is small enough, we define the map  $p_0(x, t, \xi', \tau) : E^\bullet \rightarrow \mathbb{C}^2$  by the formula

$$(2.24) \quad p_0[w] = \chi_0 \cdot \oint_{\mathbb{T}^d} \left\{ \frac{1}{2i\pi} \int_{\partial D} \frac{dz}{z - \mathbb{M}^0} \right\} [w] \quad w \in E^s, \quad s \in [-\sigma_0, \sigma_0]$$

(where  $D \subset \mathbb{C}$  is a small disk with center  $z = 0$ ).

It satisfies the estimates

$$(2.25) \quad \exists C \forall s \in [-\sigma_0, \sigma_0] \forall w \in E^s \quad \|p_0(w) - \chi_0 \oint_{\mathbb{T}^d} w\|_{\mathbb{C}^2} \leq C\tau^2 \|w\|_{E^s}$$

and there exists  $L^0(x, t, \xi', \tau) \in C^\infty(\mathit{t}T^*X; M_2(\mathbb{C}))$ , defined near  $\xi' = \tau = 0$  such that

$$(2.26) \quad p_0 \circ \mathbb{M}^0 = L^0 \circ p_0.$$

By a Taylor expansion near  $\xi' = \tau = 0$ , one gets

$$(2.27) \quad L^0 = \begin{bmatrix} a_0^{-1}(x)a_1(x, i\xi') & -a_0^{-1}(x) \\ a_2(x, i\xi') + \hat{\rho}(x)\tau^2 & 0 \end{bmatrix} + O(\tau^4)$$



We then suitably quantize the above construction and we obtain tangential pseudo differential operators

$$(2.28) \quad \begin{cases} \Pi_0(\varepsilon, t, x, \varepsilon\partial_t, \varepsilon\partial_{x'}) & : L^2(X; E^s) \rightarrow L^2(X, \mathbb{C}^2), s \in [-\sigma_0, \sigma_0] \\ L(\varepsilon, t, x, \varepsilon\partial_t, \varepsilon\partial_{x'}) & : L^2(X; \mathbb{C}^2) \rightarrow L^2(X, \mathbb{C}^2) \end{cases}$$

with principal symbol  $\sigma(\Pi_0) = p_0$ ,  $\sigma(L) = L^0$ , which satisfy the relation

$$(2.29) \quad \Pi_0(\varepsilon\partial_{x_d} + \mathbb{M}) = (\varepsilon\partial_{x_d} + L)\Pi_0 + R(\varepsilon, t, x, \varepsilon\partial_t, \varepsilon\partial_{x'})$$

In (2.29), the error term  $R : L^2(X; E^s) \rightarrow L^2(X, \mathbb{C}^2)$  will be a tangential pseudo differential operator such that for any tangential o.p.d.  $Q$  with essential support in  $W_1$  and any  $s \in [-\sigma_0, \sigma_0]$ , one has

$$(2.30) \quad \|Q \circ R; L^2(X; E^s) \rightarrow L^2(X, \mathbb{C}^2)\| \in O(\varepsilon^\infty)$$

**Definition 2.1** For  $u^\varepsilon \in I_h^\varepsilon$ , we define the Bloch-wave  $\Gamma(u^\varepsilon) \in L^2(X; \mathbb{C}^2)$  by the formula

$$(2.31) \quad \Gamma(u^\varepsilon) = \begin{bmatrix} \Gamma_0(u^\varepsilon) \\ \Gamma_1(u^\varepsilon) \end{bmatrix} = \Pi_0 \mathcal{T}(u^\varepsilon) \quad (\mathcal{T} = \mathcal{A} \circ T)$$

■

Let  $\gamma_0, \varepsilon_0, h_0$  be given small enough,  $\varepsilon \in ]0, \varepsilon_0]$ ,  $h \in [\varepsilon/\gamma_0, h_0]$ . For  $u_\varepsilon \in I_h^\varepsilon$ ,  $u^\varepsilon = \sum_{0.9 \leq \omega_n^\varepsilon h \leq 2.1} (u_{+,n} e^{it\omega_n^\varepsilon} + u_{-,n} e^{-it\omega_n^\varepsilon}) e_n^\varepsilon(x)$ , we define  $\|u^\varepsilon\|^2 \left( \simeq \int_0^{T_0} \int_\Omega |u^\varepsilon|^2 \right)$

by

$$(2.32) \quad \|u^\varepsilon\|^2 = \sum_{0.9 \leq \omega_n^\varepsilon h \leq 2.1} |u_{+,n}|^2 + |u_{-,n}|^2$$

Let  $X_{T_0} = \partial\Omega \times [-T_0, 2T_0] \times [0, r_0]$ , and let  $K$  be the compact subset of  ${}^tT^*X$ ,  $K = \partial\Omega \times [0, T_0] \times [0, r_0/2] \times \{\xi' = 0, \tau = 0\}$ . By a localization argument near the zero section ( $\gamma_0$  small), and a propagation argument in the interior, we first verify

**Proposition 2.1** Let  $Q(\varepsilon, t, x, \varepsilon\partial_{x'}, \varepsilon\partial_t)$  be a zero order tangential opd on  $X$ , equal to  $Id$  near  $K$ . If  $\gamma_0, \varepsilon_0, h_0$  are small enough, there exists a constant  $C > 0$ , such that for any  $\varepsilon \in ]0, \varepsilon_0]$ ,  $h \in [\varepsilon/\gamma_0, h_0]$ , one has

$$(2.33) \quad \|u^\varepsilon\|^2 \leq C \left[ \|Q\Gamma_0(u^\varepsilon)\|_{L^2(X_{T_0})}^2 + \|u^\varepsilon\|_{L^2((0, T_0) \times V)}^2 \right] \quad \forall u^\varepsilon \in I_h^\varepsilon$$

**3.** By Proposition 2.1, we shall obtain the inequality (1.10), if we are able to estimate the  $L^2$  norm of the first component  $\Gamma_0(u^\varepsilon)$  of the Bloch-wave near the set  $K$ .

