# Boundary Trace of Positive Solutions of Nonlinear Elliptic Inequalities 

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#### Abstract

We develop a new method for proving the existence of a boundary trace, in the class of Borel measures, of nonnegative solutions of $-\Delta u+g(x, u) \geq 0$ in a smooth domain $\Omega$ under very general assumptions on $g$. This new definition which extends the previous notions of boundary trace is based upon a sweeping technique by solutions of Dirichlet problems with measure boundary data. We also prove a boundary pointwise blow-up estimate of any solution of such inequalities in terms of the Poisson kernel. If the nonlinearity is very degenerate near the boundary, for example if $g(x, u) \approx \exp \left(-\rho_{\partial \Omega}^{-1}(x)\right) u^{q}$, we exhibit a new full boundary blow-up phenomenon.


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## Introduction

Let $\Omega$ be a bounded open domain in $\mathbb{R}^{N}$ with a $C^{2}$ boundary $\partial \Omega$. This paper is concerned with the study of the generalized boundary value problem for the equation

$$
\begin{equation*}
-\Delta u+g(x, u)=0 \quad \text { in } \Omega \tag{0.1}
\end{equation*}
$$

where $(x, r) \mapsto g(x, r)$ is a continuous function defined on $\Omega \times \mathbb{R}$, nondecreasing in the $r$ variable, and nonnegative if $r \geq 0$. When $g(r)=r^{q}$ this problem has been thoroughly investigated with a probabilistic approach by Le Gall [19], [20] in the case $N=2=q$, then by Marcus and Véron [21], [22], [23] in the general case $q>1, N>1$ by analytic tools. Related studies were carried on by Dynkin and Kuznetsov [10], [11] with a mixing of probabilistic and analytic methods. In [16] the same problem is investigated with $g(r)=\exp (r)$. In all those cases, the boundary trace dichotomy argument is settled upon duality techniques which were first introduced by Baras and Pierre [1], but in the case of general nonlinearity, this method fails.

In [27] an new approach of the boundary trace is developed for positive solutions of (0.1). This approach is settled upon two ingredients:
I - The coerciveness, which asserts that the set of nonnegative solutions of (0.1) is bounded in the local uniform topology upon $C(\Omega)$.

II - The strong-barrier property which is the property that for any boundary point $z$ and for any $r>0$, small enough, there exist supersolutions $\varphi$ of (0.1) in $\Omega \cap B_{r}(z)$ with infinite value on $\partial B_{r}(z) \cap \Omega$, and zero value on $\partial \Omega \cap B_{r}(z)$.
When $g$ depends only of $r$, those two notions coincide thanks to the OssermanKeller condition,

$$
\begin{equation*}
\int_{a}^{\infty} \frac{d s}{\sqrt{G(s)}}<\infty, \quad \text { for any } a>0 \tag{0.2}
\end{equation*}
$$

where $G(s)=\int_{0}^{s} g(t) d t$. The same equivalence holds if $\inf _{x \in \Omega} g(x, r)=g(r)$, and $g$ satisfies $(0.2)$. Moreover, in these cases, the local uniform upper bound of any positive solution of (0.1) achieves the following form

$$
\begin{equation*}
u(x) \leq \psi_{g}\left(\rho_{\partial \Omega}(x)\right), \quad \forall x \in \Omega \tag{0.3}
\end{equation*}
$$

where $\psi_{g}(t)=\int_{t}^{\infty} \frac{d s}{\sqrt{2 G(s)}}$, and $\left.\rho_{\partial \Omega}(x)\right)=\operatorname{dist}(x, \partial \Omega)$.
If the strong barrier property is uniform with respect to $z \in \partial \Omega$, it implies the coerciveness, but when $\lim _{\left.\rho_{\partial \Omega}(x)\right) \rightarrow 0} g(x, r)=0$ for any $r>0$, (we say that the nonlinearity degenerates near the boundary), the reverse implication may not hold. However, if

$$
g(x, r) \geq \rho_{\partial \Omega}^{\alpha}(x) r^{q} \quad \forall(x, r) \in \Omega \times \mathbb{R}_{+},
$$

for some $\alpha>0$ and $q>1$, it is proved in [27] that the equivalence still holds.
We adopt here a different point of view in connecting the existence of a boundary trace and the question of solving a Dirichlet problem with measure data. If $\mu \in \mathfrak{M}(\partial \Omega)$, the set of Radon measures on $\partial \Omega$, and $(x, r) \mapsto g(x, r)$ is a continuous function defined on $\Omega \times \mathbb{R}$, a function $u$ defined in $\Omega$ is a solution of

$$
\begin{align*}
-\Delta u+g(x, u) & =0 \text { in } \Omega, \\
u & =\mu \text { on } \partial \Omega, \tag{0.4}
\end{align*}
$$

if $u \in L^{1}(\Omega), g(., u) \in L^{1}\left(\Omega, \rho_{\partial \Omega} d x\right)$ and

$$
\begin{equation*}
\int_{\Omega}(-u \Delta \zeta+g(x, u) \zeta) d x=-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \mu(y) \tag{0.5}
\end{equation*}
$$

for any $\zeta \in C_{c}^{1,1}(\bar{\Omega})$, the subspace of $C^{1}(\bar{\Omega})$ functions with Lipschitz continuous gradient and zero value on $\partial \Omega$. If $g(., x)$ is nondecreasing, this solution is unique whenever it exists and we denote it by $u=u_{\mu}$, since

$$
\begin{equation*}
\left\|u_{\mu}-u_{\mu^{*}}\right\|_{L^{1}(\Omega)}+\left\|\rho_{\partial \Omega}\left(g\left(u_{\mu}, .\right)-g\left(u_{\mu^{*}}, .\right)\right)\right\|_{L^{1}(\Omega)} \leq C\left\|\mu-\mu^{*}\right\|_{\mathfrak{M}(\partial \Omega)} . \tag{0.6}
\end{equation*}
$$

The mapping $\mu \mapsto u_{\mu}$ is nondecreasing, moreover if $g(x, 0)=0, \mu \geq 0 \Longrightarrow$ $u_{\mu} \geq 0$. Conditions for existence are various.

Let $\mathcal{G}_{0}$ be the set of continuous functions $g$ defined in $\Omega \times \mathbb{R}$ such that $g(x, 0)=0$ and $r \mapsto g(r, x)$ is nondecreasing for any $x \in \Omega$, and $(x, y) \mapsto$ $P(x, y)$ be the Poisson kernel in $\Omega \times \partial \Omega$. If $\mu \in \mathfrak{M}(\partial \Omega)$, we denote by $\mathbb{P}_{\mu}$ its Poisson's potential. If $g \in \mathcal{G}_{0}$ we say that $\mu$ is $g$-admissible if

$$
\begin{equation*}
\int_{\Omega} g\left(x, \mathbb{P}_{|\mu|}(x)\right) \rho_{\partial \Omega}(x)<\infty \tag{0.7}
\end{equation*}
$$

It is proved in [27] that problem (0.4) is uniquely solvable if $\mu$ is $g$-admissible.
However to check this condition on every measure might be far out of reach and a more tractable condition is introduced. We denote $\mathcal{H} \mathcal{G}_{0}$ the subset of $g \in \mathcal{G}_{0}$ such that there exist two continuous, nondecreasing and nonnegative functions $h$ and $f$ defined on $\mathbb{R}_{+}$, such that

$$
\begin{gather*}
0 \leq|g(x, r)| \leq h\left(\rho_{\partial \Omega}(x)\right) f(|r|), \quad \forall(x, r) \in \Omega \times \mathbb{R}, \\
\int_{0}^{1} h(s) f\left(\sigma s^{1-N}\right) s^{N} d s<\infty, \quad \forall \sigma \geq 0,  \tag{0.8}\\
\text { either } h(s)=s^{\alpha}, \text { for some } \alpha \geq 0, \quad \text { or } f \text { is convex } .
\end{gather*}
$$

In the first section of this article we prove the following.
If $g \in \mathcal{H} \mathcal{G}_{0}$, then for any $\mu \in \mathfrak{M}(\partial \Omega)$, problem (0.4) admits a unique solution $u_{\mu}$. Moreover the problem is stable, in the sense that if $\left\{\mu_{n}\right\} \subset \mathfrak{M}(\partial \Omega)$ converges to $\mu$ in the weak sense of measures on $\partial \Omega$, the corresponding solutions $\left\{u_{\mu_{n}}\right\}$ converge to $u_{\mu}$, locally uniformly in $\Omega$.

In the second section we introduce a new definition of the boundary trace for nonnegative solutions of elliptic inequalities.

$$
\begin{equation*}
-\Delta u+g(x, u) \geq 0 \quad \text { in } \Omega \tag{0.9}
\end{equation*}
$$

which extends the previous results concerning equations. A key observation for defining this notion is a supremum technique introduced by Richard and Véron [30] in the study of isolated singularities of elliptic inequalities. Following [27] we say that a function $g \in \mathcal{G}_{0}$ is positively subcritical if for any $\mu \in \mathfrak{M}_{+}(\partial \Omega)$, problem (0.4) admits a solution $u_{\mu}$ (unique and nonnegative). If $u \in C(\Omega)$ such that $\Delta u \in L_{\mathrm{loc}}^{1}(\Omega)$ is a nonnegative solution of (0.9), then $w_{\mu}=\min \left\{u, u_{\mu}\right\}$ satisfies ( 0.9 ), and it admits a boundary trace $\gamma_{u}(\mu) \in \mathfrak{M}_{+}(\partial \Omega)$. Furthermore
$\mu \mapsto \gamma_{u}(\mu)$ is nondecreasing, concave if $r \mapsto g(x, r)$ is convex. Therefore the formula

$$
\begin{equation*}
v=\sup _{\mu \in \mathfrak{M}_{+}(\partial \Omega)} \gamma_{u}(\mu) \tag{0.10}
\end{equation*}
$$

defines a Borel measure $v=\operatorname{Tr}_{\partial \Omega}^{e}(u)$ on $\partial \Omega$, that we call the extended boundary trace of $u$. This measure may not be an outer regular one except in some particular cases. A particularly important case deals with the choice $\mu=\lambda \delta_{a}$ for $\lambda>0, a \in \partial \Omega$. The corresponding solution $u_{\lambda \delta_{a}}$ is called a fundamental solution. In such a case the boundary trace of $w_{\lambda \delta_{a}}$ is a measure concentrated at $a$, that we denote $\tilde{\gamma}_{u}(a, \lambda) \delta_{a}$. The mapping $\lambda \mapsto \tilde{\gamma}_{u}(a, \lambda)$ is a nondecreasing on $\mathbb{R}_{+}$, and satisfies

$$
0 \leq \tilde{\gamma}_{u}(a, \lambda) \leq \lambda, \quad \forall \lambda \geq 0, \forall a \in \partial \Omega .
$$

We define

$$
\tilde{\gamma}_{u}(a)=\lim _{\lambda \rightarrow \infty} \tilde{\gamma}_{u}(a, \lambda),
$$

and denote by $\mathcal{A}(u)$ the set of atoms of $u$,

$$
\mathcal{A}(u)=\left\{a \in \partial \Omega: \gamma_{u}(a)>0\right\} .
$$

The regular set $\mathcal{R}(u)$ of $u$ is the relatively open subset of the boundary points $a$ with the property that there exists a relatively neighborhood of $a, \mathcal{O} \subset \partial \Omega$ such that

$$
\sum_{\omega \in \mathcal{O}} \tilde{\gamma}_{u}(\omega)<\infty .
$$

The singular set $\mathcal{S}(u)$ of $u$ is the closed subset of the boundary points $a$ with the property that for any relatively open neighborhood $\mathcal{O} \subset \partial \Omega$ of $a$, there holds

$$
\sum_{\omega \in \mathcal{O}} \tilde{\gamma}_{u}(\omega)=\infty
$$

Those two definitions extend the classical notions of regular or singular sets of the boundary trace of the solution of an equation (see [22], [24]).
We prove in particular

$$
\nu(a)=\tilde{\gamma}_{u}(a), \quad \forall a \in \partial \Omega,
$$

and the equivalence between
(i) $\nu(\mathcal{O})=\infty$, for any relatively open neighborhood $\mathcal{O} \subset \partial \Omega$ of $a$, and
(ii) $u_{\infty, a} \leq u$,
under a general stability condition which holds in particular if $g \in \mathcal{H} \mathcal{G}_{0}$.

As for $u_{\infty, a}$, different features may occur, in particular,

- $u_{\infty, a} \equiv+\infty$, the full blow-up case.
- $u_{\infty, a}(x)<\infty$ for any $x \in \Omega$, but $\lim _{\rho_{\partial \Omega}(x) \rightarrow 0} u_{\infty, a}(x)=\infty$, the uniform boundary blow-up case.
- $u_{\infty, a}(x)<\infty$ for any $x \in \bar{\Omega} \backslash\{a\}$, and $\lim _{x \rightarrow a} u_{\infty, a}(x)=\infty$ (non-tangential limit), the strong isolated singularity case.

Using precise pointwise estimates of positive super-harmonic functions near the boundary and the sweeping of any positive solution $u$ of (0.9) by the solutions of (0.4) with Dirac masses as boundary data, we prove that for any $a \in \partial \Omega$,

$$
x \mapsto|x-a|^{N-1} u(x)
$$

converges in measure on the set $\{\sigma=(y-a) /|y-a|: y \in \Omega\}$ to $C(N) \tilde{\gamma}_{u}(a)$ as $x \rightarrow a$, where $C(N)$ is some positive constant depending only on the dimension.

If $g(x, r)$ satisfies

$$
\begin{equation*}
g(x, r) \geq \tilde{h}(x) \tilde{g}(r), \forall(x, r) \in \Omega \times \mathbb{R}_{+}, \tag{0.11}
\end{equation*}
$$

where $\tilde{h} \in C(\Omega)$ takes positive values, and $\tilde{g}$ is nondecreasing and satisfies (0.2), there exists a maximal solution $U_{M}$ to ( 0.1 ) in $\Omega$ (actually the global positivity of $\tilde{h}$ can be weakened, since the positivity near $\partial \Omega$ is sufficient for the existence of $U_{M}$ ). In that case (ii) implies

$$
u_{\infty, a}(x) \leq u(x) \leq U_{M}(x)
$$

which rules out the full blow-up case. The nature of $u_{\infty, a}$ depends strongly on $\tilde{h}$ and $\tilde{g}$. For example if it is assumed that $\tilde{h}$ is a positive constant, it follows from the method of construction of maximal solutions that the uniform boundary blow-up case does not hold, and we are left with the strong isolated singularity case. However, this situation also holds even if $\tilde{h}$ depends truly of $x$. It is proved in [27] that if
(0.12) $h(x)=\rho_{\partial \Omega}^{\alpha}(x)$, with $\alpha>-2$ and $1<q<(N+\alpha+1) /(N-1)=q_{c}(\alpha)$, the strong isolated singularity case occurs. We prove here that if

$$
\begin{equation*}
g(x, r)=\exp \left(-1 / \rho_{\partial \Omega}(x)\right) r^{q}, \quad \text { with } q>1 \tag{0.13}
\end{equation*}
$$

the uniform boundary blow-up case occurs for any $a \in \partial \Omega$. In such a case, either the boundary trace is a bounded Borel measure, or $u \equiv U_{M}$.
A parabolic version of this phenomenon has been observed in [28].

