

## THE WAVE EQUATION WITH OSCILLATING DENSITY: OBSERVABILITY AT LOW FREQUENCY

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**Abstract.** We prove an observability estimate for a wave equation with rapidly oscillating density, in a bounded domain with Dirichlet boundary condition.

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### 0. INTRODUCTION AND RESULTS

Let  $\Omega$  be a smooth bounded domain in  $\mathbb{R}^d$ , and  $\rho(x, y)$  a smooth function on  $\mathbb{R}^d \times \mathbb{R}^d$ , such that

$$0 < \rho_{\min} \leq \rho(x, y) \leq \rho_{\max} \quad \forall (x, y) \quad (0.1)$$

$\rho$  is  $2\pi$ -periodic with respect to the second variable, *i.e.*

$$\rho(x, y) = \rho(x, y + 2\pi\ell) \quad \forall \ell \in \mathbb{Z}^d. \quad (0.2)$$

For  $\varepsilon > 0$ , let  $(\omega_n^\varepsilon, e_n^\varepsilon(x))$  be the spectrum of the Dirichlet problem for the operator  $-\rho^{-1}(x, x/\varepsilon)\Delta_g$  on  $L^2(\Omega; \rho(x, x/\varepsilon)d_g x)$  normalized in the form

$$\begin{cases} \rho(x, x/\varepsilon)(\omega_n^\varepsilon)^2 e_n^\varepsilon(x) = -\Delta_g e_n^\varepsilon(x) & \text{in } \Omega \\ e_n^\varepsilon(x) = 0 & \text{on } \partial\Omega \\ \int_{\Omega} e_n^\varepsilon(x) \overline{e_m^\varepsilon(x)} \rho(x, x/\varepsilon) d_g x = \delta_{n,m}; & 0 < \omega_1^\varepsilon \leq \omega_2^\varepsilon \leq \dots \end{cases} \quad (0.3)$$

Here,  $\Delta_g$  denotes the Laplace operator for some fixed smooth metric  $g$  on  $\overline{\Omega}$ , and  $d_g x$  is the volume form associated to  $g$ .

For any given  $\gamma_0 > 0$ , we shall denote by  $J_{\gamma_0}^\varepsilon$  the space of solutions  $u^\varepsilon(t, x)$  of the wave equation with oscillating density  $\rho$

$$\begin{cases} (\rho(x, x/\varepsilon)\partial_t^2 - \Delta_g) u^\varepsilon(t, x) = 0 & \text{in } \mathbb{R}_+ \times \Omega \\ u^\varepsilon(t, x)|_{x \in \partial\Omega} = 0 \end{cases} \quad (0.4)$$

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with maximum frequency less than  $\gamma_0/\varepsilon$ .

In other words,  $J_{\gamma_0}^\varepsilon$  is the set

$$J_{\gamma_0}^\varepsilon = \left\{ u^\varepsilon(t, x) = \sum_{\varepsilon\omega_n^\varepsilon \leq \gamma_0} \left( u_{+,n} e^{it\omega_n^\varepsilon} + u_{-,n} e^{-it\omega_n^\varepsilon} \right) e_n^\varepsilon(x) \right\}. \quad (0.5)$$

Let  $\{u_k^{\varepsilon_k}\}$  be a bounded sequence (in  $L_{\text{loc}}^2(\mathbb{R}_t, L^2(\Omega))$ ), of solutions of (0.4), with  $\lim \varepsilon_k = 0$ . It is well known that any weak limit of this sequence will satisfy the homogenized wave equation in  $\Omega$

$$\begin{cases} (\underline{\rho}(x)\partial_t^2 - \Delta_g)u(t, x) = 0 & \text{in } \mathbb{R} \times \Omega \\ u(t, x)|_{x \in \partial\Omega} = 0 \end{cases} \quad (0.6)$$

where  $\underline{\rho}(x) = \oint \rho(x, y) dy$  is the mean value of  $\rho$ .

Let  $\bar{V}$  be an open subset of  $\Omega$ , and  $T_0 > 0$ .

One says that waves solution of (0.6) are observable from  $V$  in time  $T_0$  if there exists a constant  $C_0$  s.t for any  $L^2$ -solution of (0.6) one has

$$\int_0^{T_0} \int_\Omega |u(t, x)|^2 \underline{\rho}(x) dt d_g x \leq C_0 \int_0^{T_0} \int_V |u(t, x)|^2 \underline{\rho}(x) dt d_g x. \quad (0.7)$$

If  $u = \sum_{\pm, n} u_{\pm, n} e^{\pm it\omega_n} e_n(x)$  is the Fourier series of  $u$  in the spectral decomposition of  $(-\underline{\rho})^{-1}(x)\Delta_g$ , we deduce from the elementary fact

$$\forall T > 0, \forall \omega_0 > 0, \exists C > 0 \text{ such that } \forall \omega \geq \omega_0, |c_+|^2 + |c_-|^2 \leq C \int_0^T |c_+ e^{it\omega} + c_- e^{-it\omega}|^2 dt$$

that the condition (0.7) is equivalent to the following

$$\begin{cases} \exists C_0 \text{ s.t. } \forall (u_{+,n}, u_{-,n})_n \in \ell^2 \times \ell^2 \\ \sum_n |u_{+,n}|^2 + |u_{-,n}|^2 \leq C_0 \int_0^{T_0} \int_V |u(t, x)|^2 \underline{\rho}(x) dt d_g x. \end{cases} \quad (0.8)$$

It is proved in [4] that (0.7) holds true under the geometric-control hypothesis

$$\begin{cases} 1) & \text{there is no infinite order of contact between the boundary} \\ & \partial\Omega \text{ and the bicharacteristics of } \underline{\rho}(x)\partial_t^2 - \Delta_g \\ 2) & \text{any generalized bicharacteristic of } \underline{\rho}(x)\partial_t^2 - \Delta_g \\ & \text{parameterized by } t \in ]0, T_0[ \text{ meets } \bar{V}. \end{cases} \quad (0.9)$$

Here the generalized bicharacteristic flow is the one defined by Melrose and Sjöstrand in [11].

The main result of this paper is the following theorem, which asserts that the estimate (0.7) remains true under the hypothesis (0.9) for  $\underline{\rho}(x)$ , for solutions of (0.4) in  $J_{\gamma_0}^\varepsilon$ , if  $\gamma_0$  is small enough.

**Theorem 0.1.** *Let the hypothesis (0.9) be satisfied. There exist small positive constants  $\gamma_0, \varepsilon_0$  and a constant  $C_0$ , such that for any  $\varepsilon \in ]0, \varepsilon_0[$  and any  $u^\varepsilon \in J_{\gamma_0}^\varepsilon$*

$$\int_0^{T_0} \int_\Omega |u^\varepsilon(t, x)|^2 \rho(x, x/\varepsilon) dt d_g x \leq C_0 \int_0^{T_0} \int_V |u^\varepsilon(t, x)|^2 \rho(x, x/\varepsilon) dt d_g x. \quad (0.10)$$

This is clearly a stability result of the observability estimate (0.7) under the singular perturbation  $\underline{\rho}(x) \rightarrow \rho(x, x/\varepsilon)$ . Let us recall that Theorem 0.1 has been proved in the 1-d case by Castro and Zuazua [6], and that in the 1-d case, the counter-example of Avellaneda *et al.* [1] shows that (0.10) fails for  $\gamma_0$  large. Indeed, in the 1-d case, when  $\rho = \rho(x/\varepsilon)$ , Castro [5] has shown that the greatest value of  $\gamma_0$  such that (0.10) holds true for some  $T_0$  (when  $V \Subset [a, b] = \Omega$ ) is related with the first instability interval of the Hill equation on the line  $\left(\frac{d}{dy}\right)^2 + \omega^2 \rho(y)$ . In the multi-d case, the understanding of the best value of  $\gamma_0$  such that (0.10) holds true will clearly involve the understanding of the localization and propagation of Bloch waves for the boundary value problem (0.4): this highly difficult problem is out of the scope of the present paper.

The conserved energy for solutions of (0.4) is

$$E(u^\varepsilon) = \frac{1}{2} \int_{\Omega} \{ |\partial_t u^\varepsilon|^2 \rho(x, x/\varepsilon) + |\nabla_g u^\varepsilon|^2 \} d_g x. \tag{0.11}$$

Applying the estimate (0.10) to  $\partial_t u^\varepsilon$ , one easily gets the energy observability estimate

**Corollary 0.1.** *Under the hypothesis and with the notations of Theorem 0.1, there exists a constant  $C_0$  s.t. for any  $\varepsilon \in ]0, \varepsilon_0[$  and any  $u^\varepsilon \in J_{\gamma_0}^\varepsilon$  one has*

$$E(u^\varepsilon) \leq C_0 \int_0^{T_0} \int_V |\partial_t u^\varepsilon|^2 \rho(x, x/\varepsilon) dt d_g x. \tag{0.12}$$

The paper is organized as follows:

1. reduction to a semi-classical estimate;
2. the Bloch wave;
3. Lopatinski estimate;
4. propagation estimate;
5. Appendix A: semi-classical o.p.d with operators values;
6. Appendix B: proofs of Lemmas 3.4–3.6.

**1.** In the first part, using a Littlewood-Paley decomposition, we reduce the proof of the inequality (0.10) to the assertion

$$\left\{ \begin{array}{l} \text{there exist } \gamma_0, \varepsilon_0, h_0, C_0 \text{ such that for any } \varepsilon \in ]0, \varepsilon_0[, \text{ and} \\ h \in [\varepsilon/\gamma_0, h_0] \text{ the inequality (0.10) holds true for any } u^\varepsilon \in I_h^\varepsilon, \\ \text{where } I_h^\varepsilon = \left\{ u^\varepsilon = \sum_{0.9 \leq \omega_n^\varepsilon h \leq 2.1} (u_{+,n} e^{it\omega_n^\varepsilon} + u_{-,n} e^{-it\omega_n^\varepsilon}) e_n^\varepsilon(x) \right\}. \end{array} \right. \tag{0.13}$$

**2.** In the second part, we introduce the Bloch wave at the boundary  $\Gamma(u^\varepsilon)$ . We refer to [2] and [7] for the study of Bloch waves in equations with oscillating coefficients. We choose a coordinate system

$$\left\{ \begin{array}{l} \partial\Omega \times [0, r_0] \xrightarrow{\Theta} \mathbb{R}^d \\ (x', x_d) \mapsto \Theta(x', x_d) \end{array} \right. \tag{0.14}$$

which satisfies

$$\left\{ \begin{array}{l} i) \Theta(\partial\Omega \times [0, r_0]) \subset \overline{\Omega} \\ ii) \text{ for } x_d \text{ small, } x_d \mapsto \Theta(x', x_d) \text{ is the geodesic normal to the} \\ \text{boundary at } x' \in \partial\Omega, \text{ for the metric } g \text{ on } \overline{\Omega}. \end{array} \right. \tag{0.15}$$

In these coordinates, the Laplace operator takes the form

$$\left\{ \begin{array}{l} \Delta_g = \frac{\partial}{\partial x_d} \left( A_0(x) \frac{\partial}{\partial x_d} + A_1(x, \partial_{x'}) \right) + A_2(x, \partial_{x'}); \\ x = (x', x_d), x' \in \partial\Omega \end{array} \right. \quad (0.16)$$

where  $A_j(x, \partial_{x'})$  are differential operators of order  $j$  on  $\partial\Omega$ , with  $x_d$  as parameter. Let  $a_j(x, \xi')$  be the principal symbol of  $A_j$ . The dual metric  $g^{-1}(x, \xi) \stackrel{\text{def}}{=} \|\xi\|_x^2$  on the cotangent bundle  $T^*\Omega$  is

$$\|\xi\|_x^2 = a_0(x)\xi_d^2 + a_1(x, \xi')\xi_d + a_2(x, \xi'). \quad (0.17)$$

Let  $\mathbb{T}^d = (\mathbb{R}/2\pi\mathbb{Z})^d$  be the  $d$ -dimensional torus and for  $\varepsilon > 0$ ,  $S_\varepsilon \subset \partial\Omega \times [0, r_0] \times \mathbb{T}_y^d$  the submanifold

$$S_\varepsilon = \{(x, y); y = \Theta(x)/\varepsilon \bmod (2\pi\mathbb{Z})^d\}. \quad (0.18)$$

Let  $f(x)$  be a function on  $\partial\Omega \times [0, r_0]$ . We define a distribution  $T(f)$  on  $\partial\Omega \times [0, r_0] \times \mathbb{T}_y^d$  by the formula

$$T(f) = \sum_{\ell \in \mathbb{Z}^d} e^{i\ell(y - \Theta(x)/\varepsilon)} f(x) = (2\pi)^d \delta_{y = \Theta(x)/\varepsilon} \otimes f(x). \quad (0.19)$$

If  $X$  is a vector field on  $\partial\Omega \times [0, r_0]$ , we shall denote by  $X_\varepsilon^*$  the lift of  $X$  on  $S_\varepsilon$ . If  $x' = (x_1, \dots, x_{d-1})$  is a local coordinate system on  $\partial\Omega$ , and  $(\Theta_1(x), \dots, \Theta_d(x)) = \Theta(x)$  are the Cartesian coordinates of  $\Theta$ , one has

$$\left( \frac{\partial}{\partial x_k} \right)_\varepsilon^* = \frac{\partial}{\partial x_k} + \frac{1}{\varepsilon} \sum_{j=1}^d \frac{\partial \Theta_j}{\partial x_k}(x) \frac{\partial}{\partial y_j} \quad \text{for } 1 \leq k \leq d \quad (0.20)$$

and

$$\left( \frac{\partial}{\partial x_k} \right)_\varepsilon^* T(f) = T \left( \frac{\partial}{\partial x_k} f \right) \quad \text{for } 1 \leq k \leq d. \quad (0.21)$$

The Bloch operator on  $\partial\Omega \times [0, r_0] \times \mathbb{T}^d$  is defined by

$$\left\{ \begin{array}{l} \mathbb{B}_\varepsilon(x, \varepsilon \partial_x, \varepsilon \partial_t; y, \partial_y) = \hat{\rho}(x, y)(\varepsilon \partial_t)^2 - \varepsilon^2 (\Delta_g)_\varepsilon^*; \quad \hat{\rho}(x, y) = \rho(\Theta(x), y) \\ (\Delta_g)_\varepsilon^* = \left( \frac{\partial}{\partial x_d} \right)_\varepsilon^* \left( A_0(x) \left( \frac{\partial}{\partial x_d} \right)_\varepsilon^* + A_1(x, (\partial_{x'})_\varepsilon^*) \right) + A_2(x, (\partial_{x'})_\varepsilon^*). \end{array} \right. \quad (0.22)$$

It satisfies the identity

$$\mathbb{B}_\varepsilon(T(u(x, t))) = \varepsilon^2 T((\rho(\Theta(x), \Theta(x)/\varepsilon) \partial_t^2 - \Delta_g)(u(x, t))). \quad (0.23)$$

Let  $\tilde{A}_j$  be the operators

$$\tilde{A}_j = A_j(x, (\partial_{x'})_\varepsilon^*) \quad (0.24)$$

and let  $e_k(x)$   $1 \leq k \leq d$  be the vectors of  $\mathbb{R}^d$

$$e_k(x) = \frac{\partial \Theta}{\partial x_k}(x). \quad (0.25)$$

If  $v(t, x, y)$  is a distribution on  $\overset{\circ}{X} \times \mathbb{T}^d$ , with  $X = \mathbb{R}_t \times (\partial\Omega \times [0, r_0])$ , we shall write the equation  $\mathbb{B}_\varepsilon(v) = 0$  as a  $2 \times 2$  system for the vector  $w = \mathcal{A}(v)$ .

$$\mathcal{A}(v) = \begin{bmatrix} w_0 \\ w_1 \end{bmatrix} \stackrel{\text{def}}{=} \begin{bmatrix} v \\ (A_0(x)(\varepsilon \frac{\partial}{\partial x_d})^* + \varepsilon \tilde{A}_1)v \end{bmatrix}. \quad (0.26)$$

This system takes the form

$$\begin{cases} \varepsilon \frac{\partial}{\partial x_d} w + \mathbb{M}w = 0 \\ \mathbb{M} = \begin{bmatrix} e_d(x) \cdot \partial_y + \varepsilon A_0^{-1}(x) \tilde{A}_1 & -A_0^{-1}(x) \\ \varepsilon^2 \tilde{A}_2 - \hat{\rho}(x, y)(\varepsilon \partial_t)^2 & e_d(x) \cdot \partial_y \end{bmatrix} \end{cases}. \quad (0.27)$$

The operator  $\mathbb{M}$  will be seen as a semi-classical operator in  $t, x, \frac{\varepsilon}{i} \partial_{x'} = \xi', \frac{\varepsilon}{i} \partial_t = \tau$  with operator values in the fiber  $\mathbb{T}^d$

$$\mathbb{M} = \sum_{j=0}^2 \left( \frac{\varepsilon}{i} \right)^j \mathbb{M}^j(x, \xi', \tau; y, \partial_y). \quad (0.28)$$

The differential degree in  $y$  of  $\mathbb{M}^j$  is at most  $2 - j$  and the principal symbol  $\mathbb{M}^0$  is the matrix

$$\mathbb{M}^0(x, \xi', \tau; y, \partial_y) = \begin{bmatrix} e_d(x) \cdot \partial_y + a_0^{-1}(x) a_1(x, i\xi' + e'(x) \cdot \partial_y) & -a_0^{-1}(x) \\ a_2(x, i\xi' + e'(x) \cdot \partial_y) + \hat{\rho}(x, y) \tau^2 & e_d(x) \cdot \partial_y \end{bmatrix}. \quad (0.29)$$

Let  $E^\bullet = \{E^s, s \in \mathbb{R}\}$  be the scale of Hilbert spaces on the torus

$$E^s = H^s(\mathbb{T}^d) \oplus H^{s-1}(\mathbb{T}^d). \quad (0.30)$$

For any  $\rho = (x, \xi', \tau)$ ,  $\mathbb{M}^j(\rho, y, \partial_y)$  maps  $E^s$  into  $E^{s-1+j}$  and  $\mathbb{M}^0$  is an elliptic operator. Let  $\mathbb{M}_0^0$  be the restriction of  $\mathbb{M}^0$  to the zero section  $\xi' = \tau = 0$ .

$$\mathbb{M}_0^0(x, \partial_y) = \mathbb{M}^0(x, 0, 0, y, \partial_y) = \begin{bmatrix} e_d(x) \cdot \partial_y + a_0^{-1} a_1(x, e'(x) \cdot \partial_y) & -a_0^{-1}(x) \\ a_2(x, e'(x) \cdot \partial_y) & e_d(x) \cdot \partial_y \end{bmatrix}. \quad (0.31)$$

The eigenvalues  $\lambda_{\pm, \ell}^0$  of  $\frac{1}{i} \mathbb{M}_0^0(x, \partial_y)$  on the space  $e^{i\ell y} \mathbb{C}^2$ , for  $\ell \in \mathbb{Z}^d$  are the complex roots of the equation

$$a_0(x)(-\lambda + e_d \cdot \ell)^2 + (-\lambda + e_d \cdot \ell) a_1(x, e' \cdot \ell) + a_2(x, e' \cdot \ell) = 0 \quad (0.32)$$

which is equivalent to

$$\|{}^t d\Theta(x)(\ell) - \lambda(0, \dots, 0, 1)\|_x^2 = 0. \quad (0.33)$$

In particular we have

$$\inf_x \min_{\ell \neq 0} |\lambda_{\pm, \ell}^0(x)| > 0 \quad (0.34)$$

so the double eigenvalue  $\lambda_{\pm, 0}^0(x) = 0$  is isolated in the spectrum of  $\mathbb{M}_0^0(x, \partial_y)$ .

In the sequel, we shall restrict the values of the Sobolev index of regularity  $s$  on the torus to some fixed large interval,  $s \in [-\sigma_0, \sigma_0]$ ,  $\sigma_0 \gg \frac{d}{2}$ .

Let  $X = \partial\Omega \times \mathbb{R}_t \times [0, r_0]$ . We denote by  ${}^tT^*X$  the tangential cotangent bundle

$${}^tT^*X = T^*(\partial\Omega \times \mathbb{R}_t) \times [0, r_0]. \quad (0.35)$$

Let  $W_1 \Subset W_0$  be two small neighborhoods of the set  $\{\xi' = \tau = 0\} \times \{t \in [-T_0, 2T_0]\}$  in  ${}^tT^*X$ .

We choose a non-negative function  $\chi_0 \in C_0^\infty(W_0)$ , such that  $\chi_0 \equiv 1$  on  $W_1$ .

If  $W_0$  is small enough, we define the map  $p_0(x, t, \xi', \tau) : E^\bullet \rightarrow \mathbb{C}^2$  by the formula

$$p_0[w] = \chi_0 \cdot \oint_{\mathbb{T}^d} \left\{ \frac{1}{2i\pi} \int_{\partial D} \frac{dz}{z - \mathbb{M}^0} \right\} [w] \quad w \in E^s, \quad s \in [-\sigma_0, \sigma_0] \quad (0.36)$$

(where  $D \subset \mathbb{C}$  is a small disk with center  $z = 0$ ).

It satisfies the estimates

$$\exists C \forall s \in [-\sigma_0, \sigma_0] \quad \forall w \in E^s \quad \|p_0(w) - \chi_0 \oint_{\mathbb{T}^d} w\|_{\mathbb{C}^2} \leq C\tau^2 \|w\|_{E^s} \quad (0.37)$$

and there exists  $L^0(x, t, \xi', \tau) \in C^\infty({}^tT^*X; M_2(\mathbb{C}))$ , defined near  $\xi' = \tau = 0$  such that (see (2.29–2.31))

$$p_0 \circ \mathbb{M}^0 = L^0 \circ p_0. \quad (0.38)$$

By a Taylor expansion near  $\xi' = \tau = 0$ , one gets

$$L^0 = \begin{bmatrix} a_0^{-1}(x)a_1(x, i\xi') & -a_0^{-1}(x) \\ a_2(x, i\xi') + \hat{\rho}(x)\tau^2 & 0 \end{bmatrix} + O(\tau^4). \quad (0.39)$$

We then suitably quantize the above construction and we obtain tangential pseudo differential operators (see Append. A1)

$$\begin{cases} \Pi_0(\varepsilon, t, x, \varepsilon\partial_t, \varepsilon\partial_{x'}) & : L^2(X; E^s) \rightarrow L^2(X, \mathbb{C}^2), \quad s \in [-\sigma_0, \sigma_0] \\ L(\varepsilon, t, x, \varepsilon\partial_t, \varepsilon\partial_{x'}) & : L^2(X; \mathbb{C}^2) \rightarrow L^2(X, \mathbb{C}^2) \end{cases} \quad (0.40)$$

with principal symbol  $\sigma(\Pi_0) = p_0$ ,  $\sigma(L) = L^0$ , which satisfy the relation

$$\Pi_0(\varepsilon\partial_{x_d} + \mathbb{M}) = (\varepsilon\partial_{x_d} + L)\Pi_0 + R(\varepsilon, t, x, \varepsilon\partial_t, \varepsilon\partial_{x'}). \quad (0.41)$$

In (0.41), the error term  $R : L^2(X; E^s) \rightarrow L^2(X, \mathbb{C}^2)$  will be a tangential pseudo differential operator such that for any tangential o.p.d.  $Q$  with essential support in  $W_1$  and any  $s \in [-\sigma_0, \sigma_0]$ , one has

$$\|Q \circ R; L^2(X; E^s) \rightarrow L^2(X, \mathbb{C}^2)\| \in \mathcal{O}(\varepsilon^\infty). \quad (0.42)$$

**Definition 0.1.** For  $u^\varepsilon \in I_h^\varepsilon$ , we define the Bloch wave  $\Gamma(u^\varepsilon) \in L^2(X; \mathbb{C}^2)$  by the formula

$$\Gamma(u^\varepsilon) = \begin{bmatrix} \Gamma_0(u^\varepsilon) \\ \Gamma_1(u^\varepsilon) \end{bmatrix} = \Pi_0 \mathcal{T}(u^\varepsilon) \quad (\mathcal{T} = \mathcal{A} \circ T). \quad (0.43)$$

Let  $\gamma_0, \varepsilon_0, h_0$  be given small enough,  $\varepsilon \in ]0, \varepsilon_0]$ ,  $h \in [\varepsilon/\gamma_0, h_0]$ . For  $u_\varepsilon \in I_h^\varepsilon$ ,  $u^\varepsilon = \sum_{0.9 \leq \omega_n^\varepsilon h \leq 2.1} (u_{+,n} e^{it\omega_n^\varepsilon}$

$+ u_{-,n} e^{-it\omega_n^\varepsilon}) e_n^\varepsilon(x)$ , we define  $\|u^\varepsilon\|^2 \left( \simeq \int_0^{T_0} \int_\Omega |u^\varepsilon|^2 \right)$  by

$$\|u^\varepsilon\|^2 = \sum_{0.9 \leq \omega_n^\varepsilon h \leq 2.1} |u_{+,n}|^2 + |u_{-,n}|^2. \quad (0.44)$$

Let  $X_{T_0} = \partial\Omega \times [-T_0, 2T_0] \times [0, r_0]$ , and let  $K$  be the compact subset of  ${}^tT^*X$ ,  $K = \partial\Omega \times [0, T_0] \times [0, r_0/2] \times \{\xi' = 0, \tau = 0\}$ . The following proposition will be proven in Section 2.

