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# FUNCTIONAL EQUATION SATISFIED BY CERTAIN L-FUNCTIONS 

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## Introduction

The purpose of this paper is to establish the existence of a functional equation for certain $L$-functions of degree 4 attached to cusp forms on $P G L_{2}$ over a number field.

More precisely, let ${ }^{\circ} \pi$ be an irreducible (admissible) constituent of the space of cusp forms on $P G L_{2}$ over a number field (same definition as for $G L_{2}$, cf. [8]). We can write ${ }^{\circ} \pi$ as a tensor product of the local representations ${ }^{\circ} \pi_{v}$ which are class-one for almost all $v$.

Suppose ${ }^{\circ} \pi_{v}$ is a class-one representation. Let $\rho$ be a four dimensional irreducible representation of $S L_{2}(\mathrm{C})$, the corresponding associated group, and let $\alpha_{v}$ be the semi-simple conjugacy class in $S L_{2}(\mathbb{C})$ determined by ${ }^{\circ} \pi_{v}$ (cf. [10], also see §5). Then in [10], Langlands defines a local $L$-function attached to $\rho,{ }^{\circ} \pi_{v}$ and a complex number $s$ as follows:

$$
L\left(s, \rho,{ }^{\circ} \pi_{v}\right)=\operatorname{det}\left(I-\rho\left(\alpha_{v}^{\hat{v}}\right) q_{v}^{-s}\right) .
$$

Clearly $L\left(s, \rho,{ }^{\circ} \pi_{v}\right)$ depends only upon the classes of $\rho$ and ${ }^{\circ} \pi_{v}$.
Now, let $S$ be the finite set of places for which the corresponding representations are ramified. We put

$$
L_{S}\left(s, \rho,{ }^{\circ} \pi\right)=\prod_{v \notin S} L\left(s, \rho,{ }^{\circ} \pi_{v}\right)
$$

We also define a global coefficient $\gamma_{S}\left(s, \rho,{ }^{\circ} \pi\right)$ (see §5) which is closely related to the global root number $\epsilon\left(s, \rho,{ }^{\circ} \pi\right)$ introduced in [10]. The

[^0]main result of this paper is the functional equation
$$
L_{S}\left(1-s, \tilde{\rho},{ }^{\circ} \pi\right)=\gamma_{S}\left(s, \rho,{ }^{\circ} \pi\right) L_{S}\left(s, \rho,{ }^{\circ} \pi\right)
$$
where $\tilde{\rho}$ denotes the contragradient representation of $\rho$ (Theorem 5.9). This is in fact the functional equation for the $L$-function attached to $\rho$ and ${ }^{\circ} \pi$ which has been conjectured in the context of an arbitrary reductive group by Langlands [10]. This result together with a result of Langlands [11] concerning the meromorphicity of these $L$-functions, will prove the first conjecture in [10] for any cusp form on $P G L_{2}$, when $\rho$ is as above. We extend this result to any cusp form on $G L_{2}$ for which the center acts according to an unramified quasicharacter of the corresponding idele class group (Corollary 5.10).

We should mention that the same result will essentially follow from the works of P. Deligne, S. Gelbart, H. Jacquet and J.A. Shalika [3, 9, 16] (see §2.3 of Gelbart's paper 'Automorphic Forms and Artin's Conjecture' which is based on a talk given at the 1976 Bonn Conference on Modular Forms). However the approach is entirely different from the one explained in this paper.

Except for certain technical difficulties (particularly at infinite places), the same methods can be used to establish certain functional equations for the adjoint groups of type $A$, following the nonvanishing of some of their Fourier coefficients [15].

Interwoven with the proof of the functional equation is a description of the unique non-degenerate quotient of a certain induced representation (cf. [14]) of a simple algebraic group of type $G_{2}$ over a $p$-adic field (see Appendix).

In this paper, we have used the principle of applying the Eisenstein series to $L$-functions which is due to Langlands [11] (see also: R. Godement, Formes automorphes et produits eulérien; d'après R.P. Langlands; Séminaire Bourbaki, no. 349, 1968). For this reason we shall recall the important facts of the theory of Eisenstein series in §2.

In §3, we shall study certain properties of the Whittaker models for the algebraic groups of type $G_{2}$. Needless to say these results are true and can be proved by the same methods for an arbitrary split group over a $p$-adic field. This is based on certain unpublished results of W . Casselman and J.A. Shalika.

Most of the local computations are carried out in §4. In loose terms, we shall compute 'the local coefficients' which are basic to the definition of $\gamma_{S}\left(s, \rho,{ }^{\circ} \pi\right)$, using a key lemma (lemma 4.4) and certain results of Jacquet [7]. Notice that these local coefficients differ from the Langlands' root numbers by a ratio of local $L$-functions. Again I
should mention that these computations can be done for an arbitrary Chevalley group and it has been only for the sake of simplicity that we have limited ourselves to the case of $G_{2}$.

It is in §5 that we prove the functional equation (Theorem 5.9). We observe that the local coefficients are in fact defined by the local multiplicity one theorem of J.A. Shalika [15], together with a result of F. Rodier [14].

The problem of finding the poles of these $L$-functions still remains open.

I would like to express my gratitude to Professor Robert Langlands for his suggestion of the problem and for many helpful discussions during my stay at the Institute for Advanced Study, where the bulk of this paper was prepared.

I would like to thank Professor Joseph Shalika for useful discussions and suggestions during the last year.

## 1. Notation and terminology

Let $G$ be a simple algebraic group of type $G_{2}$ and let $\mathfrak{g}$ denote its Lie algebra. We assume that $G$ splits over $Q$. We fix a Cartan subgroup $T$ of $G$ with Lie algebra $\mathfrak{h}$. We use $B$ to denote a fixed Borel subgroup containing $T$. Put $U$ for its unipotent radical.

Let $\Psi$ denote the set of roots of $\mathfrak{g}$ with respect to $\mathfrak{h}$. We use $\Delta, \Psi^{+}$ and $\Psi^{-}$for simple, positive, and negative roots, respectively. Then $\Delta$ consists of two elements, $\alpha$, the short root, and $\beta$, the long one. Other positive roots are $\alpha+\beta, 2 \alpha+\beta, 3 \alpha+\beta$, and $3 \alpha+2 \beta$. We have:

$$
\mathfrak{g}=\mathfrak{h} \oplus \oplus_{\gamma \in \Psi} \mathfrak{g}_{\gamma}
$$

with root spaces $\mathfrak{g}_{\gamma}, \gamma \in \Psi$.
Let $W$ be the Weyl group of $G$. $W$ is generated by the reflections $\sigma_{\gamma}, \gamma \in \Phi$. We shall identify each $\sigma_{\gamma}$ with an element of $N_{G}(T)$ through the isomorphism between $W$ and $N_{G}(T) / C_{G}(T)$. We denote this element by $w_{r}$.

Let $F$ be a number field. For every place $v$ of $F$, we shall write $G_{v}$ for the group of $F_{v}$-rational points of $G$. We use $G_{F}$ for the group of $F$-rational points.

For each $v$, we fix a maximal compact subgroup $K_{v}$ of $G_{v}$ relative to $T_{v}$, so that

$$
G_{v}=K_{v} \cdot B_{v}
$$

To define $K_{v}$ for the finite places, we fix a Chevalley lattice $M$ for $\mathfrak{g}$ (cf. [1]). If $\mathbb{A}$ is the ring of adèles of $F$, we write $G_{A}$ for the corresponding adèlized group. We use the same index for the corresponding subgroups.

By a character of a group we shall understand a homomorphism from the group into the complex numbers of absolute value one.

We call a character $\chi$ of $U_{\boldsymbol{A}} / U_{F}$ non-degenerate if its restriction to every non-trivial subgroup $U_{A}^{W}$,

$$
U^{w}=w U w^{-1} \cap U,
$$

is nontrivial, $w \in W$. Clearly

$$
\chi=\prod_{\gamma \in \Delta} \chi_{\gamma}
$$

where each $\chi_{\gamma}$ is a non-trivial character of $U_{\wedge}^{\gamma} / U^{\gamma}$. Here $U^{\gamma}$ denotes the connected subgroup whose Lie algebra is $\mathfrak{g}_{\gamma}$, the root group for $\gamma$. Then $\chi_{y}$ can be considered as a non-trivial character of $\mathbb{A} / F$ and

$$
\chi_{\gamma}=\prod_{v} \chi_{\gamma, v}
$$

with each $\chi_{\gamma, v}$ a non-trivial character of $F_{v}$. Furthermore for almost all $v$, the largest ideal for which $\chi_{\gamma, v}$ is trivial is the ring of integers $O_{v}$ of $F_{v}$. Therefore we can write

$$
\chi=\prod_{v} \chi_{v}
$$

with

$$
\chi_{v}=\prod_{\gamma \in \Delta} \chi_{\gamma, v}
$$

a non-degenerate character of $U_{v}$.
Throughout this paper, we shall fix a non-degenerate character $\chi$ of $U$. Later in §5, we shall put certain conditions on $\chi$.

We use $G^{\wedge}$ to denote the associated complex group for $G$ which is defined in general by Langlands [10]. Then $G^{\wedge}$ is a complex group of type $G_{2}$.

Let $T^{\wedge}$ be a Cartan subgroup of $G^{\wedge}$ and let $L^{\wedge}$ be the root lattice of
$G^{\wedge}$ with respect to $T^{\wedge}$. We may identify $L^{\wedge}$ with the Z-lattice generated by

$$
\left\{H_{\gamma} \mid \gamma \in \Delta\right\}
$$

in $\mathfrak{h}_{\mathbf{R}}$. Here $H_{\gamma}$ is defined by

$$
\gamma^{\wedge}(H)=\kappa\left(H, H_{\gamma}\right) \quad(\forall H \in \mathfrak{h})
$$

where $\boldsymbol{\kappa}$ denotes the Killing form on $\mathfrak{g}$ and

$$
\gamma^{\wedge}=\frac{2 \gamma}{(\gamma, \gamma)}
$$

In fact we identify the simple roots of $G^{\wedge}$ relative to $T^{\wedge}$ with $H_{\gamma}$, $\gamma \in \Delta$.

Let $\mathfrak{h}^{\wedge}$ be the Lie algebra of $T^{\wedge}$; then $\operatorname{Hom}_{z}\left(L^{\wedge}, \mathbb{Z}\right)$ can be identified with a $\mathbb{Z}$-lattice in $\mathfrak{h}^{\wedge}$. Thus we may identify $\mathfrak{b}^{\wedge}$ with $\operatorname{Hom}_{z}\left(L^{\wedge}, \mathbb{C}\right)$, which itself is isomorphic to $\operatorname{Hom}_{c}\left(L^{\wedge} \otimes_{z} \mathbb{C}, \mathbb{C}\right)$. Since $L^{\wedge} \otimes_{z} \mathbb{C}$ is equal to $\mathfrak{b}_{c}$, the complexification of $h \otimes_{Q} \mathbb{R}$, we conclude that $\mathfrak{h}_{c}$ and $\mathfrak{b}^{\wedge}$ are dual to each other.

Finally, let $G$ be a split group defined over a local field, and let $\pi$ be an irreducible admissible representation of $G$ (or corresponding algebra) on a complex vector space $V$. Fix a Borel subgroup $B$ of $G$ and denote its unipotent radical by $U$. Let $\chi$ be a character of $U$. By a Whittaker functional on $V$, we shall mean a continuous linear functional on $V$ satisfying

$$
\lambda(\pi(u) v)=\overline{\chi(u)} \lambda(v)
$$

for all $u$ in $U$ and $v$ in $V$. Then from [15] it follows that the space of such linear functionals is at most one-dimensional. If there is such a functional, we shall say that $\pi$ is non-degenerate. Suppose $\pi$ is non-degenerate. For each $v$ in $V$, we define a complex function $w_{v}$ on $G$ by:

$$
w_{v}(g)=\lambda\left(\pi\left(g^{-1}\right) v\right) .
$$

The space of all such functions is called the Whittaker model of $\pi$. We denote this space by $W(\pi)$. The elements of $W(\pi)$ will be called the Whittaker functions of $\pi$.