When the nonlinearity is not degenerate in the sense that the function $g \in \mathcal{H} \mathcal{G}_{0}$ satisfies

$$
\begin{equation*}
0 \leq|g(x, r)| \leq f(|r|), \forall(x, r) \in \Omega \times \mathbb{R}, \text { and } \int_{0}^{1} f\left(s^{1-N}\right) s^{N} d s<\infty \tag{0.14}
\end{equation*}
$$

where $f$ is a continuous nondecreasing function defined on $\mathbb{R}_{+}$, we recover the classical definition of the boundary trace in the class of outer regular Borel measures. More precisely, if $u$ is a nonnegative solution of ( 0.9 ) with extended boundary trace $\nu$, then for any point $a \in \partial \Omega$ the following dichotomy occurs: either
(i) $a \in \mathcal{S}(u)$ and for any $\mathcal{O} \in \mathcal{N}_{a}$ (the set of its relatively open neighborhoods $\mathcal{O} \subset \partial \Omega), \nu(\mathcal{O})=\infty$. This is equivalent to

$$
\lim _{t \rightarrow 0} \int_{\mathcal{O}_{t}} u(y) d S_{t}=\infty, \quad \forall \mathcal{O} \in \mathcal{N}_{a}
$$

where $\mathcal{O}_{t}$ is the subset of points in $\Omega$ at distance $t>0$ from $\partial \Omega$, with projection in $\mathcal{O}$ and $d S_{t}$ the induced ( $N-1$ )-dimensional Hausdorff measure, or
(ii) $a \in \mathcal{R}(u)$, there exists $\mathcal{O} \in \mathcal{N}_{a}$ such that $v(\mathcal{O})<\infty$ and for any $\mathcal{O}^{\prime} \subset$ $\overline{\mathcal{O}}^{\prime} \subset \mathcal{O}$

$$
\sup _{t \in\left(0, \beta_{0}\right]} \int_{\mathcal{O}_{t}} u(y) d S<\infty
$$

Furthermore, for any $\phi \in C_{c}(\mathcal{R}(u))$

$$
\lim _{t \rightarrow 0} \int_{\mathcal{O}_{t}} u(y) \phi d S_{t}=\int_{\mathcal{R}(u)} \phi d v
$$

Our paper is organised as follows: In Section 1 we study the boundary value problem with Radon measures. In Section 2 we define and study the extended boundary trace of nonnegative solutions of inequalities. In Section 3 we give a boundary pointwise estimate for solutions of inequalities. In Section 4 we give properties of the boundary trace when the nonlinearity is not degenerate at the boundary. In Section 5 we study different examples of limit of a fundamental solution when the mass goes to infinity.

## 1. - Measure boundary data

Throughout this section, $\Omega$ is a bounded domain with a $C^{2}$ boundary $\partial \Omega$ and $\rho_{\partial \Omega}(x)=\operatorname{dist}(x, \partial \Omega)$. We put
(1.1) $\mathcal{G}_{0}=\{g \in C(\Omega \times \mathbb{R})$ s.t. $g(x, 0)=0$ and $r \mapsto g(x, r)$ nondecreasing, $\forall x \in \Omega\}$.

We denote by $C_{c}^{1,1}(\bar{\Omega})$, the subspace of $C^{1}(\bar{\Omega})$-functions with Lipschitz continuous gradient and zero value on $\partial \Omega, \mathfrak{M}(\partial \Omega)$ the space of Radon measures on $\partial \Omega$, and $\mathfrak{M}_{+}(\partial \Omega)$ its positive cone. If $P(x, y)$ is the Poisson kernel in $\Omega \times \partial \Omega$, the Poisson potential of $\mu$ denoted by $\mathbb{P}_{\mu}$ is defined by

$$
\begin{equation*}
\mathbb{P}_{\mu}(x)=\int_{\partial \Omega} P(x, y) d \mu(y), \quad \forall x \in \Omega \tag{1.2}
\end{equation*}
$$

The next variant of Herglotz' theorem is due to Brezis [5] (see [32] for a proof).
Lemma 1.1. Let $f \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$ and $\varphi \in L^{1}(\partial \Omega)$. Then there exists $a$ unique $u \in L^{1}(\Omega)$ such that

$$
\begin{equation*}
-\int_{\Omega} u \Delta \zeta=\int_{\Omega} f \zeta d x-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} \varphi d S \tag{1.3}
\end{equation*}
$$

for any $\zeta \in C_{c}^{1,1}(\bar{\Omega})$. Moreover there exists $C=C(\Omega)>0$ such that

$$
\begin{equation*}
\|u\|_{L^{1}(\Omega)} \leq C\left(\left\|\rho_{\partial \Omega} f\right\|_{L^{1}(\Omega)}+\|\varphi\|_{L^{1}(\partial \Omega)}\right) \tag{1.4}
\end{equation*}
$$

Finally u satisfies

$$
\begin{equation*}
-\int_{\Omega}|u| \Delta \zeta+\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}}|\varphi| d S \leq \int_{\Omega} f \zeta \operatorname{sgn}(u) d x \tag{1.5}
\end{equation*}
$$

and

$$
\begin{equation*}
-\int_{\Omega} u^{+} \Delta \zeta+\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} \varphi^{+} d S \leq \int_{\Omega} f \zeta \operatorname{sgn}^{+}(u) d x \tag{1.6}
\end{equation*}
$$

for any $\zeta \in C_{c}^{1,1}(\bar{\Omega}), \zeta \geq 0$.
Definition 1.2. Let $\mu \in \mathfrak{M}(\partial \Omega)$. A function $u \in L^{1}(\Omega)$ is a solution of

$$
\begin{align*}
-\Delta u+g(x, u) & =0 \text { in } \Omega \\
u & =\mu \text { on } \partial \Omega \tag{1.7}
\end{align*}
$$

if $g(., u) \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$ and

$$
\begin{equation*}
\int_{\Omega}(-u \Delta \zeta+g(x, u) \zeta) d x=-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \mu(y) \tag{1.8}
\end{equation*}
$$

for any $\zeta \in C_{c}^{1,1}(\bar{\Omega})$.
Uniqueness is a straightforward consequence of (1.5). In the case where $g(x, r)=g(r)$ and $\mu \in L^{1}(\partial \Omega)$, existence of a solution to (1.7) is due to Brezis [5]. If $g$ is continuous in $\bar{\Omega} \times \mathbb{R}$, the proof of Brezis result goes through without any difficulty. If $x \mapsto g(x, r)$ is merely continuous in $\Omega$ and unbounded near $\partial \Omega$, problem (1.7) may not have any solution, even with very regular data $\mu$.

Definition 1.2. A measure $\mu \in \mathfrak{M}(\partial \Omega)$ is g-admissible if

$$
\begin{equation*}
g\left(., \mathbb{P}_{|\mu|}\right) \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right) \tag{1.9}
\end{equation*}
$$

The two next results can be found in [27]

Proposition 1.1. Let $g \in \mathcal{G}_{0}$ and $\mu \in \mathfrak{M}(\partial \Omega)$ be $g$-admissible. Then problem (1.7) possesses a unique solution $u_{\mu}$. Moreover the mapping $\mu \mapsto u_{\mu}$ is increasing and continuous from $\mathfrak{M}(\partial \Omega)$ endowed with the total variation norm into $C(\Omega)$ with the local uniform topology.

Proposition 1.2. Let $g \in \mathcal{G}_{0}$ satisfy

$$
\begin{equation*}
g(., c) \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right), \quad \forall c \in \mathbb{R} . \tag{1.10}
\end{equation*}
$$

Then for any $\mu \in L^{1}(\partial \Omega)$, problem (0.4) admits a unique solution.
Let $d H_{N-1}$ be the ( $\mathrm{N}-1$ )-dimensional Hausdorff measure. If $\mu \in \mathfrak{M}(\partial \Omega)$ we denote by $\mu_{R}\left(\right.$ resp. $\mu_{s}$ ) its regular (resp. singular) part in the Lebesgue decomposition

$$
\mu=\mu_{R}+\mu_{s}
$$

where $\mu_{R} \prec d H_{N-1}$ ) and $\mu_{s} \perp \mu_{R}$. A variant of the next result can be found in [25].

Proposition 1.3. Assume $g \in \mathcal{G}_{0}$ satisfies (1.10) and

$$
\begin{equation*}
|g(x, 2 r)| \leq K(|g(x, r)|+\ell(x), \forall(x, r) \in \Omega \times \mathbb{R} \tag{1.11}
\end{equation*}
$$

for some fixed $K>0$ and $\ell \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. If $\mu \in \mathfrak{M}(\partial \Omega), \mu=\mu_{R}+\mu_{s}$ and $\mu_{s}$ is $g$-admissible, then the conclusions of Proposition 1.1 still hold.

Proof. First notice that relation (1.11), called the $\Delta_{2}$-condition, implies

$$
\begin{equation*}
\left|g\left(x, r+r^{\prime}\right)\right| \leq K\left(|g(x, r)|+\left|g\left(x, r^{\prime}\right)\right|\right)+\ell(x), \forall\left(x, r, r^{\prime}\right) \in \Omega \times \mathbb{R}^{2} \tag{1.12}
\end{equation*}
$$

Step 1. Assume that $\mu$ is nonnegative, and so are $\mu_{R}$ and $\mu_{s}$. Let $\left\{\mu_{R, n}\right\}$ be a sequence of smooth functions on $\partial \Omega$ converging to $\mu_{R}$ in $L^{1}(\partial \Omega)$. Since $\mathbb{P}_{\mu_{R, n}}$ is bounded, $g(x,$.$) is nondecreasing and (1.10) is satisfied, it follows from$ (1.12) that $\mu_{n}=\mu_{R, n}+\mu_{S}$ is $g$-admissible. Let $u_{n}$ and $v_{n}$ be the solutions of (0.4) with respective measure boundary data $\mu_{n}$ and $\mu_{R, n}$. Applying (1.5) with $u=v_{n}-v_{p}, f=-g\left(., v_{n}\right)+g\left(., v_{p}\right)$ and $\zeta=\mathbb{P}_{1}$, we obtain

$$
\begin{equation*}
\left\|v_{n}-v_{p}\right\|_{L^{1}(\Omega)}+\left\|\rho_{\partial \Omega}\left(g\left(., v_{n}\right)-g\left(., v_{p}\right)\right)\right\|_{L^{1}(\Omega)} \leq C\left\|\mu_{R, n}-\mu_{R, p}\right\|_{L^{1}(\partial \Omega)} . \tag{1.13}
\end{equation*}
$$

Thus $\left\{v_{n}\right\}$ and $\left\{g\left(., v_{n}\right\}\right.$ converge to $v$ and $g(., v)$, respectively in $L^{1}(\Omega)$ and $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. Furthermore $v$ is the solution of (0.4) with measure boundary data $\mu_{R}$. Because $v_{n}+\mathbb{P}_{\mu_{S}}$ is a supersolution of (0.1) with boundary data $\mu_{n}$, there also holds

$$
\begin{equation*}
0 \leq u_{n} \leq v_{n}+\mathbb{P}_{\mu_{S}} \tag{1.14}
\end{equation*}
$$

thus $\left\{u_{n}\right\}$ is uniformly integrable in $L^{1}(\Omega)$, and also locally compact in the $C_{\text {loc }}^{1}(\Omega)$ topology, by the elliptic equations regularity theory. Since inequality (1.12) leads to

$$
\begin{equation*}
0 \leq g\left(x, u_{n}\right) \leq K\left(g\left(x, \mathbb{P}_{\mu_{S}}\right)+g\left(x, v_{n}\right)\right)+\ell(x) \tag{1.15}
\end{equation*}
$$

it follows from the assumption on $\mu_{S}$ that the sequence $\left\{g\left(., u_{n}\right)\right\}$ is uniformly integrable in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. By the Vitali theorem $u_{n_{k}} \rightarrow u$ and $g\left(., u_{n_{k}}\right) \rightarrow$ $g(., u)$ respectively in $L^{1}(\Omega)$ and $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$ and $u$ is the solution of (0.4).

STEP 2. Let $\tilde{\mu}_{R, n}$ and $\bar{\mu}_{R, n}$ be smooth $L^{1}$-approximations of $\mu_{R}^{+}$and $\mu_{R}^{-}$, and denote by $u_{n}, \tilde{v}_{n}$ and $\bar{v}_{n}$ the solutions of (0.4) with respective measure boundary data $\mu_{n}=\tilde{\mu}_{R, n}+\bar{\mu}_{R, n}+\mu_{S}, \tilde{\mu}_{R, n}$ and $-\bar{\mu}_{R, n}$. By monotonicity and (1.12) there holds

$$
\bar{v}_{n}-\mathbb{P}_{\mu_{S}^{-}} \leq u_{n} \leq \tilde{v}_{n}+\mathbb{P}_{\mu_{S}^{+}}
$$

and
$K\left(g\left(x,-\mathbb{P}_{\mu_{S}^{-}}\right)+g\left(x, \bar{v}_{n}\right)\right)-\ell(x) \leq g\left(x, u_{n}\right) \leq K\left(g\left(x, \mathbb{P}_{\mu_{S}^{+}}\right)+g\left(x, \tilde{v}_{n}\right)\right)+\ell(x)$.
Since $\bar{v}_{n}, \tilde{v}_{n}, g\left(., \tilde{v}_{n}\right)$ and $g\left(., \bar{v}_{n}\right)$ inherit the uniform integrability properties of Step 1, we conclude again by the Vitali theorem.

Let $u_{\mu}$ denote the solution of (0.4) with boundary data $\mu$. The $g$ admissibility condition on $\mu$ does not imply the weak continuity of the mapping $\mu \mapsto u_{\mu}$, thus a more uniform assumption is needed.

Definition 1.3. We denote by $\mathcal{H} \mathcal{G}_{0}$ the subset of the $g \in \mathcal{G}_{0}$ such that there exist two continuous, nondecreasing and nonnegative functions $h$ and $f$ defined on $\mathbb{R}_{+}$, with the property

$$
\begin{gather*}
0 \leq|g(x, r)| \leq h\left(\rho_{\partial \Omega}(x)\right) f(|r|), \quad \forall(x, r) \in \Omega \times \mathbb{R}  \tag{1.16}\\
\int_{0}^{1} h(s) f\left(\sigma s^{1-N}\right) s^{N} d s<\infty, \quad \forall \sigma \geq 0 \tag{1.17}
\end{gather*}
$$ either $h(s)=s^{\alpha}$, for some $\alpha \geq 0$, or $f$ is convex .

The main result of this section is an existence and stability theorem which extends a previous one due to Gmira and Véron [15]. The technique involved is based upon the use of Marcinkiewicz spaces first introduced by Benilan and Brezis [3], [6] for solving semilinear equations with right-hand side measure.

Theorem 1.1. Let $g \in \mathcal{H} \mathcal{G}_{0}$. Then any measure $\mu$ on $\partial \Omega$ is $g$-admissible. Moreover, if $\left\{\mu_{n}\right\} \subset \mathfrak{M}(\partial \Omega)$ converges to $\mu$ in the weak sense of measures, the corresponding solutions $u_{\mu_{n}}$ of (0.4) with boundary data $\mu_{n}$ converge to $u_{\mu}$ locally uniformly in $\Omega$ and $g\left(., u_{\mu_{n}}\right) \rightarrow g\left(., u_{\mu}\right)$ in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$.

Proof. Step 1. The Marcinkiewicz space framework. For any nonnegative locally bounded Borel measure $\beta$ in $\Omega$ and real number $p>1$, we denote

$$
\begin{equation*}
\left.M^{p}(\Omega ; d \beta)=\left\{v \in L_{\mathrm{loc}}^{1}(\Omega ; d \beta):\|v\|_{M^{p}(\Omega ; d \beta)}\right\}<\infty\right\} \tag{1.19}
\end{equation*}
$$

where
(1.20) $\|v\|_{M^{p}(\Omega ; d \beta)}=\inf \left\{c \in[0, \infty]\right.$ s.t. $\int_{K}|v| d \beta \leq c\left(\int_{K} d \beta\right)^{1-1 / p} \forall K \subset \Omega, K$ Borel $\}$.