**Proposition 0.1.** *Let  $Q(\varepsilon, t, x, \varepsilon\partial_{x'}, \varepsilon\partial_t)$  be a zero order tangential opd on  $X$ , equal to  $Id$  near  $K$ . If  $\gamma_0, \varepsilon_0, h_0$  are small enough, there exists a constant  $C > 0$ , such that for any  $\varepsilon \in ]0, \varepsilon_0]$ ,  $h \in [\varepsilon/\gamma_0, h_0]$ , one has*

$$\|u^\varepsilon\|^2 \leq C \left[ \|Q\Gamma_0(u^\varepsilon)\|_{L^2(X_{T_0})}^2 + \|u^\varepsilon\|_{L^2((0, T_0) \times V)}^2 \right] \quad \forall u^\varepsilon \in I_h^\varepsilon. \tag{0.45}$$

**3.** By Proposition 0.1, we shall obtain the inequality (0.10), if we are able to estimate the  $L^2$  norm of the first component  $\Gamma_0(u^\varepsilon)$  of the Bloch wave near the set  $K$ .

The formula (0.41) shows that  $\Gamma(u^\varepsilon)$  satisfies the equation

$$(\varepsilon\partial_{x_d} + L)\Gamma(u^\varepsilon) \in O(\varepsilon^\infty L^2) \text{ (microlocally in } W_1). \tag{0.46}$$

By (0.39) this equation is very closed to the homogenized equation  $(\underline{\rho}(x)\partial_t^2 - \Delta_g)[\Gamma_0(u^\varepsilon)] = 0$ .

As one can see, all the difficulty in our problem is thus to obtain an estimate on the first Dirichlet data of  $\Gamma(u^\varepsilon)$  on the boundary  $x_d = 0$ , in order to apply propagation arguments to the equation (0.46). We shall prove the following proposition.

**Proposition 0.2.** *If  $\gamma_0, \varepsilon_0, h_0$  are small enough, there exists a constant  $C$  such that for any  $\varepsilon \in ]0, \varepsilon_0]$ ,  $h \in [\varepsilon/\gamma_0, h_0]$  the following estimate holds true*

$$\|\Gamma_0(u^\varepsilon)|_{x_d=0}\|_{L^2(X_{T_0} \cap x_d=0)} \leq C \varepsilon/h \|u^\varepsilon\| \quad \forall u^\varepsilon \in I_h^\varepsilon. \tag{0.47}$$

The above estimate will be obtained as a consequence of a uniform Lopatinski estimate on  $w^\varepsilon = \mathcal{T}(u^\varepsilon) = \begin{bmatrix} w_0^\varepsilon \\ w_1^\varepsilon \end{bmatrix}$ .

We shall prove

**Theorem 0.2.** *Let  $Q$  be a scalar tangential o.p.d. with essential support in  $W_0$ ; if  $W_0, \gamma_0, \varepsilon_0, h_0$  are small enough, there exist  $s_1 < 0$  and a constant  $C$  such that for any  $u^\varepsilon \in I_h^\varepsilon$  the following estimate holds true*

$$\|Q(t, x, \varepsilon\partial_{x'}, \varepsilon\partial_t)(w_1^\varepsilon)|_{x_d=0}\|_{L^2(X_{T_0} \cap x_d=0, H^{s_1}(\mathbb{T}^d))} \leq C \|u^\varepsilon\|. \tag{0.48}$$

Notice that  $w^\varepsilon$  satisfies the equation (0.27), with Dirichlet data  $w_0^\varepsilon|_{x_d=0} = 0$  on the boundary.

The weaker estimate

$$\|Q(w_1^\varepsilon)|_{x_d=0}\| \leq C \varepsilon^{-1/2} \|u^\varepsilon\| \tag{0.49}$$

is easy to obtain (it is sufficient to commute the Eq. (0.4) with the normal vector field  $\frac{\partial}{\partial n}$ ).

The proof of (0.48) is the most technical part of our work. It involves a detailed study of how the spectral theory of  $M^0(x, \xi', \tau; y, \partial y)$  (see (0.29)) depends on the parameter  $(x, \xi', \tau)$ .

**4.** This part will be devoted to the proof of the following proposition.

**Proposition 0.3.** *Let  $Q(\varepsilon, t, x, \varepsilon\partial_{x'}, \varepsilon\partial_t)$  be a zero order opd equal to  $Id$  near  $K$ , with essential support in  $W_1$ . There exist  $\gamma_0, \varepsilon_0, h_0$ , and a constant  $C_0$  such that, for any  $\varepsilon \in ]0, \varepsilon_0]$ ,  $h \in [\varepsilon/\gamma_0, h_0]$  and  $u^\varepsilon \in I_h^\varepsilon$ , the following estimate holds true*

$$\|Q\Gamma_0(u^\varepsilon)\|_{L^2(X_{T_0})}^2 \leq C_0 \left[ \|\Gamma_0(u^\varepsilon)|_{x_d=0}\|_{L^2(X_{T_0} \cap x_d=0)}^2 + \|u^\varepsilon\|_{L^2(0,T_0) \times V}^2 \right]. \tag{0.50}$$

This estimate will be obtained by rather classical arguments in the theory of control of linear waves, for the rescale equation

$$\begin{cases} \left( h \frac{\partial}{\partial x_d} + \mathcal{L} \right) \begin{bmatrix} g_0 \\ g_1 \end{bmatrix} \sim 0 & \begin{bmatrix} g_0 \\ g_1 \end{bmatrix} = \begin{bmatrix} \Gamma_0(u^\varepsilon) \\ \frac{h}{\varepsilon} \Gamma_1(u^\varepsilon) \end{bmatrix} \\ \mathcal{L} = \frac{h}{\varepsilon} \begin{pmatrix} 1 & 0 \\ 0 & h/\varepsilon \end{pmatrix} \circ L \circ \begin{pmatrix} 1 & 0 \\ 0 & \varepsilon/h \end{pmatrix}. \end{cases} \tag{0.51}$$

We shall verify that  $\mathcal{L}$  is still a  $h$ -pseudo differential operator, with  $\varepsilon/h$  as parameter. (We use this rescaling in order to be able to use propagation arguments in the range  $\varepsilon \ll h$ .)

**5.** In Appendix A.1, we recall the properties of the semi-classical calculus with operators values which is used in **2**. In Appendix A.2, we extend this calculus to a larger class of symbols; this exotic calculus will be used in **3**.

To end this introduction, we finally remark that the validity of (0.13), hence the proof of Theorem 0.1, is a direct consequence of the Propositions 0.1, 0.2 and 0.3.

### 1. SEMI-CLASSICAL REDUCTION

In this part, we verify that (0.13) implies the Theorem 1.

Let  $e_n^\varepsilon(x)$  be a normalized eigenfunction of the Dirichlet problem (0.3), and let  $\mu_1$  be the first eigenvalue of the Dirichlet problem for  $\Delta_g$  in  $\Omega$ . One has

$$\int_{\Omega} |\nabla_g e_n^\varepsilon|^2 d_g x = \int_{\Omega} \rho(x, x/\varepsilon) (\omega_n^\varepsilon)^2 |e_n^\varepsilon(x)|^2 d_g x \leq \rho_{\max} (\omega_n^\varepsilon)^2 \int_{\Omega} |e_n^\varepsilon(x)|^2 d_g x. \tag{1.1}$$

So we get the uniform lower bound

$$\omega_n^\varepsilon \geq (\rho_{\max})^{-1/2} \mu_1^{1/2}. \tag{1.2}$$

The Sobolev spaces  $L^2(\Omega), H_0^1(\Omega), H^{-1}(\Omega)$ , with norms  $(\int_{\Omega} |f|^2 \rho d_g x)^{1/2}, (\int_{\Omega} |\nabla_g f|^2 d_g x)^{1/2}, \sup\{\int_{\Omega} f \bar{h} \rho d_g x, \|h\|_{H_0^1} \leq 1\}$  are characterized in terms of Fourier series by

$$\begin{cases} f_n^\varepsilon = \int_{\Omega} f \overline{e_n^\varepsilon(x)} \rho d_g x \text{ for } f \in H^{-1}(\Omega) \\ \|f\|_{L^2}^2 = \sum_n |f_n^\varepsilon|^2; \|f\|_{H_0^1}^2 = \sum_n (\omega_n^\varepsilon)^2 |f_n^\varepsilon|^2; \|f\|_{H^{-1}}^2 = \sum_n (\omega_n^\varepsilon)^{-2} |f_n^\varepsilon|^2. \end{cases} \tag{1.3}$$

Any solution  $u^\varepsilon$  of the wave equation (0.4) with data  $(u^\varepsilon(0), \partial_t u^\varepsilon(0)) \in L^2(\Omega) \oplus H^{-1}(\Omega)$  is of the form

$$u^\varepsilon = \sum_n u_n^\varepsilon(t) e_n^\varepsilon(x) = \sum_n (u_{+,n}^\varepsilon e^{it\omega_n^\varepsilon} + u_{-,n}^\varepsilon e^{-it\omega_n^\varepsilon}) e_n^\varepsilon(x) \tag{1.4}$$

with  $(u_{\pm,n}^\varepsilon)_n \in \ell^2$ , and (1.2) implies that there exists a constant  $C$  independent of  $\varepsilon$ , s.t.

$$\frac{1}{C} \sum_{n,\pm} |u_{\pm,n}^\varepsilon|^2 \leq \int_0^{T_0} \int_{\Omega} |u^\varepsilon|^2 \rho dt d_g x \leq C \sum_{n,\pm} |u_{\pm,n}^\varepsilon|^2. \tag{1.5}$$



If the geometric hypothesis (0.9) holds true for  $T_0$ , it remains valid for  $T_0 - 2\delta$ , for  $\delta > 0$  small enough; we can therefore assume that (0.13) is valid on  $[\delta, T_0 - \delta]$ .

Take  $\varphi(t) \in C_0^\infty(]0, T_0[)$ ,  $\varphi(t) \equiv 1$  on  $[\delta, T_0 - \delta]$  and  $\psi(\sigma) \in C_0^\infty(]0.9, 2.1[)$ ,  $\psi(\sigma) \equiv 1$  on  $[1, 2]$ . Let  $\chi(\sigma) = \psi(\sigma) + \psi(-\sigma)$ . For  $u^\varepsilon \in J_{\gamma_0}^\varepsilon$ , one has  $\chi(2^{-k}D_t)u^\varepsilon \in I_{2^{-k}}^\varepsilon$ , so there exists  $C_0$  s.t.

$$\begin{cases} \forall \varepsilon \in ]0, \varepsilon_0], \forall k \in \mathbb{N} \text{ s.t. } 2^{-k} \in [\varepsilon/\gamma_0, h_0] \\ \forall u^\varepsilon = \sum_{\varepsilon\omega_n^\varepsilon \leq \gamma_0} (u_{+,n}^\varepsilon e^{it\omega_n} + u_{-,n}^\varepsilon e^{-it\omega_n}) e_n^\varepsilon(x) \in J_{\gamma_0}^\varepsilon \\ \sum_{2^k \leq \omega_n^\varepsilon \leq 2^{k+1}} |u_{+,n}^\varepsilon|^2 + |u_{-,n}^\varepsilon|^2 \leq C_0 \int_{-\infty}^{+\infty} dt \int_V d_g x |\varphi(t)\chi(2^{-k}D_t)u^\varepsilon|^2. \end{cases} \quad (1.6)$$

On the other hand, using classical estimates as in ([9], Sect. 4), one gets  $\exists C_1, C_2, k_0$  s.t. for any  $k_1 \geq k_0$ , and any  $u^\varepsilon \in J_{\gamma_0}^\varepsilon$

$$\sum_{k \geq k_1} \int_{-\infty}^{+\infty} dt \int_V |\varphi(t)\chi(2^{-k}D_t)u^\varepsilon|^2 d_g x \leq C_1 \int_0^{T_0} \int_V |u^\varepsilon|^2 d_g x + C_2 2^{-2k_1} \left( \sum_n |u_{\pm,n}^\varepsilon|^2 \right). \quad (1.7)$$

Let  $\gamma_1 = \gamma_0/2$ ; for  $u^\varepsilon \in J_{\gamma_1}^\varepsilon$  and  $2^{-k} < \varepsilon/\gamma_0$  one has  $\chi(2^{-k}D_t)u^\varepsilon \equiv 0$ , so putting together (1.6) and (1.7) we get

$$\begin{cases} \exists n_0, \exists C_3, \forall \varepsilon \in ]0, \varepsilon_0], \forall u^\varepsilon \in J_{\gamma_1}^\varepsilon \\ \sum_{n \geq n_0, \varepsilon\omega_n^\varepsilon \leq \gamma_1} |u_{+,n}^\varepsilon|^2 + |u_{-,n}^\varepsilon|^2 \leq C_3 \left( \int_0^{T_0} dt \int_V d_g x |u^\varepsilon|^2 + \sum_{n \leq n_0} |u_{\pm,n}^\varepsilon|^2 \right) \end{cases} \quad (1.8)$$

and (1.8) is equivalent to

$$\begin{cases} \exists n_0, \exists C_4, C_5, \forall \varepsilon \in ]0, \varepsilon_0], \forall u^\varepsilon \in J_{\gamma_1}^\varepsilon \\ \int_0^T \int_\Omega |u^\varepsilon|^2 \rho dt d_g x \leq C_3 \int_0^{T_0} \int_V |u^\varepsilon|^2 \rho dt d_g x \\ + C_4 \left( \sum_{n \leq n_0} |u_{+,n}^\varepsilon|^2 + |u_{-,n}^\varepsilon|^2 \right). \end{cases} \quad (1.9)$$

It is now easy to conclude the proof of Theorem 1 by a uniqueness argument. In fact if (0.10) is untrue, there exist a sequence  $\varepsilon_k \rightarrow 0$  and  $u^{\varepsilon_k} \in J_{\gamma_1}^{\varepsilon_k}$  such that  $\int_0^{T_0} \int_\Omega |u^{\varepsilon_k}|^2 \rho dt d_g x = 1$  and  $\int_0^{T_0} \int_V |u^{\varepsilon_k}|^2 \rho dt d_g x \rightarrow 0$ ; let  $u$  be a weak limit in  $L^2$  of  $\{u^{\varepsilon_k}\}$ ;  $u$  satisfies

$$\begin{cases} \rho(x)\partial_t^2 u - \Delta_g u = 0 \text{ on } \mathbb{R}_t \times \Omega \\ u|_{\partial\Omega} = 0; \quad u|_{]0, T_0[ \times V} = 0 \end{cases} \quad (1.10)$$

and from the observability inequality (0.7), we get  $u \equiv 0$ . Then (1.9) implies that  $u \equiv 0$  is the strong limit in  $L^2$  of  $u^{\varepsilon_k}$ , which contradicts  $\int_0^{T_0} \int_\Omega |u^{\varepsilon_k}|^2 \rho dt d_g x \equiv 1$ .

## 2. THE BLOCH WAVE

We shall now recall how one can quantize the principal symbols maps  $p_0, L^0$  defined in (0.36, 0.38) in order to obtain the pseudo differential relation (0.41).

Let

$$I = [-\sigma_0, \sigma_0]. \quad (2.1)$$

For any  $s \in I$ , we split  $E^s = H^s(\mathbb{T}^d) \oplus H^{s-1}(\mathbb{T}^d)$  into the decomposition

$$\begin{cases} E^s &= E_0 \oplus E_\perp^s & E_0 = \mathbb{C}^2 \\ w &= w_{(0)} + w_\perp & w_{(0)} = \oint_{\mathbb{T}^d} w. \end{cases} \quad (2.2)$$

In other words we write  $w = \sum_{\ell} w_{(\ell)} e^{i\ell y}$  and  $w_\perp = \sum_{\ell \neq 0} w_{(\ell)} e^{i\ell y}$ .

We then construct tangential pseudo differential operators defined near  $\varepsilon \partial_t = i\tau = 0$ ,  $\varepsilon \partial_{x'} = i\xi' = 0$ , semi-classical in  $\varepsilon$

$$A_0(\varepsilon, x, \varepsilon \partial_t, \varepsilon \partial_{x'}) : L^2(X, E_0) \rightarrow L^2(X, \bigcap_{s \in I} E^s) \quad (2.3)$$

$$A_\perp(\varepsilon, x, \varepsilon \partial_t, \varepsilon \partial_{x'}) : L^2(X, E_\perp^s) \rightarrow L^2(X, E^s) \quad (\forall s \in I) \quad (2.4)$$

$$L(\varepsilon, x, \varepsilon \partial_t, \varepsilon \partial_{x'}) : L^2(X, E_0) \rightarrow L^2(X, E_0) \quad (2.5)$$

$$L_\perp(\varepsilon, x, \varepsilon \partial_t, \varepsilon \partial_{x'}) : L^2(X, E_\perp^s) \rightarrow L^2(X, E_\perp^{s-1}) \quad (\forall s \in I) \quad (2.6)$$

with symbols admitting asymptotic expansions

$$\sum_k \left(\frac{\varepsilon}{i}\right)^k A_0^k(x, \tau, \xi') \quad A_0^k \text{ bounded from } E_0 \text{ to } \bigcap_{s \in I} E^s \quad (2.7)$$

$$\sum_k \left(\frac{\varepsilon}{i}\right)^k A_\perp^k(x, \tau, \xi') \quad A_\perp^k \text{ bounded from } E_\perp^s \text{ to } E^s \quad (\forall s \in I) \quad (2.8)$$

$$\sum_k \left(\frac{\varepsilon}{i}\right)^k L^k(x, \tau, \xi') \quad L^k \text{ bounded from } E_0 \text{ to } E_0 \quad (2.9)$$

$$\sum_k \left(\frac{\varepsilon}{i}\right)^k L_\perp^k(x, \tau, \xi') \quad L_\perp^k \text{ bounded from } E_\perp^s \text{ to } E_\perp^{s-1} \quad (\forall s \in I) \quad (2.10)$$

such that near the zero section  $\tau = \xi' = 0$ , the two following identities hold true, in the algebra of tangential pseudo differential operators

$$\begin{cases} \left(\varepsilon \frac{\partial}{\partial x_d} + \mathbb{M}\right) A_0 = A_0 (\varepsilon \partial_{x_d} + L) \\ \left(\varepsilon \frac{\partial}{\partial x_d} + \mathbb{M}\right) A_\perp = A_\perp (\varepsilon \partial_{x_d} + L_\perp). \end{cases} \quad (2.11)$$

Using the formula (0.28)  $\mathbb{M} = \sum_{j=0}^2 \left(\frac{\varepsilon}{i}\right)^j \mathbb{M}^j(x, \xi', \tau; y, \partial y)$ , and the rules of composition of pseudo differential operators, one gets that (2.11) is equivalent to the following set of equations (2.12, 2.13)

$$k=0 \quad \begin{cases} \mathbb{M}^0 A_0^0 = A_0^0 L^0 \\ \mathbb{M}^0 A_\perp^0 = A_\perp^0 L_\perp^0 \end{cases} \quad (2.12)$$

$$\begin{cases} \sum_{j+\ell+|\alpha|=k} \frac{1}{\alpha!} \partial_{\xi'}^{\alpha} \mathbb{M}^j \partial_{x'}^{\alpha} A_0^{\ell} + i \partial_{x_d} A_0^{k-1} & = \sum_{j+\ell+|\alpha|=k} \frac{1}{\alpha!} \partial_{\xi'}^{\alpha} A_0^j \partial_{x'}^{\alpha} L^{\ell} \\ \sum_{j+\ell+|\alpha|=k} \frac{1}{\alpha!} \partial_{\xi'}^{\alpha} \mathbb{M}^j \partial_{x'}^{\alpha} A_{\perp}^{\ell} + i \partial_{x_d} A_{\perp}^{k-1} & = \sum_{j+\ell+|\alpha|=k} \frac{1}{\alpha!} \partial_{\xi'}^{\alpha} A_{\perp}^j \partial_{x'}^{\alpha} L_{\perp}^{\ell}. \end{cases} \quad (2.13)$$

Let  $j_0$  and  $j_{\perp}$  be the inclusion maps

$$E_0 \xrightarrow{j_0} E^s \quad E_{\perp}^s \xrightarrow{j_{\perp}} E^s \quad (2.14)$$

and let  $\pi_0 = \pi_0(x, \xi', \tau)$  be the spectral projector of  $\mathbb{M}^0$ , which is defined near  $(\xi', \tau) = (0, 0)$ , by

$$\pi_0 = \frac{1}{2i\pi} \int_{\partial D} \frac{dz}{z - \mathbb{M}^0} \quad (2.15)$$

where  $D \subset \mathbb{C}$  is a small disk with center at  $z = 0$ .

The range of  $\pi_0$  is a two-dimensional invariant subspace of  $\mathbb{M}^0$ , and by the definition formula (0.29) of  $\mathbb{M}^0$ , one gets for  $|\tau|$  small enough

$$\left\| \oint_{\mathbb{T}^d} \pi_0 j_0 - Id_{E_0} \right\| \leq \text{Cte } \tau^2; \quad \|\pi_0 j_{\perp}; E_{\perp}^s \rightarrow E^s\| \leq \text{Cte } \tau^2. \quad (2.16)$$

In order to obtain the relations (2.12), it is clearly sufficient to select isomorphisms

$$\begin{cases} A_0^0 : E_0 \xrightarrow{\sim} \text{range}(\pi_0) \\ A_{\perp}^0 : E_{\perp}^s \xrightarrow{\sim} \text{range}(Id - \pi_0). \end{cases} \quad (2.17)$$

We can choose in view of (2.16), for  $|\tau|$  small enough

$$\begin{cases} A_{\perp}^0 & = (Id - \pi_0) j_{\perp} \\ A_0^0 & = \pi_0 j_0 \alpha \end{cases} \quad (2.18)$$

where  $\alpha = \alpha(x, \tau, \xi')$  is the unique endomorphism of  $E_0$ , such that

$$\oint_{\mathbb{T}^d} A_0^0 = \oint_{\mathbb{T}^d} \pi_0 j_0 \alpha = Id_{E_0}. \quad (2.19)$$

(This choice of  $A_0^0$  will insure the consistency with the definition (0.36) of  $p_0$ .)