## 2. Eisenstein series and Fourier coefficients

Let $P$ be a maximal parabolic subgroup of $G$ containing $B$. We put $M$ for a fixed Levi factor of $P$ and we write:

$$
P=M N,
$$

where $N$ is the unipotent radical of $P$ with Lie algebra $\mathfrak{n}$. We shall identify $M$ with the quotient $P / N$. We assume that $\mathfrak{n}$ is generated by the root spaces $\mathfrak{g}_{\alpha}, \mathfrak{g}_{\alpha+\beta}, \mathfrak{g}_{2 \alpha+\beta}, \mathfrak{g}_{3 \alpha+\beta}$, and $\mathfrak{g}_{3 \alpha+2 \beta}$. We also assume that $M$ contains $T$. Let $A$ be the center of $M$. As in [11] we put

$$
{ }^{\circ} G=M / A .
$$

Then ${ }^{\circ} G$ is the adjoint group of a split Lie algebra ${ }^{\circ} \mathfrak{g}$ of type $A_{1}$. More precisely ${ }^{\circ} G$ is isomorphic to $P S L_{2}$.

Let ${ }^{\circ} T$ be the image of $T$ in ${ }^{\circ} G$. Put

$$
{ }^{\circ} U=U \cap M .
$$

We may consider ${ }^{\circ} U$ as unipotent radical of a Borel subgroup of ${ }^{\circ} G$. Then the characters of ${ }^{\circ} U$ are the restriction of those of $U$.

As usual we use the index $v$ for the $F_{v}$-rational points of each of the groups mentioned here. We put

$$
{ }^{\circ} K_{v}=\overline{P_{v} \cup K_{v}},
$$

where $\overline{P_{v} \cap K_{v}}$ denotes the image of $P_{v} \cap K_{v}$ under the natural projection,

$$
{ }^{\circ} \boldsymbol{K}=\prod_{v}{ }^{\circ} \boldsymbol{K}_{v},
$$

and

$$
K=\prod_{v} K_{v}
$$

From now on, we shall identify ${ }^{\circ} G$ with the group $P S L_{2}$.
Let $\varphi^{*}$ be a cusp form on $G L_{2}(\mathbb{A})$ as defined in [8]. More precisely $\varphi^{*}$ is a continuous function on $G L_{2}(\mathbb{A}) / G L_{2}(F)$ which under the action of the global Hecke algebra generates an irreducible (admissible) constituent of $A_{0}(\omega)$, the subspace of the cusp forms for which the center of $G L_{2}(\mathbb{A})$ acts according to the quasi-character of $\omega$ of $\mathbb{A}^{*} / F^{*}$. We shall assume that $\omega$ is unramified.
$\varphi^{*}$ is assumed to be slowly increasing at infinite places (cf. [8]).
Put

$$
\varphi(g)=\varphi^{*}(g) \omega^{-1 / 2}(\operatorname{det} g) \quad g \in G L_{2}(\mathbb{A}) .
$$

Then $\varphi$ is a cusp form on $P G L_{2}(\mathbb{A})$ and by restriction on $P S L_{2}(\mathbb{A})$. In our notation, $\varphi$ is a function in $\mathbb{A}_{0}(\mathbb{1})$, where $\mathbb{1}$ denotes the trivial character of the group of the ideles of $\mathbb{A}$. In fact, every cusp form on $P G L_{2}(\mathbb{A})$ is a constituent of $A_{0}(1)$.

For every place $v$, we shall identify ${ }^{\circ} K_{v}$ with the maximal compact subgroup of $P S L_{2}\left(F_{v}\right)$ induced from the standard maximal compact subgroup of $G L_{2}\left(F_{v}\right)$. We also identify ${ }^{\circ} U$ with the unipotent radical of the standard parabolic subgroup of $G L_{2}$.

The restriction of the character $\chi$ of $U_{A} / U_{F}$ to ${ }^{\circ} U_{A} /{ }^{\circ} U_{F}$ is a non-degenerate character of ${ }^{\circ} U_{A} /{ }^{\circ} U_{F}$ which we still denote by $\chi$. We consider

$$
\begin{equation*}
{ }^{\circ} w(g)=\int_{{ }^{\circ} U_{A} / U_{F}} \varphi(g u) \overline{\chi(u)} d u \tag{2.1}
\end{equation*}
$$

with $g$ in ${ }^{\circ} G_{A}$. In §5, we shall use ${ }^{\circ} \chi$ to denote the restriction of $\chi$ to ${ }^{\circ} U_{\mathrm{A}} /{ }^{\circ} \boldsymbol{U}_{\mathrm{F}}$.

We shall assume that

$$
{ }^{\circ} w={\underset{v}{ }}^{\circ} w_{v},
$$

where for every $v,{ }^{\circ} w_{v}$ denotes a local Whittaker function which for almost all $v$ is a class-one function, i.e.

$$
{ }^{\circ} w_{v}\left(k_{v}\right)=1 \quad k_{v} \in{ }^{\circ} K_{v} .
$$

Then for $g=\left(g_{v}\right)$ in ${ }^{\circ} G_{\wedge}$, we have

$$
{ }^{\circ} w(g)=\prod_{v}{ }^{\circ} w_{v}\left(g_{v}\right)
$$

where almost all the factors are equal to 1 . The function ${ }^{\circ} w$ is in fact defined for $g$ in $\mathrm{GL}_{2}(\mathbb{A})$.

Let $\sigma_{v}$ be the finite dimensional representation of ${ }^{\circ} K_{v}$ on the span of ${ }^{\circ} \boldsymbol{w}_{v}$ by the elements of ${ }^{\circ} K_{v}$.

The representation $\sigma_{v}$ is in fact the restriction of the same represen-
tation when ${ }^{\circ} w_{v}$ is considered as a function on $G L_{2}\left(F_{v}\right) . \sigma_{v}$ can be extended to a representation of $P_{v} \cap K_{v}$ trivial on $A_{v}$ and $N_{v}$.

For every place $v$, we shall fix a finite dimensional representation of $K_{v}$ whose restriction to $P_{v} \cap K_{v}$ contains $\sigma_{v}$. When $\sigma_{v}$ is the (one dimensional) trivial representation of ${ }^{\circ} K_{v}$, we take this representation also to be trivial. We denote this representation by $\tilde{\boldsymbol{\sigma}}_{v}$. We put

$$
\sigma=\bigotimes_{v} \sigma_{v}
$$

and

$$
\tilde{\boldsymbol{\sigma}}=\bigotimes_{v} \tilde{\boldsymbol{\sigma}}_{v}
$$

which are representations of ${ }^{\circ} K$ and $K$, respectively.
Let $P_{\sigma}$ denote the projection onto the space of $\sigma$. Clearly

$$
P_{\sigma}=\bigotimes_{v} P_{\sigma_{v}} .
$$

We define the following well-defined operator valued function (projection)

$$
\begin{equation*}
\tilde{\varphi}(m)=\int_{\sigma_{K}} \varphi(k \bar{m}) \tilde{\sigma}\left(k^{-1}\right) d k \cdot p_{\sigma} \tag{2.2}
\end{equation*}
$$

for $m$ in $M_{A}$. Here $d k$ denotes the Haar measure on ${ }^{\circ} K$ which is a product of local measures $d k_{v}$ for which

$$
\int_{\bullet_{K_{v}}} d k_{v}=1
$$

In the terminology of [5], this is a $\tilde{\sigma}$-function on $M_{A}$. It can be extended to a $\tilde{\sigma}$-function on $G_{A}$ as follows:

$$
\tilde{\varphi}(g)=\tilde{\sigma}(k) \tilde{\varphi}(m)
$$

where

$$
g=k m n
$$

with $k$ in $K, m$ in $M_{\wedge}$, and $n$ in $N_{A}$.
Let $\delta_{P, v}$ denote the modulus character of $M$ with respect tont. More precisely, at each place $v$

$$
\delta_{P, v}(m n)=\left|\operatorname{det} A d_{n}(m)\right|_{v .} .
$$

We define the global $\delta_{P}$ by

$$
\delta_{P}(m)=\prod_{v} \delta_{P, v}\left(m_{v}\right)
$$

for $m=\left(m_{v}\right)$ in $M_{A}$.
For a complex number $s$, we define

$$
\begin{equation*}
\Phi_{s}(g)=\delta_{P}^{s-1 / 2}(b) \tilde{\varphi}(g) \tag{2.3}
\end{equation*}
$$

for

$$
g=k b
$$

with $k$ in $K$ and $b$ in $B_{\text {a }}$. Then $\Phi_{s}$ is a 3 -finite function (cf. [5]), where 3 is the center of the universal enveloping algebra of

$$
G_{\infty}=\prod_{v=\infty} G_{v} .
$$

Now we define the Eisenstein series attached to $\varphi$ as follows:

$$
\begin{equation*}
E(s ; \tilde{\varphi} ; g ; P)=\sum_{\gamma \in G_{F} P_{F}} \Phi_{s}(g \gamma) . \tag{2.4}
\end{equation*}
$$

Here $g$ is in $G_{A}$.
The series converges absolutely for $\operatorname{Re}(s)<-\frac{1}{2}$ and defines a function which is holomorphic in $s$ (whenever it converges absolutely). As a function on $G_{A}$, it is a $\tilde{\sigma}$-function. Its restriction to $G_{\infty}$ is smooth and 3 -finite (cf. [5] and [12]). Clearly (3.4) is a right $G_{F}$-invariant function.

As a function of $s, E(s ; \tilde{\varphi} ; g ; P)$ can be continued to a meromorphic function on the whole complex plane.

We use $W_{M}$ to denote $N_{M}(T) / C_{M}(T)$. Then $W_{M}$ can be considered as a subgroup of $W$. It is equal to the Weyl group for ${ }^{\circ} G$. We shall assume that each $w_{\gamma}$ (see §1) has been chosen to lie in $G_{Z} \cap K_{\infty}$, where

$$
K_{\infty}=\prod_{v=\infty} K_{v} .
$$

We use $w_{0}$ to denote $w_{2 \alpha+\beta}$. Modulo $W_{M}, w_{0}$ is the longest element in $W$.

For $\operatorname{Re}(s)<-\frac{1}{2}$, the integral

$$
\begin{equation*}
\int_{N_{A}} \Phi_{s}\left(g n w_{0}\right) d n \tag{2.5}
\end{equation*}
$$

converges absolutely and can be continued to a meromorphic function (in $s$ ) on the whole complex plane (cf. [5]).

Now, we define the function $M(s) \tilde{\varphi}$ on $M_{A}$ by

$$
\begin{equation*}
\delta_{P}^{-s-1 / 2}(m)(M(s) \tilde{\varphi})(m)=\int_{N_{\Lambda}} \Phi_{s}\left(m n w_{0}\right) d n \tag{2.6}
\end{equation*}
$$

for $m$ in $M_{A}$. The function $M(s) \tilde{\varphi}$ is in fact a cuspidal function on ${ }^{\circ} G_{A}$. It is a $\tilde{\sigma}$-function and ${ }^{\circ} 3$-finite $\left({ }^{\circ} 0\right.$ denotes the center of the universal enveloping algebra of ${ }^{\circ} G_{\infty}$ ).

The functional equation for (2.4) can be written as follows (cf. [5] and [12]):

$$
\begin{equation*}
E(-s ; M(s) \tilde{\varphi} ; g ; P)=E(s ; \tilde{\varphi} ; g ; P) \tag{2.7}
\end{equation*}
$$

where $g$ is in $G_{A}$.