Besides the classical imbedding of $M^{p}(\Omega ; d \beta)$ into $L_{\text {loc }}^{\tilde{p}}(\Omega ; d \beta)$ for any $1 \leq$ $\tilde{p}<p$, the next inequality plays an important role

$$
\begin{equation*}
\left.C(p)\|v\|_{M^{p}(\Omega ; d \beta)}\right\} \leq \sup _{\lambda>0}\left\{\lambda^{p} \int_{\{|u|>\lambda\}} d \beta\right\} \leq\|v\|_{M^{p}(\Omega ; d \beta)} . \tag{1.21}
\end{equation*}
$$

Moreover the following estimates are proved in [16]: there exists $K=K(\Omega)>0$ such that for any $v \in \mathfrak{M}(\partial \Omega)$,

$$
\begin{align*}
\left\|\mathbb{P}_{\nu}\right\|_{M^{(N+1) /(N-1)}\left(\Omega ; \rho_{\partial \Omega} d x\right)} & \leq K\|v\|_{L^{1}(\partial \Omega)}  \tag{1.22}\\
\left\|\mathbb{P}_{\nu}\right\|_{M^{N /(N-1)}(\Omega)} & \leq K\|v\|_{L^{1}(\partial \Omega)}  \tag{1.23}\\
\left\|\mathbb{P}_{\nu}\right\|_{L^{\infty}\left(\Omega_{r}^{c}\right)} & \leq K r^{1-N}\|\nu\|_{L^{1}(\partial \Omega)} \tag{1.24}
\end{align*}
$$

where $\Omega_{r}=\left\{x \in \Omega: \rho_{\partial \Omega}(x) \leq r\right\}$, and $\Omega_{r}^{c}=\Omega \backslash \bar{\Omega}_{r}=\left\{x \in \Omega: \rho_{\partial \Omega}(x)>r\right\}$.
Step 2. We claim that there exist two positive constants $C_{1}=C_{1}(\Omega)$ and $C_{2}=C_{2}(N)$ such that for any $a \in \partial \Omega$ and $\lambda>0$

$$
\begin{equation*}
\beta_{a}(\lambda)=\int_{\Gamma_{a}(\lambda)} h(\rho(x)) \rho(x) d x \leq C_{2} \int_{0}^{\left(C_{1} / \lambda\right)^{1 /(N-1)}} h(s) s^{N} d s \tag{1.25}
\end{equation*}
$$

where

$$
\Gamma_{a}(\lambda)=\{x \in \Omega: P(x, a)>\lambda\} .
$$

Since

$$
\begin{equation*}
C_{1}^{-1} \rho_{\partial \Omega}(x)|x-a|^{-N} \leq P(x, a) \leq C_{1} \rho_{\partial \Omega}(x)|x-a|^{-N} \tag{1.26}
\end{equation*}
$$

for some $C_{1}>0$ independent of $(x, a) \in \Omega \times \partial \Omega$,

$$
\Gamma_{a}(\lambda) \subset\left\{x \in \Omega: \rho_{\partial \Omega}(x)|x-a|^{-N}>\lambda / C_{1}\right\} \subset \Omega \cap B_{r_{\lambda}}(a)
$$

with $r_{\lambda}=\left(C_{1} / \lambda\right)^{1 /(N-1)}$. Since $h$ is nondecreasing,

$$
\int_{\Gamma_{a}(\lambda)} h(\rho(x)) \rho(x) d x \leq \int_{B_{r_{\lambda}}(a)} h(|x|) \rho(|x|) d x=\left|S^{N-1}\right| \int_{0}^{r_{\lambda}} h(s) s^{N} d s,
$$

which implies (1.25).
Step 3. Let $G \subset \Omega$ be a Borel subset, then for any $m>0, \lambda>0$, and $a \in \partial \Omega$ there holds

$$
\begin{align*}
& \int_{G} h\left(\rho_{\partial \Omega}\right) f\left(m P(., a) \rho_{\partial \Omega} d x \leq f(\lambda) \int_{G} h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x\right. \\
& \quad+C_{3} m^{(N+1) /(N-1)} \int_{\lambda}^{\infty} f(s) h\left(\left(m C_{1} / s\right)^{1 /(N-1)}\right) s^{-2 N /(N-1)} d s \tag{1.27}
\end{align*}
$$

with $C_{3}=C_{3}(N)>0$. Actually,

$$
\begin{aligned}
\int_{G} h\left(\rho_{\partial \Omega}(x)\right) f\left(m P(x, a) \rho_{\partial \Omega}(x) d x=\right. & \int_{G \cap\{P(x, a) \leq \lambda / m\}} h\left(\rho_{\partial \Omega}(x)\right) f\left(m P(x, a) \rho_{\partial \Omega}(x) d x\right. \\
& +\int_{G \cap \Gamma_{a}(\lambda / m)} h\left(\rho_{\partial \Omega}(x)\right) f\left(m P(x, a) \rho_{\partial \Omega}(x) d x .\right.
\end{aligned}
$$

Since $f$ is nondecreasing,

$$
\int_{G \cap\{P(x, a) \leq \lambda / m\}} f\left(m P(x, a) h\left(\rho_{\partial \Omega}(x)\right) \rho_{\partial \Omega}(x) d x \leq f(\lambda) \int_{G} h\left(\rho_{\partial \Omega}(x)\right) \rho_{\partial \Omega}(x) d x .\right.
$$

Moreover

$$
\int_{G \cap \Gamma_{a}(\lambda / m)} h\left(\rho_{\partial \Omega}(x)\right) f\left(m P(x, a) \rho_{\partial \Omega}(x) d x \leq-\int_{\lambda / m}^{\infty} f(m s) d \beta_{a}(s)\right.
$$

But

$$
-\int_{\lambda / m}^{\infty} f(m s) d \beta_{a}(s)=f(\lambda) \beta_{a}(\lambda)+\int_{\lambda / m}^{\infty} \beta_{a}(s) d f(m s)
$$

Using (1.25) in Step 2 infers

$$
-\int_{\lambda / m}^{\infty} f(m s) d \beta_{a}(s) \leq f(\lambda) \beta_{a}(\lambda)+C_{2} \int_{\lambda / m}^{\infty} \int_{0}^{\left(C_{1} / s\right)^{1 /(N-1)}} h(\tau) \tau^{N} d \tau d f(m s)
$$

Since

$$
\begin{aligned}
& \int_{\lambda / m}^{\infty} \int_{0}^{\left(C_{1} / s\right)^{1 /(N-1)}} h(\tau) \tau^{N} d \tau d f(m s)=-f(\lambda) \int_{0}^{\left(m C_{1} / \lambda\right)^{1 /(N-1)}} h(s) s^{N} d s \\
& \quad+\frac{C_{1}^{(N+1) /(N-1)}}{N-1} \int_{\lambda / m}^{\infty} h\left(\left(C_{1} / s\right)^{1 /(N-1)}\right) s^{-2 N /(N-1)} f(m s) d s
\end{aligned}
$$

(1.27) follows by change of variable, with $C_{3}=C_{2} C_{1}^{(N+1) /(N-1)}$. It is important to notice that this integral is convergent because of (1.16).

Step 4. Suppose $f$ is convex, then for any $\mu \in \mathfrak{M}^{+}(\partial \Omega)$ with total mass $m$ and any Borel subset $G \subset \Omega$, there holds

$$
\begin{align*}
& \int_{G} f\left(\mathbb{P}_{\mu}\right) h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x \leq f(\lambda) \int_{G} h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x \\
& \quad+C_{3} m^{(N+1) /(N-1)} \int_{\lambda}^{\infty} f(s) h\left(\left(m C_{1} / s\right)^{1 /(N-1)}\right) s^{-2 N /(N-1)} d s \tag{1.28}
\end{align*}
$$

First, let us assume that

$$
\mu=m \sum_{i=1}^{k} \theta_{i} \delta_{a_{i}}
$$

for some $a_{i} \in \partial \Omega$ and $\theta_{i}>0$ with $\sum_{i=1}^{k} \theta_{i}=1$. Then

$$
\mathbb{P}_{\mu}(x)=m \sum_{i=1}^{k} \theta_{i} P\left(x, a_{i}\right)
$$

Since

$$
\begin{gathered}
f\left(\mathbb{P}_{\mu}(x)\right)=f\left(m \sum_{i=1}^{k} \theta_{i} P\left(x, a_{i}\right)\right) \leq \sum_{i=1}^{k} \theta_{i} f\left(m P\left(x, a_{i}\right)\right. \\
\int_{G} f\left(\mathbb{P}_{\mu}\right) h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x \leq \sum_{i=1}^{k} \theta_{i} \int_{G} f\left(m P\left(x, a_{i}\right) h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x .\right.
\end{gathered}
$$

Therefore (1.28) follows from (1.27).
For a general nonnegative measure $\mu$ with total mass $m$, there exists a sequence of finite combinations of positive Dirac measures $\mu_{n}$ with same total mass converging to $\mu$ in the weak sense of measures. Then $\mathbb{P}_{\mu_{n}}$ converges to $\mathbb{P}_{\mu}$ locally uniformly in $\Omega$ and in $L^{p}(\Omega)$ for any $1 \leq p<N /(N-1)$. Thus (1.28) follows by the Fatou's lemma.

Step 5. Suppose $h(s)=s^{\alpha}$, then for any $\mu \in \mathfrak{M}^{+}(\partial \Omega)$ with total mass $m$ and any Borel subset $G \subset \Omega$, there holds

$$
\begin{align*}
& \int_{G} f\left(\mathbb{P}_{\mu}\right) h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x \\
& \quad \leq f(\lambda) \int_{G} \rho_{\partial \Omega}^{1+\alpha} d x+C_{7} m^{(N+1+\alpha) /(N-1)} \int_{\lambda}^{\infty} s^{-(2 N+\alpha) /(N-1)} f(s) d s \tag{1.29}
\end{align*}
$$

By Step 2 there exists $C_{4}=C_{4}(\Omega, \alpha)>0$ such that

$$
\beta_{a}(\lambda) \leq C_{4} \lambda^{(N+1+\alpha) /(N-1)}
$$

for any $\lambda>0$. Thus we derive an estimate of $P(., a)$ in $M^{(N+1+\alpha) /(N-1)}$,

$$
\begin{equation*}
\int_{G} P(x, a) \rho_{\partial \Omega}^{1+\alpha} d x \leq C_{5}\left(\int_{G} \rho_{\partial \Omega}^{1+\alpha} d x\right)^{(2+\alpha) /(N+1+\alpha)} \tag{1.30}
\end{equation*}
$$

Therefore

$$
\begin{aligned}
\left.\int_{G} \mathbb{P}_{\mu}(x)\right) \rho_{\partial \Omega}^{1+\alpha} d x & =\int_{\partial \Omega} d \mu(a) \int_{G} P(x, a) \rho_{\partial \Omega}^{1+\alpha} d x \\
& \leq\left(\int_{\partial \Omega} d \mu(a)\right) \max _{a \in \partial \Omega} \int_{G} P(x, a) \rho_{\partial \Omega}^{1+\alpha} d x
\end{aligned}
$$

From this estimate follows

$$
\begin{equation*}
\left.\int_{G} \mathbb{P}_{\mu}(x)\right) \rho_{\partial \Omega}^{1+\alpha} d x \leq C_{5}\|\mu\|_{L^{1}(\partial \Omega)}\left(\int_{G} \rho_{\partial \Omega}^{1+\alpha} d x\right)^{(2+\alpha) /(N+1+\alpha)} \tag{1.31}
\end{equation*}
$$

Now

$$
\begin{align*}
\int_{G} f\left(\mathbb{P}_{\mu}\right) h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x & =\int_{G} f\left(\mathbb{P}_{\mu}\right) \rho_{\partial \Omega}^{1+\alpha} d x \\
& \leq f(\lambda) \int_{G} \rho_{\partial \Omega}^{1+\alpha} d x+\int_{\left\{\mathbb{P}_{\mu}>\lambda\right\}} f\left(\mathbb{P}_{\mu}\right) \rho_{\partial \Omega}^{1+\alpha} d x \tag{1.32}
\end{align*}
$$

But

$$
\int_{\left\{\mathbb{P}_{\mu}>\lambda\right\}} f\left(\mathbb{P}_{\mu}\right) \rho_{\partial \Omega}^{1+\alpha} d x=-\int_{\lambda}^{\infty} f(s) d \beta^{\mu}(s)
$$

where

$$
\beta^{\mu}(s)=\int_{\Gamma^{\mu}(s)} \rho_{\partial \Omega}^{1+\alpha} d x \quad \text { with } \quad \Gamma^{\mu}(s)=\left\{x \in \Omega: \mathbb{P}_{\mu}(x)>s\right\}
$$

Moreover

$$
\beta^{\mu}(s) \leq C_{6} m^{(N+1+\alpha) /(N-1)} \lambda^{-(N+1+\alpha) /(N-1)}
$$

by (1.30), with $G=\Gamma^{\mu}(\lambda)$ and $C_{6}=C_{5}^{(N+1+\alpha) /(N-1)}$. Therefore

$$
\begin{align*}
-\int_{\lambda}^{\infty} f(s) d \beta^{\mu}(s)= & f(\lambda) \beta^{\mu}(\lambda)+\int_{\lambda}^{\infty} \beta^{\mu}(s) d f(s) \\
\leq & C_{6} m^{(N+1+\alpha) /(N-1)} \lambda^{-(N+1+\alpha) /(N-1)} f(\lambda) \\
& +C_{6} m^{(N+1+\alpha) /(N-1)} \int_{\lambda}^{\infty} s^{-(N+1+\alpha) /(N-1)} d f(s),  \tag{1.33}\\
\leq & C_{7} m^{(N+1+\alpha) /(N-1)} \int_{\lambda}^{\infty} s^{-(2 N+\alpha) /(N-1)} f(s) d s
\end{align*}
$$

where $C_{7}=(N+1+\alpha) C_{6} m^{(N+1+\alpha) /(N-1)}$. Combining (1.32) and (1.33) yields (1.29).

If we take $G=\Omega$ in (1.27) and (1.29) we derive that (1.9) holds with $\mu$ and we conclude by Proposition 1.1. However those two estimates are much more powerfull since they leads to uniform-integrability properties.

Step 6. Put $v_{n}=\mathbb{P}_{\left|\mu_{n}\right|}$ and $v=\mathbb{P}_{|\mu|}$. Then

$$
\begin{equation*}
0 \leq\left|u_{n}\right| \leq v_{n} \text { and } 0 \leq|u| \leq v . \tag{1.34}
\end{equation*}
$$

The fact that $u_{n}$ is locally bounded in $\Omega$ independently of $n$ follows from (1.24). Since $g$ is continuous, $g\left(., u_{n}\right)$ remains also locally bounded in $\Omega$. By the elliptic equations regularity theory there exists a subsequence $\left\{u_{n_{k}}\right\}$ and a $C^{1}(\Omega)$-function $u$ such that $u_{n_{k}} \rightarrow u$ in the $C_{\mathrm{loc}}^{1}(\Omega)$-topology. This clearly implies that $u$ solves (0.1) in $\Omega$.

