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CHARACTERISTIC CAUCHY PROBLEM.

by Anders Melin

Introduction.

We shall consider a real-valued function $v \in C^{\infty}(\mathbf{R}^n)$ when n > 1 is odd. In order to have sufficiently regular scattering data associated to the Schrödinger operator $H_v = -\Delta_x + v(x)$ we shall assume that v satisfies the following short-range condition:

$$\int_{\mathbf{R}^n} (1+|x|)^{2-n+|\alpha|} |v^{(\alpha)}(x)| \, dx < \infty \quad \text{for any } \alpha. \tag{1}$$

The class of such potentials will be denoted V. By using polar coordinates in the frequency variables one may write the Lippmann-Schwinger equation on the form

$$(-\Delta_x + v(x))\phi(x,\theta,k) = k^2\phi(x,\theta,k), \quad x \in \mathbf{R}^n, \ \theta \in S^{n-1}, \ k \in \mathbf{R}.$$
 (2)

One has also to impose some condition on $\phi(x, \theta, k)$ as $|x| \to \infty$ in order to obtain a unique solution of (2). We shall always consider ϕ as a perturbation of the function $\phi_0(x, \theta, k) = e^{ik\langle x, \theta \rangle}$ which solves (2) when v = 0. Moreover, ϕ will be a continuous function of $k \in \mathbb{R} \setminus 0$ with a meromorhic extension to the upper half-plane. If $0 < \Im k$ is small then

$$\phi = \phi_0 - (H_0 - k^2)^{-1}(v\phi),$$

where $(H_0 - k^2)^{-1}$ is the L^2 - bounded inverse of $H_0 - k^2$. In the case of a compactly supported potential v this leads to the formula

$$\phi(x,\theta,k) - \phi_0(x,\theta,k) = 2^{-1} \left(\frac{4\pi}{ik|x|} \right)^{(n-1)/2} e^{ik|x|} T(k,x/|x|,\theta) + O(|x|^{-(n+1)/2}), \quad (3)$$

where T is the scattering amplitude. We also remark that ϕ can be defined in terms of the distribution kernels of the wave operators $W_{\pm} = \lim_{t \to \pm \infty} e^{itH_v} e^{-itH_0}$, and one often calls the solutions of (2) generalized eigenfunctions.

In this note we show how ϕ , or rather its Fourier transform w.r.t. the variable k, can be obtained as the solution of a characteristic Cauchy problem for the differential operator $\Delta_x - \partial_t^2 - v(x)$. This viewpoint will give us extra information about ϕ and enables us to prove that

$$e^{-ik\langle x,\theta\rangle}\phi(x,\theta,k) = 1 + \int_0^\infty w_\theta(x,t)e^{ikt} dt, \tag{4}$$

where $w_{\theta}(x,t)$ is a smooth function. In particular we shall recover an identity which is usually referred to as the miracle (cf [N1, N2, C]). We also remark that part of the discussions here can be carried over to the case of more general short range potentials.

Construction of ϕ by means of intertwining operators.

We shall first consider the equation

$$(\Delta_x - \Delta_y - v(x))A_\theta(x, y) = 0. \tag{5}$$

In [M5] it was proved that this equation has a solution which is supported in the set $(y-x,\theta) \ge 0$ and given by a series

$$\sum_{0}^{\infty} U_{N,\theta}(x,y),$$

where

$$(\Delta_x - \Delta_y)U_{N+1,\theta}(x,y) = v(x)U_N(x,y), \quad U_0(x,y) = \delta(x-y).$$

In order to describe the regularity of the solution one introduces the set \mathcal{P}_{λ} of all seminorms

$$p(U) = \sup_{x} \int_{\mathbf{R}^{n}} e^{-\lambda \langle y - x, \theta \rangle} |(\partial_{x} + \partial_{y})^{\alpha} (\langle x, \partial_{x} \rangle + \langle y, \partial_{y} \rangle)^{\beta} U(x, y)| \, dy.$$

Then for each v which satisfies (1) there is a $\lambda = \lambda_v \geq 0$ so that

$$\sum_{1}^{\infty} p(U_{N,\theta}) < \infty, \quad p \in \mathcal{P}_{\lambda}. \tag{6}$$

Moreover, for each $m \geq 0$ there is a positive integer N(m) so that

$$\sum_{N(m)}^{\infty} p(\partial_x^{\alpha} \partial_y^{\beta} U_{N,\theta}) < \infty, \quad |\alpha + \beta| \le m, \ p \in \mathcal{P}_{\lambda}.$$
 (7)

In order to make the $U_{N,\theta}$ unique one also has to introduce some conditions at infinity which will exclude from the considerations functions which are constant in the direction of (θ, θ) . We shall not discuss these details here.