Remark: If we denote the function defined by the integral in (2.5) by $\tilde{\Psi}_{s}$, then the left hand side of (2.7) is equal to

$$
\begin{equation*}
\sum_{\gamma \in G_{F} P_{F}} \tilde{\Psi}_{s}(g \gamma) . \tag{2.8}
\end{equation*}
$$

We need the following notations.
For $g$ in $G_{\wedge}$ (resp. $G_{v}$ ), we write (Iwasawa decomposition)

$$
g=k(g) b(g)
$$

with $k(g)$ in $K$ (resp. $\left.K_{v}\right)$ and $b(g)$ in $B_{A}\left(\right.$ resp. $\left.B_{v}\right)$. Also for $b$ in $B_{A}$ (resp. $B_{v}$ ), we write

$$
b=t(b) u(b)
$$

with $t(b)$ in $T_{A}\left(\operatorname{resp} . T_{v}\right)$ and $u(b)$ in $U_{A}\left(\right.$ resp. $\left.U_{v}\right)$. Hence we can write

$$
g=k(g) t(g) u(g)
$$

where

$$
t(g)=t(b(g))
$$

and

$$
u(g)=u(b(g))
$$

We also write

$$
\begin{aligned}
g & =k(g) p(g) \\
& =k(g) m(g) n(g)
\end{aligned}
$$

with $k(g)$ in $K\left(\operatorname{resp}, K_{v}\right), p(g)$ in $P_{A}\left(\right.$ resp. $\left.P_{v}\right), m(g)$ in $M_{A}$ (resp. $\left.M_{v}\right)$, and $n(g)$ in $N_{A}\left(\right.$ resp. $\left.N_{v}\right)$.

The convergent integral

$$
\begin{equation*}
\int_{U_{\Lambda} / U_{F}} E(s ; \tilde{\varphi} ; g u ; P) \overline{\chi(u)} d u \tag{2.9}
\end{equation*}
$$

is called a Fourier coefficient of $E(s ; \tilde{\varphi} ; g ; P)$. We denote (2.9) by $\underline{E}_{x}(s ; \tilde{\varphi} ; g ; P)$. Then

$$
\underline{E}_{x}(s ; \tilde{\varphi} ; g u ; P)=\chi(u) \underline{E}_{x}(s ; \tilde{\varphi} ; g ; P)
$$

for $u$ in $U_{\mathrm{A}}$.

Lemma 2.1: For $\operatorname{Re}(s)<-\frac{1}{2}, \underline{E_{X}}(s ; \tilde{\varphi} ; g ; P)$ is equal to

$$
\begin{equation*}
\int_{u \in^{\circ} U^{\prime}{ }^{\circ} U_{F}} \int_{u^{\prime} \in N_{\Lambda}} \Phi_{s}\left(g u^{\prime} w_{0} u\right) \overline{\chi\left(u u^{\prime}\right)} d u^{\prime} d u . \tag{2.10}
\end{equation*}
$$

Proof: For $\operatorname{Re}(s)<-\frac{1}{2}$, we have
(2.10.1) $\int_{U_{\wedge} / U_{F}} E(s ; \varphi ; g u ; P) \overline{\chi(u)} d u=\int_{u \in U_{A} / U_{F}} \sum_{\gamma \in G_{F} P_{F}} \Phi_{s}(g u \gamma) \overline{\chi(u)} d u$.

Then (2.10.1) is equal to

$$
\begin{gather*}
\int_{U_{A} / U_{F}} \sum_{\gamma \in U_{F}\left|G_{F}\right| P_{F}} \sum_{\delta \in U_{F} / \gamma P_{F \gamma} \gamma^{-1} \cap U_{F}} \Phi_{s}(g u \delta \gamma) \overline{\chi(u)} d u  \tag{2.10.2}\\
\quad=\sum_{\gamma \in U_{F}\left|G_{F}\right| P_{F}} \int_{U_{A} / \gamma P_{F \gamma^{-1} \cap U_{F}}} \Phi_{s}(g u \gamma) \overline{\chi(u)} d u .
\end{gather*}
$$

Using the Bruhat decomposition in (2.10.2) we have

$$
\begin{equation*}
\sum_{w \in w / w_{M}} \int_{U_{A} / w P_{F w^{-1} \cap U_{F}}} \Phi_{s}(g u w) \overline{\chi(u)} d u . \tag{2.10.3}
\end{equation*}
$$

In fact, $w$ has been chosen to lie in $G_{Z} \cap K_{\infty}$. We put

$$
V=w P w^{-1} \cap U
$$

and

$$
V_{w}=V \cap w N w^{-1}
$$

Then

$$
V_{w}=U \cap w N w^{-1} .
$$

Also let $N^{V}$ be the quotient of $U$ by $V$.
A single term in (2.10.3) is equal to

$$
\int_{u^{\prime} \in N_{\AA}^{v}} \int_{u \in V_{\Lambda} V_{F}} \Phi_{s}\left(g u^{\prime} u w\right) \overline{\chi(u) \chi\left(u^{\prime}\right)} d u d u^{\prime}
$$

for some $w$. The integral over $V_{\mathrm{A}} / V_{F}$ factors through the integral

$$
\int_{\left(V_{w}\right)_{A}\left(V_{w}\right)_{F}} \overline{x(u)} d u
$$

which vanishes if $V_{w}$ is non-trivial. Hence there is only one non-zero term which corresponds to the class of $w_{0}$ modulo $W_{M}$. This completes the proof of the lemma.

## 3. Whittaker models

The purpose of this section is to study the Whittaker models of certain classes of induced representations. As we shall see later, these are crucial in the definition of the corresponding local root numbers.

Let ${ }^{\circ} \pi$ be the irreducible (admissible) representation of the global Hecke algebra of $G L_{2}$ on the space generated by $\varphi$. Then ${ }^{\circ} \pi$ can be written as

$$
{ }^{\circ} \pi=\bigotimes_{v}^{\circ} \pi_{v}
$$

where the tensor product is defined as in [8]. Let ${ }^{\circ} V_{v}$ be the space of ${ }^{\circ} \pi_{v}$. For almost all $v$, the representation ${ }^{\circ} \pi_{v}$ is a class one representation. For every $v$, let $W\left({ }^{\circ} \pi_{v}\right)$ be the Whittaker model of ${ }^{\circ} \pi_{v}$.

The same is true for ${ }^{\circ} \pi^{*}$, the representation generated by $\varphi^{*}$, in particular,

$$
{ }^{\circ} \pi^{*}=\bigotimes_{v}{ }^{\circ} \pi_{v}^{*}
$$

In this section we shall assume that $v$ is finite.
${ }^{\circ} \pi_{v}$ can be considered as a representation of $P_{v}$ (trivial on $A_{v}$ and $N_{v}$ ). We shall study the Whittaker model of the following induced representation

$$
\Pi_{\mathrm{v}}=\operatorname{ind}_{P_{v} \uparrow G_{v}}^{\circ} \pi_{v} \otimes \delta_{P, v}^{s}
$$

The space $V_{v}$ of this representation consists of all the left $K_{v}$-finite functions $f_{v}$ from $G_{v}$ into the space $W\left({ }^{\circ} \pi_{v}\right)$ which satisfy

$$
f_{v}(g p)={ }^{\circ} \pi_{v}\left(p^{-1}\right) \delta_{P, v}^{s-1 / 2}(p) f_{v}(g) \quad\left(p \in P_{v}\right)
$$

The representation $\Pi_{v}$ is given by left inverse translations.
We use $\left(f_{v}(g), m\right)$ to denote the value of the Whittaker function $f_{v}(g)$ at a point $m$ in $M_{v}$.

We need the following lemma.

Lemma 3.1: Let $N$ be a unipotent group over a non-archimedean field. Then there exists an increasing sequence $\left\{N_{i}\right\}$ of open compact subgroups of $N$ which exhausts $N$.

Proof: $N$ can be imbedded in the subgroup of $n \times n$ upper triangular matrices (some $n>0$ ) for which the lemma holds.

Let $\left\{N_{v, i}\right\}$ be a filtration of $N_{v}$ as in Lemma 3.1. The following proposition is essentially due to Casselman-Shalika [2], and I am indebted to H . Jacquet for mentioning it to me.

Proposition 3.2: Given $f_{v}$ in $V_{v}$, there exists an integer $i\left(f_{v}\right)$ so that the integral

$$
\begin{equation*}
\int_{N_{v, i}}\left(f_{v}\left(n w_{0}\right), e\right) \overline{\chi_{v}(n)} d n \tag{3.1}
\end{equation*}
$$

has a value $\lambda\left(f_{v}\right)$ independent of $i$ if $i \geq i\left(f_{v}\right)$.
Proof: Let $V_{\chi}\left(U_{v}\right)$ denote the subspace generated by all the functions in $V_{v}$ of the form

$$
\Pi_{v}(u) f_{v}-\overline{\chi_{v}(u)} f_{v}
$$

with $u$ in $U_{v}$ and $f_{v}$ in $V_{v}$. Then every $f_{v}$ in $V_{v}$ can be written as

$$
f_{v}=f_{1, v}+f_{2, v}
$$

where $f_{1, v}$ is in $V_{\chi}\left(U_{v}\right)$ and $f_{2, v}$ has support in $N_{v} w_{0} P_{v}$ (cf. [2], see also $\S 2$ in [9]). We may assume that

$$
f_{1, v}=\Pi_{v}\left(u_{0}\right) f_{0, v}-\overline{\chi_{v}\left(u_{0}\right)} f_{0, v}
$$

where $u_{0}$ is in $U_{v}$ and $f_{0, v}$ is in $V_{v}$. Then for the large values of $i$, the integral vanishes for $f_{1, v}$ and converges for $f_{2, v}$ since $f_{2, v}$ has compact support modulo $P$. This completes the proof of the proposition.

Corollary 3.3: The linear functional $\lambda$ defined by

$$
\begin{equation*}
\lambda\left(f_{v}\right)=\int_{N_{v}}\left(f_{v}\left(n w_{0}\right), e\right) \overline{\chi_{v}(n)} d n \tag{3.2}
\end{equation*}
$$

(in the sense of Proposition 3.2) is a Whittaker functional for the space of $\Pi_{v}$.

Corollary 3.4: As a function of $s$, the integral (3.2) is entire.
Corollary 3.5: There exists a function $f_{v}$ in $V_{v}$ for which $\lambda\left(f_{v}\right)$ is non-zero.

Let $\left({ }^{\circ} \tilde{\pi}_{v},{ }^{\circ} \tilde{V}_{v}\right)$ denote the contragradient representation of $\left({ }^{\circ} \pi_{v},{ }^{\circ} V_{v}\right)$. We use

$$
\langle\langle,\rangle\rangle
$$

to denote the pairing on ${ }^{\circ} V_{v} \times{ }^{\circ} \tilde{V}_{v}$. Let $\tilde{v}$ be an arbitrary vector in ${ }^{\circ} \tilde{V}_{v}$. Then:

Proposition 3.6: The integral

$$
F_{v}(s)=\int_{N_{v}}\left\langle\left\langle f_{v}\left(g n w_{0}\right), \tilde{v}\right\rangle\right\rangle d n \quad\left(f_{v} \in V_{v}\right)
$$

converges absolutely for $\operatorname{Re}(s)$ small enough and it can be extended to a meromorphic function on the whole complex plane. Furthermore there exists a polynomial $P\left(q_{v}^{s}\right)\left(q_{v}\right.$ denotes the number of the elements
in the residual field) so that as a function of $s$

$$
P\left(q_{v}^{s}\right) F_{v}(s)
$$

is entire.

Proof: In the case in hand, this can be proved using the same technique as in [4]. In general this is a result of Harish-Chandra [6].