Step 7. We claim that $u$ is a solution of (0.4) with $\mu$ as boundary data. By the definition of the Marcinkiewicz norm,

$$
\begin{aligned}
\int_{\Omega}\left|u_{n_{k}}-u\right| d x & =\int_{\Omega_{r}^{c}}\left|u_{n_{k}}-u\right| d x+\int_{\Omega_{r}}\left|u_{n_{k}}-u\right| d x \\
& \leq \int_{\Omega_{r}^{c}}\left|u_{n_{k}}-u\right| d x+2\left\|u_{n_{k}}-u\right\|_{M^{N /(N-1)}(\Omega)}\left(\text { meas } \Omega_{r}\right)^{1 / N}
\end{aligned}
$$

But $\left\|u_{n_{k}}-u\right\|_{M^{N /(N-1)}(\Omega)}$ remains bounded independently of $n_{k}$ by (1.22). Thus $u_{n_{k}} \rightarrow u$ in $L^{1}(\Omega)$. In order to prove that $g\left(., u_{n_{k}}\right) \rightarrow g(., u)$ in $L^{1}\left(\Omega, \rho_{\partial \Omega} d x\right)$ put $m_{n}=\int_{\partial \Omega} d\left|\mu_{n}\right|$ and let $G \subset \Omega$ be a Borel set. Because of (1.16) and (1.34) there holds

$$
\left|g\left(., u_{n}\right)\right| \leq f\left(v_{n}\right) h\left(\rho_{\partial \Omega}\right)
$$

If $f$ is convex (Step 4) it follows

$$
\begin{align*}
& \int_{G}\left|g\left(., u_{n}\right)\right| \rho_{\partial \Omega} d x \leq f(\lambda) \int_{G} h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x \\
& \quad+C_{3} m_{n}^{(N+1) /(N-1)} \int_{\lambda}^{\infty} f(s) h\left(\left(m_{n} C_{1} / s\right)^{1 /(N-1))} s^{-2 N /(N-1)} d s\right. \tag{1.35}
\end{align*}
$$

If $h(s)=s^{\alpha}($ Step 5$)$, then

$$
\begin{align*}
& \int_{G}\left|g\left(., u_{n}\right)\right| \rho_{\partial \Omega} d x  \tag{1.36}\\
& \quad \leq f(\lambda) \int_{G} \rho_{\partial \Omega}^{1+\alpha} d x+C_{7} m_{n}^{(N+1+\alpha) /(N-1)} \int_{\lambda}^{\infty} s^{-(2 N+\alpha) /(N-1)} f(s) d s .
\end{align*}
$$

Because $\left\{m_{n}\right\}$ is bounded, for any $\varepsilon>0$, we first choose $\lambda>0$ large enough so that, for any $n \in \mathbb{N}$,

$$
C_{3} m_{n}^{(N+1) /(N-1)} \int_{\lambda}^{\infty} f(s) h\left(\left(m_{n} C_{1} / s\right)^{1 /(N-1)}\right) s^{-2 N /(N-1)} d s \leq \varepsilon / 2
$$

in the case $f$ is convex, or

$$
C_{7} m_{n}^{(N+1+\alpha) /(N-1)} \int_{\lambda}^{\infty} s^{-(2 N+\alpha) /(N-1)} f(s) d s \leq \varepsilon / 2
$$

in the case $h(s)=s^{\alpha}$. Then we take meas. $G$ small enough so that

$$
f(\lambda) \int_{G} h\left(\rho_{\partial \Omega}\right) \rho_{\partial \Omega} d x, \leq \varepsilon / 2
$$

and we conclude that

$$
\int_{G}\left|g\left(., u_{n}\right)\right| \rho_{\partial \Omega} d x \leq \varepsilon
$$

independently of $n$. Therefore $\left\{g\left(., u_{n}\right\}\right.$ is uniformly integrable for the measure $\rho_{\partial \Omega} d x$. Since $g$ is continuous and $u_{n} \rightarrow u$ in $\Omega$, it follows that $g\left(., u_{n}\right) \rightarrow$ $g(., u)$ in $L^{1}\left(\Omega, \rho_{\partial \Omega} d x\right)$. Letting $n \rightarrow \infty$ in the integral formulation of $u_{n}$ for (0.4) implies that $u$ solves (0.4) with $\mu$ as boundary data.

The next stability result is a straightforward extension of the previous result
Proposition 1.4. Let $(x, r) \mapsto g_{n}(x, r)$ be a sequence of functions in $\in C(\Omega \times$ $\mathbb{R})$, nondecreasing with respect to $r$, vanishing at $r=0$ for any $x \in \Omega$, and satisfying (1.16)-(1.18) uniformly with respect to $n$. If there exists $g \in C(\Omega \times \mathbb{R})$ such that $g_{n}(x, r) \rightarrow g(x, r)$ pointwise in $\Omega \times \mathbb{R}$, then $g$ satisfies (1.16). Moreover, if $\mu_{n} \in \mathfrak{M}(\partial \Omega)$ converges to $\mu$ in the weak sense of measures, the sequence of solutions $u_{\mu_{n}, g_{n}}$ of

$$
\begin{aligned}
-\Delta u_{n}+g_{n}\left(x, u_{n}\right) & =0 \quad \text { in } \Omega \\
u_{n} & =\mu_{n} \text { on } \partial \Omega
\end{aligned}
$$

converges locally uniformly in $\Omega$ to the solution $u_{\mu}$ of (0.4).
Proof. By assumption

$$
0 \leq g_{n}(x, r) \leq h\left(\rho_{\partial \Omega}(x)\right) f(r), \forall n \in \mathbb{N}^{*}, \forall(x, r) \in \bar{\Omega} \times \mathbb{R}
$$

and (1.16) holds. Then $g$ satisfies the same upper bound. Since $0 \leq\left|u_{n}\right| \leq \mathbb{P}_{\left|\mu_{n}\right|}$, the inequalities (1.35)-(1.36) hold with $g$ replaced by $g_{n}$ which infers the uniform integrability of $\left\{g_{n}\left(., u_{n}\right)\right\}$ in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. The rest of the proof is similar to the one of Proposition 1.1.

## 2. - The extended boundary trace

If $\Omega$ is a $C^{3}$ bounded domain in $\mathbb{R}^{N}$ and $x \in \partial \Omega$, let $\mathbf{n}_{x}$ denote the normal pointing outward unit vector at $x$. Let us recall some notations and definitions from [25]. The mapping $\Pi$ from $\partial \Omega \times(0, \infty)$ into $\mathbb{R}^{N}$ is defined by

$$
\begin{equation*}
\Pi(x, t)=x-t \mathbf{n}_{x} \quad \forall(x, t) \in \partial \Omega \times(0,1) \tag{2.1}
\end{equation*}
$$

It is known that there exists $0<\beta_{0}$ such that $\Pi$ is a diffeomorphism from $\partial \Omega \times\left[0, \beta_{0}\right)$ onto

$$
\begin{equation*}
\Omega_{\beta_{0}}=\left\{x \in \Omega: \rho_{\partial \Omega}(x)<\beta_{0}\right\} . \tag{2.2}
\end{equation*}
$$

In particular, for any $t \in\left[0, \beta_{0}\right)$ the set

$$
\begin{equation*}
\Sigma_{t}=\left\{y=x-t \mathbf{n}_{x}: x \in \partial \Omega\right\} \tag{2.3}
\end{equation*}
$$

is diffeomorphic to $\partial \Omega=\Sigma_{0}$ (for the sake of simplicity, we shall denote $\left.\Sigma_{0}=\Sigma\right)$, and for any $y \in \Sigma_{t}, \rho_{\partial \Omega}(y)=t$. If $U \subset \partial \Omega$, we denote

$$
U_{t}=\left\{y=x-t \mathbf{n}_{x}: x \in U\right\}
$$

and if $\zeta$ is a function defined in $U_{t}$,

$$
\zeta_{t}(y)=\zeta(x) \text { for any } y=x-t \mathbf{n}_{x} \in U_{t}
$$

We denote by $\mathfrak{H}_{t}$ the mapping from $\Sigma_{t}$ to $\Sigma$ defined by $\mathfrak{H}_{t}(x)=\sigma(x)$ for $x \in \Sigma_{t}$. Thus $\mathfrak{H}_{t}^{-1}(x)=\Pi^{-1}(., t)$.

Given $t \in\left[0, \beta_{0}\right)$, a Borel measure $\mu$ and a function $f$ on $\Sigma_{t}$, we define a corresponding measure $\mu^{t}$ and function $f^{t}$ on $\Sigma$ by

$$
\begin{cases}\mu^{t}(E)=\mu\left(\mathfrak{H}_{t}^{-1}(E)\right), & \forall E \subset \Sigma, E \text { Borel, }  \tag{2.4}\\ f^{t}(\sigma)=f\left(\sigma-t \mathbf{n}_{\sigma}\right), & \forall \sigma \in \Sigma\end{cases}
$$

Then

$$
\mu \in \mathfrak{M}\left(\Sigma_{t}\right), f \in L^{1}\left(\Sigma_{t},|\mu|\right) \Longrightarrow\left\{\begin{array}{l}
f^{t} \in L^{1}\left(\Sigma,\left|\mu^{t}\right|\right)  \tag{2.5}\\
\int_{\Sigma_{t}} f d \mu=\int_{\Sigma} f^{t} d \mu^{t}
\end{array}\right.
$$

This section is devoted to the definition and properties of the notion of extended boundary trace for nonnegative solutions of

$$
\begin{equation*}
-\Delta u+g(x, u) \geq 0 \quad \text { in } \Omega \tag{2.6}
\end{equation*}
$$

and of the associated equation

$$
\begin{equation*}
-\Delta u+g(x, u)=0 \quad \text { in } \Omega \tag{2.7}
\end{equation*}
$$

Definition 2.1. I- We say that a function $g \in \mathcal{G}_{0}$ is positively subcritical if for any $\mu \in \mathfrak{M}_{+}(\partial \Omega)$, problem (0.4) admits a solution $u_{\mu}$.
II- The function $g$ is said positively subcritical and stable if $\mu_{n} \rightarrow \mu$ weakly in $\mathfrak{M}_{+}(\partial \Omega)$ implies
$u_{\mu_{n}} \rightarrow u_{\mu}$ locally uniformly in $\Omega$ and $g\left(., u_{\mu_{n}}\right) \rightarrow g\left(., u_{\mu}\right)$ in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$.

Remark 2.1. I- The convergence of $\mu_{n}$ to $\mu$ in the weak sense of measures implies that $\mathbb{P}_{\mu_{n}}$ converges to $\mathbb{P}_{\mu}$ in $L^{s}(\Omega)$ for any $1 \leq s<N /(N-1)$, by classical potential analysis. Because $0 \leq u_{\mu_{n}} \leq \mathbb{P}_{\mu_{n}}, u_{\mu_{n}} \rightarrow u_{\mu}$ in $L^{s}(\Omega)$ (by Vitali's theorem) and in the local uniform topology of $C(\Omega)$ ). If $\zeta \in C_{c}^{2}(\bar{\Omega})$, $\zeta \geq 0$, then

$$
\int_{\Omega}\left(-u_{\mu_{n}} \Delta \zeta+g\left(x, u_{\mu_{n}}\right) \zeta\right) d x=-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \mu_{n}
$$

Since $u_{\mu_{n_{k}}} \rightarrow \tilde{u}$ and $g\left(., u_{\mu_{n_{k}}}\right) \rightarrow g(., \tilde{u})$ locally uniformly in $\Omega$, Fatou's lemma infers

$$
\int_{\Omega}(-\tilde{u} \Delta \zeta+g(x, \tilde{u}) \zeta) d x \leq-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \mu
$$

and there always holds $\tilde{u} \leq u_{\mu}$. If it is assumed that $\tilde{u}=u_{\mu}$, then

$$
\lim _{n \rightarrow \infty} \int_{\Omega} g\left(x, u_{\mu_{n}}\right) \zeta d x=\int_{\Omega} g\left(x, u_{\mu}\right) \zeta d x
$$

and this convergence holds for any $\zeta \in C_{c}(\bar{\Omega})$ such that $\zeta / \rho_{\partial \Omega}$ is bounded. However it does not imply weak convergence in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. Notice that in such a case, weak convergence in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$ implies strong convergence by Dunford-Pettis and Vitali's theorems.
II- It follows by Theorem 1.1 that any $g \in \mathcal{H} \mathcal{G}_{0}$ is positively subcritical and stable.

Definition 2.2. Let $U$ be a relatively open subset of $\Sigma$ and $\mu \in \mathfrak{M}(U)$. We say that a function $v \in C(\Omega)$ admits $\mu$ for trace on $\mathcal{O}$, and we denote it by $\operatorname{Tr}_{U}(v)$, if

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{U_{t}} v(x) \phi(\sigma(x)) d S=\int_{U} \phi d \mu, \quad \forall \phi \in C_{c}(U) \tag{2.8}
\end{equation*}
$$

It is proved in [27, Corollary 1.3] that the solution $u_{\mu}$ of (0.4) admits $\mu$ for trace on $\partial \Omega$. The role of nonnegative super-solutions is enlighted by the next result.

Proposition 2.1. Let $g \in \mathcal{G}_{0}$ and $u \in C(\Omega)$ be a nonnegative solution of (2.6) in $\Omega$ such that $g(., u) \in L^{1}\left(\Omega, \rho_{\partial \Omega} d x\right)$. Then $u$ admits a boundary trace $\mu \in \mathfrak{M}_{+}(\partial \Omega)$. Moreover, if

$$
\begin{equation*}
u^{*}=\sup \{v \in C(\Omega): 0 \leq v \leq u, v \text { solution of }(2.7)\} \tag{2.9}
\end{equation*}
$$

then $u^{*}$ is a solution of (2.7) and

$$
\operatorname{Tr}_{\partial \Omega}(u)=\operatorname{Tr}_{\partial \Omega}\left(u^{*}\right) .
$$

Proof. Let $\psi=\mathbb{G}_{g(., u)}$ be the Green potential of the function $g(., u)$. Then $\psi+u$ is nonnegative and super-harmonic. Therefore $\psi+u$ admits a boundary trace belonging to $\mathfrak{M}_{+}(\partial \Omega)$. Thus the same holds for $u$ since $\psi$ vanishes on $\partial \Omega$ and the boundary trace of $u$ is a nonnegative Radon measure. The construction of $u^{*}$ is performed in considering the sequence of smooth domains $\Omega_{n}(n \geq 1)$, defined by

$$
\Omega_{n}=\left\{x \in \Omega: \rho_{\partial \Omega}(x)>\beta_{n}\right\}
$$

where $0<\beta_{n+1}<\beta_{n}<\beta_{0}$ and $\beta_{n} \rightarrow 0$ as $n \rightarrow \infty$. For each $n$ we denote by $u_{n}$ the solution of

$$
\begin{aligned}
-\Delta u_{n}+g\left(x, u_{n}\right) & =0 \text { in } \Omega_{n}, \\
u_{n} & =u \text { on } \partial \Omega_{n} .
\end{aligned}
$$

Since $u$ is a super solution, $u \geq u_{n} \geq 0$ and consequently the sequence $\left\{u_{n}\right\}$ is decreasing. Therefore $u^{*}$ exists as the decreasing limit of the $u_{n}$. By the regularity theory of elliptic equations, the convergence holds in $C_{\text {loc }}^{1}(\Omega)$ and $u^{*}$ satisfies (2.7). If $\tilde{u}$ is any nonnegative solution of (2.7) dominated by $u$, then $\tilde{u} \leq u^{*}=u_{n}$ on $\partial \Omega_{n}$, thus $\tilde{u} \leq u_{n}$ in $\Omega_{n}$. Letting $n \rightarrow \infty$ yields $\tilde{u} \leq u^{*}$. Clearly the correspondence $u \mapsto u^{*}$ which associates to a solution $u$ of (2.6) the largest solution of (2.7) dominated by $u$ inherits the following properties:

$$
\begin{equation*}
u_{1} \leq u_{2} \Longrightarrow u_{1}^{*} \leq u_{2}^{*} \tag{2.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(u^{*}\right)^{*}=u^{*} \tag{2.11}
\end{equation*}
$$

In the above construction, the boundary trace of $u$ plays no role. In order to prove that $\operatorname{Tr}_{\partial \Omega}\left(u^{*}\right)=\mu$, let $\zeta \in C_{c}^{1,1}(\bar{\Omega})$ and $\zeta_{n}$ be the solution of

$$
\begin{array}{cl}
-\Delta \zeta_{n}=-\Delta \zeta & \text { in } \Omega_{n} \\
\zeta_{n}=0 & \\
\text { on } \partial \Omega_{n}
\end{array}
$$