The maps  $L^0(x, \tau, \xi') : E_0 \rightarrow E_0$  and  $L_{\perp}^0(x, \tau, \xi') : E_{\perp}^s \rightarrow E_{\perp}^{s-1}$  are then uniquely determined by (2.12).  $L_{\perp}^0$  is a smooth function of  $(x, \tau, \xi')$  defined near  $\tau = \xi' = 0$ , taking its values in the set of pseudo-differential operators of order 1 for the scale  $\{E_{\perp}^s\}$  on the torus: for any  $w_{\perp} \in \cup_s E_{\perp}^s$  one has

$$\mathbb{M}^0 j_{\perp}(w_{\perp}) - j_{\perp} L_{\perp}^0(w_{\perp}) = \mathbb{M}^0 \pi_0 j_{\perp}(w_{\perp}) - \pi_0 j_{\perp} L_{\perp}^0(w_{\perp}) \in \cap_s E^s. \quad (2.20)$$

The map  $A^0 = A_0^0 \oplus A_{\perp}^0$

$$E^s = E_0 \oplus E_{\perp}^s \xrightarrow{A^0} E^s \quad (2.21)$$

is an isomorphism; by (2.16) it satisfies

$$\|A^0 - Id\|_{E^s} \leq \text{Cte } \tau^2 \quad (\forall s \in I). \quad (2.22)$$

The equation (2.13) is equivalent to

$$\begin{cases} \mathbb{M}^0 A_0^k - A_0^k L^0 - A_0^0 L^k = R_0^k, & R_0^k \text{ bounded from } E_0 \text{ to } E^s \\ \mathbb{M}^0 A_\perp^k - A_\perp^k L_\perp^0 - A_\perp^0 L_\perp^k = R_\perp^k, & R_\perp^k \text{ bounded from } E_\perp^s \text{ to } E^{s-1} \end{cases} \quad (2.23)$$

where the right hand side is given by induction by the formula

$$k \geq 1 \quad R_{0,\perp}^k = \sum_{\substack{j+\ell+|\alpha|=k \\ j \neq k, \ell \neq k}} \frac{1}{\alpha!} \partial_{\xi'}^\alpha A_{0,\perp}^j \partial_{x'}^\alpha L_{\cdot,\perp}^\ell - \sum_{\substack{j+\ell+|\alpha|=k \\ \ell \neq k}} \frac{1}{\alpha!} \partial_{\xi'}^\alpha \mathbb{M}^j \partial_{x'}^\alpha A_{0,\perp}^\ell - i \partial_{x_d} A_{0,\perp}^{k-1}. \quad (2.24)$$

Let  $A^k = A_0^k \oplus A_\perp^k$ ,  $\tilde{A}^k = (A^0)^{-1} A^k$ ,  $\mathcal{L}^k = L^k \oplus L_\perp^k$ ,  $R^k = R_0^k \oplus R_\perp^k$  and  $\tilde{R}^k = (A^0)^{-1} R^k$ . The equation (2.23) can be rewritten  $\mathbb{M}^0 A^k - A^k \mathcal{L}^0 - A^0 \mathcal{L}^k = R^k$ , which is equivalent by (2.12) [ $\mathbb{M}^0 A^0 = A^0 \mathcal{L}^0$ ] to  $\mathcal{L}^0 \tilde{A}^k - \tilde{A}^k \mathcal{L}^0 - \mathcal{L}^k = \tilde{R}^k$ . The matrix form of this equation on  $E_0 \oplus E_\perp^s$  is

$$\begin{cases} L^0 (\tilde{A}^k)_{1,1} - (\tilde{A}^k)_{1,1} L^0 & = L^k + (\tilde{R}^k)_{1,1} \\ L_\perp^0 (\tilde{A}^k)_{2,2} - (\tilde{A}^k)_{2,2} L_\perp^0 & = L_\perp^k + (\tilde{R}^k)_{2,2} \end{cases} \quad (2.25)$$

$$\begin{cases} L^0 (\tilde{A}^k)_{1,2} - (\tilde{A}^k)_{1,2} L_\perp^0 & = (\tilde{R}^k)_{1,2} \\ L_\perp^0 (\tilde{A}^k)_{2,1} - (\tilde{A}^k)_{2,1} L^0 & = (\tilde{R}^k)_{2,1}. \end{cases} \quad (2.26)$$

The choice  $(\tilde{A}^k)_{1,1} = 0$ ,  $(\tilde{A}^k)_{2,2} = 0$  gives then  $L^k, L_\perp^k$  by (2.25). The unique solvability of (2.26) is a consequence of (0.34) which implies for  $|\tau| + |\xi'|$  small enough

$$\begin{cases} L_\perp^0 \text{ is invertible and } \|(L_\perp^0)^{-1}; E_\perp^{s-1} \rightarrow E_\perp^s\| \leq C & (\forall s \in I) \\ \text{Spectrum } (L^0) \subset \{z \in \mathbb{C}; |z| \leq \text{Cte}(|\tau| + |\xi'|)\}. \end{cases} \quad (2.27)$$

Thus, solving the second equation in (2.26) is equivalent to find a linear map  $u : E_0 = \mathbb{C}^2 \rightarrow \bigcap_{s \in I} E_\perp^s = E_\perp^{\sigma_0}$  such that

$$u - (L_\perp^0)^{-1} \circ u \circ L^0 = v \quad (2.28)$$

where  $v : E_0 \rightarrow E_\perp^{\sigma_0}$  is given and (2.27) implies for  $|\tau| + |\xi'|$  small the existence of a unique solution  $u$  to (2.28). The first equation in (2.26) can be reduced to the second one by taking adjoints.

**Remark.** We have chosen to work with a fixed interval of regularity on the torus,  $s \in [-\sigma_0, \sigma_0] = I$  in order to work in the classical theory of semi-classical (in  $\varepsilon$ ) pseudo-differential operators with values in bounded operators between Hilbert spaces. On the other hand, the neighborhood of the zero section  $\tau = \xi' = 0$  where the above construction applies may depends on  $I$ .

In view of (2.22), the tangential pseudo-differential operator  $A = A_0 \oplus A_\perp$  is elliptic near the zero section  $\tau = 0, \xi' = 0$ . Let  $A^{-1}$  be a pseudo-differential inverse and  $\mathcal{L} = L \oplus L_\perp$ .

By construction we have  $(\varepsilon \partial_{x_d} + \mathbb{M}) A \equiv A (\varepsilon \partial_{x_d} + \mathcal{L})$  near the zero section, and  $\mathcal{L}$  is diagonal in the decomposition  $E_0 \oplus E_\perp$ . Therefore, one deduces that the following identity holds true near the zero section

$$\oint_{\mathbb{T}^d} A^{-1} \left( \varepsilon \frac{\partial}{\partial x_d} + \mathbb{M} \right) \equiv \left( \varepsilon \frac{\partial}{\partial x_d} + L \right) \oint_{\mathbb{T}^d} A^{-1}. \quad (2.29)$$

If we choose  $W_1 \Subset W_0$  two sufficiently small neighborhoods of the set  $\{\xi' = 0, \tau = 0\} \times \{t \in [-T_0, 2T_0]\}$  in  $T^*X$ , and  $Q_0, Q_1$  two scalars tangential o.p.d., with  $SE(Q_j) \subset W_j, j = 0, 1$ , such that  $Q_0 \equiv Id$  on  $\overline{W_1}$ , we then get, with

$$\Pi_0 \stackrel{\text{def}}{=} Q_0 \oint_{\mathbb{T}^d} A^{-1}. \tag{2.30}$$

$$\Pi_0 \left( \varepsilon \frac{\partial}{\partial x_d} + \mathbb{M} \right) = \left( \varepsilon \frac{\partial}{\partial x_d} + L \right) \Pi_0 + R \tag{2.31}$$

where  $R$  is such that  $\|Q_1 R; L^2(X, E^s) \rightarrow L^2(X, \mathbb{C}^2)\| \in O(\varepsilon^\infty)$  for any  $s \in I$ .

The principal symbol  $p_0$  of  $\Pi_0$  is easy to compute:

If  $w = (Id - \pi_0)j_\perp(w_\perp) + \pi_0 j_0 \alpha(w_{(0)})$ , one has  $\pi_0(w) = \pi_0 j_0 \alpha(w_{(0)})$  and  $(A^0)^{-1}(w) = w_{(0)} \oplus w_\perp$ , so we get  $\oint_{\mathbb{T}^d} (A^0)^{-1}(w) = w_{(0)}$  (using (2.19))  $\oint_{\mathbb{T}^d} \pi_0 j_0 \alpha(w_{(0)}) = \oint_{\mathbb{T}^d} \pi_0(w)$  and we recover the definition formula (0.36) of  $p_0$ , if one takes  $\chi_0$  equal to the principal symbol of  $Q_0$ .

**Lemma 2.1.** *The tangential o.p.d.  $L \simeq \sum_k (\frac{\varepsilon}{i})^k L^k(x, \tau, \xi')$  satisfies*

$$\begin{aligned} \text{i)} \quad L &\equiv \oint_{\pi^d} (\mathbb{M}|_{\tau=0}) j_0 \quad \text{modulo } \tau^2 \\ \text{ii)} \quad L^0 &= \begin{bmatrix} a_0^{-1}(x) a_1(x, i\xi') & -a_0^{-1}(x) \\ a_2(x, i\xi') + \hat{\rho}(x)\tau^2 & 0 \end{bmatrix} + O(\tau^4). \end{aligned} \tag{2.32}$$

*Proof.* For i), we observe that  $\tau^2$  is a smooth parameter in the above construction, and that by formulas (0.27, 0.28), the restriction  $\mathbb{M}|_{\tau=0}$  is a constant coefficient operator on the torus  $\mathbb{T}_y^d$ .

We thus get  $\pi_0|_{\tau=0} = \oint_{\mathbb{T}^d}$ ,  $\alpha|_{\tau=0} = Id$ ,  $A_0|_{\tau=0} = j_0$ ,  $A_\perp|_{\tau=0} = j_\perp$ ,  $L|_{\tau=0} = \oint_{\mathbb{T}^d} (\mathbb{M}|_{\tau=0}) j_0$ ,  $L_\perp|_{\tau=0} = (Id - \oint_{\mathbb{T}^d})|_{\tau=0} (\mathbb{M}|_{\tau=0}) j_\perp$ .

One has  $\oint_{\mathbb{T}^d} A_0^0 = Id_{E_0}$  and  $A_0^0 = j_0 + O(\tau^2)$ , so there exists a map  $\theta(x, \tau^2, \xi') : E_0 \rightarrow E_\perp$  such that  $A_0^0 = j + \tau^2 \theta$ . Using (2.12), we get

$$L^0 = \oint_{\mathbb{T}^d} \mathbb{M}^0 A_0^0 \tag{2.33}$$

so for any  $w \in E_0$

$$L^0(w) = \oint_{\mathbb{T}^d} \mathbb{M}^0 j_0(w) + \tau^2 \oint_{\mathbb{T}^d} \mathbb{M}^0 \theta(w). \tag{2.34}$$

The definition formula (0.29) of  $\mathbb{M}^0$  and (2.34) gives the second part of the lemma. □

For  $u^\varepsilon \in I_h^\varepsilon$ , we define  $\underline{u}^\varepsilon$  by

$$\underline{u}^\varepsilon = \begin{bmatrix} u_0^\varepsilon \\ u_1^\varepsilon \end{bmatrix} = \begin{bmatrix} u^\varepsilon \\ A_0(x) \varepsilon \partial_{x_d} u^\varepsilon + \varepsilon A_1(x, \partial_{x'}) u^\varepsilon \end{bmatrix} \tag{2.35}$$

and  $w^\varepsilon = \mathcal{T}(u^\varepsilon) = \mathcal{T}(\underline{u}^\varepsilon)$  by

$$w^\varepsilon = \begin{bmatrix} w_0^\varepsilon \\ w_1^\varepsilon \end{bmatrix} = \begin{bmatrix} \mathcal{T}(u_0^\varepsilon) \\ \mathcal{T}(u_1^\varepsilon) \end{bmatrix} \tag{2.36}$$

where  $T$  is the transformation (0.19):

$$T(f)(t, x, y) = \sum_{\ell \in \mathbb{Z}^d} e^{i\ell(y - \Theta((x)/\varepsilon))} f(t, x). \tag{2.37}$$

Then  $w^\varepsilon$  satisfies, for  $s_0 < -d/2$ .

$$w^\varepsilon(t, x, y) \in L^2([t_1, t_2] \times \partial\Omega \times [0, r_0]; [H^{s_0}(\mathbb{T}^d)]^2) \quad \forall t_1, t_2 \in \mathbb{R} \tag{2.38}$$

$$\begin{cases} \left(\varepsilon \frac{\partial}{\partial x_d} + \mathbb{M}\right) w^\varepsilon = 0 & \text{on } \mathbb{R}_t \times \partial\Omega \times ]0, r_0[ \times \mathbb{T}_y^d \\ w_0^\varepsilon|_{x_d=0} = 0. \end{cases} \tag{2.39}$$

We recall that we define the Bloch wave  $\Gamma(u^\varepsilon) \in L^2(X; \mathbb{C}^2)$  by

$$\Gamma(u^\varepsilon) = \begin{bmatrix} \Gamma_0(u^\varepsilon) \\ \Gamma_1(u^\varepsilon) \end{bmatrix} = \Pi_0 \circ T(\underline{u}^\varepsilon). \tag{2.40}$$

*Proof of Proposition 1.*

(We denote by  $C$  various constants which are independent of  $\varepsilon, h$ .)

For  $u^\varepsilon = \sum_{0.9 \leq \omega_n^\varepsilon h \leq 2.1} (u_{+,n} e^{it\omega_n^\varepsilon} + u_{-,n} e^{it\omega_n^\varepsilon}) e_n^\varepsilon(x)$  we put  $\|u^\varepsilon\|^2 = \sum |u_{+,n}|^2 + |u_{-,n}|^2$ . For any  $t_1 < t_2$ , there exists a constant  $C$  such that for any  $\varepsilon, h$  and  $u^\varepsilon \in I_h^\varepsilon$  one has

$$\int_\Omega \int_{t_1}^{t_2} |h\nabla u^\varepsilon|^2 + |h\partial_t u^\varepsilon|^2 dt dx \leq C \|u^\varepsilon\|^2. \tag{2.41}$$

Let  $\gamma = \varepsilon/h$ ; we rewrite (2.41) on the form

$$\int_\Omega \int_{t_1}^{t_2} |\varepsilon\nabla u^\varepsilon|^2 + |\varepsilon\partial_t u^\varepsilon|^2 dt dx \leq C\gamma^2 \|u^\varepsilon\|^2. \tag{2.42}$$

Let  $K = \partial\Omega \times [0, T_0] \times [0, r_0/2] \times \{\xi' = 0, \tau = 0\}$  and  $Q(\varepsilon, t, x, \varepsilon\partial_{x'}, \varepsilon\partial_t)$  be a scalar tangential o.p.d. on  $X$ , equal to  $Id$  near  $K$ .

Let  $\alpha$  small such that the geometric control hypothesis (0.9) holds true for  $T_0 - 4\alpha$ , and let  $Y = \partial\Omega \times [\alpha, T_0 - \alpha] \times [0, r_0/2]$ . By (2.42), for  $\gamma$  small, the  $L^2$  norm of  $u^\varepsilon$  on  $Y$  is concentrated near the set  $\xi' = 0, \tau = 0$  where  $Q$  is equal to  $Id$ ; so we get

$$\|u^\varepsilon\|_{L^2(Y)}^2 \leq C \left[ \|Q(u^\varepsilon)\|_{L^2(X_{T_0})}^2 + (\gamma + \varepsilon)^2 \|u^\varepsilon\|^2 \right]. \tag{2.43}$$

By construction of  $\Pi_0$ , one has

$$\Pi_0 = Q_0 \left[ \oint_{\mathbb{T}^d} Id + R_0(\varepsilon\partial_t) + \varepsilon R_1 \right] \tag{2.44}$$

where  $R_{0,1}$  are tangential o.p.d. from  $L^2(X; E^s)$  in  $L^2(X; E_0)$  ( $s \in I$ ). Therefore we get

$$\left\| \Gamma(u^\varepsilon) - Q_0 \begin{pmatrix} u_0^\varepsilon \\ u_1^\varepsilon \end{pmatrix} \right\|_{L^2(X; E_0)} \leq C[\gamma + \varepsilon] \|u^\varepsilon\| \tag{2.45}$$

(here we have used the fact that  $\varepsilon\partial_t$  commutes with  $T$  and is bounded by  $O(\gamma = \varepsilon/h)$  on  $I_h^\varepsilon$ ). Since  $QQ_0$  is equal to  $\text{Id}$  near  $K$ , we deduce from (2.43, 2.45), for  $\gamma_0, \varepsilon_0$  small enough

$$\|u^\varepsilon\|_{L^2(Y)}^2 \leq C \left[ \|Q\Gamma_0(u^\varepsilon)\|_{L^2(X_{T_0})}^2 + (\gamma + \varepsilon)^2 \|u^\varepsilon\|^2 \right]. \tag{2.46}$$

We are now ready to prove (0.45) by a contradiction argument. If (0.45) is untrue, there exist sequences  $\varepsilon_k \rightarrow 0, \gamma_k \rightarrow 0, h_k \rightarrow 0$ ,  $h_k \geq \varepsilon_k/\gamma_k$ ,  $u^k \in I_{h_k}^{\varepsilon_k}$  such that

$$\begin{cases} \|u^k\| = 1 \\ \lim_{k \rightarrow \infty} \|Q\Gamma_0(u^k)\|_{L^2(X_{T_0})}^2 + \|u^k\|_{L^2((0, T_0) \times V)}^2 = 0. \end{cases} \tag{2.47}$$

Moreover, we can suppose that the weak limit  $u = \text{weak} - \lim(u^k)$  exist. Then  $u$  satisfies (0.6) and is equal to 0 on  $(0, T_0) \times V$ . By the geometric control hypothesis (0.9) of [4], the estimate (0.7) holds true for  $u$ , so we get  $u = 0$ . We deduce from (2.46)

$$\lim_{k \rightarrow \infty} \|u^k\|_{L^2(Y)} = 0. \tag{2.48}$$

We are thus reduced to an interior problem in  $\Omega$ .

Let  $Z = \{x \in \Omega ; \text{dist}(x, \partial\Omega) > r_0/4\} \times \mathbb{R}_t$ . We denote by  $\widetilde{M} = \rho(x, y)(\varepsilon\partial_t)^2 - \varepsilon^2(\Delta_g)_\varepsilon^*$  the Bloch operator on  $Z$ , and  $G^s = H^s(\mathbb{T}^d)$ . By the same construction as above, there exist a  $\varepsilon$ -pseudo-differential operator  $\widetilde{\Pi}_0(x, \xi, \tau, y, \partial_y) : L^2(Z, G^\bullet) \rightarrow L^2(Z, \mathbb{C})$  and a scalar  $\varepsilon$ -o.p.d.  $\widetilde{L}(x, \xi, \tau) : L^2(Z; \mathbb{C}) \rightarrow L^2(Z, \mathbb{C})$ , defined near the zero section  $\xi = \tau = 0$ , such that

$$\widetilde{\Pi}_0 \widetilde{M} = \widetilde{L} \widetilde{\Pi}_0 + \widetilde{R}. \tag{2.49}$$

The principal symbol of  $\widetilde{\Pi}_0$  is  $\tilde{\chi}_0 \int_{\mathbb{R}^d} \frac{1}{2i\pi} \int_{\partial D} \frac{dz}{z - i\mathbb{M}^0}$  with  $\tilde{\chi}_0 \in C_0^\infty(\widetilde{W}_0), \tilde{\chi}_0 \equiv 1$  on  $\widetilde{W}_1$ , where  $\widetilde{W}_1 \Subset \widetilde{W}_0$  are two small neighborhood of the set  $\{\xi = \tau = 0\} \times \{t \in [-T_0, 2T_0]\}$  in  $T^*Z$ . The scalar operator  $\widetilde{L}$  satisfies

$$\begin{cases} \widetilde{L} \simeq \sum_k (\frac{\varepsilon}{i})^k \widetilde{L}^k(x, \tau, \xi) \\ \widetilde{L}|_{\tau=0} = -\varepsilon^2 \Delta_g \text{ modulo } \tau^2 \\ \widetilde{L}^0(x, \tau, \xi) = -\rho(x)\tau^2 + \|\xi\|^2 + o(\tau^4). \end{cases} \tag{2.50}$$

The error terms  $\widetilde{R}$  in (2.49) is such that for any  $\varepsilon$ -o.p.d.  $\widetilde{Q}$  with essential support in  $\widetilde{W}_1$ , one has

$$\|\widetilde{Q} \circ \widetilde{R} ; L^2(Z; G^s) \rightarrow L^2(Z; \mathbb{C})\| \in \mathcal{O}(\varepsilon^\infty) \quad \forall s \in [-\sigma_0, \sigma_0]. \tag{2.51}$$

Let  $v^k(t, x, y)$  be the distribution on  $Z \times \mathbb{T}^d$

$$v^k = T(u^k) = \sum_{\ell \in \mathbb{Z}^d} e^{i\ell(y-x/\varepsilon_k)} u^k(t, x). \tag{2.52}$$

We deduce from (2.50) that  $(\frac{h}{\varepsilon})^2 \widetilde{L} \stackrel{\text{def}}{=} \widetilde{\mathcal{L}}$  is an  $h$ -o.p.d.; writing  $\frac{\varepsilon}{i} \partial_x = \frac{\varepsilon}{h} (\frac{h}{i} \partial_x)$ , and using  $\frac{h_k}{\varepsilon_k} \geq \frac{1}{\gamma_k} \rightarrow \infty$  (2.49, 2.51) we get, for any  $h$ -o.p.d.  $Q$  compactly supported in  $\{\xi, \tau\}$  and with support in  $Z \times \{t \in (-T_0, 2T_0)\}$

$$\|Q \widetilde{\mathcal{L}} \widetilde{\Pi}_0 v^k\|_{L^2(Z)} \in o(h_k^\infty). \tag{2.53}$$

By the analogue of (2.45) in the interior case, we also have

$$\|\tilde{\Pi}_0 v^k - \tilde{Q}_0 u^k\|_{L^2(Z \cap \{t \in [-T_0, 2T_0]\})} \leq C[\gamma_k + \varepsilon_k] \tag{2.54}$$

where  $\tilde{Q}_0$  is an  $\varepsilon$ -o.p.d. with principal symbol  $\chi_0$ , with essential support in  $\tilde{W}_0$ .