Now, for $f_{v}$ in $V_{v}$ we shall define

$$
\begin{equation*}
f_{v}^{\prime}(g)=\int_{N_{v}} f_{v}\left(g n w_{0}\right) d n \tag{3.3}
\end{equation*}
$$

We put

$$
{ }^{\circ} \pi_{v}^{\prime}(m)={ }^{\circ} \pi_{v}\left(w_{0}^{-1} m w_{0}\right) ;
$$

then $f_{v}^{\prime}$ belongs to the space $V_{v}^{\prime}$ of the representation

$$
\Pi_{v}^{\prime}=\operatorname{ind}_{P_{v} \uparrow G_{v}}^{\circ} \pi_{v}^{\prime} \otimes \delta_{P, v}^{-s} \quad\left(\forall^{\prime} s \in \mathbb{C}\right)
$$

Let $\lambda^{\prime}$ denote the primed analogue of the Whittaker functional defined by (3.2). More precisely,

$$
\begin{equation*}
\lambda^{\prime}\left(f_{v}^{\prime}\right)=\int_{N_{v}}\left(f_{v}^{\prime}\left(n w_{0}\right), e\right) \overline{\chi_{v}(n)} d n \tag{3.4}
\end{equation*}
$$

with $f_{v}^{\prime}$ in $V_{v}^{\prime}$.
Now, let $f_{v}^{\prime}$ be a function defined by (3.3). Then (3.4) defines another Whittaker functional (which we still denote by $\lambda^{\prime}$ ) on $V_{v}$ by

$$
\begin{equation*}
\lambda^{\prime}\left(f_{v}\right)=\lambda^{\prime}\left(f_{v}^{\prime}\right) \tag{3.5}
\end{equation*}
$$

Later in $\S 5$, we shall show that $\lambda$ is in fact a non-zero Whittaker functional on $\Pi_{v}$ (Theorem 5.5). Therefore from [14] and [15] it follows that $\lambda$ and $\lambda^{\prime}$ are proportional and the computations of the next section will show that the coefficient of proportionality is directly related to the Langlands' root number attached to a four dimensional representation of $S L_{2}$.

More precisely, let $\rho$ be a four dimensional representation of $S L_{2}$, the associated group for ${ }^{\circ} G$. Assume at each place $v$ the local $L$-function $L\left(s, \rho,{ }^{\circ} \pi_{v}\right)$ is defined. As we shall see later these coefficients of proportionality will lead to certain local coefficients which we denote by $\gamma\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right)$ (see §5). Then the Langlands' root numbers are defined by

$$
\epsilon\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right)=\gamma\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right) \cdot \frac{L\left(1-s, \tilde{\rho},{ }^{\circ} \pi_{v}\right)}{L\left(s, \rho,{ }^{\circ} \pi_{v}\right)}
$$

where $\tilde{\rho}$ denotes the contragredient representation of $\rho$.

## 4. Local coefficients

In this section, we shall explicitly compute the coefficients of proportionality, which was mentioned at the end of $\S 3$, for certain classes of representations.

Let $v$ and $\tilde{v}$ be two fixed vectors in the space of $\sigma_{v}$ and that of the contragredient representation of $\tilde{\sigma}_{v}$, respectively. We use

$$
\langle,\rangle
$$

for the natural pairing between $\tilde{\sigma}_{v}$ and its contragredient. We shall start with the following trivial lemma.

## Lemma 4.1: The function

$$
\begin{align*}
f_{s, v}(g)=\delta_{P, v}^{s-1 / 2}\left(p_{0}\right) \int_{\cdot K_{v}} & \left\langle\tilde{\sigma}_{v}\left(k_{0}\right) \sigma_{v}\left(k_{v}^{-1}\right) v, \tilde{v}\right\rangle  \tag{4.1}\\
& \cdot{ }^{\circ} \pi_{v}\left(p_{0}^{-1} k_{v}^{-1}\right)^{\circ} w_{v} d k_{v}
\end{align*}
$$

with

$$
g=k_{0} p_{0}
$$

is a well defined function on $G_{v}$. Furthermore $f_{s, v}$ belongs to $V_{v}$ and every function in $V_{v}$ is a finite linear combination of the functions of this type (different $\sigma_{v}, \tilde{\sigma}_{v}$ and ${ }^{\circ} w_{v}$ ).

For $f_{s, v}$ as above, we put

$$
w_{s, v}(g)=\lambda\left(\Pi_{v}\left(g^{-1}\right) f_{s, v}\right) .
$$

Then

$$
\begin{gather*}
w_{s, v}(g)=\int_{N_{v}} \int_{{ }_{K_{v}}} \delta_{P, v}^{s-1 / 2}\left(b\left(g n w_{0}\right)\right)\left\langle\tilde{\sigma}_{v}\left(k\left(g n w_{0}\right)\right) \sigma_{v}\left(k_{v}^{-1}\right) v, \tilde{v}\right\rangle  \tag{4.2}\\
\cdot{ }^{\circ} w_{v}\left(k_{v} \cdot \overline{b\left(g n w_{0}\right)}\right) \overline{\chi_{v}(n)} d k_{v} d n
\end{gather*}
$$

with the same notation as in §2, i.e.

$$
g n w_{0}=k\left(g n w_{0}\right) \cdot b\left(g n w_{0}\right)
$$

with the obvious meanings for $k\left(g n w_{0}\right)$ and $b\left(g n w_{0}\right)$.
We shall use $\underline{w}_{s, v}$ to denote the corresponding operator valued function, i.e.

$$
\begin{aligned}
\underline{w}_{s, v}(g)= & \int_{N_{v}} \int_{K_{i}} \delta_{P, v}^{s-1 / 2}\left(b\left(g n w_{0}\right)\right)^{\circ} w_{v}\left(k_{v} \cdot \overline{b\left(g n w_{0}\right)}\right) \overline{\chi_{v}(n)} \\
& \times \tilde{\sigma}_{v}\left(k\left(g n w_{0}\right)\right) \sigma_{v}\left(k_{v}^{-1}\right) p_{\sigma_{v}} \cdot d k_{v} d n .
\end{aligned}
$$

For each $f_{s, v}$, the function $f_{s, v}^{\prime}$ defined by (3.3) is equal to

$$
\begin{align*}
f_{s, v}^{\prime}(g)= & \int_{N_{v}} \int_{\sigma_{K_{v}}} \delta_{P, v}^{s-1 / 2}\left(b\left(g n w_{0}\right)\right)\left\langle\tilde{\sigma}_{v}\left(k\left(g n w_{0}\right)\right) \sigma_{v}\left(k_{v}^{-1}\right) v, \tilde{v}\right\rangle  \tag{4.3}\\
& \cdot{ }^{\circ} \pi_{v}\left(b\left(g n w_{0}\right)^{-1} \cdot k_{v}^{-1}\right)^{\circ} w_{v} d k_{v} d n .
\end{align*}
$$

We still use $f_{s, v}^{\prime}$ to denote the analytic continuation of (4.3) by means of Proposition 3.6. We put

$$
w_{s, v}^{\prime}(g)=\lambda^{\prime}\left(\Pi_{v}^{\prime}\left(g^{-1}\right) f_{s, v}^{\prime}\right)
$$

for $f_{s, v}^{\prime}$ defined by (4.3).
Now we shall assume that ${ }^{\circ} \pi_{v}$ is not supercuspidal (when $v$ is a finite place). Then there is a quasi-character $\eta_{v}$ of ${ }^{\circ} T_{v}$ so that ${ }^{\circ} \pi_{v}$ can be realized as a quotient of the space of the left ${ }^{\circ} K_{v}$-finite functions ${ }^{\circ} f_{v}$ on ${ }^{\circ} G_{v}$ which satisfy

$$
\begin{equation*}
{ }^{\circ} f_{v}(g t u)=\eta_{v}(t)^{\circ} \delta_{v}^{-1 / 2}(t)^{\circ} f_{v}(g) \tag{4.4}
\end{equation*}
$$

with $t$ in ${ }^{\circ} T_{v}, u$ in ${ }^{\circ} U_{v}$, and $g$ in ${ }^{\circ} G_{v}$. Here ${ }^{\circ} \delta_{v}$ denotes the modulus character of ${ }^{\circ} T_{v}$ and the corresponding Hecke algebra acts by convolutions (cf. [8]).

We use $w_{1}$ to denote $w_{\beta}$ which we realize as the longest element in
$W_{M}$. Then, there is a function ${ }^{\circ} f_{v}$ so that

$$
\begin{equation*}
{ }^{\circ} w_{v}(g)=\int_{\circ U_{v}}{ }^{\circ} f_{v}\left(g u w_{1}\right) \overline{\chi_{v}(u)} d u \tag{4.5}
\end{equation*}
$$

where $g$ is in ${ }^{\circ} G_{v}$.
We shall consider $\eta_{v}$ as a character of $T_{v}$ in the obvious manner (i.e. trivial on $A_{v}$ ).

The goal of this section is to compute explicitly $w_{s, v}^{\prime}$ in terms of $\boldsymbol{w}_{s, v}$.

Following Jacquet [7], we define the following operator valued function on $\boldsymbol{G}_{\boldsymbol{v}}$

$$
\begin{equation*}
\underline{h}_{s, v}(g)=\eta_{v} \cdot{ }^{\circ} \delta_{v}^{-1 / 2} \delta_{P, v}^{s-1 / 2}(t(g)) \tilde{\sigma}_{v}(k(g)) P\left(\tilde{\sigma}_{v}, \eta_{v}\right) \tag{4.6}
\end{equation*}
$$

Where $P\left(\tilde{\sigma}_{v}, \eta_{v}\right)$ denotes the projection onto the subspace of the vectors $v$ which satisfy

$$
\tilde{\sigma}_{v}(t u) v=\eta_{v}(t) v
$$

with $t$ in $K_{v} \cap T_{v}$ and $u$ in $K_{v} \cap U_{v}$. The corresponding Whittaker function is defined by

$$
\begin{equation*}
\underline{w}_{s, v}(g)=\int_{U_{v}} \underline{h}_{s, v}\left(g u w_{2}\right) \overline{\chi_{v}(u)} d u \tag{4.7}
\end{equation*}
$$

with

$$
w_{2}=w_{1} w_{0},
$$

the longest element in $W$. As a function of $s, \underline{w}_{s, v}$ is entire (cf. [7]).
We shall consider the following operator valued function

$$
\begin{align*}
\underline{f}_{s, v}^{\prime}(g)= & \int_{N_{v}} \int_{v_{K_{v}}} \delta_{P, v}^{s-1 / 2}\left(b\left(g n w_{0}\right)\right)^{\circ} w_{v}\left(k_{v} \overline{b\left(g n w_{0}\right)}\right) \tilde{\sigma}_{v}\left(k\left(g n w_{0}\right)\right)  \tag{4.8}\\
& \cdot \tilde{\sigma}_{v}\left(k_{v}^{-1}\right) \cdot p_{\sigma_{v}} d k_{v} d n
\end{align*}
$$

where the integral is convergent for $\operatorname{Re}(s)$ sufficiently small. Then (4.3) implies that

$$
\left(f_{s, v}^{\prime}(g), e\right)=\left\langle f_{s, v}^{\prime}(g) v, \tilde{v}\right\rangle
$$

We need the following lemma.