Although $\zeta_{n} \notin C_{c}^{1,1}\left(\bar{\Omega}_{n}\right), \zeta_{n}$ remains uniformly bounded in $C_{c}^{1, \alpha}\left(\bar{\Omega}_{n}\right)$ for any $\alpha \in(0,1), \Delta \zeta_{n}$ is bounded and

$$
\int_{\Omega_{n}}\left(-u_{n} \Delta \zeta_{n}+g\left(x, u_{n}\right) \zeta_{n}\right) d x=-\int_{\Sigma_{\beta_{n}}} \frac{\partial \zeta_{n}}{\partial \mathbf{n}} u(y) d S
$$

by approximation. We extend $u_{n}$ by putting the zero value outside $\bar{\Omega}_{n}$ and call $\tilde{u}_{n}$ this extension. For $\zeta_{n}$ we perform an extension by reflexion following the normal direction and define $\tilde{\zeta}_{n}$ thanks to the following formula

$$
\begin{aligned}
& \forall x \in \bar{\Omega} \backslash \Omega_{n} \text { with } \rho_{\partial \Omega}(x)=t_{x} \text { and } x=\sigma(x)-t_{x} \mathbf{n}_{\sigma(\mathbf{x})}, \\
& \tilde{\zeta}_{n}(x)=-\zeta_{n}\left(\sigma(x)-\left(2 \beta_{n}-t_{x}\right) \mathbf{n}_{\sigma(\mathbf{x})}\right)
\end{aligned}
$$

Notice that we have to assume $2 \beta_{n} \leq \beta_{0}$. It follows by the elliptic equations regularity theory that there exist a subsequence $\zeta_{n_{k}}$ and some $\tilde{\zeta} \in C_{c}^{1, \alpha}(\bar{\Omega})$ such that $\zeta_{n_{k}} \rightarrow \tilde{\zeta}$ in $C^{1}(\bar{\Omega})$. By the uniqueness of the solution of the Dirichlet problem $\tilde{\zeta}=\zeta$, and the wole sequence $\left\{\zeta_{n}\right\}$ is convergent. Moreover

$$
0 \leq \tilde{u}_{n}(x) \leq u(x) \Longrightarrow 0 \leq g\left(x, \tilde{u}_{n}\right) \leq g(x, u), \quad \forall x \in \Omega
$$

Thus the sequence $\left\{g\left(., \tilde{u}_{n}\right) \rho_{\partial \Omega}\right\}$ is uniformly integrable in $\Omega$. By Vitali's theorem

$$
\int_{\Omega}\left(-u^{*} \Delta \zeta+g\left(x, u^{*}\right) \zeta\right) d x=-\lim _{n_{k} \rightarrow 0} \int_{\Sigma_{\beta_{n_{k}}}} \frac{\partial \zeta_{n_{k}}}{\partial \mathbf{n}} u(y) d S=\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \mu
$$

This indicates that $u^{*}$ is a solution of (0.4) with boundary $\mu$ and therefore $\operatorname{Tr}_{\partial \Omega}\left(u^{*}\right)=\mu$.

The key observation on which is based the definition of the boundary trace is the following

Proposition 2.2. Let $g \in \mathcal{G}_{0}$ is positively subcritical and $u$ a nonnegative solution of (2.6) in $\Omega$. If $\mu \in \mathfrak{M}_{+}(\partial \Omega)$ set $w_{\mu}=\min \left(u, u_{\mu}\right)$. Then $w_{\mu}$ satifies

$$
\begin{equation*}
-\Delta w_{\mu}+g\left(., w_{\mu}\right) \geq 0 \quad \text { in } \Omega \tag{2.12}
\end{equation*}
$$

and there exists $\gamma_{u}(\mu) \in \mathfrak{M}_{+}(\partial \Omega)$ such that $\operatorname{Tr}_{\partial \Omega}\left(w_{\mu}\right)=\gamma_{u}(\mu)$. The mapping $\mu \mapsto \gamma_{u}(\mu)$ is nondecreasing and $0 \leq \gamma_{u}(\mu) \leq \mu$. Moreover, if for any $x \in \Omega$ the function $r \mapsto g(x, r)$ is convex on $\mathbb{R}_{+}$, the mapping $\mu \mapsto \gamma_{u}(\mu)$ is concave on $\mathfrak{M}_{+}(\partial \Omega)$.

Proof. Let $\delta>0$ and $p$ be the $C^{1,1}$ even convex function defined on $\mathbb{R}$ by

$$
p(t)= \begin{cases}|t|-\delta / 2 & \text { for }|t| \geq \delta \\ t^{2} / 2 \delta & \text { for }|t| \leq \delta\end{cases}
$$

Then $\omega_{\delta}=\frac{1}{2}\left(u+u_{\mu}-p\left(u-u_{\mu}\right)\right)$ satisfies

$$
\begin{aligned}
\Delta \omega_{\delta} & =\frac{1}{2}\left(\Delta u+\Delta u_{\mu}-p^{\prime}\left(u+u_{\mu}\right) \Delta\left(u-u_{\mu}\right)-p^{\prime \prime}\left(u-u_{\mu}\right)\left|D\left(u-u_{\mu}\right)\right|^{2}\right), \\
& \leq \frac{1}{2}\left(\Delta u+\Delta u_{\mu}-p^{\prime}\left(u+u_{\mu}\right) \Delta\left(u-u_{\mu}\right)\right)=F .
\end{aligned}
$$

Put

$$
\begin{align*}
& G_{1}=\left\{x \in \Omega:\left(u-u_{\mu}\right)(x)>\delta\right\}, \\
& G_{2}=\left\{x \in \Omega:\left(u-u_{\mu}\right)(x)<-\delta\right\},  \tag{2.13}\\
& G_{3}=\left\{x \in \Omega:\left|u-u_{\mu}\right|(x) \leq \delta\right\} .
\end{align*}
$$

On $G_{1}, p^{\prime}\left(u-u_{\mu}\right)=1$ and

$$
F=\Delta u_{\mu}=g\left(., u_{\mu}\right)=g\left(., \omega_{\delta}-\delta / 4\right)
$$

On $G_{2}, p^{\prime}\left(u-u_{\mu}\right)=-1$ and

$$
F=\Delta u \leq g\left(., u_{\mu}\right)=g\left(., \omega_{\delta}-\delta / 4\right) .
$$

On $G_{3}, p^{\prime}\left(u-u_{\mu}\right)=\delta^{-1}\left(u-u_{\mu}\right)$ and

$$
\begin{aligned}
F & =\frac{1}{2}\left(1-\frac{u-u_{\mu}}{\delta}\right) \Delta u+\frac{1}{2}\left(1+\frac{u-u_{\mu}}{\delta}\right) \Delta u_{\mu} \\
& \leq\left(1-\frac{u-u_{\mu}}{\delta}\right) g(., u)+\frac{1}{2}\left(1+\frac{u-u_{\mu}}{\delta}\right) g\left(., u_{\mu}\right) .
\end{aligned}
$$

By continuity of $r \mapsto g(x, r)$ there exists $\theta=\theta(x) \in[0,1]$ such that

$$
F \leq g\left(., \theta u+(1-\theta) u_{\mu}\right) \leq g\left(., \omega_{\delta}+3 \delta / 4\right) \leq g(., v+\delta)
$$

Combining those inequalities infers

$$
\begin{equation*}
\Delta \omega_{\delta} \leq g\left(., \omega_{\delta}+3 \delta / 4\right) \leq g\left(., u_{\mu}+\delta\right) \tag{2.14}
\end{equation*}
$$

If we let $\delta \rightarrow 0, \omega_{\delta} \rightarrow w_{\mu}=\min \left(u, u_{\mu}\right)$ and (2.12) holds in the sense of distributions in $\Omega$. Since $0 \leq g\left(., w_{\mu}\right) \leq g\left(., u_{\mu}\right), g\left(., w_{\mu}\right) \in L^{1}\left(\Omega ; \rho_{\partial \Omega}\right) d x$. For the last assertion let $\mu_{i} \in \mathfrak{M}_{+}(\partial \Omega)(i=1,2), \theta \in[0,1], \mu_{\theta}=\theta \mu_{1}+(1-\theta) \mu_{2}$ and $u_{\theta}=\theta u_{\mu_{1}}+(1-\theta) u_{\mu_{2}}$. Since

$$
g\left(x, u_{\theta}\right) \leq \theta g\left(x, u_{\mu_{1}}\right)+(1-\theta) g\left(x, u_{\mu_{2}}\right),
$$

there holds

$$
-\Delta u_{\theta}+g\left(x, u_{\theta}\right) \leq 0
$$

and $u_{\theta} \leq u_{\mu_{\theta}}$ by the comparison principle between solutions of (0.4). Moreover,

$$
\begin{aligned}
u+u_{\theta}-\left|u-u_{\theta}\right| & =\theta\left(u+u_{\mu_{1}}\right)+(1-\theta)\left(u+u_{\mu_{2}}\right)-\left|\theta\left(u-u_{\mu_{1}}\right)+(1-\theta)\left(u-u_{\mu_{2}}\right)\right| \\
& \geq \theta\left(u+u_{\mu_{1}}-\left|u-u_{\mu_{1}}\right|\right)+(1-\theta)\left(u+u_{\mu_{2}}-\left|u-u_{\mu_{2}}\right|\right) \\
& =\theta \min \left\{u, u_{\mu_{1}}\right\}+(1-\theta) \min \left\{u, u_{\mu_{2}}\right\} .
\end{aligned}
$$

Thus

$$
\min \left\{u, u_{\mu_{\theta}}\right\} \geq \min \left\{u, u_{\theta}\right\} \geq \theta \min \left\{u, u_{\mu_{1}}\right\}+(1-\theta) \min \left\{u, u_{\mu_{2}}\right\}
$$

which implies

$$
\gamma_{u}\left(\theta \mu_{1}+(1-\theta) \mu_{2}\right) \geq \theta \gamma_{u}\left(\mu_{1}\right)+(1-\theta) \gamma_{u}\left(\mu_{2}\right)
$$

Remark 2.2. It follows also from [8], [9] that $\Delta w_{\mu} \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. Moreover

$$
\begin{equation*}
\int_{\Omega}\left(-w_{\mu} \Delta \zeta+g\left(x, w_{\mu}\right)\right) d x=\int_{\Omega} \Phi \zeta d x-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \gamma_{u}(\mu), \forall \zeta \in C_{c}^{1,1}(\bar{\Omega}) \tag{2.15}
\end{equation*}
$$

where $\Phi=\left[-\Delta w_{\mu}+g\left(x, w_{\mu}\right)\right]$.
Proposition 2.3. Under the assumptions of Proposition 2.2 set

$$
\begin{equation*}
v=\sup \left\{\gamma_{u}(\mu): \mu \in \mathfrak{M}_{+}(\partial \Omega)\right\} \tag{2.16}
\end{equation*}
$$

Then $v$ is a Borel measure on $\partial \Omega$.
Proof. It is clear that $v$ is an outer measure in the sense that

$$
\begin{equation*}
v(\emptyset)=0, \quad \text { and } v(A) \leq \sum_{k=1}^{\infty} v\left(A_{k}\right), \quad \text { whenever } A \subset \bigcup_{k=1}^{\infty} A_{k} \tag{2.17}
\end{equation*}
$$

Let $A$ and $B \subset \partial \Omega$ be disjoint Borel subsets. In order to prove that

$$
\begin{equation*}
v(A \cup B)=v(A)+v(B) \tag{2.18}
\end{equation*}
$$

we first notice that the relation holds if $\max \{v(A), \nu(B)\}=\infty$. Therefore we assume that $\nu(A)$ and $\nu(B)$ are finite. For $\varepsilon>0$ there exist two bounded positive measures $\mu_{1}$ and $\mu_{2}$ such that

$$
\gamma_{u}\left(\mu_{1}\right)(A) \leq \nu(A) \leq \gamma_{u}\left(\mu_{1}\right)(A)+\varepsilon / 2
$$

and

$$
\gamma_{u}\left(\mu_{1}\right)(B) \leq \nu(B) \leq \gamma_{u}\left(\mu_{1}\right)(B)+\varepsilon / 2
$$

Hence

$$
\begin{aligned}
v(A)+v(B) & \leq \gamma_{u}\left(\mu_{1}\right)(A)+\gamma_{u}\left(\mu_{2}\right)(B)+\varepsilon \\
& \leq \gamma_{u}\left(\mu_{1}+\mu_{2}\right)(A)+\gamma_{u}\left(\mu_{1}+\mu_{2}\right)(B)+\varepsilon \\
& =\gamma_{u}\left(\mu_{1}+\mu_{2}\right)(A \cup B)+\varepsilon \\
& \leq v(A \cup B)+\varepsilon .
\end{aligned}
$$

Therefore $v$ is a finitely additive measure. If $\left\{A_{k}\right\}(k \geq 0)$ is a sequence of of disjoint Borel sets and $A=\cup A_{k}$, then

$$
v(A) \geq v\left(\bigcup_{1 \leq k \leq n} A_{k}\right)=\sum_{k=1}^{n} v\left(A_{k}\right) \Longrightarrow v(A) \geq \sum_{k=1}^{\infty} v\left(A_{k}\right) .
$$

By (2.17), it implies that $v$ is a countably additive measure.
Remark 2.3. The measure $v$ may not be regular. If $v(B)=\infty$ then $\nu(\mathcal{O})=\infty$ for any relatively open subset $\mathcal{O}$ containing $B$. On the other hand, if $v(B)<\infty$, there exists a sequence of positive Radon measures $\mu_{n}$ such that

$$
\gamma_{u}\left(\mu_{n}\right)(B) \uparrow v(B) \text { as } n \rightarrow \infty .
$$

Even if for each $n \in \mathbb{N}_{*}$ and $\epsilon>0$ there exists a relatively open subset $\mathcal{O}_{n, \epsilon}$ containing $B$ such that

$$
\gamma_{u}\left(\mu_{n}\right)\left(\mathcal{O}_{n, \epsilon}\right) \leq \gamma_{u}\left(\mu_{n}\right)(B)+\epsilon,
$$

there is no reason that there exists some open subset containing $B$ such that $\gamma_{u}\left(\mu_{n}\right)(\mathcal{O})$ would remain bounded independently of $n$.

Definition 2.3. The outer Borel measure $v$ defined by the above process is called the extended boundary trace of $u$ and denoted by $\operatorname{Tr}_{\partial \Omega}^{e}(u)$.

The next result shows that in the study of the extended boundary trace, it is always possible to replace the inequation by an equation.