Let  $\mu$  be a  $h$ -semi classical measure associated to  $\{u^k\}$  (see [8]). (The hypothesis  $u^k \in I_{h_k}^{\varepsilon_k}$  implies that  $\mu$  is supported in  $|\tau| \in [0.9, 2.1]$ .) From (2.47) and (2.46) we deduce that

$$\mu|_{Y \cap Z} \equiv 0 \text{ and } \mu|_{]0, T_0[ \times V} \equiv 0. \tag{2.55}$$

Let  $\nu$  be a  $h$ -semiclassical measure associated to  $\tilde{\Pi}_0 v^k$ . Using (2.54) and  $\lim_{k \rightarrow \infty} \varepsilon_k/h_k = 0$  we get

$$\nu = \tilde{\chi}_0^2(t, x; \xi' = 0, \tau = 0)\mu. \tag{2.56}$$

The principal symbol of  $\tilde{\mathcal{L}}$  is  $-\underline{\rho}(x)\tau^2 + \|\xi\|^2 + \gamma_k^2 0(\tau^4)$ . In the equation (2.53) we view  $\gamma_k = \varepsilon_k/h_k$  as a small parameter. We can then use the proof of the interior propagation theorem (see [8]) with the additional parameter  $\gamma_k$  going to zero. We get from (2.53) that the support of  $\nu$  is contained in the set  $\underline{\rho}(x)\tau^2 - \|\xi\|^2 = 0$ , and that the support of  $\nu$  propagates along the bicharacteristic flow of  $\underline{\rho}(x)\tau^2 - \|\xi\|^2$ . Using (2.55, 2.56), and the hypothesis (0.9) we obtain for  $\beta$  small

$$\mu|_{T_0/2-\beta, T_0/2+\beta[} \equiv 0. \tag{2.57}$$

Using (2.41), we get that the sequence  $u^k$  is  $h$ -oscillatory (see [7]), so from (2.57) we deduce

$$\lim_{k \rightarrow \infty} \|u^k\|_{L^2(Z \times ]T_0/2-\beta, T_0/2+\beta[)} \equiv 0.$$

Then from (2.48), we obtain  $\lim_{k \rightarrow \infty} \|u^k\|_{L^2(\Omega \times (T_0/2-\beta, T_0/2+\beta))} = 0$  which contradicts  $\|u^k\| \equiv 1$ . □

### 3. LOPATINSKI ESTIMATE

#### 3.1. Proof of Proposition 2

We first verify the implication Theorem 2  $\Rightarrow$  Proposition 2. For  $u^\varepsilon \in I_h^\varepsilon$ , we have

$$w^\varepsilon = \begin{bmatrix} w_0^\varepsilon \\ w_1^\varepsilon \end{bmatrix} = \begin{bmatrix} T(u^\varepsilon) \\ T(A_0(\varepsilon \partial_{x_d} u^\varepsilon) + \varepsilon A_1(x, \partial_{x'}) u^\varepsilon) \end{bmatrix} \tag{3.1}$$

and by (2.44)

$$\Gamma(u^\varepsilon) = Q_0 \left[ \oint_{\mathbb{T}^d} w^\varepsilon + R_0(\varepsilon \partial_t) w^\varepsilon + \varepsilon R_1 w^\varepsilon \right]. \tag{3.2}$$

The Dirichlet boundary condition  $u^\varepsilon|_{x_d=0}$  implies  $w_0^\varepsilon|_{x_d=0} = 0$ , so we get

$$\Gamma_0(u^\varepsilon)|_{x_d=0} = Q_0 \left[ \oint_{\mathbb{T}^d} (R_0(\varepsilon \partial_t) + \varepsilon R_1) \begin{bmatrix} 0 \\ w_1^\varepsilon|_{x_d=0} \end{bmatrix} \right]_{1^{st} \text{ component}}. \tag{3.3}$$



If one multiplies the equation (0.4) by  $\varepsilon^3 \varphi(x_d) \frac{\partial}{\partial x_d}$  where  $\varphi \in C_0^\infty(-r_0/2, r_0/2[)$  is equal to 1 near the boundary  $x_d = 0$ , and integrates by part, one gets

$$\left\{ \begin{array}{l} \text{For any } t_1, t_2, \text{ there exist } C \text{ s.t. } \quad \forall \varepsilon \\ \|\varepsilon \partial_n u^\varepsilon\|_{L^2((t_1, t_2) \times \partial\Omega)} \leq C \varepsilon^{-1/2} \|u^\varepsilon\| \quad \forall u^\varepsilon \in I_h^\varepsilon. \end{array} \right. \quad (3.4)$$

Therefore, by (3.1) we get for  $s_0 < -d/2$

$$\|w_1^\varepsilon|_{x_d=0}; L^2((t_1, t_2) \times \partial\Omega; H^{s_0}(\mathbb{T}_y^d))\| \leq C \varepsilon^{-1/2} \|u^\varepsilon\|. \quad (3.5)$$

If  $R$  is an o.p.d from  $L^2(X_{T_0 \cap x_d=0}; H^{s_0}(\mathbb{T}^d))$  in  $L^2(X_{T_0 \cap x_d=0})$ , using the classical calculus of Appendix A.1, we get from the *a priori* bound (3.5) on the trace  $w_1^\varepsilon|_{x_d=0}$

$$\|[Q_0, R]w_1^\varepsilon|_{x_d=0}; L^2(X_{T_0 \cap x_d=0})\| \leq C \varepsilon^{1/2} \|u^\varepsilon\|. \quad (3.6)$$

If Theorem 2 holds true, we have

$$\|Q_0 w_1^\varepsilon|_{x_d=0}; L^2(X_{T_0 \cap x_d=0}; H^{s_1}(\mathbb{T}^d))\| \leq C \|u^\varepsilon\|. \quad (3.7)$$

Now using the fact that  $\varepsilon \partial_t$  commutes with  $T$  and is bounded by  $\mathcal{O}(\gamma = \varepsilon/h)$  on  $I_h^\varepsilon$ , (3.3, 3.6, 3.7) and  $\varepsilon \leq h_0 \varepsilon/h$ , we get (0.47), *i.e.*

$$\|\Gamma_0(u^\varepsilon)|_{x_d=0}; L^2(X_{T_0} \cap x_d = 0)\| \leq C \frac{\varepsilon}{h} \|u^\varepsilon\|.$$

### 3.2. Proof of Theorem 2

In this part, we work with a family  $\{u^\varepsilon\}_\varepsilon, u^\varepsilon \in I_h^\varepsilon$  with  $\varepsilon \in ]0, \varepsilon_0], h \in [\varepsilon/\gamma_0, h_0]$ ; we always assume  $\|u^\varepsilon\| \leq 1$ . We first remark that the Theorem 2 is local near any  $\rho_0 = (t_0, x'_0, \tau_0 = 0, \xi'_0 = 0) \in T^*(\mathbb{R}_t \times \partial\Omega)$ . Let  $Q_1$  be a tangential scalar o.p.d equal to  $Id$  near  $\rho_0$ , and with essential support close to  $\rho_0$ , and contained in  $W_0$ . By (2.38, 2.39) we get (see (0.30) for the definition of  $E^s$ )

$$\left\{ \begin{array}{l} \left( \varepsilon \frac{\partial}{\partial x_d} + \mathbb{M}^0 \right) Q_1 w^\varepsilon = \tilde{g}^\varepsilon \\ \tilde{g}^\varepsilon = \left[ \varepsilon \frac{\partial}{\partial x_d} + \mathbb{M}, Q_1 \right] w^\varepsilon - \frac{\varepsilon}{i} \sum_{j=1}^2 \mathbb{M}^j Q_1 w^\varepsilon \end{array} \right. \quad (3.8)$$

and for any  $s_0 + 1 < -d/2$  and any  $t_1, t_2$

$$\sup_\varepsilon \|Q_1 w^\varepsilon; L^2([t_1, t_2] \times \partial\Omega \times [0, r_0]; E^{s_0+1})\| < +\infty \quad (3.9)$$

$$\sup_\varepsilon \varepsilon^{-1} \|\tilde{g}^\varepsilon; L^2([t_1, t_2] \times \partial\Omega \times [0, r_0]; E^{s_0})\| < +\infty. \quad (3.10)$$

We define  $f^\varepsilon, g^\varepsilon$  by

$$Q_1 w^\varepsilon = \begin{pmatrix} f_0^\varepsilon \\ i f_1^\varepsilon \end{pmatrix}, f^\varepsilon = \begin{pmatrix} f_0^\varepsilon \\ f_1^\varepsilon \end{pmatrix}, g^\varepsilon = \begin{pmatrix} \tilde{g}_0^\varepsilon \\ -\tilde{g}_1^\varepsilon \end{pmatrix}. \quad (3.11)$$

We may assume that  $f^\varepsilon$  is supported in a small neighborhood  $U = U_0 \times [0, r_1[$  of  $(t_0, x'_0)$  in  $\mathbb{R}_t \times \partial\Omega \times [0, r_0]$ , and we denote by  $(x_1, \dots, x_{d-1})$  a local coordinate system near  $x'_0$  in  $\partial\Omega$ . Near the boundary by the choice of coordinates (0.15), we have  $a_0(x) \equiv 1$  and  $a_1(x, \xi') \equiv 0$ , so equation (3.8) may be rewritten as

$$\begin{cases} \frac{\varepsilon}{i} \frac{\partial}{\partial x_d} f^\varepsilon + \mathbb{N} f^\varepsilon = g^\varepsilon \\ \mathbb{N} = \begin{pmatrix} & e_d(x) \cdot D_y & & -1 \\ a_2(x, \frac{\varepsilon}{i} \partial_{x'} + e'(x) D_y) - \hat{\rho}(x, y) \left( \frac{\varepsilon \partial_t}{i} \right)^2 & & e_d(x) \cdot D_y & \end{pmatrix} \end{cases} \quad (3.12)$$

with  $D_y = \frac{1}{i} \frac{\partial}{\partial y}$ ,  $e'(x) D_y = (e_1(x) \cdot D_y, \dots, e_{d-1}(x) \cdot D_y)$ , we define the trace operators  $Tr_0, Tr_1$  by

$$Tr_0(f^\varepsilon) = f_0^\varepsilon|_{x_d=0} \quad Tr_1(f^\varepsilon) = f_1^\varepsilon|_{x_d=0}. \quad (3.13)$$

We have  $Tr_0(f^\varepsilon) \equiv 0$  and we have to prove

$$\begin{cases} \text{If } W_0 \subset \{|\xi'| + |\tau| < \alpha_0\}, \text{ with } \alpha_0 \text{ small enough, there exist } s_1, C, \text{ s.t.} \\ \sup_\varepsilon \|Tr_1(f^\varepsilon); L^2(U_0; H^{s_1}(\mathbb{T}^d))\| \leq C. \end{cases} \quad (3.14)$$

For any  $\ell \in \mathbb{Z}^d$ , we define  $\ell_x^\perp$  and  $\ell_x''$

$$\ell_x^\perp = e_d(x) \cdot \ell, \quad \ell_x'' = (e_1(x) \cdot \ell, \dots, e_{d-1}(x) \cdot \ell). \quad (3.15)$$

We have by (32), with  $\|\ell_x''\|^2 = a_2(x, \ell_x'')$

$$\|{}^t d\theta(x)(\ell)\|_x^2 = (\ell_x^\perp)^2 + \|\ell_x''\|^2. \quad (3.16)$$

Let  $\mathbb{N}_0(x)$  be the restriction of  $\mathbb{N}$  to the zero section  $\xi' = \tau = 0$ . We have

$$\begin{cases} \mathbb{N}_{0,\ell}(x) = \begin{pmatrix} \ell_x^\perp & -1 \\ \|\ell_x''\|^2 & \ell_x^\perp \end{pmatrix} \in \mathcal{M}_2(\mathbb{R}) \\ \mathbb{N}_0(x) \left( \sum_\ell z_\ell e^{i\ell y} \right) = \sum_\ell \mathbb{N}_{0,\ell}(x) (z_\ell) e^{i\ell y} \end{cases} \quad (3.17)$$

and the eigenvalues of  $\mathbb{N}_{0,\ell}(x)$  are

$$\lambda_{\pm,\ell}^0(x) = \ell_x^\perp \pm i \|\ell_x''\|. \quad (3.18)$$

Our strategy of proof of the estimate (3.14) is to split  $f^\varepsilon$  into two pieces. The first one will be concentrate near  $\|\ell_x''\|$  small, where the spectrum of  $\mathbb{N}$  is close to the real axis; we shall treat this part by a perturbation argument on the spectral theory of  $\mathbb{N}$ . The second one  $\|\ell_x''\| \geq c^{te} > 0$  will be handle by elliptic estimates on  $\mathbb{N}$ .

To achieve this program, we shall use the ‘‘exotic’’ pseudo-differential calculus of Appendice A.2, with  $Z = \mathbb{R}_t \times \mathbb{R}_{x'}^{d-1} \times [0, r_0]_{x_d}$ ; to simplify notation we denote by  $\mathcal{S}^{t,m}$  (resp.  $\mathcal{B}^{t,m}$ ) the class of symbols (resp. operators) defined in (A.15) (resp. (A.17)). The restriction on  $x_d = 0$  of these class of symbols and operators will be denoted by  $\mathcal{S}^m, \mathcal{B}^m$ .

We first conjugate the equation (3.12) so that the natural scale of space on the torus will be

$$\mathcal{H}^s \stackrel{\text{def}}{=} [H^s(\mathbb{T}^d)]^2. \quad (3.19)$$

Let  $\langle \ell''_x \rangle = (1 + \|\ell''_x\|^2)^{1/2}$ . We define the operators  $\Lambda = \Lambda(x)$  and  $\mathbb{E}_0 = \mathbb{E}_0(x)$  on the torus by

$$\Lambda(\sum_{\ell} z_{\ell} e^{i\ell y}) = \sum_{\ell} \begin{pmatrix} 1 & 0 \\ 0 & \langle \ell''_x \rangle \end{pmatrix} (z_{\ell}) e^{i\ell y} \quad (3.20)$$

$$\mathbb{E}_0(\sum_{\ell} z_{\ell} e^{i\ell y}) = \sum_{\ell} \begin{pmatrix} \ell_x^{\perp} & -\langle \ell''_x \rangle \\ \frac{\|\ell''_x\|^2}{\langle \ell''_x \rangle} & \ell_x^{\perp} \end{pmatrix} (z_{\ell}) e^{i\ell y}. \quad (3.21)$$

Let  $F^{\varepsilon}$  be

$$F^{\varepsilon} = \Lambda^{-1}(f^{\varepsilon}). \quad (3.22)$$

We have  $Tr_0(F^{\varepsilon}) = 0$  and by (3.9), and the fact that  $\Lambda^{-1}$  maps clearly  $E^{s+1}$  in  $\mathcal{H}^s$ , we get

$$\sup_{\varepsilon} \|F^{\varepsilon}; L^2(U; \mathcal{H}^{s_0})\| < \infty. \quad (3.23)$$

**Lemma 3.1.** *There exist  $q \in \mathcal{S}^{t,0}$ , with*

$$q|_{\xi'=0, \tau=0} \equiv 0 \quad (3.24)$$

such that, for any scalar tangential symbol  $\theta(t, x, \tau, \xi')$  equal to  $Id$  near the essential support of  $Q_1$  and with support in  $\{|\xi'| + |\tau| \leq \alpha_0\}$ ,  $F^{\varepsilon}$  satisfies

$$G^{\varepsilon} = \frac{\varepsilon}{i} \frac{\partial}{\partial x_d} F^{\varepsilon} + \left( \mathbb{E}_0 + \begin{pmatrix} 0 & 0 \\ Op(q\theta) & 0 \end{pmatrix} \right) F^{\varepsilon} \quad (3.25)$$

$$\sup_{\varepsilon} \varepsilon^{-1} \|G^{\varepsilon}; L^2(U; \mathcal{H}^{s_0-1})\| < +\infty. \quad (3.26)$$

*Proof.* We conjugate (3.12) by  $\Lambda$  and we obtain

$$\frac{\varepsilon}{i} \frac{\partial}{\partial x_d} F^{\varepsilon} + \Lambda^{-1} \mathbb{N} \Lambda F^{\varepsilon} = \Lambda^{-1} g^{\varepsilon} - \Lambda^{-1} \frac{\varepsilon}{i} \left( \frac{\partial}{\partial x_d} \Lambda \right) F^{\varepsilon}. \quad (3.27)$$

We have

$$\left( \frac{\partial}{\partial x_d} \Lambda \right) \left( \sum_{\ell} z_{\ell} e^{i\ell y} \right) = \sum_{\ell} \begin{pmatrix} 0 & 0 \\ 0 & \frac{\partial}{\partial x_d} \langle \ell''_x \rangle \end{pmatrix} (z_{\ell}) e^{i\ell y}$$

and  $|\frac{\partial}{\partial x_d} \langle \ell''_x \rangle| \leq c^{t\varepsilon} (1 + |\ell|^2)^{1/2}$ ; therefore (by (3.10, 3.19)) we get

$$\sup_{\varepsilon} \varepsilon^{-1} \|\Lambda^{-1} g^{\varepsilon} - \Lambda^{-1} \frac{\varepsilon}{i} \left( \frac{\partial}{\partial x_d} \Lambda \right) F^{\varepsilon}; L^2(U; \mathcal{H}^{s_0-1})\| < +\infty. \quad (3.28)$$

A simple computation gives

$$\begin{cases} \Lambda^{-1}\mathbb{N}\Lambda = \mathbb{E}_0 + \begin{pmatrix} 0 & 0 \\ R & 0 \end{pmatrix} \\ R = Op(\oplus_{\ell}\langle\ell''_x\rangle)^{-1} \left[ a_2 \left( x, \frac{\varepsilon}{i} \frac{\partial}{\partial x'} \right) + \sum_{j=1}^{d-1} \frac{\partial a_2}{\partial \xi'_j} \left( x, \frac{\varepsilon}{i} \frac{\partial}{\partial x'} \right) (e_j(x).D_y) + \hat{\rho}(x, y)(\varepsilon\partial_t)^2 \right] \end{cases} \quad (3.29)$$

with

$$Op \left( \oplus_{\ell}\langle\ell''_x\rangle \right)^{-1} \left( \sum_{\ell} z_{\ell} e^{i\ell y} \right) = \sum_{\ell} \langle\ell''_x\rangle^{-1} z_{\ell} e^{i\ell y}.$$

Let  $\theta(t, x, \tau, \xi')$  be a classical tangential *o.p.d.* with support in  $\{|\xi'| + |\tau| \leq \alpha_0\}$  and equal to  $Id$  near the essential support of  $Q_1$ . By (3.11) we have

$$\|Op(\theta)F^{\varepsilon} - F^{\varepsilon}; L^2(U, \mathcal{H}^{s_0})\| \in \mathcal{O}(\varepsilon^{\infty}). \quad (3.30)$$

Therefore we can move  $R((1 - Op(\theta)F^{\varepsilon}))$  from the left to the right of (3.27). So we just have to verify

$$R \circ Op(\theta) = Op(q\theta) + \varepsilon Op(\oplus_{\ell}\langle\ell''_x\rangle)^{-1} \circ Op(b) \quad (3.31)$$

with  $q \in \mathcal{S}^{t,0}$  so that (3.24) holds true, and  $b \in \mathcal{S}^{t,1}$ . The  $b$  term in (3.31) is defined by  $[Op(a_2(x, \xi') + \dots) \circ Op(\theta) = Op(\theta(a_2 + \dots))] + \varepsilon Op(b)$  and belongs clearly to  $\mathcal{S}^{t,1}$  (there is no loose in the  $x$  derivatives of  $b$  in (A.15)). Let  $\chi(\tau, \xi') \in C_0^{\infty}$  equal to 1 for  $(|\tau| + |\xi'|) \leq 2\alpha_0$ . We define  $q$  by

$$q = \left( \oplus_{\ell}\langle\ell''_x\rangle^{-1} \right) \left[ a_2(x, \xi') + \sum_{j=1}^{d-1} \frac{\partial a_2}{\partial \xi'_j}(x, \xi')(e_j(x).D_y) - \hat{\rho}(x, y)\tau^2 \right] \cdot \chi(\tau, \xi'). \quad (3.32)$$

The estimates  $|e_j(x).\ell| \leq C^{te}\langle\ell''_x\rangle, j \leq d-1$ , and

$$\forall \alpha \exists C_{\alpha} |\partial_x^{\alpha}(\langle\ell''_x\rangle^{-1})| \leq C_{\alpha}(1 + |\ell|)^{|\alpha|}(\langle\ell''_x\rangle^{-1}) \quad (3.33)$$

implies  $q \in \mathcal{S}^{t,0}$ . The function  $a_2(x, \xi')$  is quadratic in  $\xi'$  so (3.24) follows from (3.32).  $\square$

The eigenvalues of  $\mathbb{E}_0 = \Lambda^{-1}\mathbb{N}_0\Lambda$  are  $\lambda_{\pm, \ell}^0(x) = \ell_x^{\perp} \pm i\|\ell''_x\|$ . For any  $x$ , the set  $(e_1(x), \dots, e_d(x))$  is a basis of  $\mathbb{R}^d$ , so by the definition (3.15) of  $\ell_x^{\perp}$  and  $\ell''_x$ , there exist  $c_1 > 0$  such that

$$|\ell_x^{\perp} - k_x^{\perp}| + \|\ell''_x - k''_x\| \geq 4c_1|\ell - k| \quad \forall x, \forall k, \ell \in \mathbb{Z}^d. \quad (3.34)$$

This implies the following separation property for the spectrum of  $\mathbb{E}_0$  near the real axis

**Lemma 3.2.** *For any  $x, \ell \in \mathbb{Z}^d$  such that  $\|\ell''_x\| \leq c_1$ , one has*

$$\text{dist}(\{\lambda_{\pm, \ell}^0(x)\}, \{\lambda_{\pm, k}^0(x)\}) \geq c_1 \quad \forall k \neq \ell. \quad (3.35)$$

*Proof.* If (3.35) is false, one has  $|\ell_x^{\perp} - k_x^{\perp}| < c_1$  and

$$\|\|\ell''_x\| - \|k''_x\|\| < c_1, \text{ so we get } |\ell_x^{\perp} - k_x^{\perp}| + \|\ell''_x - k''_x\| < c_1 + 3c_1$$

in contradiction with (3.34).  $\square$

Let  $Sp_0(x)$  be the spectrum of  $\mathbb{E}_0(x)$

$$Sp_0(x) = \bigcup_{\pm, \ell} \lambda_{\pm, \ell}^0(x). \tag{3.36}$$

By (3.21), for  $\lambda \notin Sp_0(x)$  the resolvent  $(\lambda - \mathbb{E}_0(x))^{-1}$  is diagonal with respect to the decomposition  $\bigoplus_{\ell} e^{i\ell y} \mathbb{C}^2$ ,  $(\lambda - \mathbb{E}_0(x))^{-1} = \bigoplus_{\ell} (\lambda - \mathbb{E}_{0, \ell}(x))^{-1}$  with

$$(\lambda - \mathbb{E}_{0, \ell}(x))^{-1} = \frac{1}{(\lambda - \lambda_{+, \ell}^0)(\lambda - \lambda_{-, \ell}^0)} \begin{pmatrix} \lambda - \ell_x^\perp & -\langle \ell_x'' \rangle \\ \frac{\|\ell_x''\|^2}{\langle \ell_x'' \rangle} & \lambda - \ell_x^\perp \end{pmatrix}. \tag{3.37}$$

**Lemma 3.3.** *For any  $c_0 > 0$ , there exist  $M$  such that for any  $x$ ,  $\text{dist}(\lambda, Sp_0(x)) \geq c_0$  implies*

$$\|(\lambda - \mathbb{E}_{0, \ell}(x))^{-1}\| \leq M \quad \forall \ell. \tag{3.38}$$

*Proof.* We may suppose  $\text{Im} \lambda \geq 0$ . Then we have  $|\lambda - \lambda_{+, \ell}^0| |\lambda - \lambda_{-, \ell}^0| \geq c_0 |\lambda - \ell_x^\perp + i \|\ell_x''\| |$ , so

$$\frac{|\lambda - \ell_x^\perp|}{|\lambda - \lambda_{+, \ell}^0| |\lambda - \lambda_{-, \ell}^0|} \leq \frac{1}{c_0}$$

and

$$\frac{c_0 \langle \ell_x'' \rangle}{|\lambda - \lambda_{+, \ell}^0| |\lambda - \lambda_{-, \ell}^0|} \leq \frac{\sqrt{1 + \|\ell_x''\|^2}}{\max(c_0, \|\ell_x''\|)}.$$

The lemma follows from these two inequalities by (3.37). □

Let us define  $\mathbb{E} = \mathbb{E}(t, x, \tau, \xi')$  by (see (3.25))

$$\mathbb{E} = \mathbb{E}_0 + \begin{pmatrix} 0 & 0 \\ q\theta & 0 \end{pmatrix}. \tag{3.39}$$

For  $\beta > 0$ , let  $\Sigma_\beta(x) \subset \mathbb{Z}^d$  be the set

$$\Sigma_\beta(x) = \{\ell \in \mathbb{Z}^d, \|\ell_x''\| < \beta\} \tag{3.40}$$

and for  $\ell \in \mathbb{Z}^d$ , let  $\gamma_\ell(x)$  be the circle

$$\gamma_\ell(x) = \{z \in \mathbb{C}, |z - \ell_x^\perp| = c_1/4\} \tag{3.41}$$

where  $c_1$  is the constant of Lemma 3.2.