The formula (2.29) shows that  $\Gamma(u^\varepsilon)$  satisfies the equation

$$(2.34) \quad (\varepsilon \partial_{x_d} + L)\Gamma(u^\varepsilon) \in O(\varepsilon^\infty L^2) \text{ (microlocally in } W_1).$$

By (2.27) this equation is very closed to the homogenized equation  $(\underline{\rho}(x)\partial_t^2 - \Delta_g)[\Gamma_0(u^\varepsilon)] = 0$ .

As one can see, all the difficulty in our problem is thus to obtain an estimate on the first Dirichlet data of  $\Gamma(u^\varepsilon)$  on the boundary  $x_d = 0$ , in order to apply propagation arguments to the equation (2.34).

**Proposition 2.2** *If  $\gamma_0, \varepsilon_0, h_0$  are small enough, there exists a constant  $C$  such that for any  $\varepsilon \in ]0, \varepsilon_0[$ ,  $h \in [\varepsilon/\gamma_0, h_0]$  the following estimate holds true*

$$(2.35) \quad \|\Gamma_0(u^\varepsilon)|_{x_d=0}\|_{L^2(X_{T_0} \cap x_d=0)} \leq C \varepsilon/h \|u^\varepsilon\| \quad \forall u^\varepsilon \in I_h^\varepsilon.$$

■

The above estimate is obtain as a consequence of a uniform Lopatinski estimate on  $w^\varepsilon = \mathcal{T}(u^\varepsilon) = \begin{bmatrix} w_0^\varepsilon \\ w_1^\varepsilon \end{bmatrix}$ . More precisely, we have

**Theorem 2.1** *Let  $Q$  be a scalar tangential o.p.d. with essential support in  $W_0$  ; if  $W_0, \gamma_0, \varepsilon_0, h_0$  are small enough, there exist  $s_1 < 0$  and a constant  $C$  such that for any  $u^\varepsilon \in I_h^\varepsilon$  the following estimate holds true*

$$(2.36) \quad \|Q(t, x, \varepsilon \partial_{x'}, \varepsilon \partial_t)(w_1^\varepsilon)|_{x_d=0}\|_{L^2(X_{T_0} \cap x_d=0, H^{s_1}(\mathbb{T}^d))} \leq C \|u^\varepsilon\|$$

■

Notice that  $w^\varepsilon$  satisfies the equation (2.15), with Dirichlet data  $w_0^\varepsilon|_{x_d=0} = 0$  on the boundary.

The weaker estimate

$$(2.37) \quad \|Q(w_1^\varepsilon)|_{x_d=0}\| \leq C \varepsilon^{-1/2} \|u^\varepsilon\|$$

is easy to obtain (it is sufficient to commute the equation (1.4) with the normal vector field  $\frac{\partial}{\partial n}$ ).

The proof of (2.36) is the most technical part of our work. It involves a detail study of how the spectral theory of  $\mathbb{M}^0(x, \xi', \tau; y, \partial y)$  (see (2.15)) depends on the parameter  $(x, \xi', \tau)$ . This involves both the real and the complex part of the Bloch variety.

**4.** The last part is devoted to the proof of the following proposition.

**Proposition 2.3** *Let  $Q(\varepsilon, t, x, \varepsilon \partial_{x'}, \varepsilon \partial_t)$  be a zero order opd equal to  $Id$  near  $K$ , with essential support in  $W_1$ . There exist  $\gamma_0, \varepsilon_0, h_0$ , and a constant  $C_0$  such that, for any  $\varepsilon \in ]0, \varepsilon_0]$ ,  $h \in [\varepsilon/\gamma_0, h_0]$  and  $u^\varepsilon \in I_h^\varepsilon$ , the following estimate holds true*

$$(2.38) \quad \|Q\Gamma_0(u^\varepsilon)\|_{L^2(X_{T_0})}^2 \leq C_0 \left[ \|\Gamma_0(u^\varepsilon)|_{x_d=0}\|_{L^2(X_{T_0} \cap x_d=0)}^2 + \|u^\varepsilon\|_{L^2(0, T_0) \times V}^2 \right]$$

■

This estimate is obtain by rather classical arguments in the theory of control of linear waves, for the rescale equation

$$(2.39) \quad \begin{cases} \left( h \frac{\partial}{\partial x_d} + \mathcal{L} \right) \begin{bmatrix} g_0 \\ g_1 \end{bmatrix} \sim 0 & \begin{bmatrix} g_0 \\ g_1 \end{bmatrix} = \begin{bmatrix} \Gamma_0(u^\varepsilon) \\ \frac{h}{\varepsilon} \Gamma_1(u^\varepsilon) \end{bmatrix} \\ \mathcal{L} = \frac{h}{\varepsilon} \begin{pmatrix} 1 & 0 \\ 0 & h/\varepsilon \end{pmatrix} \circ L \circ \begin{pmatrix} 1 & 0 \\ 0 & \varepsilon/h \end{pmatrix} \end{cases}$$

We verify that  $\mathcal{L}$  is still a  $h$ -pseudo differential operator, with  $\varepsilon/h$  as parameter. (We use this rescaling in order to be able to use propagation arguments in the range  $\varepsilon \ll h$ ).

We finally remark that the validity of (2.1), hence the proof of Theorem 0.1, is a direct consequence of the Proposition 2.1, 2.2 and 2.3.

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