Proposition 2.4. Let $g \in \mathcal{G}_{0}$ be positively subcritical. If $u$ is a nonnegative solution of (2.6) and $u^{*}$ is the largest solution of $(0.1)$ dominated by $u$, then

$$
\begin{equation*}
\operatorname{Tr}_{\partial \Omega}^{e}(u)=\operatorname{Tr}_{\partial \Omega}^{e}\left(u^{*}\right) . \tag{2.19}
\end{equation*}
$$

If, in addition, there exist an open domain $\mathcal{O} \subset \mathbb{R}^{N}$ and a nonnegative Radon measure $\tilde{\mu}$ on $\mathcal{O} \cap \partial \Omega$ such that

$$
\begin{equation*}
\lim _{\beta \rightarrow 0} \int_{\Sigma_{\beta} \cap \mathcal{O}} u(x) \phi(\sigma(x)) d S=\int_{\mathcal{O} \cap \partial \Omega} \phi d \tilde{\mu} . \tag{2.20}
\end{equation*}
$$

for any $\phi \in C(\mathcal{O} \cap \partial \Omega)$, then

$$
\begin{equation*}
\left.\operatorname{Tr}_{\partial \Omega}^{e}(u)\right|_{\partial \Omega \cap \mathcal{O}}=\tilde{\mu} . \tag{2.21}
\end{equation*}
$$

Proof. Because of the definition of the extended boundary trace, it is sufficient to prove that for any $\mu \in \mathfrak{M}_{+}(\partial \Omega)$,

$$
\begin{equation*}
\gamma_{u}(\mu)=\gamma_{u^{*}}(\mu) \tag{2.22}
\end{equation*}
$$

Because $u^{*} \leq u$, then $\min \left\{u^{*}, u_{\mu}\right\} \leq \min \left\{u, u_{\mu}\right\}$ and $\gamma_{u^{*}}(\mu) \leq \gamma_{u}(\mu)$. Conversely

$$
u \geq \min \left\{u, u_{\mu}\right\} \Longrightarrow u^{*} \geq\left[\min \left\{u, u_{\mu}\right\}\right]^{*}
$$

and

$$
u_{\mu} \geq \min \left\{u, u_{\mu}\right\} \Longrightarrow u_{\mu}^{*}=u_{\mu} \geq\left[\min \left\{u, u_{\mu}\right\}\right]^{*}
$$

by (2.10). Therefore

$$
\min \left\{u^{*}, u_{\mu}\right\} \geq\left[\min \left\{u, u_{\mu}\right\}\right]^{*}
$$

and

$$
\left[\min \left\{u^{*}, u_{\mu}\right\}\right]^{*} \geq\left(\left[\min \left\{u, u_{\mu}\right\}\right]^{*}\right)^{*}=\left[\min \left\{u, u_{\mu}\right\}\right]^{*}
$$

by (2.11). By Proposition 2.1, $\min \left\{u, u_{\mu}\right\}$ and $\left[\min \left\{u, u_{\mu}\right\}\right]^{*}$ have the same boundary trace, and the same holds $\min \left\{u^{*}, u_{\mu}\right\}$ and $\left[\min \left\{u^{*}, u_{\mu}\right\}\right]^{*}$. Therefore

$$
\gamma_{u^{*}}(\mu) \geq \gamma_{u}(\mu)
$$

which implies (2.22).
For the second assertion, we assume that $u$ admits $\tilde{\mu}$ for boundary trace on $\mathcal{O} \cap \partial \Omega$. Let $\lambda \in \mathfrak{M}_{+}(\partial \Omega)$ and $\phi \in C_{c}(\partial \Omega \cap \mathcal{O}), \phi \geq 0$. Since

$$
\int_{\Sigma_{\beta} \cap \mathcal{O}} u(x) \phi(\sigma(x)) d S \geq \int_{\Sigma_{\beta} \cap \mathcal{O}} \min \left\{u(x), u_{\lambda}(x)\right\} \phi(\sigma(x)) d S,
$$

there holds, as $\beta \rightarrow 0$,

$$
\int_{\partial \Omega \cap \mathcal{O}} \phi d \tilde{\mu} \geq \int_{\partial \Omega \cap \mathcal{O}} \phi d \gamma_{u}(\lambda)
$$

thus

$$
\begin{equation*}
\tilde{\mu} \geq\left.\operatorname{Tr}_{\partial \Omega}^{e}(u)\right|_{\partial \Omega \cap \mathcal{O}} \tag{2.23}
\end{equation*}
$$

Conversely, by reducing the set $\mathcal{O}$, we first suppose that $\tilde{\mu}$ is bounded and we extend it by 0 outside $\mathcal{O} \cap \partial \Omega$. We can also suppose that $\mathcal{O} \cap \Omega$ is $C^{2}$ and that $u \in L^{1}(\Omega \cap \partial \mathcal{O})$ by (2.20) and Fubini's theorem. Let $v=v_{\tilde{\mu}}^{\mathcal{O}}$ be the solution of

$$
\begin{align*}
&-\Delta v+g(x, v)=0 \text { in } \mathcal{O} \cap \Omega \\
& v=0  \tag{2.24}\\
& \text { in } \Omega \cap \partial \mathcal{O} \\
& v=\tilde{\mu}
\end{align*} \text { in } \partial \Omega \cap \mathcal{O} .
$$

Since $\left.u\right|_{\mathcal{O} \cap \Omega}$ satisfies the same equation, with the exception of the data on $\Omega \cap \partial \mathcal{O}$ which is an integrable nonnegative function, $\left.g(., u)\right|_{\mathcal{O} \Omega \Omega} \in L^{1}\left(\mathcal{O} \cap \Omega ; \rho_{\partial(\mathcal{O} \cap \Omega)} d x\right)$ and

$$
v_{\tilde{\mu}}^{\mathcal{O}} \leq u \quad \text { in } \mathcal{O} \cap \Omega
$$

by the maximum principle and [5]. If $\tilde{v}=\tilde{v}_{\tilde{\mu}}^{\mathcal{O}}$, is the extension of $v$ by zero in $\Omega \backslash(\mathcal{O} \cap \Omega)$, then

$$
\tilde{v} \leq u_{\tilde{\mu}} \leq \min \left\{u, u_{\tilde{\mu}}\right\}
$$

Therefore

$$
\int_{\mathcal{O}} \phi d \tilde{\mu} \leq \int_{\mathcal{O}} \phi d \gamma_{u}(\tilde{\mu})
$$

for any $\phi \in C_{c}(\partial \Omega \cap \mathcal{O})$. Clearly the same relation holds even if we no longer assume that $\tilde{\mu}$ is bounded. Thus

$$
\tilde{\mu}(E) \leq \gamma_{u}(\tilde{\mu})(E), \quad \forall E \subset \partial \Omega \cap \mathcal{O}, E \text { Borel }
$$

and consequently

$$
\begin{equation*}
\tilde{\mu} \leq\left.\operatorname{Tr}_{\partial \Omega}^{e}(u)\right|_{\partial \Omega \cap \mathcal{O}} . \tag{2.25}
\end{equation*}
$$

Remark 2.4. The relation (2.20) means that $u$ admits a boundary trace in the usual sense on $\mathcal{O} \cap \partial \Omega$ which is precisely $\tilde{\mu}$. The reverse implication " $(2.21) \Longrightarrow(2.20)$ " holds under an additional strong stability assumption which will be developped in Section 4. However we can give a weaker form of this implication if $u^{*}$ is dominated by the minimal large solution of (2.7), whenever it exists.

Let us denote by $\lambda_{n}(n \geq 0)$ the measure $n \chi_{\partial \Omega} d S$ and

$$
\begin{equation*}
u_{m}=\sup \left\{u_{\lambda_{n}}: n \in \mathbb{N}\right\} \tag{2.26}
\end{equation*}
$$

If $u_{m}$ is locally bounded in $\Omega$, it is a solution of (2.7) which blows up on the boundary. In such a case it is called the minimal large solution. Depending upon the nonlinearity, $u_{m}$ may also be infinite in whole $\Omega$ or in part of $\Omega$. Moreover, by the maximum principle, it dominates any solution $u$ of the same equation which is obtained by approximation by solutions with finite values as boundary data.

Proposition 2.5. Let $g \in \mathcal{H}_{0}$ be positively subcritical. If $\operatorname{Tr}_{\partial \Omega}^{e}(u)$ is a bounded Borel measure and $u^{*} \leq u_{m}$, then

$$
\begin{equation*}
\operatorname{Tr}_{\partial \Omega}^{e}\left(u^{*}\right)=\operatorname{Tr}_{\partial \Omega}\left(u^{*}\right) \in \mathfrak{M}_{+}(\partial \Omega) \tag{2.27}
\end{equation*}
$$

Proof. By assumption $v=\operatorname{Tr}_{\partial \Omega}^{e}(u)$ is bounded. Thus there exists a sequence $\left\{\mu_{n}\right\} \subset \mathfrak{M}_{+}(\partial \Omega)$ such that

$$
\gamma_{u}\left(\mu_{n}\right)(1)=\gamma_{u^{*}}\left(\mu_{n}\right)(1)=\operatorname{Tr}_{\partial \Omega}\left(\left[\min \left\{u^{*}, u_{\mu_{n}}\right\}\right]^{*}\right)(1) \uparrow \nu(1),
$$

since $u$ and $u^{*}$ have the same extended boundary trace. Because the extended boundary trace is defined by a supremum over all measures, it can also be
assumed that the regular part of $\mu_{n}$ is a.e. bounded from below by $n$. Let $\zeta \in C_{c}^{1,1}(\bar{\Omega})$ be the solution of

$$
\begin{aligned}
-\Delta \zeta & =1 \text { in } \Omega \\
\zeta & =0 \text { on } \partial \Omega
\end{aligned}
$$

By the definition of the boundary trace in $\mathfrak{M}_{+}(\partial \Omega)$,

$$
\begin{aligned}
\int_{\Omega}\left(\left[\min \left\{u^{*}, u_{\lambda_{n}}\right\}\right]^{*}+\zeta g\left(x,\left[\min \left\{u^{*}, u_{\lambda_{n}}\right\}\right]^{*}\right) d x\right. & =-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \operatorname{Tr}_{\partial \Omega}\left(\left[\min \left\{u^{*}, u_{\lambda_{n}}\right\}\right]^{*}\right), \\
& \leq C v(1) .
\end{aligned}
$$

Since $u_{m}$ dominates $u^{*}, \lim _{n \rightarrow \infty} \min \left\{u^{*}, u_{\lambda_{n}}\right\}=u^{*}$ and

$$
\begin{equation*}
u^{*} \in L^{1}(\Omega), \text { and } g\left(., u^{*}\right) \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right) \tag{2.28}
\end{equation*}
$$

It follows from [27, Corollary 1.3] that $u^{*}$ has a boundary trace, say $\mu^{*}$ in $\mathfrak{M}_{+}(\partial \Omega)$, and $u^{*}=u_{\mu^{*}}$. Consequently for any measure $\mu$ larger than $\mu^{*}$, $\min \left\{u_{\mu}, u_{\mu^{*}}\right\}=u_{\mu^{*}}$ and

$$
\operatorname{Tr}_{\partial \Omega}^{e}\left(u^{*}\right)=\operatorname{Tr}_{\partial \Omega}^{e}\left(u_{\mu^{*}}\right)=\sup _{\mu \geq \mu^{*}} \gamma_{u^{*}}(\mu)=\mu^{*}=\operatorname{Tr}_{\partial \Omega}\left(u^{*}\right)
$$

Remark 2.5. The previous result still holds if, in the domination assumption $u^{*} \leq u_{m}$, the function $u_{m}$ is no longer the minimal large solution, but any $\sigma$ moderate solution in the sense of Dynkin and Kuznetsov [12], that is a solution of (2.7) which is an increasing limit of solutions $u_{\mu_{n}}$ for some $\mu_{n} \in \mathfrak{M}_{+}(\partial \Omega)$

Proposition 2.6. Let $g \in \mathcal{G}_{0}$ be positively subcritical and stable, and let $u$ be a nonnegative solution of (0.9). If $\left\{\mu_{n}\right\} \subset \mathfrak{M}_{+}(\partial \Omega)$ converges weakly to $\mu$, then $\limsup _{n \rightarrow \infty} \gamma_{u}\left(\mu_{n}\right) \leq \gamma_{u}(\mu)$. If we assume moreover that the sequence $\left\{\mu_{n}\right\}$ is nonincreasing, $\lim _{n \rightarrow \infty} \gamma_{u}\left(\mu_{n}\right)=\gamma_{u}(\mu)$.

Proof. Since $\mu_{n} \rightarrow \mu$ in the weak sense of measures on $\partial \Omega, u_{\mu_{n}} \rightarrow u_{\mu}$ locally uniformly in $\Omega$ by definition of the positive subcriticality and stability. Thus

$$
w_{\mu_{n}}=\min \left\{u^{*}, u_{\mu_{n}}\right\} \rightarrow w_{\mu}=\min \left\{u^{*}, u_{\mu}\right\} \text { in the } C_{\text {loc }}(\Omega) \text {-topology. }
$$

Since $u_{\mu_{n}} \leq \mathbb{P}_{\mu_{n}}$ and $P_{\mu_{n}} \rightarrow P_{\mu}$ in $L^{1}(\Omega)$, the convergence of $w_{\mu_{n}}$ to $w_{\mu}$ holds also in $L^{1}(\Omega)$. Moreover

$$
g\left(., w_{\mu_{n}}\right) \rightarrow g\left(., w_{\mu}\right)
$$

in $C_{\text {loc }}(\Omega)$. Since

$$
0 \leq g\left(., w_{\mu_{n}}\right) \leq g\left(., u_{\mu_{n}}\right)
$$

and

$$
g\left(., u_{\mu_{n}}\right) \rightarrow g\left(., u_{\mu}\right)
$$

in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$,

$$
\begin{equation*}
g\left(., w_{\mu_{n}}\right) \rightarrow g\left(., w_{\mu}\right) \tag{2.29}
\end{equation*}
$$

also in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. Put $v_{n}^{*}=\left[w_{\mu_{n}}\right]^{*}$. By the elliptic equations regularity theory, $\left\{v_{n}^{*}\right\}$ remains bounded in $C_{\text {loc }}^{1}(\Omega)$. Since $\gamma_{u}\left(\mu_{n}\right)$ is dominated by $\mu_{n}$ which is bounded let us consider a subsequence $\gamma_{u}\left(\mu_{n_{k}}\right)$ weakly convergent to some nonnegative measure $\lambda$. Up to an extraction of a subsequence, it is always possible to assume that $v_{n_{k}}^{*}$ converges (in the $C_{\text {loc }}(\Omega)$-topology) to $\bar{v}$. Clearly $\bar{v}$ is a solution of (0.1) in $\Omega$ and

$$
\begin{equation*}
v_{n}^{*} \leq w_{\mu_{n}} \Longrightarrow \bar{v} \leq \lim _{n \rightarrow \infty} w_{\mu_{n}}=w_{\mu} \tag{2.30}
\end{equation*}
$$

Therefore

$$
\begin{equation*}
\bar{v}=\bar{v}^{*} \leq v^{*}=\left[w_{\mu}\right]^{*} \tag{2.31}
\end{equation*}
$$

Inasmuch

$$
v_{n_{k}}^{*} \rightarrow \bar{v} \text { in } L^{1}(\Omega) \text { and } g\left(., v_{n_{k}}^{*}\right) \rightarrow g(., \bar{v}) \text { in } L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right),
$$

(for the second relation we use $0 \leq g\left(., v_{n_{k}}^{*}\right) \leq g\left(., w_{\mu_{n}}\right)$, together with (2.29) and Vitali's theorem) and

$$
v_{n_{k}}+\mathbb{G}_{g\left(., v_{n_{k}}^{*}\right)}=\mathbb{P}_{\gamma_{u}\left(\mu_{n}\right)}
$$

it follows

$$
\bar{v}+\mathbb{G}_{g(., \bar{v})}=\mathbb{P}_{\lambda}
$$

But

$$
v^{*}+\mathbb{G}_{g\left(., v^{*}\right)}=\mathbb{P}_{\gamma_{u}(\mu)} .
$$

As $\bar{v} \leq v^{*}$ and $g(., \bar{v}) \leq g\left(., v^{*}\right)$,

$$
\begin{equation*}
\mathbb{P}_{\lambda} \leq \mathbb{P}_{\gamma_{u}(\mu)} \Longrightarrow \lambda \leq \gamma_{u}(\mu) \tag{2.32}
\end{equation*}
$$

which is the first assertion.
If we assume that $\left\{\mu_{n}\right\}$ is nonincreasing, the same holds with $\left\{u_{\mu_{n}}\right\},\left\{w_{\mu_{n}}\right\}$, $\left\{\gamma_{u}\left(\mu_{n}\right)\right\}$ and $\left\{v_{n}^{*}\right\}$. If $\mu_{n} \downarrow \mu$, any solution of (0.1) dominated by $w_{\mu}$ is dominated by $w_{\mu_{n}}$. Thus $v^{*} \leq v_{n}^{*}$ and $\bar{v}=\lim _{n \rightarrow \infty} v_{n}^{*}=v^{*}$ by (2.30) and (2.31).