Next we introduce

$$V_{N,\theta}(x,t) = \int_{\langle y-x,\theta\rangle=t} U_{N,\theta}(x,y) \, dy,$$

and we let Q_{λ} be the family of semi-norms

$$q(V) = \sup_{x} \int_{0}^{\infty} e^{-\lambda t} |\partial_{x}^{\alpha}(\langle x, \partial_{x} \rangle + t \partial_{t})^{\beta} V(x, t)| dt.$$

It follows from (6) and (7) then that

$$\sum_{1}^{\infty} q(V_{N,\theta}) < \infty, \quad q \in \mathcal{Q}_{\lambda}, \tag{6}$$

and

$$\sum_{N(m)}^{\infty} q(\partial_x^{\alpha} \partial_t^{\beta} V_{N,\theta}) < \infty, \quad |\alpha| + \beta \le m, \ q \in \mathcal{Q}_{\lambda}, \tag{7}$$

It was proved in [M5] that

$$\phi(x,\theta,k) = \int A_{\theta}(x,y)e^{ik\langle y,\theta\rangle} dy.$$

Hence if $V_{\theta}(x,t) = \sum_{0}^{\infty} V_{N,\theta}(x,t)$ we must (in view of the definition of $V_{N,\theta}$) have

$$e^{-ik\langle x,\theta\rangle}\phi(x,\theta,k) = \int e^{itk}V_{\theta}(x,t)\,dt.$$
 (8)

It follows from (6)' that the integrand is continuous w.r.t. x and integrable w.r.t. t when $\Im k$ is large enough. Moreover, $t \ge 0$ in the support of V_{θ} , and $V_{\theta}(x,t) = \delta(t)$ if v = 0.

The main result.

It follows immediately from (6)' that $V_{\theta}(x,t)$ is a smooth function of x and t when t > 0, and the next result implies that one may write $V_{\theta}(x,t) = \delta(t) + Y_{+}(t)w_{\theta}(x,t)$, where Y_{+} is the Heaviside function and $w_{\theta}(x,t)$ is smooth when $x \in \mathbb{R}^{n}$ and $t \geq 0$.

Theorem 1. There is a positive number λ such that

$$\sup_{x,\theta} \int_{+0}^{\infty} e^{-\lambda t} |\partial_x^{\alpha} \partial_t^{\beta} V_{\theta}(x,t)| \langle x \rangle^{-\beta} dt < \infty \tag{9}$$

for any α and β . (Here $\langle x \rangle = (1 + |x|^2)^{1/2}$.)

We have already seen that $V_{\theta}(x,t)$ is smooth when t>0, and it follows from (6)' also that we need only consider the integral over the interval (0,1) in (9). Moreover, the estimates (7)' imply that it it suffices to prove a similar result for each of the $V_{N,\theta}$. Hence Theorem 1 results from the following

Theorem 1'. If $N \geq 1$, then

$$\sup_{x,\theta} \int_{t=0}^{1} \langle x \rangle^{-\beta} |\partial_x^{\alpha} \partial_t^{\beta} V_{N,\theta}(x,t)| dt < \infty.$$
 (10)

We have now come to the point where one has to study the wave equation. In fact, the equation $(\Delta_x - \Delta_y)U_{N+1,\theta}(x,y) = v(x)U_{N,\theta}(x,y)$ implies that

$$\mathcal{L}_{\theta}V_{N+1,\theta}(x,t) = v(x)V_{N,\theta}(x,t), \tag{11}$$

if $\mathcal{L}_{\theta} = \Delta_x - 2\langle \theta, \partial_x \rangle \partial_t$. We observe that \mathcal{L}_{θ} is obtained from the wave operator $\Delta_x - (\partial_{t_0})^2$, after the substitution $t_0 = t + \langle x, \theta \rangle$. Hence $t \geq 0$ is a characteristic half-space for \mathcal{L}_{θ} . We let G_{θ} be the image of the fundamental solution of $\Delta_x - 2\langle \theta, \partial_x \rangle \partial_t$ with support in the set $t_0 \geq 0$ under the substitution above. Then $t \geq 0$ in the support of G_{θ} .