Now, we fixe  $\beta$ ,  $0 < \beta \ll c_1/4$ . Then for any  $x$  and  $\ell \in \Sigma_\beta(x)$ , one has  $|\lambda_{\pm, \ell}^0(x) - \ell_x^\perp| \leq \beta \ll c_1/4$ , so the eigenvalues  $\lambda_{\pm, \ell}^0(x)$  are the only ones inside the circle  $\gamma_\ell(x)$ . By Lemma 3.3, one gets

$$\left\{ \begin{array}{l} (\lambda - \mathbb{E}_0(x))^{-1} \in \mathcal{A}^0 \\ \|(\lambda - \mathbb{E}_0(x))^{-1}; \mathcal{H}^0 \rightarrow \mathcal{H}^0\| \leq M \end{array} \right\} \quad \begin{array}{l} \forall \lambda \in \bigcup_{\ell \in \Sigma_\beta(x)} \gamma_\ell(x) \\ \forall x. \end{array} \tag{3.42}$$

We then apply Lemma A.1:  $q$  vanishes on  $\xi' = 0, \tau = 0$  and  $\theta$  is supported in  $|\xi'| + |\tau| \leq \alpha_0$ . Therefore, if  $\alpha_0$  is small enough, the resolvent  $(\lambda - \mathbb{E}(t, x, \tau, \xi'))^{-1}$  exist for any  $(t, x, \tau, \xi', \lambda)$  for  $\lambda \in \bigcup_{\ell \in \Sigma_\beta(x)} \gamma_\ell(x)$ . Obviously,

one has

$$\frac{1}{2i\pi} \int_{\gamma_{\ell}(x)} (\lambda - \mathbb{E}_0(x))^{-1} d\lambda \left( \sum_k z_k e^{iky} \right) = z_{\ell} e^{i\ell y}. \quad (3.43)$$

We choose  $\psi \in C_0^{\infty}(\cdot - 1, 1]$  equal to 1 on  $[-1/2, 1/2]$  and we define  $pr_0(x), pr(t, x, \tau, \xi')$  by the formulas

$$pr_0(x) \left[ \sum_{\ell} z_{\ell} e^{i\ell y} \right] = \sum_{\ell} \psi \left( \frac{\|\ell''_x\|^2}{\beta^2} \right) z_{\ell} e^{i\ell y} \quad (3.44)$$

$$pr(t, x, \tau, \xi') = \sum_{\ell} \psi \left( \frac{\|\ell''_x\|^2}{\beta^2} \right) \frac{1}{2i\pi} \int_{\gamma_{\ell}(x)} (\lambda - \mathbb{E}(t, x, \tau, \xi'))^{-1} d\lambda. \quad (3.45)$$

The next lemma shows that  $pr$  is well defined.

**Lemma 3.4.** *There exists a  $2 \times 2$  matrix  $\delta pr(t, x, \tau, \xi')$  with entries in  $\mathcal{S}^{t,0}$ , such that*

$$\begin{cases} pr = pr_0 + \delta pr \\ \delta pr|_{\xi'=0, \tau=0} = 0. \end{cases} \quad (3.46)$$

*Proof.* See Appendix B.

Let  $\varphi(t, x) \in C_0^{\infty}(U)$  equal to 1 near  $(t_0, x'_0)$ . We next define  $Q_0(t, x)$  and  $Q(t, x, \tau, \xi')$  by the formulas, where  $\langle \ell_x^{\perp} \rangle = \sqrt{1 + |\ell_x^{\perp}|^2}$ , and  $\sigma = 2|s_0| + 2$

$$Q_0(t, x) \left[ \sum z_{\ell} e^{i\ell y} \right] = \varphi \sum_{\ell} \psi \left( \frac{\|\ell''_x\|^2}{4\beta^2} \right) \frac{1}{\langle \ell_x^{\perp} \rangle^{\sigma}} \begin{pmatrix} 0 & -\langle \ell''_x \rangle \\ \frac{\|\ell''_x\|^2}{\langle \ell''_x \rangle} & 0 \end{pmatrix} (z_{\ell}) e^{i\ell y} \quad (3.47)$$

$$Q(t, x, \tau, \xi') = \varphi \sum_{\ell} \psi \left( \frac{\|\ell''_x\|^2}{4\beta^2} \right) \frac{1}{\langle \ell_x^{\perp} \rangle^{\sigma}} \frac{1}{2i\pi} \int_{\gamma_{\ell}(x)} (\lambda - \mathbb{E}(t, x, \tau, \xi'))^{-1} (\mathbb{E}(t, x, \tau, \xi') - \ell_x^{\perp}) d\lambda. \quad (3.48)$$

**Lemma 3.5.** *There exist a  $2 \times 2$  matrix  $\delta Q(t, x, \tau, \xi')$  with entries in  $\mathcal{S}^{t,-\sigma}$ , such that*

$$\begin{cases} Q = Q_0 + \delta Q \\ \delta Q|_{\xi'=0, \tau=0} = 0. \end{cases} \quad (3.49)$$

*Proof.* See Appendix B.

We then define  $F^{\varepsilon, \mathbb{R}}$  and  $F^{\varepsilon, I}$  by

$$\begin{cases} F^{\varepsilon, \mathbb{R}} = Op(pr) F^{\varepsilon} \\ F^{\varepsilon, I} = F^{\varepsilon} - F^{\varepsilon, \mathbb{R}}. \end{cases} \quad (3.50)$$

The Lemmas 3.4, A.2, the estimates (3.23) and (3.5), and the assumption  $\|u^{\varepsilon}\| \leq 1$  imply

$$\sup_{\varepsilon} \|F^{\varepsilon, \mathbb{R}, I}; L^2(U; \mathcal{H}^{s_0})\| < +\infty \quad (3.51)$$

$$\sup_{\varepsilon} \varepsilon^{1/2} \|Tr_{0,1}(F^{\varepsilon, \mathbb{R}, I}); L^2(U_0; H^{s_0}(\mathbb{T}^d))\| < +\infty. \quad (3.52)$$

Moreover,  $F^{\varepsilon, I}$  satisfies the following elliptic estimate

**Lemma 3.6.** *There exist  $D(t, x', \tau, \xi') \in \mathcal{S}^0$  such that*

$$\sup_{\varepsilon} \varepsilon^{-1/2} \|Tr_1(F^{\varepsilon, I}) - Op(D)Tr_0(F^{\varepsilon, I}); L^2(U_0; H^{s_0-1}(\mathbb{T}^d))\| < +\infty.$$

*Proof.* See Appendix B.

To simplify notations, for  $A \in \mathcal{S}^{t,*}$  we define  $\tilde{A}$  by  $\tilde{A} = Op(A)$ , and for  $g^\varepsilon$ , a family depending on  $\varepsilon$  in a norm space  $B$ ,  $g^\varepsilon \in \varepsilon^\alpha B$  means  $\sup_{\varepsilon} \varepsilon^{-\alpha} \|g^\varepsilon; B\| < +\infty$ . We denote also by  $\delta$  various symbols in  $\mathcal{S}^0$  such that  $\delta|_{\xi'=0, \tau=0} = 0$ . We first notice that  $Tr_0(F^\varepsilon) = 0$  and (3.44) imply  $Tr_0(\tilde{p}r_0(F^\varepsilon)) = 0$ , so by Lemma 3.4 we get

$$Tr_0(F^{\varepsilon, \mathbb{R}}) = \tilde{\delta} Tr_1(F^{\varepsilon, \mathbb{R}} + F^{\varepsilon, I}). \quad (3.53)$$

By Lemma 3.6, and the Lemmas A.2 and A.3 on the symbolic calculus, we deduce from (3.53)

$$Tr_0(F^{\varepsilon, \mathbb{R}}) + \tilde{\delta} Tr_1(F^{\varepsilon, \mathbb{R}}) + \tilde{\delta} Tr_0(F^{\varepsilon, I}) \in \varepsilon^{1/2} L^2(U_0, H^{s_0-1}). \quad (3.54)$$

We have  $Tr_0(F^{\varepsilon, I}) = -Tr_0(F^{\varepsilon, \mathbb{R}})$ , so (3.54) may be rewrite as a boundary condition for  $F^{\varepsilon, \mathbb{R}}$

$$(1 - \tilde{\delta}) Tr_0(F^{\varepsilon, \mathbb{R}}) + \tilde{\delta} Tr_1(F^{\varepsilon, \mathbb{R}}) \in \varepsilon^{1/2} L^2(U_0, H^{s_0-1}). \quad (3.55)$$

By Lemma 3.1,  $F^{\varepsilon, \mathbb{R}}$  satisfy the equation

$$\begin{cases} \frac{\varepsilon}{i} \partial_{x_d} F^{\varepsilon, \mathbb{R}} + \tilde{\mathbb{E}} F^{\varepsilon, \mathbb{R}} = G^{\varepsilon, \mathbb{R}} \\ G^{\varepsilon, \mathbb{R}} = \tilde{p}r(G^\varepsilon) + [\tilde{E}, \tilde{p}r] F^\varepsilon + \frac{\varepsilon}{i} (\partial_{x_d} \tilde{p}r) F^\varepsilon. \end{cases} \quad (3.56)$$

By construction, we have  $[\mathbb{E}, pr] \equiv 0$ , so by Lemma A.3  $[\tilde{\mathbb{E}}, \tilde{p}r] \in \varepsilon \mathcal{S}^{t,2}$  and from (3.23, 3.26) and Lemma A.2 we deduce

$$G^{\varepsilon, \mathbb{R}} \in \varepsilon L^2(U, \mathcal{H}^{s_0-2}). \quad (3.57)$$

For  $u(x, y) \in L^2(U, \mathcal{H}^s)$ ,  $v(x, y) \in L^2(U, \mathcal{H}^{-s})$  let  $\langle u|v \rangle$  be the duality

$$\langle u|v \rangle = \int_U \left( \int_{\mathbb{T}^d} u(x, y) \bar{v}(x, y) dy \right) dx \quad (3.58)$$

and let us define  $J : L^2(U, \mathcal{H}^s) \rightarrow L^2(U, \mathcal{H}^s)$  by

$$J \begin{pmatrix} u_0(x, y) \\ u_1(x, y) \end{pmatrix} = \begin{pmatrix} u_1(x, y) \\ u_0(x, y) \end{pmatrix}. \quad (3.59)$$

By the choice  $\sigma = 2|s_0| + 2$  and Lemma 3.6 (A.2) we have

$$\begin{cases} J \tilde{Q} F^{\varepsilon, \mathbb{R}} \in L^2(U; \mathcal{H}^{|s_0|+2}) \\ J \tilde{Q} F^{\varepsilon, \mathbb{R}} \text{ is compactly supported in } U. \end{cases} \quad (3.60)$$

Multiplying (3.56) by  $J\tilde{Q}F^{\varepsilon,\mathbb{R}}$ , we obtain (where  $(\cdot|\cdot)$  is the duality on  $x_d = 0$ )

$$\left\{ \begin{aligned} \langle G^{\varepsilon,\mathbb{R}} | J\tilde{Q}F^{\varepsilon,\mathbb{R}} \rangle &= \left\langle \frac{\varepsilon}{i} \partial_{x_d} F^{\varepsilon,\mathbb{R}} | J\tilde{Q}F^{\varepsilon,\mathbb{R}} \right\rangle + \left\langle J\tilde{\mathbb{E}}F^{\varepsilon,\mathbb{R}} | \tilde{Q}F^{\varepsilon,\mathbb{R}} \right\rangle \\ &= -\frac{\varepsilon}{i} \left( F^{\varepsilon,\mathbb{R}} |_{x_d=0} | J\tilde{Q}F^{\varepsilon,\mathbb{R}} |_{x_d=0} \right) + \mathcal{J}_1 + \mathcal{J}_2 + \mathcal{J}_3 \\ \mathcal{J}_1 &= \left\langle F^{\varepsilon,\mathbb{R}} | \left\{ (J\tilde{\mathbb{E}})^* - J\tilde{\mathbb{E}} \right\} \tilde{Q}F^{\varepsilon,\mathbb{R}} \right\rangle \\ \mathcal{J}_2 &= \left\langle JF^{\varepsilon,\mathbb{R}} | \left[ \frac{\varepsilon}{i} \partial_{x_d} + \tilde{\mathbb{E}}, \tilde{Q} \right] F^{\varepsilon,\mathbb{R}} \right\rangle \\ \mathcal{J}_3 &= \left\langle JF^{\varepsilon,\mathbb{R}} | \tilde{Q}G^{\varepsilon,\mathbb{R}} \right\rangle. \end{aligned} \right. \quad (3.61)$$

By (3.57) and (3.60), both  $|\langle G^{\varepsilon,\mathbb{R}} | J\tilde{Q}F^{\varepsilon,\mathbb{R}} \rangle|$  and  $\mathcal{J}_3$  are  $\mathcal{O}(\varepsilon)$ . By construction of  $Q$  (see (3.50)) we have  $[\mathbb{E}, Q] \equiv 0$  so by Lemma A.3, we get  $\mathcal{J}_2 \in \mathcal{O}(\varepsilon)$ . Finally, we have

$$J\mathbb{E} = J\mathbb{E}_0 + \begin{pmatrix} q\theta & 0 \\ 0 & 0 \end{pmatrix};$$

$$J\mathbb{E}_0 \left( \sum_{\ell} z_{\ell} e^{i\ell y} \right) = \sum_{\ell} \begin{pmatrix} \frac{\|\ell''_x\|^2}{\langle \ell''_x \rangle} & \ell_x^{\perp} \\ \ell_x^{\perp} & -\langle \ell''_x \rangle \end{pmatrix} (z_{\ell}) e^{i\ell y}$$

so we obtain

$$(J\tilde{\mathbb{E}})^* - (J\tilde{\mathbb{E}}) = \begin{pmatrix} (\tilde{q}\theta)^* - q\theta & 0 \\ 0 & 0 \end{pmatrix}. \quad (3.62)$$

By formula (3.32) and if we choose  $\theta(t, x, \tau, \xi')$  real,  $q\theta$  is self adjoint, so from Lemma A.4 we deduce  $\mathcal{J}_1 \in \mathcal{O}(\varepsilon)$ . Summing up, we have thus

$$\sup_{\varepsilon} \left| \left( F^{\varepsilon,\mathbb{R}} |_{x_d=0} | J\tilde{Q}F^{\varepsilon,\mathbb{R}} |_{x_d=0} \right) \right| < +\infty. \quad (3.63)$$

We now remark that if  $\delta_1 \in \mathcal{S}^0$  vanishes on  $\xi' = 0, \tau = 0$ , there exist  $\delta_2 \in \mathcal{S}^0$ , vanishing on  $\xi' = \tau = 0$  such that  $(1 + \delta_2)(1 - \delta_1) = 1 - p$ , where  $p(t, x, \tau, \xi') \in \mathcal{S}^0$ , is supported in  $c_0 \leq |\tau| + |\xi'| \leq 1/c_0$  for some  $c_0 > 0$ . Decreasing  $\alpha_0$ , hence  $W_0$ , if necessary, we then will have  $\tilde{p}Tr_{0,1}(F^{\varepsilon,\mathbb{R}}) \in \varepsilon^{1/2}L^2(U_0, H^{s_0-1})$ . Using once more Lemma A.3, we can thus rewrite the boundary condition (3.55) on the form

$$Tr_0(F^{\varepsilon,\mathbb{R}}) - \tilde{\delta}Tr_1(F^{\varepsilon,\mathbb{R}}) \in \varepsilon^{1/2}L^2(U_0, H^{s_0-1}). \quad (3.64)$$

Let  $Q = \begin{pmatrix} Q^1 & Q^2 \\ Q^3 & Q^4 \end{pmatrix}$ ; inserting (3.64) in (3.63) and taking in account the *a priori* estimate (3.52) we get

$$\left\{ \begin{aligned} \sup_{\varepsilon} |(Tr_1(F^{\varepsilon,\mathbb{R}}) | \tilde{A}Tr_1(F^{\varepsilon,\mathbb{R}}))| &< +\infty \\ \tilde{A} &= Q^2 + Q^1\delta + \delta^*Q^3\delta + \delta^*Q^4. \end{aligned} \right. \quad (3.65)$$

Let (we use Lem. 3.4 for the second equality)

$$F_0^{\varepsilon,\mathbb{R}} = \tilde{p}r_0(F^{\varepsilon}) = F^{\varepsilon,\mathbb{R}} + \delta(F^{\varepsilon,\mathbb{R}} + F^{\varepsilon,I}). \quad (3.66)$$



We know already that  $Tr_0(F^{\varepsilon, \mathbb{R}}), Tr_0(F^{\varepsilon, I})$  and  $Tr_1(F^{\varepsilon, I})$  are of the form  $\tilde{\delta}Tr_1(F^{\varepsilon, \mathbb{R}}) + \varepsilon^{1/2}L^2(U_0, H^{s_0-1})$  so we get from (3.66)

$$Tr_1(F_0^{\varepsilon, \mathbb{R}}) = (1 + \delta)Tr_1(F^{\varepsilon, \mathbb{R}}) + \varepsilon^{1/2}L^2(U_0, H^{s_0-1}). \quad (3.67)$$

Decreasing  $\alpha_0$  if necessary we get as above

$$Tr_1(F^{\varepsilon, \mathbb{R}}) = (1 + \delta)Tr_1(F_0^{\varepsilon, \mathbb{R}}) + \varepsilon^{1/2}L^2(U_0, H^{s_0-1}). \quad (3.68)$$

Therefore (3.65) and (3.68) imply

$$\begin{cases} \sup_{\varepsilon} |(Tr_1(F_0^{\varepsilon, \mathbb{R}})|\tilde{A}_0Tr_1(F_0^{\varepsilon, \mathbb{R}}))| < +\infty \\ A_0 = Q_0^2 + \delta^*A_1 + A_2\delta \end{cases} \quad (3.69)$$

with  $A_1, A_2 \in \mathcal{S}^{+\sigma}$ .

By (3.66, 3.44),  $Tr_1(F_0^{\varepsilon, \mathbb{R}})$  is of the form

$$Tr_1(F_0^{\varepsilon, \mathbb{R}}) = \sum_{\|\ell''_{(x', 0)}\| \leq \beta} z_{\ell}(t, x')e^{i\ell y} \quad (3.70)$$

and we may assume that the functions  $z_{\ell}(t, x')$  are supported in  $\{\varphi \equiv 1\}$ .

For  $\|\ell''_x\| \leq \beta$  we have  $\psi\left(\frac{\|\ell''_x\|^2}{4\beta^2}\right) = 1$ , and  $\frac{|\langle \ell''_x \rangle|}{\langle \ell''_x \rangle^{\sigma}} \sim (1 + |\ell|)^{-(2|s_0|+2)}$ ; from (3.47) we therefore get for some  $C_0 > 0$

$$|(Tr_1(F_0^{\varepsilon, \mathbb{R}})|Q_0^2Tr_1(F_0^{\varepsilon, \mathbb{R}}))| \geq C_0\|Tr_1(F_0^{\varepsilon, \mathbb{R}}); L^2(U_0, H^{s_0-1})\|^2. \quad (3.71)$$

We now remark that in (3.69), we may replace any  $\delta(t, x', \tau, \xi')$  term by  $\chi((\tau, \xi')/\alpha_0)\delta(t, x', \tau, \xi')$ , with  $\chi \in C_0^{\infty}$  equal to 1 in the unit ball, and

$$\chi((\tau, \xi')/\alpha_0)\delta = \sum_{j=1}^{d-1} \chi((\tau, \xi')/\alpha_0)\xi'_j b_j + \chi((\tau, \xi')/\alpha_0)\tau b_0$$

where  $b_* \in \mathcal{S}^0$  so we have, for some  $C_1 > 0$

$$|(Tr_1(F_0^{\varepsilon, \mathbb{R}})|(\tilde{\delta}\tilde{A}_0)Tr_1(F_0^{\varepsilon, \mathbb{R}}))| \leq C_1\alpha_0\|Tr_1(F_0^{\varepsilon, \mathbb{R}}); L^2(U_0, H^{s_0-1})\|^2. \quad (3.72)$$

From (3.69, 3.71, 3.72) we get, for  $\alpha_0$  small,

$$\sup_{\varepsilon} \|Tr_1(F_0^{\varepsilon, \mathbb{R}}); L^2(U_0; H^{s_0-1})\| < +\infty \quad (3.73)$$

so by (3.68), the same estimate holds true for  $Tr_1(F^{\varepsilon, \mathbb{R}})$ , hence also for

$$Tr_1(F^{\varepsilon, I}) = \tilde{\delta}Tr_1(F^{\varepsilon, \mathbb{R}}) + \varepsilon^{1/2}L^2(U_0, H^{s_0-1}).$$

Thus we have

$$\sup_{\varepsilon} \|Tr_1(F^{\varepsilon}); L^2(U_0, H^{s_0-1})\| < \infty. \quad (3.74)$$

This concludes the proof of Theorem 2. □

## 4. PROPAGATION ESTIMATE

This section is devoted to the proof of Proposition 0.3. We fix a zero order o.p.d.  $Q(\varepsilon, t, x, \varepsilon\partial_{x'}, \varepsilon\partial_t)$  equal to  $Id$  near  $K$ , with essential support in  $W_1$  and we argue by contradiction. If (0.50) is untrue, there exist sequences  $\varepsilon_k \rightarrow 0$ ,  $\gamma_k \rightarrow 0$ ,  $h_k \rightarrow 0$ ,  $h_k \geq \varepsilon_k/\gamma_k$ , and  $u^k \in I_{h_k}^{\varepsilon_k}$  such that

$$\begin{cases} \|u^k\| = 1 \\ \frac{1}{k} \|Q\Gamma_0(u^k)\|_{L^2(X_{T_0})}^2 \geq \left[ \|\Gamma_0(u^k)|_{x_d=0}\|_{L^2(X_{T_0} \cap x_d=0)}^2 + \|u^k\|_{L^2((0, T_0) \times V)}^2 \right]. \end{cases} \quad (4.1)$$

In particular the right hand side of the second line in (4.1) goes to zero.