A particularly important case deals with the choice $\mu=\lambda \delta_{a}$, with $a \in \partial \Omega$ and $\lambda>0$. Let $u=u_{\lambda \delta_{a}}$ be the solution of

$$
\begin{aligned}
-\Delta u+g(x, u)=0, & \text { in } \Omega \\
u=\lambda \delta_{a}, & \text { on } \partial \Omega
\end{aligned}
$$

Since $g(x,$.$) is nondecreasing, \lambda \mapsto u_{\lambda \delta_{a}}$ is increasing. Set

$$
u_{\infty, a}=\lim _{\lambda \rightarrow \infty} u_{\lambda \delta_{a}}
$$

On any open subset of $\Omega$ where it is locally finite, $u_{\infty, a}$ is a solution of (0.1).

Lemma 2.1. Let $a \in \partial \Omega, \lambda>0$ and $w_{\lambda \delta_{a}}=\min \left\{u, u_{\lambda \delta_{a}}\right\}$. Then

$$
\begin{equation*}
\operatorname{Tr}_{\partial \Omega}\left(w_{\lambda \delta_{a}}\right)=\gamma_{u}\left(\lambda \delta_{a}\right)=\tilde{\gamma}_{u}(a, \lambda) \delta_{a}, \tag{2.34}
\end{equation*}
$$

where $0 \leq \tilde{\gamma}_{u}(a, \lambda) \leq \lambda$. Moreover the mapping $\lambda \mapsto \tilde{\gamma}_{u}(a, \lambda)$ is nondecreasing, there exists

$$
\begin{equation*}
\tilde{\gamma}_{u}(a)=\lim _{\lambda \rightarrow \infty} \tilde{\gamma}_{u}(a, \lambda) \tag{2.35}
\end{equation*}
$$

and

$$
\begin{equation*}
\min \left(u, u_{\infty, a}\right) \geq u_{\tilde{\gamma}_{u}(a)} . \tag{2.36}
\end{equation*}
$$

Proof. Because $\operatorname{Tr}_{\partial \Omega}\left(w_{\lambda \delta_{a}}\right) \leq \lambda \delta_{a}$, this trace is concentrated at $a$ and achieves the form $\tilde{\gamma}_{u}(a, \lambda) \delta_{a}$. Moreover $\lambda \mapsto \tilde{\gamma}_{u}(a, \lambda)$ is nondecreasing as is the mapping $\lambda \mapsto \min \left(u, u_{\lambda \delta_{a}}\right)$. Since

$$
w_{\lambda \delta_{a}} \geq u_{\tilde{\gamma}_{u}(a, \lambda) \delta_{a}}
$$

and $u_{\lambda \delta_{a}}$ admits a limit, finite or not, when $\lambda \rightarrow \infty$, assertions (i) or (ii) follow.
The next result points out the role of $u_{\infty, a}$
Proposition 2.7. Let $g \in \mathcal{G}_{0}$ be positively subcritical. If $u$ is a nonnegative solution of (2.6) with boundary trace $v$ and $a \in \partial \Omega$, then

$$
\begin{equation*}
u \geq u_{\infty, a} \Longrightarrow v(a)=\infty \tag{2.37}
\end{equation*}
$$

If we assume moreover that $g$ is positively subcritical and stable, then

$$
\begin{equation*}
\nu(a)=\infty \Longrightarrow u \geq u_{\infty, a} \tag{2.38}
\end{equation*}
$$

Proof. Let $a \in \partial \Omega$ be such that

$$
v(a)=\infty
$$

Then for any relatively open subset $\mathcal{O} \subset \partial \Omega$ containing $a$,

$$
v(\mathcal{O})=\infty
$$

which means that there exists a sequence of positive Radon measures $\mu_{n}$ such that

$$
\lim _{n \rightarrow \infty} \gamma_{u}\left(\mu_{n}\right)(\mathcal{O})=\infty
$$

Without any restriction we can suppose that the sequence of restricted measures $\gamma_{u}^{\prime}\left(\mu_{n}\right)=\chi_{\mathcal{O}} \gamma_{u}\left(\mu_{n}\right)$ is increasing and

$$
u \geq u_{\gamma_{u}^{\prime}\left(\mu_{n}\right)}
$$

because $u \geq u_{\gamma_{u}\left(\mu_{n}\right)}$ and $\gamma_{u}\left(\mu_{n}\right) \geq \gamma_{u}^{\prime}\left(\mu_{n}\right)$. For any $k \in \mathbf{N}_{*}$, there exists $\epsilon_{n, k}>0$ such that, if we take $\mathcal{O}_{n, k}=B_{\epsilon_{n, k}}(a) \cap \partial \Omega$, there holds

$$
\gamma_{u}\left(\mu_{n}\right)\left(\mathcal{O}_{n, k}\right)=k
$$

Set $\mu_{n, k}=\chi_{\mathcal{O}_{n, k}} \gamma_{u}\left(\mu_{n}\right)$. Then

$$
\gamma_{u}\left(\mu_{n}\right) \geq \mu_{n, k} \Longrightarrow u_{\gamma u\left(\mu_{n}\right)} \geq u_{\mu_{n, k}}
$$

Since $\lim _{n \rightarrow \infty} \epsilon_{n, k}=0, \lim _{n \rightarrow \infty} \mu_{n, k}=k \delta_{a}$. Consequently

$$
u(x) \geq u_{k \delta_{a}}(x)
$$

Letting $k \rightarrow \infty$ yields to the following implication

$$
v(a)=\infty \Longrightarrow u \geq u_{\infty a} .
$$

Conversely, assume $u \geq u_{\infty a}$, then for any $k>0, u \geq u_{k \delta_{a}}$. On one hand the boundary trace of $w_{k \delta_{a}}=\min \left\{u, u_{k \delta_{a}}\right\}$ is the measure $\gamma_{u}(a, k) \delta_{a}$. But $\min \left\{u, u_{k \delta_{a}}\right\}=u_{k \delta_{a}}$ implies $\gamma_{u}(a, k)=k$ and therefore $\gamma_{u}\left(k \delta_{a}\right)=k \delta_{a}$. By the definition of $\nu$,

$$
v(a) \geq \gamma_{u}\left(k \delta_{a}\right)(a)=k
$$

Since this holds for any $k>0, v(a)=\infty$.
The characterisation of Borel subsets on which the boundary trace of $u$ takes finite values is less complete, however there holds

Proposition 2.8. Assume the assumptions of Proposition 2.2 are fulfilled and $u$ is a nonnegative solution of (2.6) with boundary trace $\nu$. If $\mathcal{O} \subset \partial \Omega$ is a relatively open subset of $\partial \Omega$ such that

$$
\int_{\mathcal{O}_{t}} u(y) d S_{t}
$$

remains bounded independently of $t \in\left(0, \beta_{0}\right]$, then $\nu(\mathcal{O})$ is finite.
Proof. Let $\mu$ be a nonnegative measure. Since $u_{\gamma u(\mu)} \leq u$ in $\Omega$,

$$
\int_{\mathcal{O}_{t}} u_{\gamma_{u}(\mu)}(y) d S_{t} \leq \int_{\mathcal{O}_{t}} u(y) d S_{t}
$$

Thus

$$
\gamma_{u}(\mu)(\mathcal{O}) \leq \sup _{0<\beta \leq \beta_{0}} \int_{\mathcal{O}_{t}} u(y) d S_{t}=M
$$

Therefore

$$
\nu(\mathcal{O}))=\sup _{\mu \in \mathcal{M}_{+}(\partial \Omega)} \gamma_{u}(\mu)(\mathcal{O}) \leq M .
$$

Remark 2.6. The reverse implication

$$
\nu(\mathcal{O})<\infty \Longrightarrow \int_{\mathcal{K}_{t}} u(y) d S_{t} \leq M, \quad \forall t \in\left(0, \beta_{0}\right]
$$

for any compact subset $K \subset \mathcal{O}$, where $M=M(K)>0$, may not hold in the case of general inequalities. However, $g \in \mathcal{H} \mathcal{G}_{0}$ with $\alpha=0$, that is $g$ satisfies (0.14), or if

$$
g(x, r) \leq \rho_{\partial \Omega}^{\alpha}(x) r^{q}=0, \quad \text { in } \Omega \times \mathbb{R}_{+},
$$

with $\alpha>-2$, and $1<q<(N+1+\alpha) /(N-1)$, such a result is still valid. Under both assumptions the proof is much intricated : in the first one it is given in next section, and in the second one, in [27]. In both cases the proof is settled on the notion of stability from inside approximations of the Dirichlet problems ( 0.4 ) which means that if a sequence of measures $\lambda_{n} \in \mathfrak{M}_{+}(\Omega)$ converges weakly to a measure $\mu \in \mathfrak{M}_{+}(\partial \Omega)$ the solutions $v_{n}$ of the semilinear equation with forcing term

$$
\begin{align*}
-\Delta v_{n}+g\left(x, v_{n}\right)=\lambda_{n} & \text { in } \Omega  \tag{2.39}\\
v_{n}=0 & \text { on } \partial \Omega \tag{2.40}
\end{align*}
$$

converges to $u_{\mu}$ locally uniformly in $\Omega$.
The real number $\gamma_{u}(a)$ plays an important role in the study of the boundary behavior of $u$ at $a$. If $a \in \partial \Omega$, we denote by $\mathcal{N}_{a}$ the set of relatively open neighborhoods of $a$ in $\partial \Omega$.

Definition 2.4. We define by $\mathcal{A}(u)$ the set of atoms of $u$,

$$
\mathcal{A}(u)=\left\{a \in \partial \Omega: \gamma_{u}(a)>0\right\},
$$

by $\mathcal{S}(u)$ the singular set of $u$,

$$
\mathcal{S}(u)=\left\{a \in \mathcal{A}: \forall N_{a} \in \mathcal{N}_{a}, \sum_{\omega \in N_{a}} \gamma_{u}(\omega)=\infty\right\},
$$

the symbol $\sum$ being taken in the sense of summable family, and by $\mathcal{R}(u)$ the regular set of $u$,

$$
\mathcal{R}(u)=\partial \Omega \backslash \mathcal{S}(u)=\left\{a \in \mathcal{A}: \exists N_{a} \in \mathcal{N}_{a}, \sum_{\omega \in N_{a}} \gamma_{u}(\omega)<\infty\right\}
$$

The set $\mathcal{S}(u)$ is closed and $\mathcal{R}(u)$ relatively open. Moreover, if $a \in \mathcal{R}(u)$, there exists a relatively open neighborhood $N_{a} \in \mathcal{N}_{a}$ such that $\mathcal{A}(u) \cap N_{a}$ is at most countable.

The next result complements Propositions 2.7 and 2.8

Theorem 2.1. Assume the assumptions of Proposition 2.2 are fulfilled, $u$ is a nonnegative solution of (2.6) with boundary trace $\nu$, and $\mathcal{O}$ is a relatively open subset of $\partial \Omega$ such that $\nu(\mathcal{O})<\infty$. Then

$$
\sum_{a \in \mathcal{O}} \tilde{\gamma}_{u}(a)<\infty .
$$

If we assume moreover that $g \in \mathcal{H} \mathcal{G}_{0}$, then, for any $\omega \in \partial \Omega$, there holds

$$
\nu(\omega)=\tilde{\gamma}_{u}(\omega),
$$

and the measure $\chi_{\mathcal{O}} \nu-\sum_{\omega \in \mathcal{O}} \tilde{\gamma}_{u}(\omega) \delta_{\omega}$ has no atom.
Proof. Let $K$ be a finite subset of $\mathcal{R}(u) \cap \mathcal{O}$ and put $\mu_{K}=\sum_{a \in K} \delta_{a}$. Then for any $\lambda>0$

$$
\sum_{a \in K} \tilde{\gamma}_{u}(a, \lambda) \leq \gamma_{u}\left(\lambda \mu_{K}\right)(\mathcal{O}) \leq \nu(\mathcal{O}) .
$$

Therefore the following family $\left\{\tilde{\gamma}_{u}(a)\right\}_{a \in \mathcal{O}}$ is summable, and

$$
\sum_{a \in \mathcal{O}} \tilde{\gamma}_{u}(a) \leq v(\mathcal{O})
$$

For the next statement, for any $\lambda>0$ and $\omega \in \partial \Omega$, there holds

$$
\nu(\omega) \geq \tilde{\gamma}_{u}(\omega, \lambda) \Longrightarrow \nu(\omega) \geq \tilde{\gamma}_{u}(\omega)
$$

Conversely, for any relatively open neighborhood of $\omega, \mathcal{O}_{\omega}$, there exists a sequence of Radon measures $\mu_{n} \in \mathfrak{M}_{+}(\partial \Omega)$ such that

$$
\int_{\mathcal{O}_{\omega}} d \mu_{n} \uparrow \nu\left(\mathcal{O}_{\omega}\right), \quad \text { as } n \rightarrow \infty
$$

If we assume that $v(\omega)=\infty$, we know from Proposition 2.7 that $\tilde{\gamma}_{u}(\omega)=\infty$. Thus we assume $\nu(\omega)<\infty$. For $\epsilon>0$, there exists $\mu_{\epsilon} \in \mathfrak{M}_{+}(\partial \Omega)$ such that

$$
\gamma_{u}\left(\mu_{\epsilon}\right)(\omega) \leq \nu(\omega) \leq \gamma_{u}\left(\mu_{\epsilon}\right)(\omega)+\epsilon,
$$

and there exists $\eta_{0}>0$ such that $0<\eta \leq \eta_{0}$ implies

$$
\int_{\Gamma_{\eta}(\omega)} d \gamma_{u}\left(\mu_{\epsilon}\right)-\epsilon \leq \nu(\omega) \leq \int_{\Gamma_{\eta}(\omega)} d \gamma_{u}\left(\mu_{\epsilon}\right)+\epsilon,
$$

where $\Gamma_{\eta}(\omega)=B_{\eta}(\omega) \cap \partial \Omega$, which yields to

$$
\left|\int_{\Gamma_{\eta}(\omega)} d \gamma_{u}\left(\mu_{\epsilon, \eta}\right)-v(\omega)\right| \leq 2 \epsilon
$$

If we take $\epsilon=1 / n$, then $\eta_{0}=\eta_{0}(n) \rightarrow 0$ and $\chi_{\Gamma_{\eta}(\omega)} \gamma_{u}\left(\mu_{\epsilon}\right) \rightarrow \nu(\omega) \delta_{\omega}$ as $n \rightarrow \infty$. But

$$
u_{\chi_{\Gamma_{\eta}(\omega)}^{\gamma_{u}}\left(\mu_{\epsilon}\right)} \leq u_{\gamma_{u}\left(\mu_{\epsilon}\right)} \leq w_{\gamma_{u}\left(\mu_{\epsilon}\right)} \leq u .
$$

Letting $n \rightarrow \infty$ and using the fact that $u_{\chi_{\Gamma_{\eta}(\omega)} \gamma_{u}\left(\mu_{\epsilon}\right)} \rightarrow u_{\nu(\omega) \delta_{\omega}}$ implies

$$
u_{\nu(\omega) \delta_{\omega}} \leq u .
$$

Therefore

$$
u_{\nu(\omega) \delta_{\omega}}=\min \left\{u, u_{\nu(\omega) \delta_{\omega}}\right\}=w_{\nu(\omega) \delta_{\omega}}=u_{\tilde{\gamma}_{u}(\omega, \nu(\omega)) \delta_{\omega}} \leq u_{\tilde{\gamma}_{u}(\omega) \delta_{\omega}} .
$$

This implies $\nu(\omega) \leq \gamma_{u}(\omega)$ and the equality follows. Consequently $\chi_{\mathcal{O}} v-$ $\sum_{\omega \in \mathcal{O}} \tilde{\gamma}_{u}(\omega) \delta_{\omega}$ has no atom.