We shall consider approximate solutions of the equation $\mathcal{L}_{\theta}V(x,t) = v(x)V_{N,\theta}(x,t)$, which is solved by $V_{N+1,\theta}$. The construction of such a solution will be similar to the methods of geometrical optics used in microlocal analysis, and an exact solution will then be obtained after convolving the error term $\mathcal{L}_{\theta}V(x,t)-v(x)V_{N,\theta}(x,t)$ with some fundamental solution Q_{θ} of \mathcal{L}_{θ} with support in the set $t \geq 0$. The following result shows that one has to take $Q_{\theta} = G_{\theta}$ if one hopes to obtain good bounds for the solutions.

Proposition 2. Assume that $\mathcal{L}_{\theta}u(x,t) = 0$ and that $t \geq 0$ in the support of u. If u(x,t) is temperate w.r.t. x then

$$u(x,t) = \sum_{|\alpha| \le \mu(t)} f_{\alpha}(t) x^{\alpha},$$

where $f_{\alpha} \in \mathcal{D}'(\mathbf{R})$ and the integer valued function $\mu(t)$ is locally bounded.

PROOF: We may assume that $\theta = e_n$. The function $G(x,y) = u(x,y_n - x_n)$ then solves the equation $(\Delta_x - \Delta_y)G(x,y) = 0$, and $y_n \geq x_n$ in its support. The proof of Theorem 3.5 of [M5] shows then that

$$G(x,y) = G(x,y_n) = \sum_{n=0}^{\infty} g_j(x',y_n - x_n)x_n^j,$$

where $x' = (x_1, \ldots, x_{n-1})$ and $g_j \in \mathcal{D}'(\mathbb{R}^n)$. Then

$$u(x,t) = \sum_{0}^{\infty} g_j(x',t)x_n^j,$$

where only finitely many of the g_j are $\neq 0$ when t stays in any bounded open set ω . The equation $(\Delta_x - 2\partial_{x_n}\partial_t)u = 0$ implies that

$$\Delta_{x'}g_j(x',t) - 2(j+1)\partial g_{j+1}(x',t)/\partial t + (j+2)(j+1)g_{j+2}(x',t) = 0, \quad j = 0,1,\ldots$$

Hence, when t is in any ω as above, then $\Delta_{x'}^{N}g_{j}(x',t)=0$ for any j if N is large enough. Since the g_{j} are temperate in x', this implies that they are polynomials in this variable and the proposition follows.

Let Γ_0 be the cone $|x| = t_0$ and Γ be its image under the substitution $t_0 = t + \langle x, \theta \rangle$, i.e. Γ is defined by $t = |x| - \langle x, \theta \rangle$. The half-plane $B: t \geq 0$ corresponds to $B_0: t_0 \geq \langle x, \theta \rangle$, which intersects Γ_0 only along the ray $\{(t_0\theta, t_0); t_0 \geq 0\}$. Hence $-\Gamma_0$ intersects B_0 only along the opposite ray $\{(t_0\theta, t_0); t_0 \leq 0\}$, and any distribution u_0 supported in B_0 and vanishing over a conic neighbourhood of $\gamma = \mathbf{R}_-\theta$ has to vanish identically if it satisfies the wave equation $(\Delta_x - (\partial_{t_0})^2)u_0(x, t_0) = 0$. In fact, if (x, s) is in the wave cone Γ_0 , (y, t) belongs to the support of u_0 and x + y, s + t belong to bounded sets, then |y| can not tend to infinity unless y/|y| tends to $-\theta$. In (x, t) space this implies that $w = G_\theta * u$ is defined if $t \geq 0$ in the support of u and u vanishes over a conic neighbourhood of the ray $\gamma = \mathbf{R}_-\theta$, and w is the unique solution of the equation $\mathcal{L}_\theta w = u$ in the space of such distributions.