Let  $\mathcal{L}$  and  $\begin{bmatrix} g_0 \\ g_1 \end{bmatrix}$  be defined by the formula (0.51) with  $u^\varepsilon = u^k$ . We have

$$L \sim \sum_n \left(\frac{\varepsilon}{i}\right)^n L^n, \quad L^n = \begin{pmatrix} L_1^n & L_2^n \\ L_3^n & L_4^n \end{pmatrix} \quad (4.2)$$

so we get

$$\mathcal{L} \sim \begin{pmatrix} h/\varepsilon L_1^0 & L_2^0 \\ h^2/\varepsilon^2 L_3^0 & h/\varepsilon L_4^0 \end{pmatrix} + \frac{h}{i} \begin{pmatrix} L_1^1 & \varepsilon/h L_2^1 \\ h/\varepsilon L_3^1 & L_4^1 \end{pmatrix} + \sum_{n \geq 2} \left(\frac{h}{i}\right)^n \left(\frac{\varepsilon}{h}\right)^{n-1} \begin{pmatrix} L_1^n & \varepsilon/h L_2^n \\ h/\varepsilon L_3^n & L_4^n \end{pmatrix}. \quad (4.3)$$

By Lemma 2.1, i)  $\frac{h}{\varepsilon} L_3^1$  is a smooth function of  $x, \xi' = \frac{h}{i} \partial_{x'}, \tau = \frac{h}{i} \partial_t$ , defined for  $\frac{h}{\varepsilon}(|\xi'| + |\tau|)$  small. Therefore,  $\mathcal{L}(h, x, \frac{h}{i} \partial_{x'}, \frac{h}{i} \partial_t)$  is a  $h$ -o.p.d. defined for  $\frac{h}{\varepsilon}(|\xi'| + |\tau|)$  small, with asymptotic development

$$\mathcal{L} \sim \sum_{n \geq 0} \left(\frac{h}{i}\right)^n \mathcal{L}^n \quad (4.4)$$

and by Lemma 2.1, ii) we get

$$\mathcal{L}_0 = \begin{pmatrix} a_0^{-1}(x) a_1(x, i\xi') & -a_0^{-1}(x) \\ a_2(x, i\xi') + \hat{\rho}(x)\tau^2 & 0 \end{pmatrix} + 0 \left( \left(\frac{\varepsilon}{h}\right)^2 \tau^4 \right). \quad (4.5)$$

Let  $\underline{u}^k$  be the extension of  $u^k$  by zero outside  $\Omega$ . Let  $\mu$  be a  $h$ -semiclassical measure associated to  $\{\underline{u}^k\}$  (see [8]). Let  $\underline{g}_0^k$  the extension of  $g_0^k = \Gamma_0(u^k)$  by zero on  $x_d < 0$  and let  $\nu$  be a  $h$ -semiclassical measure associated to  $\{\underline{g}_0^k\}$ . Using (2.45) and  $\lim_{k \rightarrow \infty} \varepsilon_k/h_k = 0$  we get

$$\nu = \chi_0^2(t, x; \xi' = 0, \tau = 0) \mu \quad (\text{for } x \in \partial\Omega \times ]-r_0, r_0[). \quad (4.6)$$

We have  $u^k \in I_{h_k}^{\varepsilon_k}$  so we know that  $\mu$  is supported in  $|\tau| \in [0.9, 2.1]$ ; moreover, by the proof of Proposition 1 Section 2, the support of  $\mu|_\Omega$  is contained in the set  $\hat{\rho}(x)\tau^2 - \|\xi\|^2 = 0$ , and  $\mu|_\Omega$  propagates on the bicharacteristic flow of  $\hat{\rho}(x)\tau^2 - \|\xi\|^2$ . Let  $g^k = \Gamma(u^k)$ , and let  $A(h, t, x, h\partial_{x'}, h\partial_t)$  be any  $h$ -o.p.d. compactly supported in  $T^*(X_{T_0})$ . Using (0.41, 0.42) and  $\lim_{k \rightarrow \infty} \varepsilon_k/h_k = 0$  we get with  $h = h_k$

$$A \left[ \left( h \frac{\partial}{\partial x_d} + \mathcal{L} \right) g^k \right] \in \mathcal{O}(h^\infty L^2). \quad (4.7)$$

Writing  $\mathcal{L} = \begin{pmatrix} \mathcal{L}_1 & \mathcal{L}_2 \\ \mathcal{L}_3 & \mathcal{L}_4 \end{pmatrix}$ , we observe that the principal symbol of  $\mathcal{L}_1$  vanishes near  $x_d = 0$ . Using (4.1, 4.7) we get that  $g_0^k$  satisfies near the boundary the second order tangential  $h$ -pseudo differential equation, with  $h = h_k$

$$\begin{cases} A \left[ (h\partial_{x_d})^2 g_0^k + \left( R_2 + hR_1 h \frac{\partial}{\partial x_d} \right) g_0^k \right] \in \mathcal{O}(h^\infty L^2) \\ \lim_{k \rightarrow \infty} \|g_0^k|_{x_d=0}\|_{L^2} = 0 \end{cases} \tag{4.8}$$

where  $R_{1,2}$  are  $h$ -tangential o.p.d. defined for  $\frac{h}{\varepsilon}(|\xi'| + |\tau|)$  small, and the principal symbol of  $R_2$ ,  $R_2^0$  is given by

$$R_2^0 = a_2(x, i\xi') + \underline{\rho}(x)\tau^2 + 0 \left( \left( \frac{\varepsilon}{h} \right)^2 \tau^4 \right). \tag{4.9}$$

We can now use the propagation theorem at the boundary for second order Dirichlet problem (see [8] for the localization and propagation at hyperbolic point and [10], Append. or [3], Th. 1 for the propagation result near the glancing set; here we view  $\gamma_k = \varepsilon_k/h_k$  as a small parameter in equation (4.8), and we notice that the proof of the propagation theorem allows this additional parameter going to zero). We get that the support of  $\nu$  is contained in the set  $\underline{\rho}(x)\tau^2 - \|\xi\|^2 = 0$ , and that the support of  $\nu$  propagates along the generalized bicharacteristic flow of  $\underline{\rho}(x)\tau^2 - \|\xi\|^2$ ; but (4.1) implies  $\mu_{\llbracket 0, T_0 \times V} \equiv 0$ , so from (0.9) and (4.6) we get  $\mu_{|t \in ]T_0/2 - \alpha, T_0/2 + \alpha[} \equiv 0$  for  $\alpha$  small. This is in contradiction with  $\|u^k\| = 1$  by (2.41).  $\square$

### A. SEMI-CLASSICAL O.P.D. WITH OPERATOR VALUES

#### A.1. Classical calculus

We recall here some classical properties of semi-classical tangential pseudo differential operators. Let  $Z = \mathbb{R}_z^p \times [0, r_0]_{x_d}$  and  $H_1, H_2$  two separable Hilbert spaces.

We denote by  $S_Z^t(H_1 \rightarrow H_2)$  the vector space of functions  $q(\varepsilon, z, \zeta, x_d)$  defined for  $\varepsilon \in ]0, \varepsilon_0]$  ( $\varepsilon_0$  small) smooth in  $(z, \zeta) \in T^*\mathbb{R}_z^p$ ,  $x_d \in [0, r_0]$ , compactly supported in  $z$ , with values in bounded operators from  $H_1$  to  $H_2$  which satisfies the estimates

$$\begin{aligned} \forall \alpha, k \exists C_{\alpha, k} \forall \varepsilon, z, \zeta, x_d \\ \|(1 + |\zeta|)^k \partial_{z, \zeta, x_d}^\alpha q(\varepsilon, z, \zeta, x_d); H_1 \rightarrow H_2\| \leq C_{\alpha, k} \end{aligned} \tag{A.1}$$

and admitting classical asymptotic expansions in  $\varepsilon$

$$q \sim \sum_{n=0}^{\infty} \left( \frac{\varepsilon}{i} \right)^n q_n(z, \zeta, x_d) \Leftrightarrow \forall N \quad q - \sum_{n < N} \left( \frac{\varepsilon}{i} \right)^n q_n \in \varepsilon^N S_Z^t. \tag{A.2}$$

For  $f(z, x_d) \in L^2(Z, H_1)$  with compact support in  $z$ , the Fourier transform  $\hat{f}_\varepsilon(\zeta, x_d)$  is defined by

$$\hat{f}_\varepsilon(\zeta, x_d) = \int e^{-iz\zeta/\varepsilon} f(z, x_d) dz \in L^2(\mathbb{R}_\zeta^p \times [0, x_d], H_1) \tag{A.3}$$

and for  $q \in S_Z^t(H_1, H_2)$ ,  $Op(q)(f)$  is defined by

$$Op(q)(f)(\varepsilon, z, x_d) = (2\pi\varepsilon)^{-p} \int e^{iz\zeta/\varepsilon} q(\varepsilon, z, \zeta, x_d) [\hat{f}_\varepsilon(\zeta, x_d)] d\zeta. \tag{A.4}$$

We define the set  $\mathcal{E}_Z^t(H_1 \rightarrow H_2)$  of tangential pseudo differential operators from  $L^2(Z, H_1)$  to  $L^2(Z, H_2)$  by

$$\left\{ \begin{array}{l} Q = Q_\varepsilon \in \mathcal{E}_Z^t(H_1 \rightarrow H_2) \text{ iff there exist } \varphi(z) \in C_0^\infty(\mathbb{R}_z^p) \\ \text{and } \tilde{q} \in S_Z^t(H_1 \rightarrow H_2) \text{ such that} \\ Q_\varepsilon(f)(z) = Op(\tilde{q})[\varphi(z)f] \quad \forall f \in L^2(Z, H_1). \end{array} \right. \quad (\text{A.5})$$

For  $Q \in \mathcal{E}_Z^t(H_1 \rightarrow H_2)$ , one has  $Q = Op(q)$  with

$$q(\varepsilon, z, \zeta, x_d) = (2\pi)^{-p} \int e^{-it\theta} \tilde{q}(\varepsilon, z, \zeta + \varepsilon\theta, x_d) \varphi(z+t) dt d\theta$$

and  $Q$  is bounded on  $L^2$ , *i.e.*

$$\exists C \quad \forall \varepsilon \quad \|Q_\varepsilon(f); L^2(Z, H_2)\| \leq C \|f; L^2(Z, H_1)\|. \quad (\text{A.6})$$

For  $Q_1 = Op(q_1) \in \mathcal{E}_Z^t(H_1 \rightarrow H_2)$  and  $Q_2 = Op(q_2) \in \mathcal{E}_Z^t(H_2 \rightarrow H_3)$ , one has  $Q_1 \circ Q_2 = Op(q) = Q \in \mathcal{E}_Z^t(H_1 \rightarrow H_3)$  with

$$q(\varepsilon, z, \zeta, x_d) = (2\pi)^{-p} \int e^{-it\theta} q_1(\varepsilon, z, \zeta + \varepsilon\theta, x_d) \circ q_2(\varepsilon, z+t, \zeta, x_d) dt d\theta$$

and the asymptotic expansion of  $q$  is given by the rule

$$q \sim \sum_{\alpha} \left(\frac{\varepsilon}{i}\right)^{|\alpha|} \frac{1}{\alpha!} \partial_\zeta^\alpha q_1 \circ \partial_z^\alpha q_2. \quad (\text{A.7})$$

The set of operators  $\mathcal{E}_Z^t(H_1 \rightarrow H_2)$  is free of coordinates, *i.e.*, if  $z \mapsto \phi(z)$  is a smooth diffeomorphism of  $\mathbb{R}_z^p$ , and  $Q \in \mathcal{E}_Z^t$ , then  $\phi \circ Q \circ \phi^{-1} \in \mathcal{E}_Z^t$ . Thus, in the definition of  $\mathcal{E}_Z^t$ , we can replace  $\mathbb{R}_z^p$  by a smooth manifold  $M$ . For  $Q = Op(q) \in \mathcal{E}_Z^t(H_1 \rightarrow H_2)$  its principal symbol,  $q_0(z, \zeta, x_d)$  is then defined as a smooth function of  $(z, \zeta, x_d) \in T^*M \times [0, r_0]$ , with values in bounded operators from  $H_1$  to  $H_2$ . For  $Q = Op(q) \in \mathcal{E}_Z^t$ , the essential support of  $Q$   $SE(Q)$  is the closed subset of  $T^*M \times [0, r_0]$  defined by

$$\left\{ \begin{array}{l} \rho_0 = (z_0, \zeta_0, x_{d,0}) \notin SE(Q) \text{ iff there exists a neighborhood} \\ W \text{ of } \rho_0 \text{ such that } q|_W \sim 0. \end{array} \right. \quad (\text{A.8})$$

Let  $K$  be a compact subset of  $T^*M \times [0, r_0]$ . One says that  $Q_1 \equiv Q_2$  near  $K$  if  $SE(Q_1 - Q_2) \cap K = \emptyset$  and if  $u : H_1 \rightarrow H_2$  is bounded,  $Q \equiv u$  near  $K$  means  $Q - \varphi(y, x_d)u \equiv 0$  for some  $\varphi \in C_0^\infty(M \times [0, r_0])$  equal to 1 near the projection of  $K$  on  $M \times [0, r_0]$ . If  $Q \equiv 0$  near  $K$ , for any scalar tangential o.p.d.  $P \in \mathcal{E}_Z^t(\mathbb{C} \rightarrow \mathbb{C})$ , such that  $SE(P) \subset K$  one has

$$\forall N, \exists C_N \|QP \text{ or } PQ; L^2(Z, H_1) \rightarrow L^2(Z, H_2)\| \leq C_N \varepsilon^N. \quad (\text{A.9})$$

One says that  $Q = Op(q) \in \mathcal{E}_Z^t$  is elliptic on  $K$  if for any  $\rho = (z, \zeta, x_d) \in K$ , the principal symbol  $q_0(\rho)$  is an isomorphism from  $H_1$  onto  $H_2$ . In that case, there exist  $E \in \mathcal{E}_Z^t(H_2 \rightarrow H_1)$  with principal symbol  $e_0$  equal to  $q_0^{-1}$  near  $K$  such that  $E \circ Q \equiv Id_{H_1}$  and  $Q \circ E \equiv Id_{H_2}$  near  $K$ .

## A.2. An exotic calculus

Let  $\mathbb{T}_Y^d$  be the  $d$ -dimensional torus, and for  $s \in \mathbb{R}$ ,  $H^s$  the usual Sobolev space

$$H^s = \left\{ \sum_{\ell \in \mathbb{Z}^d} a_\ell e^{i\ell y}, \sum_{\ell} (1 + |\ell|^2)^s |a_\ell|^2 < \infty \right\}. \quad (\text{A.10})$$

For any operator  $A : \cap_s H^s \rightarrow \cup_s H^s$ , we denote by  $A_{\ell,k}$  the matrix coefficient

$$A_{\ell,k} = \oint_{\mathbb{T}^d} (Ae^{iky}).e^{-ily}. \tag{A.11}$$

For  $m \in \mathbb{R}$ , let  $\mathcal{A}^m$  be the following class of operators on the torus

$$\mathcal{A}^m = \left\{ A ; \forall N, \exists C_N \quad |A_{\ell,k}| \leq C_N \frac{(1 + |\ell|)^m}{(1 + |\ell - k|)^N} \quad \forall \ell, k \in \mathbb{Z}^d \right\}. \tag{A.12}$$

One has  $\mathcal{A}^m \circ \mathcal{A}^{m'} \subset \mathcal{A}^{m+m'}$ , and for  $A \in \mathcal{A}^m$ ,  $A$  is bounded from  $H^s$  to  $H^{s-m}$  for any  $s \in \mathbb{R}$ . The identity

$$[D_j, A]_{\ell,k} = (\ell_j - k_j)A_{\ell,k} \quad D_j = \frac{1}{i} \frac{\partial}{\partial y_j} \tag{A.13}$$

shows that  $\mathcal{A}^0$  is the class of bounded operators on  $L^2 = H^0$  such that all the commutators

$$[D_{j_1}, [D_{j_2}, \dots [D_{j_p}, A] \dots]] \tag{A.14}$$

are bounded on  $L^2$ . As a consequence, we get

**Lemma A.1.** *Let  $A \in \mathcal{A}^0$ , and  $\delta < (\|A; L^2 \rightarrow L^2\|)^{-1}$ . Then  $(Id + \delta A)^{-1} \in \mathcal{A}^0$ .*

*Proof.*  $B = (Id + \delta A)^{-1}$  is bounded on  $L^2$ , and all the commutators (A.14) for  $B$  can be expressed in terms of commutators for  $A$  by iteration of the formula

$$[D_j, B] = -B\delta[D_j, A]B.$$

□

Let  $Z = \mathbb{R}_z^p \times [0, r_0]$ .

We denote by  $\mathcal{S}_Z^{t,m}$  the vector space of functions  $A(\varepsilon, z, \zeta, x_d)$  defined for  $\varepsilon \in ]0, \varepsilon_0]$  smooth in  $(z, \zeta) \in T^*\mathbb{R}_z^p, x_d \in [0, r_0]$ , with values operators on the torus, which satisfy the estimates

$$\left\{ \begin{array}{l} \forall \alpha, \beta, \gamma, N, \exists C, \quad \forall \varepsilon, \ell, k, z, \zeta, x_d \\ |(1 + |\zeta|)^\gamma \partial_{z, x_d}^\alpha \partial_\zeta^\beta A_{\ell,k}(\varepsilon, z, \zeta, x_d)| \leq C \frac{(1 + |\ell|)^{m+|\alpha|}}{(1 + |\ell - k|)^N}. \end{array} \right. \tag{A.15}$$

In other words,  $A \in \mathcal{S}_Z^{t,m}$  means

$$\forall \alpha, \beta, \gamma \quad (1 + |\zeta|)^\gamma \partial_{z, x_d}^\alpha \partial_\zeta^\beta A \in \mathcal{A}^{m+|\alpha|}$$

uniformly in  $\varepsilon, z, \zeta, x_d$ .

Leibniz formula implies

$$\mathcal{S}_Z^{t,m} \circ \mathcal{S}_Z^{t,m'} \subset \mathcal{S}_Z^{t,m+m'}. \tag{A.16}$$

We denote by  $\mathcal{B}_Z^{t,m}$  the class of operators

$$\left\{ \begin{array}{l} Op(A) ; A \in \mathcal{S}_Z^{t,m} \\ Op(A)[f](z, x_d) = (2\pi\varepsilon)^{-p} \int e^{iz\zeta/\varepsilon} A(\varepsilon, z, \zeta, x_d)[\hat{f}_\varepsilon(\zeta, x_d)]d\zeta \end{array} \right. \tag{A.17}$$

where  $f \in L^2(Z; H^s)$  for some  $s$ ,  $Z = \mathbb{R}_z^p \times (0, r_0)_{x_d}$ , and  $\hat{f}_\varepsilon$  is the partial Fourier transform

$$\hat{f}_\varepsilon(\zeta, x_d) = \int e^{-iz\zeta/\varepsilon} f(z, x_d) dz \in L^2(\mathbb{R}_\zeta^p \times (0, r_0); H^s). \quad (\text{A.18})$$

**Lemma A.2.** *For any  $A \in \mathcal{S}_Z^{t,m}$ ,  $Op(A)$  is bounded from  $L^2(Z; H^s)$  in  $L^2(Z; H^{s-m})$  for any  $s$ , uniformly in  $\varepsilon \in ]0, \varepsilon_0]$ .*

*Proof.* To avoid the loose of derivative in  $z$ , we use the fact that  $A$  is a Schwartz function in  $\zeta$ , so we can write, by Fourier inversion formula

$$A(\varepsilon, z, \zeta, x_d) = (2\pi)^{-p} \int e^{i\zeta\theta} B(\varepsilon, z, \theta, x_d) d\theta \quad (\text{A.19})$$

with  $B \in \mathcal{S}_Z^{t,m}$ ; we obtain

$$Op(A)(f) = (2\pi)^{-p} \int B(\varepsilon, z, \theta, x_d) [f(z + \varepsilon\theta, x_d)] d\theta. \quad (\text{A.20})$$

The bounds (A.15) for  $B$  (with  $\alpha = \beta = 0$ ,  $|\gamma| = p + 1$ ) imply

$$\forall s, \exists C_s \sup_\varepsilon \|B(\varepsilon, \cdot, \theta, \cdot); L^2(Z; H^s) \rightarrow L^2(Z; H^{s-m})\| \leq \frac{C_s}{(1 + |\theta|)^{p+1}} \quad (\text{A.21})$$

and the lemma follows from (A.21) and

$$\|f(z + \varepsilon\theta, x_d)\| = \|f(z, x_d)\| \text{ in } L^2(Z; H^s).$$

□

The next lemma gives the principal part of the symbolic calculus

**Lemma A.3.** *For  $A_1 \in \mathcal{S}_Z^{t,m_1}$ ,  $A_2 \in \mathcal{S}_Z^{t,m_2}$ , one has*

$$\begin{cases} Op(A_1) \circ Op(A_2) = Op(B) \\ B = A_1 \circ A_2 + \varepsilon R \\ B \in \mathcal{S}_Z^{t,m_1+m_2}, R \in \mathcal{S}_Z^{t,m_1+m_2+1}. \end{cases} \quad (\text{A.22})$$

*Proof.* We have  $Op(A_1) \circ Op(A_2) = Op(B)$  with

$$B(\varepsilon, z, \zeta, x_d) = (2\pi)^{-p} \int \int e^{-it\eta} A_1(\varepsilon, z, \zeta + \varepsilon\eta, x_d) \circ A_2(\varepsilon, z + t, \zeta, x_d) d\eta dt. \quad (\text{A.23})$$

Using the Taylor formula  $f(\zeta + \varepsilon\eta) = f(\zeta) + \sum_j \varepsilon\eta_j \int_0^1 \frac{\partial f}{\partial \zeta_j}(\zeta + \varepsilon s\eta) ds$  and integrating by part with respect to  $t_j$ , we get  $B = A_1 \circ A_2 + \varepsilon R$  with

$$R = \frac{1}{i} \sum_j \int_0^1 ds \int \int (2\pi)^{-p} e^{-it\eta} \frac{\partial A_1}{\partial \zeta_j}(\varepsilon, z, \zeta + \varepsilon s\eta, x_d) \circ \frac{\partial A_2}{\partial z_j}(\varepsilon, z + t, \zeta, x_d) d\eta dt. \quad (\text{A.24})$$

We shall verify  $R \in \mathcal{S}_Z^{t,m_1+m_2+1}$  (the proof of  $B \in \mathcal{S}_Z^{t,m_1+m_2}$  is similar).

If we define  $B_2 = \frac{\partial A_2}{\partial z_j} \in \mathcal{S}_Z^{t,m_2+1}$  and  $B_1 \in \mathcal{S}_Z^{t,m_1}$  by

$$\frac{\partial A_1}{\partial \zeta_j}(\varepsilon, z, \zeta, x_d) = (2\pi)^{-p} \int e^{i\theta\zeta} B_1(\varepsilon, z, \theta, x_d) d\theta$$

we are reduce to prove

$$\int_0^1 ds \int e^{i\theta\zeta} B_1(\varepsilon, z, \theta, x_d) \circ B_2(\varepsilon, z + \varepsilon s\theta, \zeta, x_d) d\theta \in \mathcal{S}_Z^{t, m_1 + m_2 + 1}. \quad (\text{A.25})$$

The verification of (A.25) is now easy using (A.16), the Leibniz rule for derivatives and the fact that  $B_1$  (resp.  $B_2$ ) is in the Schwartz space with respect to  $\theta$  (resp.  $\zeta$ ).  $\square$

**Lemma A.4.** *Let  $\psi \in \mathcal{S}(\mathbb{R}^p)$  and  $A \in \mathcal{S}_Z^{t, m}$ . One has*

$$Op(A)^* \circ \psi\left(\frac{\varepsilon}{i}\partial z\right) = Op(A^* \psi(\zeta)) + \varepsilon Op(R) \quad (\text{A.26})$$

with  $R \in \mathcal{S}_Z^{t, m+1}$ .