Lemma 4.2: The non-zero operator

$$
\int_{\odot_{K_{v}}}{ }^{\circ} f_{v}\left(k_{v}\right) \tilde{\sigma}_{v}\left(k_{v}^{-1}\right) d k_{v} \cdot p_{\sigma_{v}}
$$

sends the space of $\tilde{\sigma}_{v}$ into the range of

$$
p\left(\tilde{\sigma}_{v}, \eta_{v}\right) \cdot p_{\sigma_{v}}
$$

We use $H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right)$ to denote the operator introduced in Lemma 4.2.
Lemma 4.3: For $\operatorname{Re}(s)$ sufficiently small, we have

$$
\begin{equation*}
w_{s, v}^{\prime}(g)=\underline{w}_{-s, v}(g) \int_{N_{v}} \underline{h}_{s, v}\left(n w_{0}\right) d n \cdot H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right) \tag{4.9}
\end{equation*}
$$

Proof: Substitution of (4.5) into (4.8), followed by a simple change of variables, shows that for $\operatorname{Re}(s)$ sufficiently small

$$
\begin{align*}
f_{s, v}^{\prime}(g)= & \int_{N_{v}} \int_{U_{v}} \delta_{P, v}^{s-1 / 2}\left(b_{0}\right) \eta_{v}{ }^{\circ} \delta_{v}^{-1 / 2}\left(t\left(b_{0} u w_{1}\right)\right) \overline{\chi_{v}(u)}  \tag{4.9.1}\\
& \cdot \tilde{\sigma}_{v}\left(k_{0}\right) \tilde{\sigma}_{v}\left(k\left(\bar{b}_{0} u w_{1}\right)\right) d u d n \cdot H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right)
\end{align*}
$$

with

$$
b_{0}=b\left(g n w_{0}\right)
$$

and

$$
k_{0}=k\left(g n w_{0}\right) .
$$

Using

$$
\begin{gathered}
k\left(g n w_{0} u w_{1}\right)=k_{0} \cdot k\left(b_{0} u w_{1}\right), \\
\delta_{P, v}\left(u w_{1}\right)=1 \quad\left(u \in{ }^{\circ} U_{v}\right)
\end{gathered}
$$

and Lemma 4.2, (4.9.1) reduces to

$$
\int_{N_{v}} \int_{U_{v}} \underline{h}_{s, v}\left(g n w_{0} u w_{1}\right) \overline{\chi_{v}(u)} d u d n \cdot H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right)
$$

which is equal to

$$
\begin{equation*}
\int_{N_{v}} \int_{U_{U_{v}}} \underline{h}_{s, v}\left(g u w_{1} n w_{0}\right) \overline{\chi_{v}(u)} d u d n \cdot H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right) \tag{4.9.2}
\end{equation*}
$$

Therefore for $\operatorname{Re}(s)$ sufficiently small (4.9.2) implies that

$$
\begin{equation*}
f_{s, v}^{\prime}(g)=\int_{{ }^{\circ} U_{v}} \underline{h}_{-s, v}\left(g u w_{1}\right) \overline{\chi_{v}(u)} d u \cdot \int_{N_{v}} \underline{h}_{s, v}\left(n w_{0}\right) d n \cdot H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right) \tag{4.9.3}
\end{equation*}
$$

Then using Proposition 3.6 and (4.9.3), we can define $f_{s, v}^{\prime}$ for (almost) all $s$ and the lemma follows immediately.

To simplify the right hand side of (4.9), we need the following results from the first part of [7].

For a local field $K$ (archimedean or non-archimedean), let $H$ denote the group $S L_{2}(K)$ and let $M_{H}$ be the maximal compact subgroup of $H$ as in [7]. We put $A_{H}$ and $N_{H}$ for the diagonal and the standard unipotent subgroups of $H$, respectively.

Let $\mathfrak{D}$ be a finite dimensional representation of $M_{H}$ on the Hilbert space $H(\mathfrak{D})$. We use $\eta$ to denote a character of $K^{*}$, the multiplicative subgroup of $K$. It can be considered as a character of $A_{H}$ by means of

$$
\eta\left(\left(\begin{array}{ll}
a & 0 \\
0 & a
\end{array}\right)\right)=\eta(a)
$$

We use $\mathbb{S}\left(K^{2}\right)$ to denote the space of the Schwartz-Bruhat functions on $K \times K$ and $\operatorname{End}(H(\mathfrak{D}))$ for the algebra of endomorphisms of $H(\mathfrak{D})$. Let $s$ be a complex number and let $\Psi$ be a function in

$$
\mathfrak{S}\left(K^{2}\right) \otimes \operatorname{End}(H(\mathfrak{D}))
$$

We define (Proposition 1.7 of [7])

$$
\begin{equation*}
L_{\Psi}(g, \eta, s)=\int_{K^{*}} \Psi\left(t e_{1} \cdot{ }^{t} g\right) \overline{\eta(t)}|t|^{s} d^{*} t \tag{4.10}
\end{equation*}
$$

with $g$ in $H$ and

$$
e_{1}=(1,0) .
$$

Then

$$
\begin{equation*}
L_{\Psi}(g, \eta, s)=L(g, \mathfrak{F}, \eta, s) L_{\Psi}(e, \eta, s) \tag{4.11}
\end{equation*}
$$

with

$$
\begin{equation*}
L(g, \mathfrak{D}, \eta, s)=\mathfrak{D}(m) \eta(a)|a|^{-s} P(\mathfrak{D}, \eta) \tag{4.12}
\end{equation*}
$$

if

$$
g=m a n
$$

with $m$ in $M_{H}, a$ in $A_{H}$, and $n$ in $N_{H}$. Here $P(\mathfrak{D}, \eta)$ denotes the projection into the subspace of the vectors $v$ which satisfy

$$
\mathfrak{D}(a n) v=\eta(a) v
$$

with $a$ in $A_{H} \cap M_{H}$ and $n$ in $N_{H} \cap M_{H}$. Let $\tau$ be the additive character of $K$ fixed in Section 1 of [7]. We put

$$
\hat{\Psi}(x, y)=\int_{K^{2}} \Psi(u, v) \tau(x v-y u) d u d v
$$

and

$$
w=\left(\begin{array}{rr}
0 & -1 \\
1 & 0
\end{array}\right) .
$$

Finally the operator

$$
Ч(\mathfrak{D}, \eta, s),
$$

introduced in Corollary (1.10) of [7], satisfies the following relations

$$
\begin{gather*}
L_{\psi}(e, \bar{\eta},-s+1)=Ч(\mathfrak{D}, \eta, s) L_{\psi}(e, \eta, s+1),  \tag{4.13}\\
Ч(\mathfrak{D}, \bar{\eta},-s) \cdot Ч(\mathfrak{D}, \eta, s)=P(\mathfrak{D}, \eta) \tag{4.14}
\end{gather*}
$$

and

$$
\begin{equation*}
P(\mathfrak{D}, \bar{\eta}) Ч(\mathfrak{D}, \eta, s) P(\mathfrak{D}, \eta)=Ч(\mathfrak{D}, \eta, s) . \tag{4.15}
\end{equation*}
$$

Let $\gamma(\eta, s)$ be the coefficient (cf. [17]) defined by

$$
\begin{equation*}
\int_{K^{*}} \Psi(t) \overline{\eta(t)}|t|^{1-s} d^{*} t=\gamma(\eta, s) \int_{K^{*}} \hat{\Psi}(t) \eta(t)|t|^{s} d^{*} t, \tag{4.16}
\end{equation*}
$$

where $\Psi$ is a Schwartz function on $K$ and

$$
\hat{\Psi}(t)=\int_{K} \Psi(x) \overline{\tau(t x)} d x
$$

We shall prove the following important lemma.

Lemma 4.4: For $\operatorname{Re}(s)$ sufficiently small, one has

$$
\begin{equation*}
Ч(\mathfrak{D}, \bar{\eta},-s) \int_{N_{H}} L(n w, \mathfrak{D}, \eta, s+1) d n=\gamma^{-1}(\bar{\eta}, s) P(\mathfrak{D}, \eta) . \tag{4.17}
\end{equation*}
$$

Proof: We have

$$
\begin{align*}
& \int_{N_{H}} L_{\Psi}(n w, \eta, s+1) d n=  \tag{4.17.1}\\
& \quad \int_{N_{H}} L(n w, \mathfrak{D}, \eta, s+1) d n \cdot L_{\Psi}(e, \eta, s+1)
\end{align*}
$$

Then straight forward computations imply that the left hand side of (4.17.1) is equal to

$$
\int_{K \times K^{*}} \Psi(x, t) \overline{\eta(t)}|t|^{s} d^{*} t d x
$$

Using (4.16), this is equal to

$$
\begin{gather*}
\eta(-1) \gamma(\eta, 1-s) \int_{K^{*}} \hat{\Psi}(t, 0) \eta(t)|t|^{1-s} d^{*} t  \tag{4.17.2}\\
=\gamma^{-1}(\bar{\eta}, s) L_{\varphi}(e, \bar{\eta},-s+1)
\end{gather*}
$$

The lemma follows if we apply (4.13) and (4.14) to (1.17.1) and the right hand side of (4.17.2).

Remark: The lemma is still true if we assume that $\eta$ is a quasicharacter and replace $\overline{\boldsymbol{\eta}}$ by $\boldsymbol{\eta}^{-1}$.

For every root $\gamma$, there is an isomorphism $x_{\gamma}$ from $S L_{2}$ onto the group $G^{\gamma}$ whose Lie algebra is generated by $H_{\gamma} \in \mathfrak{h}, X_{\gamma} \in \mathfrak{g}_{\gamma}$ and $X_{-\gamma} \in \mathfrak{g}_{-\gamma}(G$ is universal). It sends $\binom{1 x}{0},\binom{10}{x}$, and $\left(\begin{array}{c}t \\ 0 \\ 0\end{array} t^{0}-1\right)$ to $\exp \left(x X_{\gamma}\right), \exp \left(x X_{-\gamma}\right)$, and $h_{\gamma}(t)$, respectively. There is also an isomorphism (conjugation by $w$ ) between $G^{\gamma}$ and $G^{w(\gamma)}$ for every $w$ in $W$.

For every place $v$, we have the Iwasawa decomposition

$$
G_{v}^{\gamma}=K_{v}^{\gamma} \cdot T_{v}^{\gamma} \cdot U_{v}^{\gamma}
$$

with

$$
K_{v}^{\gamma}=G_{v}^{\gamma} \cap K_{v}
$$

and

$$
T_{v}^{\gamma}=\left\{h_{\gamma}(t) ; t \in F_{v}^{*}\right\}
$$

We put

$$
\nu=\sum_{\gamma \in \Psi^{+}-\{\beta\}} \gamma
$$

and

$$
\langle\gamma, \delta\rangle=2 \frac{(\gamma, \delta)}{(\delta, \delta)}
$$

for a pair of roots $\gamma$ and $\delta$.
As we mentioned before, there is no harm in assuming that $\boldsymbol{\eta}_{v}$ is trivial on the center and therefore $\eta_{v}$ is a quasi-character of $T_{v}$. Notice that

$$
\eta_{v}\left(h_{\gamma}(t)\right)=\eta_{v}^{\langle\beta, \gamma\rangle}(t) \quad(\forall \gamma \in \Psi)
$$

by means of which we may consider each $\eta_{v}^{\langle\beta, \gamma\rangle}$ as a quasi-character of $F_{v}^{*}$.

For a pair of roots $\gamma$ and $\delta$, we define

$$
\begin{equation*}
\underline{h}_{s, v}^{\delta, \gamma}(g)=\tilde{\sigma}_{v}(k(g)) \cdot \eta_{v}^{\langle\beta, \gamma\rangle} \cdot \alpha_{v}^{s(\nu, \gamma\rangle-1}(t) \cdot \underline{P}\left(\tilde{\sigma}_{v}, \eta_{v}\right) \tag{4.18}
\end{equation*}
$$

for $g$ in $G_{v}^{\delta}$ with

$$
g=k(g) h_{\delta}(t) u(g)
$$

Here $\alpha_{v}$ denotes the modulus character (absolute value) and $\underline{P}\left(\tilde{\sigma}_{v}, \boldsymbol{\eta}_{v}\right)$ is the corresponding projection. We put

$$
\begin{equation*}
A_{s, v}^{\delta, \gamma}=\int_{U_{v}^{\delta}} \underline{\boldsymbol{h}}_{s, v}^{\delta, \gamma}\left(u w_{\delta}\right) d u \tag{4.19}
\end{equation*}
$$

We need the following lemma.

Lemma 4.5: For $\operatorname{Re}(s)$ sufficiently small, the following relation holds.
(4.20) $\int_{N_{v}} \underline{h}_{s, v}\left(n w_{0}\right) d n=A_{s, v}^{\alpha, \alpha+\beta} \cdot A_{s, v}^{\beta, 3 \alpha+2 \beta} \cdot A_{s, v}^{\alpha, 2 \alpha+\beta} \cdot A_{s, v}^{\beta, 3 \alpha+\beta} \cdot A_{s, v}^{\alpha, \alpha}$.