In order to have $G_{\theta} * u$ defined on a larger space we introduce the following definition:

Definition 3. $\mathcal{D}'_{G_{\theta}}$ is the space of all u in $\mathcal{D}'(\mathbf{R}^n \times \mathbf{R})$ such that $\lim_{j \to \infty} G_{\theta} * (\chi_j u)$ exists for any sequence $\chi_j \in C_0^{\infty}(\mathbf{R}^n \times \mathbf{R})$ such that $\|\chi_j\|_{L^{\infty}}$ is bounded, χ_j converges pointwise and $\|\chi_j^{(\alpha)}\|_{L^{\infty}} \to 0$ as $j \to \infty$ if $\alpha \neq 0$.

If $u \in \mathcal{D}'_{G_{\theta}}$, and the sequence χ_j above tends to 1 then we define $G_{\theta} * u$ as the limit of $G_{\theta} * (\chi_j u)$. This limit is independent of the choices made. Moreover, $\mathcal{D}'_{G_{\theta}}$ is invariant

under differentiation, and $u \to G_{\theta} * u$ is a left-inverse for \mathcal{L}_{θ} on this space in the sense that $u = G_{\theta} * \mathcal{L}_{\theta} u$ when $u \in \mathcal{D}'_{G_{\theta}}$.

One can show that $u \in \mathcal{D}'_{G_{\theta}}$ if $t \geq 0$ in its support and if it decays as $|x|^{-1-\varepsilon}$ over some conic neighbourhood of γ for some positive ε . One can even allow less restrictive conditions on the decay of u, however, since we are dealing with potentials in the class \mathcal{V} , we shall only consider the following conditions:

Definition 4. Let $f \in C^{\infty}(\mathbb{R}^n)$. Then we say that $f \in \mathcal{V}_{\theta}$ if $af \in \mathcal{V}$ for any $a \in S^0(\mathbb{R}^n)$ such that the support of a is contained in some cone $\varepsilon |x| \leq -\langle x, \theta \rangle$, where $\varepsilon > 0$.

Here the condition that $a \in S^0(\mathbf{R}^n)$ means that $\langle x \rangle^{|\alpha|} a^{(\alpha)}(x)$ is bounded for any α . It is easy to see that \mathcal{V} and \mathcal{V}_{θ} are Fréchet spaces and they are also S^0 modules.

We let $C^{\infty}(\overline{\mathbf{R}_{+}}) \otimes \mathcal{V}_{\theta}$ be the space of smooth maps from $\overline{\mathbf{R}_{+}}$ to \mathcal{V}_{θ} . Set $Y_{+,j}(t) = t^{j}Y_{+}(j)/j!$. If j is a non-negative integer, p = 0 or 1, then $\mathcal{W}_{\theta,j,p}$ is the space of all functions on the form $Y_{+,j}(t)\langle x \rangle^{p}U(x,t)$, where $U \in C^{\infty}(\overline{\mathbf{R}_{+}}) \otimes \mathcal{V}_{\theta}$.

Theorem 5. \mathcal{L}_{θ} is bijective from $\mathcal{W}_{\theta,j+1,1}$ to $\mathcal{W}_{\theta,j,0}$ if $j \geq 0$.

By combining this result with some uniqueness statements obtained from Proposition 2 one can easily prove Theorem 1' now by induction over N. We leave out these details and discuss instead the proof of the theorem above.

It is clear that \mathcal{L}_{θ} maps $\mathcal{W}_{\theta,j+1,1}$ into $\mathcal{W}_{\theta,j,0}$. The injectivity of the map follows since one can show that $\mathcal{W}_{\theta,0,1}$ is contained in $\mathcal{D}'_{G_{\theta}}$. Hence convolution with G_{θ} gives us a left inverse.

In order to give the main ideas of the proof of the surjectivity of the map \mathcal{L}_{θ} in the theorem we consider the corresponding situation when j = -1 so that $Y_{+,j}(t) = \delta(t)$. This leads us to discuss the equation

$$\mathcal{L}_{\theta}V(x,t) = v(x)\delta(t) \tag{12}$$

when $v \in \mathcal{V}$. We first construct an approximate solution. We set

$$v_j(x) = (2^{j+1}j!)^{-1} \int_0^\infty \Delta_x^j v(x - s\theta) ds.$$

Then

$$2\langle \theta, \partial_x \rangle v_j(x) = \Delta_x v_{j-1}(x), \quad j > 0,$$

$$2\langle \theta, \partial_x \rangle v_0(x) = v(x),$$
(13)

and $\langle x \rangle^{-1} v_j(x) \in \mathcal{V}_{\theta}$.