*Proof.* We have  $Op(A)^* = Op(B)$  with

$$B(\varepsilon, z, \zeta, x_d) = (2\pi)^{-p} \int \int e^{-it\eta} A^*(\varepsilon, z + t, \zeta + \varepsilon\eta, x_d) d\eta dt. \quad (\text{A.27})$$

Using the Taylor formula as before, we get the identity (A.26) with

$$R = \frac{1}{i} \sum_j \int_0^1 ds \int (2\pi)^{-p} e^{-it\eta} \frac{\partial^2 A^*}{\partial z_j \partial \zeta_j}(\varepsilon, z + t, \zeta + \varepsilon s\eta, x_d) \psi(\zeta) dt d\eta. \quad (\text{A.28})$$

As in the proof of Lemma A.3, we just observe that, for  $B \in \mathcal{S}_Z^{t, m}$ , we have

$$\int_0^1 ds \int e^{i\theta\zeta} B(\varepsilon, z + \varepsilon s\theta, \theta, x_d) \psi(\zeta) d\theta \in \mathcal{S}_Z^{t, m}. \quad (\text{A.29})$$

$\square$

## B. APPENDIX

### B.1. Proof of Lemmas 3.4 and 3.5

Let  $pr_0^\ell(x)$  and  $pr^\ell(t, x, \tau, \xi')$  be the operators

$$pr_0^\ell(x) \left( \sum_k z_k e^{iky} \right) = \psi \left( \frac{\|\ell''_x\|^2}{\beta^2} \right) z_\ell e^{i\ell y} \quad (\text{B.1})$$

$$pr^\ell(t, x, \tau, \xi') = \psi \left( \frac{\|\ell''_x\|^2}{\beta^2} \right) \int_{\gamma_\ell(x)} (\lambda - \mathbb{E}(t, x, \tau, \xi'))^{-1} \frac{d\lambda}{2i\pi} \quad (\text{B.2})$$

and

$$\delta pr^\ell(t, x, \tau, \xi') = pr^\ell(t, x, \tau, \xi') - pr_0^\ell(x). \quad (\text{B.3})$$

Let  $z = (t, x)$ ,  $\zeta = (\tau, \xi')$ , and  $(\delta pr^\ell)_{j,k}$  be the matrix of  $\delta pr^\ell$  as in Appendix A.2. We shall prove:

For any  $\alpha, \beta, \gamma, N$ , there exist  $C_{\alpha, \beta, \gamma, N}$  such that

$$\begin{cases} (1 + |\zeta|)^\gamma |\partial_z^\alpha \partial_\zeta^\beta (\delta pr^\ell(z, \zeta))_{j,k}| \leq C_{\alpha, \beta, \gamma, N} \frac{(1+|\ell|)^\alpha}{(1+|\ell-j|+|\ell-k|)^N} \\ \text{for every } z, \zeta, j, k, \ell \end{cases} \quad (\text{B.4})$$

$$\delta pr^\ell(z, 0) = 0 \text{ for any } \ell, z. \quad (\text{B.5})$$

The two Lemmas 3.4 and 3.5 are consequences of the two properties (B.4, B.5), by definition of the class  $\mathcal{S}^{t,m}$  (see (A.15)). In fact (B.4) implies that the series ((3.45) and (3.48)) are convergent in the class  $\mathcal{S}^{t,0}$  and  $\mathcal{S}^{t,-\sigma}$ .

(B.5) is obvious since  $\mathbb{E}(z, 0) = \mathbb{E}_0(z)$  by (3.21) so  $pr^\ell(z, 0) = pr_0^\ell(z)$ . We have also  $\mathbb{E}(z, \zeta) = \mathbb{E}_0(z)$  for  $|\zeta| \geq \alpha_0$  since  $\mathbb{E}(z, \zeta) = \mathbb{E}_0(z) + \begin{pmatrix} 0 & 0 \\ q\theta & 0 \end{pmatrix}$  and  $\theta(z, \zeta) = 0$  for  $|\zeta| \geq \alpha_0$  so  $\delta pr^\ell(z, \zeta)$  is compactly supported in  $\zeta$  and we can forget  $\gamma$  and  $(1 + |\zeta|)^\gamma$  in the proof of (B.4). We first conjugate  $\mathbb{E}(t, x, \tau, \xi')$  by the multiplication by  $e^{i\ell y}$ . We get by (3.21, 3.32, 3.39)

$$e^{-i\ell y} \circ \mathbb{E}(z, \zeta) \circ e^{i\ell y} = \ell_x^\perp Id + \mathbb{E}^\ell(z, \zeta) \quad (\text{B.6})$$

$$\mathbb{E}^\ell(z, \zeta) = \mathbb{E}_0^\ell(z) + \begin{pmatrix} 0 & 0 \\ q^\ell \theta & 0 \end{pmatrix} \quad (\text{B.7})$$

$$\mathbb{E}_0^\ell(z) [\sum w_k e^{iky}] = \sum_k \begin{pmatrix} k_x^\perp & -\langle (k+\ell)_x'' \rangle \\ \frac{\| (k+\ell)_x'' \|^2}{\langle (k+\ell)_x'' \rangle} & k_x^\perp \end{pmatrix} (w_k) e^{iky} \quad (\text{B.8})$$

$$q^\ell(z, \zeta) = \left( \bigoplus_k \langle (k+\ell)_x'' \rangle^{-1} \right) \circ \left[ a_2(x, \xi') + \sum_{j=1}^{d-1} \frac{\partial a_2}{\partial \xi_j'}(x, \xi') (e_j(x) \cdot D_y + e_j(x) \cdot \ell) - \hat{\rho}(x, y) \tau^2 \right] \chi. \quad (\text{B.9})$$

We define  $\pi_0^\ell(z) = e^{-i\ell y} \circ pr_0^\ell(z) \circ e^{i\ell y}$ ,

$$\pi^\ell(z, \zeta) = e^{-i\ell y} \circ pr^\ell(z, \zeta) \circ e^{i\ell y}, \quad \delta \pi^\ell = \pi^\ell - \pi_0^\ell.$$

We have

$$\pi_0^\ell(z) \left[ \sum_k z_k e^{iky} \right] = \psi \left( \frac{\| \ell_x'' \|^2}{\beta^2} \right) z_0 \quad (\text{B.10})$$

$$\pi^\ell(z, \zeta) = \psi \left( \frac{\| \ell_x'' \|^2}{\beta^2} \right) \int_{|\lambda|=c_{1/4}} (\lambda - \mathbb{E}^\ell(z, \zeta))^{-1} \frac{d\lambda}{2i\pi}. \quad (\text{B.11})$$



We are now reduce to prove

$$\left\{ \begin{array}{l} \text{For any } \alpha, \beta, N, \text{ there exist } C_{\alpha, \beta, N} \text{ such that} \\ |\partial_z^\alpha \partial_\zeta^\beta (\delta \pi^\ell)_{j, k}(z, \zeta)| \leq C_{\alpha, \beta, N} \frac{(1 + |\ell|)^{|\alpha|}}{(1 + |j| + |k|)^N} \\ \text{for every } z, \zeta, j, k, \ell. \end{array} \right. \quad (\text{B.12})$$

Notice that the spectrum of  $\mathbb{E}_0^\ell(z)$ , with  $(\ell, x)$  such that  $\|\ell_x''\| \leq 2\beta \ll c_1$  can be separate in two pieces: two small eigenvalues  $\pm i\|\ell_x''\|$  with associated eigenspace  $\mathbb{C}^2 e^{i0y}$ , and the other part of the spectrum leaving outside the complex disk  $|\lambda| \geq c_1/2$ . The same is true for the spectrum  $\mathbb{E}^\ell(z, \zeta)$  (the cutt-off function  $\chi(z, \zeta)$  localize  $q^\ell$  in  $|\zeta| \leq 2\alpha_0 \ll \beta$ ). In order to prove (B.12), we use a Grushin method.

Let  $\widehat{\mathcal{A}}^m$  be the set of operators on  $(\mathcal{D}'(\mathbb{T}^d))^2 \oplus \mathbb{C}^2$  of the form

$$L = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \quad (\text{B.13})$$

with  $A$  a  $2 \times 2$  matrix with entries in  $\mathcal{A}^m$  (see Append. A.2)  $B$  a linear map from  $\mathbb{C}^2$  in  $(C^\infty(\mathbb{T}^d))^2$ ,  $C$  a continuous linear map from  $(\mathcal{D}'(\mathbb{T}^d))^2$  in  $\mathbb{C}^2$  and  $D \in \mathcal{M}^2(\mathbb{C})$ . As in Appendix A.2, we remark that  $\widehat{\mathcal{A}}^0$  is the class of bounded operators  $L$  on  $\mathcal{H} = (L^2(\mathbb{T}^d))^2 \oplus \mathbb{C}^2$  such that all the commutators  $[\widetilde{\mathcal{D}}_{j_1}, \dots, [\widetilde{\mathcal{D}}_{j_p}, L]]$  are bounded on  $\mathcal{H}$ , with

$$\widetilde{\mathcal{D}}_j = \begin{pmatrix} \frac{1}{i} \partial_{y_j} & 0 \\ 0 & 0 \end{pmatrix}.$$

In particular, Lemma A.1 remain valid. We denote by  $\widehat{\mathcal{S}}_V^m$  the vector space of functions of  $(z, \zeta) \in V$ ,  $V$  open

$$L(z, \zeta) = \begin{pmatrix} A(z, \zeta) & B(z, \zeta) \\ C(z, \zeta) & D(z, \zeta) \end{pmatrix} \quad (\text{B.14})$$

where  $B, C, D$  are as above and depends smoothly on  $(z, \zeta)$ , and  $A \in \mathcal{S}_V^m$ . In other words,  $L(z, \zeta) \in \widehat{\mathcal{S}}^m$  means

$$\forall \alpha, \beta \quad \partial_z^\alpha \partial_\zeta^\beta L \in \widehat{\mathcal{A}}^{m+|\alpha|} \text{ uniformly in } (z, \zeta) \in K \Subset V. \quad (\text{B.15})$$

Let  $j, p$  be the injection and projection

$$\left\{ \begin{array}{l} j(w) = w e^{i0y} \quad : \quad \mathbb{C}^2 \rightarrow C^\infty(\mathbb{T}^d)^2 \\ p(f) = \oint f dy \quad \mathcal{D}'(\mathbb{T}^2) \rightarrow \mathbb{C}^2 \end{array} \right. \quad (\text{B.16})$$

and

$$L^\ell(\lambda, z, \zeta) = \begin{pmatrix} \lambda - \mathbb{E}^\ell(z, \zeta) & j \\ p & 0 \end{pmatrix}. \quad (\text{B.17})$$

Then  $L^\ell(\lambda, \dots)$  is a holomorphic family in  $\lambda$  with values in  $\widehat{\mathcal{A}}_{V_\ell}^1$ , with inverse in  $\widehat{\mathcal{A}}_{V_\ell}^{-1}$  for  $|\lambda| < \frac{c_1}{2}$ , with  $V_\ell = \{(z, \zeta); \|\ell_x''\| < 2\beta\}$ .

Notice that in view of (B.8, B.9),  $\ell_x''$  can be replace by a small parameter in  $\mathbb{R}^{d-1}$  in both  $\mathbb{E}_0^\ell, q^\ell$ , so all the semi-norms of  $(L^\ell(\lambda, z, \zeta))^{-1}$  are uniform in  $(\ell, x)$  such that  $\|\ell_x''\| < 2\beta$ .

Let  $\mathcal{L}^\ell(\lambda, z, \zeta) = (L^\ell(\lambda, z, \zeta))^{-1}$

$$\mathcal{L}^\ell = \begin{pmatrix} A^\ell & B^\ell \\ C^\ell & D^\ell \end{pmatrix}. \quad (\text{B.18})$$

Then  $\lambda - \mathbb{E}^\ell(z, \zeta)$  is invertible iff  $\det(D^\ell(\lambda, z, \zeta)) \neq 0$ , and we have the algebraic identity

$$(\lambda - \mathbb{E}^\ell(z, \zeta))^{-1} = [A^\ell - B^\ell(D^\ell)^{-1}C^\ell](\lambda, z, \zeta). \tag{B.19}$$

The function  $A^\ell$  is holomorphic in  $\lambda \in \{|z| < \frac{\epsilon_0}{2}\}$  so we get by (B.11)

$$\pi^\ell(z, \zeta) = -\psi\left(\frac{\|\ell''_x\|^2}{\beta^2}\right) \int_{|\lambda|=\frac{\epsilon_0}{4}} (B^\ell(D^\ell)^{-1}C^\ell)(\lambda, z, \zeta) \frac{d\lambda}{2i\pi}. \tag{B.20}$$

This implies that the estimate (B.12) holds true for  $\pi^\ell$ , hence for  $\delta\pi^\ell$  ((B.12) is obvious for  $\pi_0^\ell$ ).

**B.2. Proof of Lemma 3.6**

One has  $[pr(t, x, \tau, \xi'), \mathbb{E}(t, x, \tau, \xi')] \equiv 0$ ,  $pr \in \mathcal{S}^{t,0}$ ,  $\mathbb{E} \in \mathcal{S}^{t,1}$ , so Lemma A.3 implies  $[Op(pr), Op(\mathbb{E})] \in \epsilon \mathcal{S}^{t,2}$ . In fact, the more precise estimate  $[Op(pr), Op(\mathbb{E})] \in \epsilon \mathcal{S}^{t,1}$  holds true. To see this, we just observe that we have  $\mathbb{E} - \mathbb{E}_0 \in \mathcal{S}^{t,0}$ ; from the definitions ((3.21) and (3.33)) we get  $\partial_\zeta \mathbb{E}_0 = 0$ ,  $\partial_z \mathbb{E}_0 \in \mathcal{S}^{t,1}$  and the result follows from the symbolic calculus formulas (A.23) and (A.24). We then deduce from Lemma 3.1 that  $F^{\epsilon,I} = Op(Id - pr)(F^\epsilon)$  satisfies the following equation

$$\frac{\epsilon}{i} \frac{\partial}{\partial x_d} F^{\epsilon,I} + Op(\mathbb{E})F^{\epsilon,I} = G^{\epsilon,I} \tag{B.21}$$

where,  $F^{\epsilon,I}$  and  $G^{\epsilon,I}$  are such that

$$\left\{ \begin{array}{l} \sup_\epsilon \|F^{\epsilon,I}; L^2(U; \mathcal{H}^{s_0})\| < +\infty \\ \sup_\epsilon \epsilon^{-1} \|G^{\epsilon,I}; L^2(U, \mathcal{H}^{s_0-1})\| < +\infty. \end{array} \right. \tag{B.22}$$

We shall first modified  $\mathbb{E}$  in (B.21) in order to work with an elliptic equation. Let us define  $\tilde{pr}_0(x)$  and  $\tilde{pr}(t, x, \tau, \xi')$  by formulas (3.44) and (3.45) with  $\psi(\frac{4\|\ell''_x\|^2}{\beta^2})$  instead of  $\psi(\frac{\|\ell''_x\|^2}{\beta^2})$ . One has

$$\tilde{pr} \circ (Id - pr)(t, x, \tau, \xi') \equiv 0 \tag{B.23}$$

and by the proof of Lemma 3.4 one gets

$$\tilde{pr} = \tilde{pr}_0 + \delta\tilde{pr}, \quad \delta\tilde{pr} \in \mathcal{M}_{2,2}(\mathcal{S}^{t,0}), \delta\tilde{pr}|_{\xi'=\tau=0} = 0. \tag{B.24}$$

Let  $\chi(\tau, \xi') \in C_0^\infty(|\tau| + |\xi'| < 2)$  equal to 1 near  $(|\tau| + |\xi'| \leq 1)$ ,  $\chi_{\alpha_0}(\tau, \xi') = \chi((\tau, \xi')/\alpha_0)$  and let  $K(x)$  be the operator on the torus

$$K(x)(\Sigma_\ell z_\ell e^{i\ell y}) = \sum_\ell \psi\left(\frac{16\|\ell''_x\|^2}{\beta^2}\right) \frac{z_\ell}{\langle \ell'_x \rangle} e^{i\ell y}. \tag{B.25}$$

Let us define  $\tilde{\mathbb{E}}$  by the formula

$$\left\{ \begin{array}{l} \tilde{\mathbb{E}} = \tilde{\mathbb{E}}_0 + \delta\tilde{\mathbb{E}} \\ \tilde{\mathbb{E}}_0 = \mathbb{E}_0 + \begin{pmatrix} 0 & 0 \\ K & 0 \end{pmatrix} \\ \delta\tilde{\mathbb{E}} = \left[ \begin{pmatrix} 0 & 0 \\ q\theta & 0 \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ K & 0 \end{pmatrix} \circ \delta\tilde{pr} \right] \chi_{\alpha_0}. \end{array} \right. \tag{B.26}$$

One has  $\theta\chi_{\alpha_0} \equiv \theta$  and  $\begin{pmatrix} 0 & 0 \\ K & 0 \end{pmatrix} \circ \tilde{p}r_0 = \begin{pmatrix} 0 & 0 \\ K & 0 \end{pmatrix}$ , which implies

$$\tilde{\mathbb{E}} = \mathbb{E} + \begin{pmatrix} 0 & 0 \\ K & 0 \end{pmatrix} \circ \tilde{p}r\chi_{\alpha_0} + \begin{pmatrix} 0 & 0 \\ K & 0 \end{pmatrix} (1 - \chi_{\alpha_0}). \quad (\text{B.27})$$

One has  $(1 - \chi_{\alpha_0})F^{\varepsilon, I} \in \varepsilon L^2(U; \mathcal{H}^{s_0-1})$  so using (B.21, B.23) and Lemma A.3, one gets

$$\begin{cases} \left[ \frac{\varepsilon}{i} \frac{\partial}{\partial x_d} + Op(\tilde{\mathbb{E}}) \right] F^{\varepsilon, I} = \tilde{G}^{\varepsilon, I} \\ \sup_{\varepsilon} \varepsilon^{-1} \|\tilde{G}^{\varepsilon, I}; L^2(U, \mathcal{H}^{s_0-1})\| < +\infty. \end{cases} \quad (\text{B.28})$$

Notice that  $\tilde{\mathbb{E}}_0$  is a diagonal operator

$$\begin{cases} \tilde{\mathbb{E}}_0(\Sigma z_{\ell} e^{i\ell y}) = \Sigma_{\ell} \tilde{\mathbb{E}}_{0, \ell}(z_{\ell}) e^{i\ell y} \\ \tilde{\mathbb{E}}_{0, \ell} = \begin{pmatrix} \ell_x^{\perp} & -\langle \ell''_x \rangle \\ \frac{\|\ell''_x\|^2 + \psi\left(\frac{16\|\ell''_x\|^2}{\beta^2}\right)}{\langle \ell''_x \rangle} & \ell_x^{\perp} \end{pmatrix}. \end{cases} \quad (\text{B.29})$$

The eigenvalues of  $\tilde{\mathbb{E}}_{0, \ell}$  are

$$\tilde{\lambda}_{\ell}^{\pm} = \ell_x^{\perp} \pm i \left( \|\ell''_x\|^2 + \psi\left(\frac{16\|\ell''_x\|^2}{\beta^2}\right) \right)^{1/2}. \quad (\text{B.30})$$

In particular, one has with  $0 < c_1 < c_2$

$$c_1 \langle \ell''_x \rangle \leq |\text{Im} \tilde{\lambda}_{\ell}^{\pm}| \leq c_2 \langle \ell''_x \rangle. \quad (\text{B.31})$$

We choose the associated eigenvectors

$$e_{\ell}^{\pm}(x) = \begin{bmatrix} \frac{-\langle \ell''_x \rangle}{\tilde{\lambda}_{\ell}^{\pm}(x) - \ell_x^{\perp}} \\ 1 \end{bmatrix}. \quad (\text{B.32})$$

The map  $J_0(x)$  defined by

$$J_0(x) \left( \sum_{\ell} \begin{pmatrix} z_{\ell}^+ \\ z_{\ell}^- \end{pmatrix} e^{i\ell y} \right) = \sum_{\ell} (z_{\ell}^+ e_{\ell}^+(x) + z_{\ell}^- e_{\ell}^-(x)) e^{i\ell y} \quad (\text{B.33})$$

is then an isomorphism of  $\mathcal{H}^s$  for any  $s$ .

Let  $D^{\pm}$  be the operators

$$D^{\pm}(\Sigma z_{\ell} e^{i\ell y}) = \Sigma \tilde{\lambda}_{\ell}^{\pm} z_{\ell} e^{i\ell y}. \quad (\text{B.34})$$

By construction, one has

$$J_0^{-1} \tilde{\mathbb{E}}_0 J_0 = \begin{pmatrix} D^+ & 0 \\ 0 & D^- \end{pmatrix}. \quad (\text{B.35})$$

**Lemma B.1.** *If  $\alpha_0$  is small enough, there exist  $\delta B, \delta C, \delta D^+, \delta D^-$  in  $\mathcal{S}^{t,0}$ , with support in  $\{|\tau| + |\xi'| \leq 2\alpha_0\}$  vanishing on  $\xi' = 0, \tau = 0$  such that the following identity holds true*

$$\begin{pmatrix} Id & \delta B \\ \delta C & Id \end{pmatrix}^{-1} J_0^{-1} \tilde{\mathbb{E}} J_0 \begin{pmatrix} Id & \delta B \\ \delta C & Id \end{pmatrix} = \begin{pmatrix} D^+ + \delta D^+ & 0 \\ 0 & D^- + \delta D^- \end{pmatrix}. \quad (\text{B.36})$$

*Proof.* By formulas (B.26, B.35), one has

$$J_0^{-1} \tilde{\mathbb{E}} J_0 = \begin{pmatrix} D^+ & 0 \\ 0 & D^- \end{pmatrix} + \delta M, \quad \delta M = \begin{pmatrix} \delta M_1 & \delta M_2 \\ \delta M_3 & \delta M_4 \end{pmatrix}$$

where  $\delta M_j \in \mathcal{S}^{t,0}$ , vanishes on  $\xi' = 0, \tau = 0$ , and has support in  $\{|\tau| + |\xi'| \leq 2\alpha_0\}$ . Equation (B.36) is then equivalent to the following system of equations

$$\begin{cases} \delta M_1 + \delta M_2 \delta C = \delta D_+ \\ \delta M_4 + \delta M_3 \delta B = \delta D_- \\ \delta M_2 + \delta M_1 \delta B + D^+ \delta B = \delta B D^- + \delta B \delta D^- \\ \delta M_3 + \delta M_4 \delta C + D^- \delta C = \delta C D^+ + \delta C \delta D^+. \end{cases} \quad (\text{B.37})$$

We are thus reduce to solve the equation, with unknown  $\delta B \in \mathcal{S}^{t,0}$

$$\begin{cases} D^+ \delta B - \delta B D^- + \delta M_1 \delta B + \delta M_2 - \delta B \delta M_4 - \delta B \delta M_3 \delta B = \phi(\delta B, \delta M) \\ \phi(\delta B, \delta M) = 0. \end{cases} \quad (\text{B.38})$$

Let  $\mathcal{E}_x, \mathcal{E}$  be the Banach space of operators on the torus:  $(A_{\ell,k} = \oint_{\mathbb{T}^d} A e^{iky} e^{-i\ell y})$

$$\mathcal{E} = \left\{ A; \|A; \mathcal{E}\| = \sup_{\ell,k} |A_{\ell,k}| (1 + |\ell - k|)^{N_0} < +\infty \right\} \quad (\text{B.39})$$

$$\mathcal{E}_x = \left\{ A; \|A; \mathcal{E}_x\| = \sup_{\ell,k} |A_{\ell,k}| (1 + |\ell - k|)^{N_0} |\tilde{\lambda}_\ell^+(x) - \tilde{\lambda}_k^-(x)| < +\infty \right\} \quad (\text{B.40})$$

where  $N_0$  is given,  $N_0 \geq d + 1$ . By (B.30, B.31), the injection  $\mathcal{E}_x \hookrightarrow \mathcal{E}$  is continuous, and the map  $(A_1, A_2) \rightarrow A_1 A_2$  is continuous on  $\mathcal{E}$  by the choice of  $N_0$ .