Proof: In fact the ordering introduced here is the one in [4] (see also [11]). Set

$$
\overline{\boldsymbol{\theta}}_{0}=\boldsymbol{\Psi}^{-}-\{-\boldsymbol{\beta}\} ;
$$

then (4.20) is equal to

$$
\tilde{\sigma}_{v}\left(w_{0}\right) \int_{N_{v}\left(\bar{\theta}_{0}\right)} \underline{h}_{s, v}(\bar{n}) d \bar{n}
$$

where $N_{v}\left(\bar{\theta}_{0}\right)$ denotes the unipotent subgroup generated by $\overline{\boldsymbol{\theta}}_{0}$. If we put

$$
\bar{\theta}_{1}=\boldsymbol{\Psi}^{-}-\{-\boldsymbol{\beta},-(\alpha+\boldsymbol{\beta})\}
$$

and denote the unipotent subgroup generated by $\bar{\theta}_{1}$ by $N_{v}\left(\bar{\theta}_{1}\right)$, we can write

$$
\begin{align*}
& \int_{N_{v}} \underline{h}_{s, v}\left(n w_{0}\right) d n=\tilde{\sigma}_{v}\left(w_{0} w_{\alpha+\beta}\right) \cdot \int_{U_{v}^{\alpha+\beta}} \tilde{\sigma}_{v}\left(k\left(u w_{\alpha+\beta}\right)\right)  \tag{4.20.1}\\
& \cdot \eta_{v} \delta_{P, v}^{s} \delta_{\alpha+\beta, v}^{-1 / 2}\left(t\left(u w_{\alpha+\beta}\right)\right) d u \\
& \cdot \int_{N_{v}\left(\bar{\theta}_{1}\right)} \underline{h}_{s, v}(\bar{n}) d \bar{n}
\end{align*}
$$

Here ${ }^{\circ} \boldsymbol{\delta}_{\alpha+\beta, v}$ is defined by

$$
{ }^{\circ} \delta_{\alpha+\beta, v}\left(h_{\alpha+\beta}(t)\right)=|t|^{2}
$$

If we conjugate $k\left(u w_{\alpha+\beta}\right)$ with $w_{0} w_{\alpha+\beta}$ and use the isomorphism between $G_{v}^{\alpha+\beta}$ and $G_{v}^{\alpha}$ (through this conjugation), we shall see that (4.20.1) is equal to

$$
A_{s, v}^{\alpha, \alpha+\beta} \cdot \sigma_{v}\left(w_{0} w_{\alpha+\beta}\right) \cdot \int_{N_{v}\left(\overline{\theta_{1}}\right)} \underline{h}_{s, v}(\bar{n}) d \bar{n}
$$

Then applying the same argument to

$$
\tilde{\sigma}_{v}\left(w_{0} w_{\alpha+\beta}\right) \cdot \int_{N_{v}\left(\bar{\theta}_{1}\right)} \underline{h}_{s, v}(\bar{n}) d \bar{n}
$$

by means of the ordering introduced in the statement of the lemma, completes the proof inductively.

At each finite place $v$, we shall fix a non-trivial character $\tau_{v}$ of $F_{v}$ so that the largest ideal for which $\tau_{v}$ is trivial is the ring of integers $O_{v}$ of $F_{v}$. Then

$$
\chi_{\gamma, v}\left(\exp \left(x X_{\gamma}\right)\right)=\tau_{v}\left(\mu_{\gamma, v} x\right) \quad(\gamma \in \Delta)
$$

with $\mu_{\gamma, v}$ in $F_{v}^{*}$. Notice that $\mu_{\gamma, v}$ is a unit for almost all $v$. When $v$ is an infinite place, we shall fix $\tau_{v}$ as in [7].

For a simple root $\delta$, we put

$$
\begin{equation*}
B_{s, v}^{\delta, \gamma}=\mathrm{Y}_{\delta}\left(\tilde{\sigma}_{v}, \eta_{v}^{-\langle\beta, \gamma\rangle},\langle\nu, \gamma\rangle \cdot s\right) . \tag{4.21}
\end{equation*}
$$

Lemma 4.6: With $B_{s, v}^{\delta, \gamma}$ as in (4.21) we have

$$
\begin{align*}
\underline{w_{-s, v}}(g)= & \left|\mu_{\alpha, v}^{2} \cdot \mu_{\beta, v}\right|^{10 s} \underline{w}_{s, v}(g) \cdot B_{s, v}^{\alpha, \alpha} \cdot B_{s, v}^{\beta, 3 \alpha+\beta}  \tag{4.22}\\
& \cdot B_{s, v}^{\alpha, 2 \alpha+\beta} \cdot B_{s, v}^{\beta, 3 \alpha+\beta} \cdot B_{s, v}^{\alpha, \alpha+\beta} .
\end{align*}
$$

Proof: Put

$$
\bar{U}=w_{2} U w_{2}^{-1} .
$$

We use the ordering introduced in Lemma 4.5 for the set

$$
\Psi^{+}-\{\boldsymbol{\beta}\} .
$$

Then if we put $\gamma_{1}$ for $\alpha+\beta$, $\gamma_{5}$ will be equal to $\alpha$.
Now, we inductively define the following quasi-characters of $T_{v}$ :

$$
\eta_{i+1}=w_{\delta(i)} \eta_{i} \quad 1 \leq i \leq 5
$$

and

$$
\lambda_{i+1}=w_{\delta(i)} \eta_{i} \quad 1 \leq i \leq 5
$$

with the obvious action of $W$ on the character group. The index $\delta(i)$ is defined by

$$
\delta(i)= \begin{cases}\alpha & i \text { odd } \\ \beta & i \text { even }\end{cases}
$$

and

$$
\begin{aligned}
\boldsymbol{\eta}_{1} & =w_{0} \boldsymbol{\eta}_{v} \\
& =\boldsymbol{\eta}_{v},
\end{aligned}
$$

and finally

$$
\begin{aligned}
\lambda_{1} & =w_{0} \delta_{P, v}^{s} \\
& =\delta_{P, v}^{s s} .
\end{aligned}
$$

Then from the functional equation of Jacquet ([7], Proposition 3.3), it
follows that

$$
\begin{align*}
E_{\bar{U}}\left(g, \tilde{\sigma}_{v}, \eta_{i}, \lambda_{i}, \chi_{v}\right)= & \eta_{v}^{\left\langle\beta, \gamma_{i}\right\rangle}\left(\mu_{\delta(i), v}\right)\left|\mu_{\delta(i), v}\right|^{\left\langle\nu, \gamma_{i}\right\rangle s}  \tag{4.22.1}\\
& \cdot E_{\bar{U}}\left(g, \tilde{\boldsymbol{\sigma}}_{v}, \boldsymbol{\eta}_{i+1}, \lambda_{i+1}, \chi_{v}\right) \\
& \cdot Ч_{\delta(i)}\left(\tilde{\sigma}_{v}, \boldsymbol{\eta}_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle},\left\langle\nu, \gamma_{i}\right\rangle s\right)
\end{align*}
$$

for $i=1, \ldots, 5$. But

$$
\begin{equation*}
E_{\bar{U}}\left(g w_{2}^{-1}, \tilde{\sigma}_{v}, \eta_{1}, \lambda_{1}, \chi_{v}\right)=\underline{w}_{-s, v}(g) \tag{4.22.2}
\end{equation*}
$$

and

$$
\begin{equation*}
E_{\bar{U}}\left(g w_{2}^{-1}, \tilde{\sigma}_{v}, \eta_{6}, \lambda_{6}, \chi_{v}\right)=\underline{w}_{s, v}(g) \tag{4.22.3}
\end{equation*}
$$

and the lemma follows if we consider (4.22.1) for $1 \leq i \leq 5$ and combine them together with (4.22.2) and (4.22.3).

If we apply Lemmas $4.4,4.5$, and 4.6 to Lemma 4.3 , we conclude
Theorem 4.7: Let $H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right)$ be defined as in Lemma 4.2; then

$$
\begin{align*}
\int_{N_{v}} f_{s, v}^{\prime}\left(g n w_{0}\right) \overline{\chi_{v}(n)} d n= & \left|\mu_{\alpha, v}^{2} \cdot \mu_{\beta, v}\right|^{10 s} \prod_{i=1}^{5} \gamma^{-1}\left(\eta_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle},-\left\langle\nu, \gamma_{i}\right\rangle s\right) \\
& \cdot \underline{w}_{s, v}(g) \cdot H\left({ }^{\circ} f_{v}, \tilde{\sigma}_{v}\right) . \tag{4.23}
\end{align*}
$$

Corollary 4.8: With notation as before, we have

$$
\begin{align*}
\underline{w}_{s, v}^{\prime}(g)= & \left|\mu_{\alpha, v}^{2} \cdot \mu_{\beta, v}\right|^{10 s} \prod_{i=1}^{s} \gamma^{-1}\left(\eta_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle},-\left\langle\nu, \gamma_{i}\right\rangle s\right)  \tag{4.24}\\
& \cdot \underline{w}_{s, v}(g)
\end{align*}
$$

We shall conclude this section by establishing a result similar to that of Corollary 3.5 for the archimedean places.

Suppose $v$ is an archimedean place. As usual, let $C_{c}^{\infty}\left(G_{v}\right)$ denote the space of smooth functions on $G_{v}$ with compact support. We assume $C_{c}^{\infty}\left(G_{v}\right)$ has the Schwartz topology.

For every function $\Psi$ in $C_{c}^{\infty}\left(G_{v}\right)$, the integral

$$
\begin{equation*}
h_{\Psi}(g)=\int_{B_{v}} \delta_{P, v}^{-s+1 / 2}(b) \eta_{v}^{-1} \cdot{ }^{\circ} \delta_{v}^{1 / 2}(t(b)) \Psi(g b) d_{\ell}(b) \tag{4.25}
\end{equation*}
$$

defines a function which satisfies

$$
\begin{equation*}
h_{\Psi}(g t u)=\eta_{v}{ }^{\circ} \delta_{v}^{-1 / 2} \delta_{P, v}^{s-1 / 2}(t) h_{\Psi}(g) \tag{4.26}
\end{equation*}
$$

with $t$ in $T_{v}$ and $u$ in $U_{v}$. Here $d_{\ell}(b)$ denotes the left invariant Haar measure on $B_{v}$.

The map sending $\Psi$ to $h_{\Psi}$ is a surjective homomorphism onto the space of smooth functions on $G_{v}$ which satisfy (4.26)

Let $h$ be such a function; then the integral

$$
\int_{U_{v}} h\left(g u w_{2}\right) \overline{\chi_{v}(u)} d u
$$

converges absolutely for $\operatorname{Re}(s)$ sufficiently small (cf. [7]), and defines a function $w(g)$ on $G_{v}$.

Lemma 4.0: There is a function $\Psi$ in $C_{c}^{\infty}\left(G_{v}\right)$ for which the function

$$
w_{\Psi}(g)=\int_{U_{v}} h_{\Psi}\left(g u w_{2}\right) \overline{\chi_{v}(u)} d u
$$

(which is defined by this integral for $\operatorname{Re}(s)$ sufficiently small) does not vanish identically.

Proof: For $\operatorname{Re}(s)$ sufficiently small we can write

$$
w_{\Psi}(g)=\int_{U_{v}} \int_{B_{v}} \delta_{P, v}^{-s+1 / 2}(b) \eta_{v}^{-10} \delta_{v}^{1 / 2}(t(b)) \cdot \overline{\chi_{v}(u)} \Psi\left(g u w_{2} b\right) d_{\ell}(b) d u .
$$

Then we may choose $\Psi$ with support in $U_{v} w_{2} B_{v}$ (which is open in $G_{v}$ ) so that $w_{\Psi}(e)$ is not zero. This completes the lemma.

Lemma 4.10: For $\operatorname{Re}(s)$ sufficiently small, the integral

$$
\begin{gathered}
w_{\Psi}(e)=\int_{U_{v}} \int_{B_{v}} \delta_{P, v}^{-s+1 / 2}(b) \cdot \eta_{v}^{-10} \delta_{v}^{1 / 2}(t(b)) \overline{\chi_{v}(u)} \Psi\left(u w_{2} b\right) d_{\ell}(b) d u \\
\left(\Psi \in C_{c}^{\infty}\left(G_{v}\right)\right)
\end{gathered}
$$

defines a distribution on $G_{v}$.