Choose $\zeta(t) \in C_0^{\infty}(\mathbb{R}^n)$ such that $\zeta(t) = 1$ in a neighbourhood of the origin. If the sequence $1 \leq L_i$ grows sufficiently fast, then it is true that the series

$$w(x,t) = \sum_{0}^{\infty} Y_{+,j}(t) \zeta(L_{j}t) v_{j}(x)$$

converges in C^{∞} and defines an element in $\mathcal{W}_{\theta,0,1}$. Moreover, it follows from (13) that

$$r(x,t) = \mathcal{L}_{\theta} w(x,t) - v(x) \delta(t)$$

is a smooth function of t with values in \mathcal{V}_{θ} . Moreover, it vanishes when $t \leq 0$. Since the dimension is odd, one has a simple explicit formula for G_{θ} which allows one to conclude that $w_r(x,t) = (G_{\theta} * r)(x,t)$ is a smooth function of t with values in \mathcal{V}_{θ} after multiplication by $\langle x \rangle^{-1}$. Hence by subtraction w_r from w we have obtained a solution V(x,t) of (12) such that $V \in \mathcal{W}_{\theta,0,1}$.

Remark. The proof shows that $V_{\theta}(x,\varepsilon) \to v_0(x)$ as $\varepsilon \to 0$. Hence it follows from (13) that

$$2\langle \theta, \partial_x \rangle V_{\theta}(x, \varepsilon) \to v(x)$$
 as $\varepsilon \to 0$.

This phenomenon was discovered by R.G. Newton and called the miracle by him ([N1, N2]).

Remarks about the case of exponentially decaying potentials.

We shall finally discuss a situation when the potential is exponentially decreasing. Let a be a positive number and assume that one has the estimates

$$|v^{(\alpha)}(x)| \le C_{\alpha} e^{-(2a+\varepsilon)|x|} \tag{14}$$

for every α and some positive ε . In this case it turns out that $V_{\theta}(x,t)$ will be exponentially decaying in the variable t except for some contributions to V_{θ} that are due to bound states and resonances:

Theorem 6. Assume that v satisfies (14). Then there is a finite set $Z \subset \{k \in \mathbb{C}; \Im k \geq -a\}$ so that

$$V_{\theta}(x,t) = \delta(t) + Y_{+}(t)a(x,\theta,t) + \sum_{z \in Z} \sum_{\mu \le \mu(z)} t^{\mu} e^{-itz} a_{z,\mu}(x,\theta), \tag{15}$$

where for some constants $C_{\alpha,\beta}$ and C_{α}

$$\int_0^\infty |\partial_x^\alpha \partial_t^\beta a(x,\theta,t)| e^{at} dt \le C_{\alpha,\beta} \langle x \rangle^\beta e^{a(|x| - \langle x,\theta \rangle)},$$

and

$$|a_{z,\mu}^{(\alpha)}(x,\theta)| \leq C_{\alpha} e^{a(|x|-\langle x,\theta\rangle)}.$$

All estimates are uniform in θ .

Remark. It is also possible to prove smoothness w.r.t. θ .

REFERENCES

[C] Cheney, M., Inverse scattering in dimension 2. J. Math. Phys. 25(1984), 94-107.

[F] Faddeev, L.D., Inverse problems in quantum scattering theory, II. J. Sov. Math., 5(1976), 334-396.

- [M1] —"—, Méthodes d'entrelacements dans le problème de scattering inverse. Journées équations aus dérivées partielles, Saint –Jean –de Monts, 1986.
- [M2] —"—, Some problems in inverse scattering theory. Journées équations aus dérivées partielles, Saint –Jean –de Monts, 1987.
- [M3] —"—, Some mathematical problems in inverse potential scattering. Séminaire sur les équations aus dérivées partielles 1986-1987, Ecole polytechnique, no XX, 1987.
- [M4] —"—, Intertwining methods in the theory of inverse scattering. Int. J. of Quantum Chemistry 31 (1987), 739-746.
- [M5] —"—,Intertwining methods in multi-dimensional scattering theory I. Univ. of Lund and Lund Inst. of Tech, Dept of Math. preprint series, 1987:13.
- [N1] Newton, R.G., Inverse scattering II. Three dimensions. J. Math. Phys. 21(1980), 1698-1715.
- [N2] —"—,An inverse spectral theorem in three dimensions. SIAM-AMS Proceedings 14, 81-90.

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