We shall first verify that (B.38) has a unique small solution  $\delta B \in \mathcal{E}_x$ , for  $(t, \tau, x, \xi')$  fixed, if  $\alpha_0$  is small enough. By construction, one has

$$(D^+ \delta B - \delta B D^-)_{\ell,k} = (\tilde{\lambda}_\ell^+ - \tilde{\lambda}_k^-) (\delta B)_{\ell,k} \quad (\text{B.41})$$

so  $\delta B \mapsto D^+ \delta B - \delta B D^-$  is an isomorphism of  $\mathcal{E}_x$  onto  $\mathcal{E}$ . The map  $(\delta B, \delta M) \rightarrow \phi(\delta B, \delta M)$  is differentiable from  $\mathcal{E}_x \times (\mathcal{E})^4$  to  $\mathcal{E}$  and satisfies

$$\frac{\partial}{\partial \delta B} \phi(0,0) = D^+(\cdot) - (\cdot)D^- \quad \phi(0,0) = 0. \quad (\text{B.42})$$

By the implicit function theorem, the equation  $\phi(\delta B, \delta M) = 0$  has thus a unique small solution  $\delta B \in \mathcal{E}_x$ , provide  $\|\delta M; \mathcal{E}\|$  is small. Using (B.26) ( $q$  and  $\delta \tilde{r}$  vanish on  $\xi' = \tau = 0$ ) and  $\chi_{\alpha_0}(\tau, \xi') = \chi(\frac{\tau, \xi'}{\alpha_0})$  one gets the estimate  $\|\delta M_j; \mathcal{E}\| \leq C^{te} \alpha_0$ . This shows the existence of  $\delta B$  solution of (B.38).

It remains to prove that for any fixed  $z = (t, x), \zeta = (\tau, \xi)$ , we have

$$\forall N, \exists C_N \quad |(\delta B)_{\ell,k}(z, \zeta)| \leq \frac{C_N}{(1 + |\ell - k|)^N} \tag{B.43}$$

and that the functions  $(z, \zeta) \mapsto (\delta B)_{\ell,k}(z, \zeta)$  are smooth and satisfy

$$\forall \alpha, \beta, \gamma, N, \exists C \quad \forall z, \zeta, \ell, k \tag{B.44}$$

$$|(1 + |\zeta|)^\gamma \partial_z^\alpha \partial_\zeta^\beta (\delta B)_{\ell,k}(z, \zeta)| \leq C \frac{(1 + |\ell|)^{|\alpha|}}{(1 + |\ell - k|)^N}.$$

Let  $\nabla_i = \frac{1}{i} \frac{\partial}{\partial y_i}$  a derivation on the torus. The commutator  $[\nabla_i, \delta B]$  satisfies the linear equation

$$\begin{cases} \mathcal{L}([\nabla_i, \delta B]) \in \mathcal{E} \\ \mathcal{L}(u) = D^+u - uD^- + \delta M_1 u - u\delta M_4 - u\delta M_3 \delta B - \delta B \delta M u. \end{cases} \tag{B.45}$$

The linear map  $\mathcal{L}$  is an isomorphism of  $\mathcal{E}_x$  onto  $\mathcal{E}$  provide  $\|\delta M_j; \mathcal{E}\|$  (hence  $\|\delta B; \mathcal{E}\|$ ) is small enough; decreasing  $\alpha_0$  if necessary, we find that (B.45) admits a unique solution  $[\nabla_i, \delta B] \in \mathcal{E}_x \hookrightarrow \mathcal{E}$ . By iteration, all the commutators  $[\nabla_{i_1}, [\nabla_{i_2}, \dots [\nabla_{i_k}, \delta B] \dots]]$  belongs to  $\mathcal{E}_x$  so (B.43) holds true. By construction, the functions  $(\delta B)_{\ell,k}(z, \zeta)$  are smooth and compactly supported in  $\{|\zeta| \leq 2\alpha_0\}$ . For any  $m \geq 0$ , let  $\mathcal{A}^m, \mathcal{A}_x^m$  be the vector spaces

$$\mathcal{A}^m = \left\{ A; \forall N, \exists C_N |A_{\ell,k}| \leq C_N \frac{(1 + |\ell|)^m}{(1 + |\ell - k|)^N} \right\} \tag{B.46}$$

$$\mathcal{A}_x^m = \left\{ A; \forall N, \exists C_N |A_{\ell,k}| \leq \frac{C_N}{|\tilde{\lambda}_\ell^+(x) - \tilde{\lambda}_k^-(x)|} \frac{(1 + |\ell|)^m}{(1 + |\ell - k|)^N} \right\}. \tag{B.47}$$

In order to prove (B.44), we differentiate (B.38) with respect to  $(z, \zeta)$  and we are reduce to verify that the following assertion holds true

$$\begin{cases} \text{There exist } \beta > 0 \text{ such that for } \delta M_j, \delta B \in \mathcal{A}^0, \\ \text{with } \sum_j \|\delta M_j; \mathcal{E}\| + \|\delta B; \mathcal{E}\| \leq \beta, \text{ the map } u \mapsto \mathcal{L}(u) \\ \text{is an isomorphism of } \mathcal{A}_x^m \text{ onto } \mathcal{A}^m \text{ for any } m \geq -1. \end{cases} \tag{B.48}$$

(Here we use the fact that  $A \mapsto (\partial_x^\alpha D^+)A - A(\partial_x^\alpha D^-)$  maps  $\mathcal{A}_x^m$  into  $\mathcal{A}^{m+|\alpha|}$  for any  $\alpha$ : it is a consequence of the estimates  $|\tilde{\lambda}_\ell^+(x) - \tilde{\lambda}_k^-(x)| \geq C^{te}(\langle \ell''_x \rangle + \langle k''_x \rangle)$  and for  $|\alpha| \geq 1 \quad |\partial_x^\alpha \tilde{\lambda}_\ell^\pm(x)| \leq C_\alpha (1 + |\ell|)^{|\alpha|} \langle \ell''_x \rangle$ .)

Let us first verify that (B.48) holds true for  $m = 0$ . We remark that  $\mathcal{A}^0$  (resp.  $\mathcal{A}_x^0$ ) is the set of operators  $A \in \mathcal{E}$  (resp.  $\mathcal{E}_x$ ) such that all the commutators  $[\nabla_{i_1}, [\nabla_{i_1}, [\nabla_{i_2}, \dots, [\nabla_{i_p}, A]]]$  belongs to  $\mathcal{E}$  (resp.  $\mathcal{E}_x$ ). For  $\beta$  small  $\mathcal{L}$  is an isomorphism between  $\mathcal{E}_x$  and  $\mathcal{E}$  and for  $u \in \mathcal{E}_x, v \in \mathcal{E}$  such that  $\mathcal{L}(u) - v = 0$ , one has  $\mathcal{L}([\nabla_i, u]) - [\nabla_i, v] \in \mathcal{E}$ . Therefore (B.48) holds true for  $m = 0$ , and by the same argument for  $m = -1$ . We now fixe  $\beta$  and we proceed by induction on  $m \geq 1$ ; let us assume that (B.48) holds true for  $-1 \leq m' \leq m - 1$ . Let  $\Lambda$  be the operator on the torus  $\Lambda(\Sigma z_\ell e^{i\ell y}) = \Sigma(1 + |\ell|)z_\ell e^{i\ell y}$ . We have  $\mathcal{L}(u) = D^+u - uD^- + pu + uq$  with  $p, q \in \mathcal{A}^0$ , so  $[\Lambda, p] \in \mathcal{A}^0$ ; from  $[\Lambda, D^\pm] = 0$ , we get  $\mathcal{L}(\Lambda w) - \Lambda \mathcal{L}(w) = [p, \Lambda]w$ . Let  $J : \mathcal{A}^m \rightarrow \mathcal{A}_x^m$  the map

$$J(v) = \Lambda \mathcal{L}^{-1}(\Lambda^{-1}v) + \mathcal{L}^{-1}([\Lambda, p]\mathcal{L}^{-1}(\Lambda^{-1}v)) \tag{B.49}$$

where  $\mathcal{L}^{-1} : \mathcal{A}^{m-1} \rightarrow \mathcal{A}_x^{m-1}$  is the inverse map of  $\mathcal{L}$ . We have  $\mathcal{L} \circ J(v) \equiv v$ , and it remains to show that  $u \in \mathcal{A}_x^m$ , and  $\mathcal{L}(u) = 0$  imply  $u = 0$  : we have  $[\Lambda^{-1}, p] \in \mathcal{A}^{-2}$  so  $\mathcal{L}(u) = 0 \Rightarrow \mathcal{L}(\Lambda^{-1}u) = [p, \Lambda^{-1}]u \in \mathcal{A}^{m-2} \Rightarrow \Lambda^{-1}u \in \mathcal{A}_x^{m-2} \Rightarrow u \in \mathcal{A}_x^{m-1}$  and we get  $u = 0$ .  $\square$

**Lemma B.2.** *Let  $U_0 = \{z \in \mathbb{R}^p, |z| \leq r_0\}$ , and  $U = U_0 \times [0, r_1]$  with  $r_0, r_1 > 0$ . For any  $\ell \in \mathbb{Z}^d$ , let  $\lambda_\ell(z, x_d) \in C^0(U; \mathbb{C})$  be given continuous functions such that*

$$\begin{aligned} & \exists c_0 > 0, \exists c_1 > 1, \forall \ell, \forall z, x_d \\ & \operatorname{Im} \lambda_\ell(z, x_d) \geq c_0 \text{ and } \frac{|\lambda_\ell(z, x_d)|}{1 + |\ell|} \in \left[ \frac{1}{c_1}, c_1 \right]. \end{aligned} \quad (\text{B.50})$$

Let  $D(z, x_d)$  be the operator on the torus

$$D(z, x_d)[\Sigma u_\ell e^{i\ell y}] = \Sigma \lambda_\ell(z, x_d) u_\ell e^{i\ell y}. \quad (\text{B.51})$$

Let  $\sigma \in \mathbb{R}$  be given, and for  $\varepsilon \in ]0, 1]$ ,  $B_\varepsilon(x_d)$  a family of bounded operator on  $E^\sigma = L^2(U_0, H^\sigma(\mathbb{T}^d))$  such that

$$\begin{cases} i) \quad \forall f \in E^\sigma, \forall \varepsilon \quad x_d \mapsto B_\varepsilon(x_d)[f] \text{ is a continuous function} \\ \quad \text{of } x_d \in [0, r_1] \text{ with values in } E^\sigma \\ ii) \quad \exists \delta, \forall \varepsilon, \forall x_d \quad \|B_\varepsilon(x_d); E^\sigma \rightarrow E^\sigma\| \leq \delta. \end{cases} \quad (\text{B.52})$$

Then, for  $\delta < c_0$  the Cauchy problem

$$\begin{cases} \left[ \frac{\varepsilon}{i} \frac{d}{dx_d} - (D + B_\varepsilon) \right] (u^\varepsilon(x_d)) = 0 \quad x_d \in ]0, r_1[ \\ u^\varepsilon(0) = u_0 \in E^\sigma \end{cases} \quad (\text{B.53})$$

admits a solution  $u^\varepsilon \in C^0([0, r_1], E^\sigma) \cap C^1([0, r_1], E^{\sigma-1})$  such that

$$\|u^\varepsilon(x_d), E^\sigma\| \leq \|u_0, E^\sigma\| e^{-(c_0 - \delta)x_d/\varepsilon}. \quad (\text{B.54})$$

*Proof.* We first observe that the assumption (B.50) implies that  $D$  maps  $C^0([0, r_1], E^\sigma)$  onto  $C^0[0, r_1], E^{\sigma-1}$  for any  $\sigma$ . We have  $\|v; E^\sigma\| = \|(1 + |D_y|^2)^{\sigma/2} v; E^0\|$  and  $[D, (1 + |D_y|^2)^{\sigma/2}] = 0$ , so if one replace  $B_\varepsilon$  by  $(1 + |D_y|^2)^{\sigma/2} B_\varepsilon (1 + |D_y|^2)^{-\sigma/2}$ , we are reduce to the case  $\sigma = 0$ . For any  $L$ , let  $\pi_L$  be the orthogonal projector  $\pi_L(\Sigma u_\ell e^{i\ell y}) = \sum_{|\ell| \leq L} u_\ell e^{i\ell y}$ . The equation

$$\begin{cases} \left( \frac{\varepsilon}{i} \frac{d}{dx_d} - \pi_L(D + B_\varepsilon)\pi_L \right) u_L^\varepsilon(x_d) = 0 \\ u_L^\varepsilon(0) = \pi_L(u_0) \end{cases} \quad (\text{B.55})$$

is an ordinary differential equation in the Hilbert space  $= L^2(U_0, \bigoplus_{|\ell| \leq L} \mathbb{C} e^{i\ell y}) = E_L \hookrightarrow E^s$ , so admit a unique solution  $u_L^\varepsilon \in C^1([0, r_1], E_L)$ . It satisfies the identity,

$$\frac{d}{dx_d} \|u_L^\varepsilon\|^2 = 2\operatorname{Re} \left( \frac{i}{\varepsilon} D u_L^\varepsilon |u_L^\varepsilon \right) + 2\operatorname{Re} (i/\varepsilon \pi_L B_\varepsilon \pi_L u_L^\varepsilon |u_L^\varepsilon) \quad (\text{B.56})$$

so we get using (B.50) and (B.52)  $\frac{d}{dx_d} \|u_L^\varepsilon\|^2 \leq \frac{-2}{\varepsilon} (c_0 - \delta) \|u_L^\varepsilon\|^2$ , which implies

$$\|u_L^\varepsilon(x_d), E^0\| \leq \|u_0, E^0\| e^{-(c_0 - \delta) \frac{x_d}{\varepsilon}}. \quad (\text{B.57})$$

Therefore  $u_L^\varepsilon$  is bounded in  $L^2([0, r_1], E^\sigma) \cap H^1([0, r_1], E^{\sigma-1}) = F$  for fixed  $\varepsilon$  so we can extract a subsequence  $u_{L_k}^\varepsilon$  so that  $u_{L_k}^\varepsilon \xrightarrow{\text{weak}} u^\varepsilon$  in  $F$  and  $u^\varepsilon$  satisfies (B.53). In particular we have  $\frac{\varepsilon}{i} \frac{d}{dx_d} u^\varepsilon - D u^\varepsilon = B_\varepsilon u^\varepsilon \in L^2([0, r_1], E^\sigma)$  so  $u^\varepsilon \in C^0([0, r_1], E^\sigma) \cap C^1([0, r_1], E^{\sigma-1})$ . The estimate (B.54) is then a consequence of (B.57).  $\square$

We can now achieve the verification of Lemma 3.6. We choose a tangential scalar *o.p.d.*  $Q_2$  equal to  $Id$  near the support of  $Q_1$  and with essential support closed to  $\rho_0$ , and we define  $T = Q_2 \delta \tilde{\mathbb{E}} Q_2$ . Then  $Op(T)(x_d)$  acts on  $L^2(U_0, H^\sigma)$ ,  $\forall \sigma$ . Using (B.28), we still have

$$\left[ \frac{\varepsilon}{i} \frac{\partial}{\partial x_d} + Op(\tilde{\mathbb{E}}_0 + T) \right] F^{\varepsilon, I} \in \varepsilon L^2(U, \mathcal{H}^{s_0-1}). \quad (\text{B.58})$$

We then apply Lemma B.1 to  $\tilde{\mathbb{E}}_0^* + T^*$  instead of  $\tilde{\mathbb{E}}$ ; let  $I_0(x)$  be the map

$$I_0(x) \left( \Sigma \begin{pmatrix} z_\ell^+ \\ z_\ell^- \end{pmatrix} e^{i\ell y} \right) = \Sigma \frac{1}{\ell} \left[ z_\ell^+ \begin{pmatrix} -\tilde{\lambda}_\ell^+(x) + \ell_x^\perp \\ \langle \ell_x'' \rangle \\ 1 \end{pmatrix} + z_\ell^- \begin{pmatrix} -\tilde{\lambda}_\ell^-(x) + \ell_x^\perp \\ \langle \ell_x'' \rangle \\ 1 \end{pmatrix} \right] e^{i\ell y}. \quad (\text{B.59})$$

We get the existence of  $\delta B, \delta C, \delta D^+, \delta D^-$  in  $\mathcal{S}^{t,0}$  such that, with

$$I = I_0 \begin{pmatrix} Id & \delta B \\ \delta C & Id \end{pmatrix}.$$

One has

$$I^{-1}(\tilde{\mathbb{E}}_0^* + T^*)I = \begin{pmatrix} D^+ + \delta D^+ & 0 \\ 0 & (D^- + \delta D^-) \end{pmatrix}. \quad (\text{B.60})$$

Moreover, by the proof of Lemma B.1, and the fact that  $\lim_{\alpha_0 \rightarrow 0} \|(|\xi'| + |\tau|)\chi_{\alpha_0}(\xi', \tau)\|_{L^\infty} = 0$ , we may suppose that the norm of the tangential operators  $\delta D^\pm(x_d)$  acting on  $L^2(U_0, H^{|\text{s}_0|+1})$  is as small as we want. Taking in account the lower bound (B.31)  $-Im\lambda_\ell^-(x) \geq c_1 \langle \ell_x'' \rangle \geq c_1$ , we can apply Lemma B.2. For every  $h \in L^2(U_0, H^{|\text{s}_0|+1})$  we get  $v^\varepsilon \in L^2(U_0 \times [0, r_1], H^{|\text{s}_0|+1})$  such that

$$\begin{cases} \left[ \frac{\varepsilon}{i} \frac{\partial}{\partial x_d} + Op(D^- + \delta D^-) \right] v^\varepsilon = 0 \\ v^\varepsilon|_{x_d=0} = \varepsilon^{-1/2} h \\ \sup_\varepsilon \|v^\varepsilon, L^2(U_0 \times [0, r_1]; H^{|\text{s}_0|+1})\| \leq C^{te} \|h; L^2(U_0; H^{|\text{s}_0|+1})\|. \end{cases} \quad (\text{B.61})$$

We put  $\underline{v}^\varepsilon = \begin{bmatrix} 0 \\ v^\varepsilon \end{bmatrix}$ .

We choose  $\theta(x_d) \in C_0^\infty([-r_1, r_1])$  equal to 1 near zero. We denote by  $\langle | \rangle$  the duality between  $L^2(V_0, H^\sigma)$  and  $L^2(V_0, H^{-\sigma})$ . We have by (B.58)

$$\int_0^\infty \left\langle \left( \frac{\varepsilon}{i} \partial_{x_d} + Op(\tilde{\mathbb{E}}_0 + T) \right) F^{\varepsilon, I} | \theta(x_d) Op(I) \underline{v}^\varepsilon \right\rangle \in 0(\varepsilon) \|h\|. \quad (\text{B.62})$$

We integrate by part, taking into account Lemma A.4 and  $\|\theta'(x_d) Op(I) \underline{v}^\varepsilon; L^2(U, \mathcal{H}^{|\text{s}_0|+1})\| \leq C^{te} \|h\|$ , we get

$$\frac{\varepsilon}{i} \langle F^{\varepsilon, I} |_{x_d=0} | Op(I) \underline{v}^\varepsilon |_{x_d=0} \rangle = \int_0^\infty \left\langle \theta(x_d) F^{\varepsilon, I} | \left( \frac{\varepsilon}{i} \partial_{x_d} + Op(\tilde{\mathbb{E}}_0^* + T^*) \right) Op(I) \underline{v}^\varepsilon \right\rangle dx_d + 0(\varepsilon \|h\|). \quad (\text{B.63})$$

We have  $\|\varepsilon [\frac{\partial}{\partial x_d} Op(I)] \underline{v}^\varepsilon; L^2(U, \mathcal{H}^{|\text{s}_0|})\| \leq C^{te} \varepsilon \|h\|$ , and the estimates

$$\begin{cases} \left| \partial_x^\alpha \left[ \frac{\tilde{\lambda}_\ell^\pm(x) - \ell_x^\perp}{\langle \ell_x'' \rangle} \right] \right| \leq C_\alpha (1 + |\ell|)^{|\alpha|} \quad \forall \alpha \\ |\partial_x^\alpha \tilde{\mathbb{E}}_{0, \ell}| \leq C_\alpha (1 + |\ell|)^{|\alpha|} \quad \forall \alpha, |\alpha| \geq 1 \end{cases} \quad (\text{B.64})$$

implies

$$\left\{ \begin{array}{l} \left\| \left[ Op(\tilde{\mathbb{E}}_0^* + T^*) \circ Op(I) - Op((\tilde{\mathbb{E}}_0^* + T^*)I) \right] \underline{v}^\varepsilon; L^2(U, \mathcal{H}^{|s_0|}) \right\| \leq C^{te} \varepsilon \|h\| \\ \left\| \left[ Op\left(I \begin{pmatrix} D^+ + \delta D^+ & 0 \\ 0 & D^- + \delta D^- \end{pmatrix}\right) - Op(I) \circ Op\left(\begin{pmatrix} D^+ + \delta D^+ & 0 \\ 0 & D^- + \delta D^- \end{pmatrix}\right) \right] \underline{v}^\varepsilon; L^2(U, \mathcal{H}^{|s_0|}) \right\| \leq C^{te} \varepsilon \|h\|. \end{array} \right. \quad (\text{B.65})$$

From (B.61, B.63, B.65), we get

$$\left| \left\langle F^{\varepsilon, I} \Big|_{x_d=0} \Big| Op(I) \begin{bmatrix} 0 \\ h \end{bmatrix} \right\rangle \right| \leq C^{te} \varepsilon^{1/2} \|h\|. \quad (\text{B.66})$$

If one use the definition of  $I$ , the fact that  $h$  is arbitrary in  $L^2(U_0, H^{|s_0|+1})$ , one get for some  $D \in \mathcal{S}^0$

$$\|Op(Id + \delta B)^* Tr_1(F^{\varepsilon, I}) - Op(D) Tr_0(F^{\varepsilon, I}); L^2(U_0, H^{s_0-1})\| \leq C \varepsilon^{1/2}. \quad (\text{B.67})$$

Lemma 3.6 is then a consequence of the estimates (3.52), Lemmas A.2–A.4, and the fact for  $\alpha_0$  small,  $Id + (\delta B)^*$  is invertible in  $\mathcal{S}^0$ .  $\square$

## REFERENCES

- [1] M. Avellaneda, C. Bardos and J. Rauch, Contrôlabilité exacte, homogénéisation et localisation d'ondes dans un milieu non-homogène. *Asymptot. Anal.* **5** (1992) 481-484.
- [2] G. Allaire and C. Conca, Bloch wave homogenization and spectral asymptotic analysis. *J. Math. Pures Appl.* **77** (1998) 153-208.
- [3] N. Burq and G. Lebeau, Mesures de défaut de compacité; applications au système de Lamé, preprint.
- [4] C. Bardos, G. Lebeau and J. Rauch, Sharp sufficient conditions for the observation, control and stabilization of waves from the boundary. *SIAM J. Control Optim.* **30** (1992) 1024-1075.
- [5] C. Castro, Boundary controllability of the one dimensional wave equation with rapidly oscillating density, preprint.
- [6] C. Castro and E. Zuazua, Contrôle de l'équation des ondes à densité rapidement oscillante à une dimension d'espace. *C. R. Acad. Sci. Paris* **324** (1997) 1237-1242.
- [7] P. Gérard, *Mesures semi-classiques et ondes de Bloch*, Séminaire X EDP, exposé 16 (1991).
- [8] P. Gérard and E. Leichtnam, Ergodic properties of eigenfunctions for the Dirichlet problem. *Duke Math. J.* **71** (1993) 559-607.
- [9] G. Lebeau, Contrôle de l'équation de Schrödinger. *J. Math. Pures Appl.* **71** (1993) 267-291.
- [10] G. Lebeau, *Équation des ondes amorties*, Algebraic and Geometric Methods in Mathematical Physics, A. Boutet de Monvel and V. Marchenko, Eds. Kluwer Academic Publishers (1996) 73-109.
- [11] R. Melrose and J. Sjöstrand, Singularities of boundary value problems I, II. *Comm. Pure Appl. Math.* **31** (1978) 593-617; **35** (1982) 129-168.