Proof: We define the following function on $G_{v}$ :

$$
f(g)=\delta_{P, v}^{-s+1 / 2}(b) \eta_{v}^{-10} \delta_{v}^{1 / 2}(t(b)) \overline{\chi_{v}(u)}
$$

if

$$
g=u w_{2} b \in U_{v} w_{2} B_{v}
$$

and zero otherwise. Then $w_{\Psi}(e)$ can be written as

$$
w_{\Psi}(e)=\int_{G_{v}} f(g) \Psi(g) d g \quad\left(\Psi \in C_{c}^{\infty}\left(G_{v}\right)\right)
$$

which is clearly a distribution since $f$ is locally integrable for $\operatorname{Re}(s)$ sufficiently small. In fact, if $\omega$ is a compact subset of $G_{v}$, there exists a positive function $h$ in $C_{c}^{\infty}\left(G_{v}\right)$ with

$$
h(x)=1 \quad(x \in \omega)
$$

Then, it is clear that

$$
\int_{\omega}|f(g)| d g \leq \int_{B_{v}} \int_{U_{v}} h\left(u w_{2} b\right)\left|\delta_{P, v}^{-s+1 / 2}(b) \cdot \eta_{v}^{-10} \delta_{v}^{1 / 2}(t(b))\right| d u d_{\ell}(b) .
$$

But for the small values of $\operatorname{Re}(s)$ the right hand side of this inequality is finite which implies the lemma.

The functions $h_{s, v}(g)$ (different $\tilde{\sigma}_{v}, v$ and $\tilde{v}$ ) will generate the subspace of left $K_{v}$-finite functions on $G_{v}$ which satisfy (4.26). Since the subspace of left $K_{v}$-finite functions with compact support on $G_{v}$ is dense in $C_{c}^{\infty}\left(G_{v}\right)$, we conclude the following corollary.

Corollary 4.11: Let $v$ be an archimedean place; then for $\operatorname{Re}(s)$ sufficiently small, there is a function $\underline{h}_{s, v}$, defined by (4.6), for which $\underline{w}_{s, v}$ is not identically zero.

## 5. Functional equation

From now on, we shall assume that the character $\chi$ has been chosen so that

$$
\mu_{\alpha, v}=\mu_{\beta, v} \quad(\forall v)
$$

The complex dual group for ${ }^{\circ} G$ (or $P G L_{2}$ ), i.e. ${ }^{\circ} G^{\wedge}$, is $S L_{2}(\mathbb{C})$.
For an unramified place (i.e. when ${ }^{\circ} \pi_{v}$ is class-one and $\mu_{\alpha, v}$ is a unit in $O_{v}$ ), $\eta_{v}$ is unramified and therefore we can write

$$
\eta_{v}\left(h_{\beta}(t)\right)=|t|_{v}^{2 s_{1}}
$$

with a complex number $s_{1}$ ( $s_{1}$ depends on $v$ ). Then the matrix

$$
t_{v}^{\wedge}=\left(\begin{array}{ll}
q_{v}^{-s_{1}} & 0 \\
0 & q_{v}^{s_{1}}
\end{array}\right)
$$

determines a semi-simple conjugacy class $\alpha_{\hat{v}}$ in $S L_{2}(\mathbb{C})$ (see the introduction). $q_{v}$ is the number of elements in the residue field.

Let $\rho$ be a four dimensional irreducible representation of $S L_{2}(\mathbb{C})$ (e.g. restriction of the third symmetric power of the standard representation of $G L_{2}(\mathbb{C})$ ). In fact up to equivalence $\rho$ is unique. The highest weight for $\rho$ is equal to $3 \delta$ where $\delta$ is the fundamental weight for $S L_{2}(\mathbb{C})$. Then the local $L$-function $L\left(s, \rho,{ }^{\circ} \pi_{v}\right)$ defined by Langlands in [10] is equal to

$$
\begin{equation*}
L\left(s, \rho,{ }^{\circ} \pi_{v}\right)=\operatorname{det}\left(I-\rho\left(t_{v}^{\wedge}\right) q_{v}^{-\delta}\right)^{-1} . \tag{5.1}
\end{equation*}
$$

In fact

$$
\begin{equation*}
L\left(s, \rho,{ }^{\circ} \pi_{v}\right)=\prod_{i=1,2,4,5}\left(I-q_{v}^{\left\langle\beta, \gamma_{i}\right\rangle s_{1}} \cdot q_{v}^{-s}\right)^{-1} \tag{5.2}
\end{equation*}
$$

with notation as in Lemma 4.6 (cf. [11]).
Let us assume that

$$
\omega=\prod_{v} \omega_{v}
$$

is an unramified quasi-character. This means that there exists a complex number $s_{2}$ such that

$$
\omega\left(\left(\begin{array}{ll}
z & 0 \\
0 & z
\end{array}\right)\right)=|z|^{s_{2}} \quad\left(z \in \mathbb{A}^{*}\right)
$$

Then

$$
\begin{equation*}
L\left(s, \rho,{ }^{\circ} \pi_{v}\right)=L\left(s+3 s_{2}, \rho^{*},{ }^{\circ} \pi_{v}^{*}\right), \tag{5.3}
\end{equation*}
$$

where $\rho^{*}$ is the third symmetric power representation of $G L_{2}(\mathbb{C})$ (the associated group for $G L_{2}$ ) and ${ }^{\circ} \pi_{v}^{*}$ is defined to be the $v$-th component of ${ }^{\circ} \pi^{*}$, the space generated by $\varphi^{*}$ (see §2).

To proceed, we need the following lemma which is a simple consequence of Lemma 2.1.

Lemma. 5.1: Let $\underline{E}_{x}(s ; \tilde{\varphi} ; g ; P)$ be the Fourier coefficient of $E(s ; \tilde{\varphi} ; g ; P)$ defined by (2.9), then for $\operatorname{Re}(s)<-1 / 2$

$$
\underline{E}_{X}(s ; \tilde{\varphi} ; g ; P)=\bigotimes_{v} w_{s, v}\left(g_{v}\right)
$$

with $g=\left(g_{v}\right)$ in $G_{A}$.

Remark: For almost all $v, g_{v}$ is in $K_{v}$ and $\underline{w}_{s, v}$ is a class-one function; thus for such places, we have

$$
\underline{w}_{s, v}\left(g_{v}\right)=\underline{w}_{s, v}(e)
$$

and the tensor product is defined without ambiguity.
The space of $\tilde{\sigma}_{v}$ (resp. its contragredient) is generated by the vectors of the form $\otimes_{v} e_{v}$ (resp. $\otimes_{v} \tilde{e}_{v}$ ) with $e_{v}$ (resp. $\tilde{e}_{v}$ ) in the space of $\tilde{\sigma}_{v}$ (resp. its contragredient) where for almost all $v,\left\langle e_{v}, \tilde{e}_{v}\right\rangle$ is equal to 1 . We put

$$
\left\langle\bigotimes_{v} e_{v}, \bigotimes_{v} \tilde{e}_{v}\right\rangle=\prod_{v}\left\langle e_{v}, \tilde{e}_{v}\right\rangle .
$$

We shall fix two such vectors $\otimes_{v} e_{v}$ and $\otimes_{v} \tilde{e}_{v}$ for which

$$
\left\langle e_{v}, \tilde{e}_{v}\right\rangle=1
$$

whenever ${ }^{\circ} \pi_{v}$ is a class-one representation. We set

$$
E_{\chi}(s ; \tilde{\varphi} ; g ; P)=\prod_{v}\left\langle\underline{w}_{s, v}\left(g_{v}\right) e_{v}, \tilde{e}_{v}\right\rangle
$$

for $g=\left(g_{v}\right)$ in $G_{A}$.
Now we shall explain the analytic continuation of $E_{\chi}(s ; \tilde{\varphi} ; g ; P)$ as a function of $s$. First we shall mention a result of W. Casselman and J.A. Shalika [2] which unfortunately has not been published yet. We use

$$
L\left(1_{v}, s\right)=\left(1-q_{v}^{-s}\right)^{-1}
$$

to denote the local Hecke $L$-function attached to the trivial character $1_{v}$ of $F_{v}^{*}$. Let $\tilde{\rho}$ denote the contragredient of $\rho$. As before put

$$
w_{s, v}(g)=\left\langle\underline{w}_{s, v}(g) e_{v}, \tilde{e}_{v}\right\rangle
$$

Theorem 5.2: (W. Casselman-J.A. Shalika). Assume $v$ is unramified; then

$$
\begin{equation*}
w_{s, v}(e) \cdot L\left(1_{v},-10 s+1\right) \cdot L\left(-5 s+1, \tilde{\rho},{ }^{\circ} \pi_{v}\right)={ }^{\circ} w_{v}(e) \tag{5.4}
\end{equation*}
$$

where $L\left(-5 s+1, \tilde{\rho},{ }^{\circ} \pi_{v}\right)$ is defined by (5.1).
Now as in introduction, let $S$ be the finite set of ramified places (including infinite places); then it follows from [10] that the product

$$
\begin{equation*}
L_{S}\left(s, \tilde{\rho},{ }^{\circ} \pi\right)=\prod_{v \in S} L\left(s, \tilde{\rho},{ }^{\circ} \pi_{v}\right) \tag{5.5}
\end{equation*}
$$

is convergent for $\operatorname{Re}(s)$ sufficiently large. We put

$$
\begin{equation*}
L_{S}(1, s)=\prod_{v \notin S} L\left(1_{v}, s\right) \tag{5.6}
\end{equation*}
$$

which is again convergent for large values of $\operatorname{Re}(s)$. We have

Corollary 5.3: As a function of $s$, the product (which converges for $\operatorname{Re}(s)$ sufficiently small)

$$
E_{x}(s ; \tilde{\varphi} ; g ; P) \cdot L_{S}(1,-10 s+1) \cdot L_{s}\left(-5 s+1, \tilde{\rho},{ }^{\circ} \pi\right)
$$

can be continued to an entire function on the whole complex plane.
Proof: This is an easy consequence of Theorem 5.2, Corollary 3.4 and the relation

$$
\prod_{v}{ }^{\circ} w_{v}(e)={ }^{\circ} w(e)
$$

Remark: Corollary 5.3 may also be considered as a consequence of the last section in [7].

As a consequence of relation (2.7) and Corollary 5.3 we have:

Theorem 5.4: As a function of $s$, the Fourier coefficient $E_{x}(s ; \tilde{\varphi} ; g ; P)$ can be continued to a meromorphic function on the whole complex plane. Furthermore

$$
\begin{equation*}
E_{\chi}(-s ; M(s) \tilde{\varphi} ; g ; P)=E_{x}(s ; \tilde{\varphi} ; g ; P) \tag{5.7}
\end{equation*}
$$

Remark: For $\operatorname{Re}(s)>\frac{1}{2}$, we have

$$
\begin{equation*}
E_{\chi}(-s ; M(s) \tilde{\varphi} ; g ; P)=\prod_{v}\left\langle\underline{w}_{s, v}^{\prime}\left(g_{v}\right) e_{v}, \tilde{e}_{v}\right\rangle \tag{5.8}
\end{equation*}
$$

with $\underline{w}_{s, v}^{\prime}$ defined as in §4.
Theorem 5.5: At each place $v$, the linear functional $\lambda^{\prime}$ defined by (3.4) is non-zero. Therefore there exists a complex function

$$
c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)
$$

meromorphic in $s$, such that

$$
\begin{equation*}
w_{s, v}(g)=c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right) w_{s, v}^{\prime}(g) \quad\left(g \in G_{v}\right) \tag{5.9}
\end{equation*}
$$

Furthermore $c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$ depends only on the class of ${ }^{\circ} \pi_{v}$.

Proof: For $\operatorname{Re}(s)>\frac{1}{2}$, it follows from Corollary 3.5, Corollary 4.11, functional equation (5.7), and the relation (5.8) that $\lambda^{\prime}$ is non-zero.

The existence of $c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$ follows from Theorem 2 of [14]. It clearly depends only upon the class of ${ }^{\circ} \pi_{v}$. Since $w_{s, v}(g)$ and $w_{s, v}^{\prime}(g)$ are at most meromorphic, it follows that $c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$ is a meromorphic function of $s$.

Corollary 5.6: As a function of $s, c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$ is holomorphic whenever $\lambda^{\prime}$ is non-zero.

Proof: For a fixed $s$, there exists a function $w_{s, v}^{\prime}$ for which $w_{s, v}^{\prime}\left(g_{0}\right)$ is not zero for some $g_{0}$ in $G_{v}$. Then if $c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$ had a pole at $s$, it would appear as a pole for $w_{s, v}\left(g_{0}\right)$ which is a contradiction to Corollary 3.4 .

Remark: Since $w_{s, v}^{\prime}$ has already some poles coming from the intertwining operators, it is not true that $c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$ has no zeros.

Let ${ }^{\circ} \chi$ be the restriction of $\chi$ to ${ }^{\circ} U_{\Lambda}$, then by the assumption on $\chi$, we may assume that $c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$ depends on ${ }^{\circ} \chi_{v}$ and write $c\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right)$ for $c\left(s, \rho,{ }^{\circ} \pi_{v}, \chi_{v}\right)$.

Proposition 5.7: Assume ${ }^{\circ} \pi_{v}$ can be realized as a quotient of the space of the left ${ }^{\circ} K_{v}$-finite functions ${ }^{\circ} f_{v}$ on ${ }^{\circ} G_{v}$ which satisfy

$$
{ }^{\circ} f_{v}(g t u)=\eta_{v}(t)^{\circ} \delta_{v}^{-1 / 2}(t) f_{v}^{\circ}(g) \quad\left(t \in{ }^{\circ} T_{v}, u \in{ }^{\circ} U_{v}, g \in{ }^{\circ} G_{v}\right),
$$

where $\eta_{v}$ is a quasi-character of ${ }^{\circ} T_{v}$ (this includes all the irreducible and admissible representations of the Hecke algebra of ${ }^{\circ} G_{v}$ except when $v$ is finite and the representation is supercuspidal). Then

$$
\begin{equation*}
c\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right)=\prod_{i=1}^{5}\left|\mu_{\alpha, v}\right|^{-\left\langle\nu, \gamma_{i}\right\rangle s} \gamma\left(\eta_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle},-\left\langle\nu, \gamma_{i}\right\rangle s\right) \tag{5.10}
\end{equation*}
$$

with $\mu_{\alpha, v}, \nu, \gamma$ and the ordering defined as before.

Proof: This is a simple consequence of Corollary 4.8.

Corollary 5.8: As a function of $s$, the product

$$
\prod_{i=1}^{5} \gamma\left(\eta_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle},-\left\langle\nu, \gamma_{i}\right\rangle s\right)
$$

is holomorphic whenever $\lambda^{\prime}$ is non-zero.
Now, for $v \in S$, put

$$
\begin{equation*}
\gamma\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right)=c\left(\frac{s-1}{5}, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right) \cdot \gamma^{-1}\left(1_{v}, 2-2 s\right), \tag{5.11}
\end{equation*}
$$

where $\gamma\left(1_{v}, 2-2 s\right)$ is defined by (4.16), and define

$$
\begin{equation*}
\gamma_{S}\left(s, \rho,{ }^{\circ} \pi\right)=\prod_{v \in S} \gamma\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right) \tag{5.12}
\end{equation*}
$$

This product, as we shall see, is independent of ${ }^{\circ} \chi$. The main result of this paper is the following theorem.

Theorem 5.9 (functional equation). Let ${ }^{\circ} \pi$ be a cusp form on $P G L_{2}(\mathbb{A})$. Denote by $S$ the corresponding set of ramified places. Define $\gamma_{S}\left(s, \rho,{ }^{\circ} \pi\right)$ and $L_{S}\left(s, \rho,{ }^{\circ} \pi\right)$ as before. Then

$$
\begin{equation*}
L_{S}\left(1-s, \tilde{\rho},{ }^{\circ} \pi\right)=\gamma_{S}\left(s, \rho,{ }^{\circ} \pi\right) L_{S}\left(s, \rho,{ }^{\circ} \pi\right) . \tag{5.13}
\end{equation*}
$$

In particular

$$
\prod_{v \in S} \gamma\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right)
$$

is independent of ${ }^{\circ} \chi$.

Proof: For $\operatorname{Re}(s)$ sufficiently small,

$$
\begin{equation*}
E_{\chi}(s ; \tilde{\varphi} ; g ; P) L_{S}(1,-10 s) \cdot L_{S}\left(-5 s, \tilde{\rho},{ }^{\circ} \pi\right) \tag{5.13.1}
\end{equation*}
$$

is equal to

$$
\begin{equation*}
\prod_{v \in S} w_{s, v}\left(g_{v}\right) \cdot \prod_{v \in S} w_{s, v}\left(g_{v}\right) \cdot L\left(1_{v},-10 s\right) \cdot L\left(-5 s, \tilde{\rho},{ }^{\circ} \pi_{v}\right) \tag{5.13.2}
\end{equation*}
$$

For any $v$, infinite or finite in which case we assume that ${ }^{\circ} \pi_{v}$ is not supercuspidal, Proposition 5.7 implies

$$
\begin{aligned}
w_{s, v}\left(g_{v}\right) & \prod_{i=1}^{s} L\left(\eta_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle s},-\left\langle\nu, \gamma_{i}\right\rangle s\right) \\
= & \prod_{i=1}^{s}\left|\mu_{\alpha, v}\right|^{-\left\langle\nu, \gamma_{i}\right\rangle s} \cdot \epsilon\left(\eta_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle},-\left\langle\nu, \gamma_{i}\right\rangle s\right) \\
& \cdot \prod_{i=1}^{s} L\left(\eta_{v}^{\left\langle\beta, \gamma_{i}\right\rangle}, 1+\left\langle\nu, \gamma_{i}\right\rangle s\right) \cdot w_{s, v}^{\prime}\left(g_{v}\right) .
\end{aligned}
$$

Here $\epsilon(\theta, s)$ (local root number), for a quasi-character $\theta$ of $F_{v}^{*}$, is defined by

$$
\epsilon(\theta, s)=\gamma(\theta, s) \frac{L(\theta, s)}{L\left(\theta^{-1}, 1-s\right)}
$$

Using (5.13.3) for unramified places, (5.13.2) can be written as follows:

$$
\begin{aligned}
\prod_{v \in S} w_{s, v}^{\prime}\left(g_{v}\right) \cdot \prod_{v \in S} c\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right) \cdot \prod_{v \in S} w_{s, v}^{\prime}\left(g_{v}\right) \cdot L_{S}(1,1+10 s) \cdot & L_{S}(1 \\
& \left.+5 s, \rho,{ }^{\circ} \pi\right)
\end{aligned}
$$

If we use analytic continuation of (5.13.1) to the large values of $\operatorname{Re}(s)$,
we conclude that
(5.13.4)

$$
\begin{gathered}
\prod_{v} w_{s, v}^{\prime}\left(g_{v}\right) \cdot \prod_{v \in S} c\left(s, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right) \cdot L_{S}(1,1+10 s) \cdot L_{S}\left(1+5 s, \rho,{ }^{\circ} \pi\right) \\
=\prod_{v} w_{s, v}^{\prime}\left(g_{v}\right) \cdot L_{S}(1,-10 s) \cdot L_{S}\left(-5 s, \tilde{\rho},{ }^{\circ} \pi\right)
\end{gathered}
$$

It is well known that

$$
L_{S}(1,-10 s)=\prod_{v \in S} \gamma\left(1_{v},-10 s\right) \cdot L_{S}(1,1+10 s)
$$

Now, if we cancel

$$
\prod_{v} w_{s, v}^{\prime}\left(g_{v}\right) \cdot L(1,-10 s)
$$

from both sides of (5.13.4), and change $s$ to $s-1 / 5$ we get (5.13) and therefore the theorem.

We put

$$
\begin{equation*}
\gamma\left(s, \rho^{*},{ }^{\circ} \pi_{v}^{*},{ }^{\circ} \chi_{v}\right)=\gamma\left(s+3 s_{2}, \rho,{ }^{\circ} \pi_{v},{ }^{\circ} \chi_{v}\right) \tag{5.14}
\end{equation*}
$$

and

$$
\begin{equation*}
\gamma_{S}\left(s, \rho^{*},{ }^{\circ} \pi^{*}\right)=\gamma_{S}\left(s+3 s_{2}, \rho,{ }^{\circ} \pi\right) . \tag{5.15}
\end{equation*}
$$

Then we have:

Corollary 5.10: Let ${ }^{\circ} \pi^{*}$ be a cusp form on $G L_{2}(\mathbb{A})$. Assume its restriction to the center of $G L_{2}(\mathbb{A})$ is unramified. Denote by $S$ the corresponding set of ramified places. Define $\gamma_{s}\left(s, \rho^{*},{ }^{\circ} \pi^{*}\right)$ as above and put

$$
L_{S}\left(s, \rho^{*},{ }^{\circ} \pi^{*}\right)=\prod_{v \in S} L\left(s, \rho^{*},{ }^{\circ} \pi_{v}^{*}\right)
$$

where the local factors are defined as before. Then

$$
\begin{equation*}
L_{S}\left(1-s, \tilde{\rho}^{*},{ }^{\circ} \pi^{*}\right)=\gamma_{S}\left(s, \rho^{*},{ }^{\circ} \pi^{*}\right) L_{S}\left(s, \rho^{*},{ }^{\circ} \pi^{*}\right) \tag{5.16}
\end{equation*}
$$

where $\tilde{\rho}^{*}$ denotes the contragredient representation of $\rho^{*}$.

Proposition 5.11: Let $v$ be any place, infinite or finite in which case we assume that ${ }^{\circ} \pi_{v}^{*}$ is not supercuspidal. Then

$$
\begin{aligned}
\gamma\left(s, \rho^{*},{ }^{\circ} \pi_{v}^{*},{ }^{\circ} \chi_{v}\right)= & \prod_{i=1,2,4,5}\left|\mu_{\alpha, v}\right|^{-\left\langle\nu, \gamma_{i}\right\rangle\left(s+3 s_{2}\right)} \\
& \cdot \gamma\left(\eta_{v}^{-\left\langle\beta, \gamma_{i}\right\rangle},-\left\langle\nu, \gamma_{i}\right\rangle\left(s+3 s_{2}\right)\right) .
\end{aligned}
$$

## Appendix

In [13], Langlands proved that, when $v$ is an infinite place and $\operatorname{Re}(s)$ is sufficiently small, the image of

$$
\Pi_{v}=\operatorname{ind}_{P_{v} \uparrow G_{v}}^{\circ} \pi_{v} \otimes \delta_{P, v}^{s}
$$

under the intertwining operator

$$
\int_{N_{v}} f_{v}\left(g n w_{0}\right) d n \quad\left(f_{v} \in V_{v}, g \in G_{v}\right),
$$

which is a subspace of

$$
\Pi_{v}^{\prime}=\operatorname{ind}_{P_{v} \uparrow G_{v}}{ }^{\circ} \pi_{v}^{\prime} \otimes \delta_{P, v}^{-s}
$$

is in fact irreducible.
Let us assume the following conjecture.

Conjecture: Assume $v$ is a finite place; then for $\operatorname{Re}(s)$ sufficiently large, the image of $\Pi_{v}$ under the above intertwining operator is irreducible (by means of analytic continuation).

Corollary to Conjecture: Assume ${ }^{\circ} \pi_{v}$ is a component of a cusp form on $\mathrm{PGL}_{2}(\mathbb{A})(v$ a finite place). Then for $\operatorname{Re}(s)$ sufficiently large, the image of the space of $\Pi_{v}$ under the above intertwining operator, is equivalent to the unique non-degenerate quotient of $\Pi_{v}$.

At any rate we should have:

Proposition: Assume ${ }^{\circ} \pi_{v}$ is a component of a cusp form on $P G L_{2}(A)$ ( $v$ finite or infinite). Then for $\operatorname{Re}(s)$ sufficiently large, the image of the space of $\Pi_{v}$ under the above intertwining operator is non-degenerate.

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