## 3. - Pointwise boundary behaviour of solutions of general inequalities

In this section, we give a precise description of the behaviour of a solution $u$ of (2.6) near an atom of its extended boundary trace. We say that the coordinates are proper at $a=\left(a_{1}, \ldots, a_{N}\right) \in \partial \Omega$ relatively to $\Omega$ if the plane $x_{1}-a_{1}=0$ is tangent to $\partial \Omega$ at $a$, and that the inward pointing vector to $\partial \Omega$ is the direction $x_{1}-a_{1}>0$.

Definition 3.1. Let $(E, \Sigma, \mu)$ be a measured space, where $\Sigma$ is $\sigma$-algebra of subsets of $E$ and $\mu$ a positive and $\sigma$-additive measure with finite mass. We recall that a set of $\mu$-measurable functions $x \mapsto \psi_{r}(x)(r>0)$, defined over $E$ converges in measure to $\psi$ when $r \rightarrow 0$, if for any $\varepsilon>0$ there holds

$$
\lim _{r \rightarrow 0} \mu\left\{x \in E:\left|\psi_{r}(x)-\psi(x)\right|>\varepsilon\right\}=0 .
$$

The functions $\psi_{r}$ converges in measure to $\infty$, if for any $k>0$,

$$
\lim _{r \rightarrow 0} \mu\left\{x \in E: \psi_{r}(x) \leq k\right\}=0
$$

The convergence is equivalent to the following statement: from any sequence $\left\{r_{n}\right\}$ converging to 0 one can extract a subsequence $\left\{r_{n_{k}}\right\}$ such that $\psi_{r_{n_{k}}}$ converges to $\psi$ (or $\infty$ ), $\mu$-a.e. in $E$.

Theorem 3.1. Assume $g \in \mathcal{G}_{0}$ is positively subcritical, $u$ is a nonnegative solution of (2.6) and $a \in \partial \Omega$. If the coordinates are proper at a relatively to $\Omega$, the following alternative holds. Either
(i) $\tilde{\gamma}_{u}(a)$ is finite and the following convergence holds

$$
\begin{equation*}
\lim _{\substack{x \rightarrow a \\\left(x_{1}-a_{1}\right) /|x-a| \rightarrow \eta_{1}}}|x-a|^{N-1} u(x)-C(N) \tilde{\gamma}_{u}(a) \eta_{1}=0, \tag{3.1}
\end{equation*}
$$

in measure on $S_{+}^{N-1}$,
or
(ii) $\tilde{\gamma}_{u}(a)$ is infinite and

$$
\begin{equation*}
\lim _{r \rightarrow 0}|x-a|^{N-1} u(x)=\infty \tag{3.2}
\end{equation*}
$$

in measure on $S_{+}^{N-1}$.
For $s>0$, put $\Omega \cap B_{s}(a)=\Omega_{s}(a), \Omega_{s}^{c}(a)=\Omega \backslash \bar{\Omega}_{s}(a)$ and $\partial B_{s}(a) \cap$ $\Omega=\Gamma_{s}^{\Omega}(a)$. The next series of results deals with the boundary behaviour of the Green potential of a weighted integrable function. In the flat case where $\Omega=\mathbb{R}_{+}^{N}=\left\{x=\left(x_{1}, \ldots, x_{N}\right): x_{1}>0\right\}$ the computation can be explicited

Lemma 3.1. Let $\Omega=\mathbb{R}_{+}^{N}, N \geq 2, \Phi \in L^{1}\left(\mathbb{R}_{+}^{N} ; x_{1} d x\right)$ and $v=\mathbb{G}_{\Phi}$. Then for any $a \in \partial \mathbb{R}_{+}^{N}$ there holds

$$
\begin{equation*}
\lim _{x \rightarrow a}|x-a|^{-1} \int_{\Gamma_{s}^{\Omega}(a)}|v| x_{1} d S=0 \tag{3.3}
\end{equation*}
$$

Proof. We can assume that $a=0, \Phi \geq 0$, and so is $v$. For $\varepsilon>0$, let $s>0$ such that

$$
\int_{B_{s}(a)} \Phi \rho_{\partial \Omega} d x \leq \varepsilon
$$

Let $(r, \sigma) \in(0,+\infty) \times S^{N-1}$ be the spherical coordinates in $\mathbb{R}^{N}, S_{+}^{N-1}=$ $S^{N-1} \cap \mathbb{R}_{+}^{N}$ and $v(x)=v(r, \sigma)$, then

$$
-\partial_{r r} v-\frac{N-1}{r} \partial_{r} v-\frac{1}{r^{2}} \Delta_{\sigma} v=\Phi
$$

where $\Delta_{\sigma}$ is the Laplace Beltami operator on $S^{N-1}$. Since $N-1$ is the first eigenvalue of $-\Delta_{\sigma}$ in $W_{0}^{1,2}\left(S_{+}^{N-1}\right)$ and $\phi_{1}(\sigma)=\left.x_{1}\right|_{S_{+}^{N-1}}$, the first eigenfunction, there holds

$$
-\bar{v}_{r r}-\frac{N-1}{r} \bar{v}_{r}+\frac{N-1}{r^{2}} \bar{v}=\bar{\Phi}
$$

where

$$
\bar{v}(r)=\int_{S_{+}^{N-1}} v(r, \sigma) \phi_{1}(\sigma) d \sigma \quad \text { and } \quad \bar{\Phi}(r)=\int_{S_{+}^{N-1}} \Phi(r, \sigma) \phi_{1}(\sigma) d \sigma
$$

Integrating the above differential equation yields to

$$
\bar{v}(r)=\alpha r^{1-N}+\beta r-\frac{r}{N} \int_{0}^{r} \bar{\Phi}(s) d s+\frac{r^{1-N}}{N} \int_{0}^{r} \bar{\Phi}(s) s^{N} d s
$$

for some constants $\alpha \geq 0$ and $\beta$. But $\alpha=0$ otherwhile $v$ would be bounded from below by $\alpha C(N) P(x, 0)$. This is impossible because $v$ admits the zero measure for trace on the boundary. Thus

$$
\limsup _{r \rightarrow 0} r^{N-1} \bar{v}(r)=0,
$$

since

$$
\int_{0}^{r} \bar{\Phi}(s) s^{N} d s=\int_{B_{r}(0)} \Phi(x) x_{1} d x
$$

and the result follows.
This result is immediately extendable for any domain which can be deduced by a conformal transformation from a half space.

Lemma 3.2. Let $\Omega$ be a ball or the complementary of a ball, $\Phi \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$ and $v=\mathbb{G}_{\Phi}$. Then for any $a \in \partial \Omega$ there holds

$$
\begin{equation*}
\lim _{x \rightarrow a}|x-a|^{-1} \int_{\Gamma_{s}^{\Omega}(a)}|v| \rho_{\partial \Omega} d S=0 \tag{3.4}
\end{equation*}
$$

In the next lemma we prove that this result is actually always valid. Our proof involves Marcinkiewicz space estimates on the Green potential of a weighted integrable function. The following estimates, similar to (1.22), (1.23), can be found in [4, Theorem 2.6]

$$
\begin{align*}
\left.\left\|\mathbb{G}_{\Phi}\right\|_{M^{(N+1) /(N-1)\left(\Omega ; \rho_{\partial \Omega}\right.}} d x\right) & \leq K\|\Phi\|_{L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)}  \tag{3.5}\\
\left\|\mathbb{G}_{\Phi}\right\|_{M^{N /(N-1)}(\Omega)} & \leq K\|\Phi\|_{L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)} \tag{3.6}
\end{align*}
$$

Actually (3.5) is obtained in [4] only in the case $N \geq 3$, but an easy adaptation of the proof fills the gap.

Lemma 3.3. Let $N \geq 2, \Phi \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$ and $v=\mathbb{G}_{\Phi}$. If $a \in \partial \Omega$, there holds

$$
\begin{equation*}
\lim _{r \rightarrow 0} r^{-1} \int_{\Gamma_{s}^{\Omega}(a)}|v| \rho_{\partial \Omega} d S=0 \tag{3.7}
\end{equation*}
$$

Proof. We still assume $\Phi \geq 0$. For $\varepsilon>0$ let $s>0$ be such that

$$
\int_{\Omega_{s}(a)} \Phi \rho_{\partial \Omega} d x \leq \varepsilon
$$

Set $\Phi_{s}=\chi_{\Omega_{s}^{c}(a)}$ and $v_{s}=\mathbb{G}_{\Phi_{s}}$. Since $v_{s}$ is harmonic in $\Omega_{s}(a)$, with zero trace on $\partial \Omega \cap B_{s}(a)$, it is continuous in a neighborhood of $a$ and

$$
\lim _{r \rightarrow 0} r^{-1} \int_{\Gamma_{s}^{\Omega}(a)}|v|_{s} \rho_{\partial \Omega} d S=0
$$

Thus there is no loss of generality in assuming that $\Phi$ has support in $\Omega_{s}(a)$ and

$$
\|\Phi\|_{L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)} \leq \varepsilon
$$

For any $r>0$ and any $\zeta \in C_{c}^{1,1}\left(\bar{\Omega}_{r}(a)\right)$, there holds

$$
\begin{equation*}
-\int_{\Omega_{r}(a)} v \Delta \zeta d x+\int_{\Gamma_{r}^{\Omega}(a)} \frac{\partial \zeta}{\partial \mathbf{n}} v d S=\int_{\Omega_{r}(a)} \Phi \zeta d x \tag{3.8}
\end{equation*}
$$

This can be established in assuming first that $\Phi=\Phi_{n}$ is regular, and then by density in $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$. By translation it can be assumed that $a=0$. We set $\zeta(x)=|x| \eta_{r}(x)$, where $\eta_{r}$ is the first eigenfunction of $-\Delta$ in $W_{0}^{1,2}\left(\Omega_{r}(0)\right)$, and let $\lambda_{r}$ be the corresponding eigenvalue. Notice that $r^{2} \lambda_{r} \approx \lambda_{1}$ where $\lambda_{1}$ is the first eigenvalue of the operator

$$
\begin{equation*}
\ell \mapsto-\ell^{\prime \prime}-\frac{N-1}{s} \ell^{\prime}+\frac{N-1}{s^{2}} \ell \quad \text { on }(0,1) \tag{3.9}
\end{equation*}
$$

subject to the limit conditions $\ell^{\prime}(0)=0, \ell(1)=0$ (thus the corresponding eigenfuction for (3.9) is a Bessel function, say $B_{1}$, and $\eta_{r}(x) \equiv B_{1}(x / r) x_{1}$ as $r \rightarrow 0$ ). Then (3.8) becomes

$$
\begin{aligned}
& \lambda_{r} \int_{\Omega_{r}(0)} v|x| \eta_{r}(x) d x-\int_{\Omega_{r}(0)} v|x|^{-1}\left\langle x . \nabla \eta_{r}\right\rangle d x \\
& \quad=(N-1) \int_{\Omega_{r}(0)} \eta_{r} u d x-\int_{\Gamma_{r}^{\Omega}(0)}\left(r+\left\langle x . \nabla \eta_{r}\right\rangle\right) v d S+\int_{\Omega_{r}(0)} \Phi|x| \eta_{r} d x
\end{aligned}
$$

Thus

$$
\begin{align*}
& \limsup _{r \rightarrow 0} \int_{\Gamma_{r}^{\Omega}(0)}\left\langle x . \nabla \eta_{r}\right\rangle v d S \leq \limsup _{r \rightarrow 0} \lambda_{r} \int_{\Omega_{r}(0)} v|x| \eta_{r}(x) d x  \tag{3.10}\\
& \quad+\limsup _{r \rightarrow 0} \int_{\Omega_{r}(0)} v|x|^{-1}\left|\left\langle x . \nabla \eta_{r}\right\rangle\right| d x .
\end{align*}
$$

But

$$
|x| \eta_{r}(x) \leq C \rho_{\partial \Omega}(x)\left|\left\langle x . \nabla \eta_{r}\right\rangle\right| \leq C|x| / r,
$$

and more precisely,

$$
\left.\lim _{r \rightarrow 0}\left\langle x . \nabla \eta_{r}\right\rangle\right|_{|x|=r}=\phi_{1}(\sigma)=x_{1} /|x| .
$$

Then

$$
\begin{align*}
\int_{\Omega_{r}(0)} v|x| \eta_{r}(x) d x & \leq C \int_{\Omega_{r}(0)} \rho_{\partial \Omega}(x) v d x, \\
& \leq C\|v\|_{M^{(N+1) /(N-1)}\left(\Omega ; \rho_{\partial \Omega} d x\right)}\left(\int_{\Omega_{r}(0)} \rho_{\partial \Omega} d x\right)^{2 /(N+1)},  \tag{3.11}\\
& \leq C C^{\prime} r^{2}\|\Phi\|_{L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)},
\end{align*}
$$

and

$$
\begin{align*}
\int_{\Omega_{r}(0)} v|x|^{-1}\left|\left\langle x . \nabla \eta_{r}\right\rangle\right| d x & \leq C r^{-1} \int_{\Omega_{r}(0)} v d x \\
& \leq C r^{-1}\|v\|_{M^{N /(N-1)}(\Omega)}\left|\Omega_{r}(0)\right|^{1 / N}  \tag{3.12}\\
& \leq C C^{\prime}\|\Phi\|_{L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)}
\end{align*}
$$

Combining (3.10), (3.11) and (3.12) yields to

$$
\begin{equation*}
\limsup _{r \rightarrow 0} r^{-1} \int_{\Gamma_{r}^{\Omega}(0)} v \rho_{\partial \Omega} d S \leq C^{\prime}\|\Phi\|_{L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)} \leq C \varepsilon \tag{3.13}
\end{equation*}
$$

which ends the proof since $\varepsilon$ is arbitrary.
Lemma 3.4. Assume the assumptions of Theorem 3.1 are fulfilled, $u$ is a nonnegative solution of (2.6), $\lambda>0$ and $a \in \partial \Omega$. If the coordinates are proper at $a$ relatively to $\Omega$, then for any $q \in[1, \infty)$,

$$
\begin{equation*}
\lim _{r \rightarrow 0} r^{N-2}\left(\int_{\Gamma_{r}^{\Omega}(a)}\left|u_{\lambda \delta_{a}}(y)-\lambda P(y, a)\right|^{q} \rho_{\partial \Omega}(y) d S\right)^{1 / q}=0 \tag{3.14}
\end{equation*}
$$

Proof. Recall that $g\left(., u_{\lambda \delta_{a}}\right) \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$, we put $v=\mathbb{G}_{g\left(., u_{\lambda \delta_{a}}\right)}$. Since $u_{\lambda \delta_{a}}=\lambda P(., a)-v$, it follows from Lemma 3.2

$$
\begin{equation*}
\lim _{r \rightarrow 0} r^{-1} \int_{\Gamma_{r}^{\Omega}(a)}\left|u_{\lambda \delta_{a}}(y)-\lambda P(y, a)\right| \rho_{\partial \Omega}(y) d S=0 \tag{3.15}
\end{equation*}
$$

Since $0 \leq u_{\lambda \delta_{a}} \leq \lambda P(., a)$,

$$
y \mapsto r^{N-2} \sup _{y \in \Gamma_{r}^{\Omega}(a)}\left|u_{\lambda \delta_{a}}(y)-\lambda P(y, a)\right| \rho_{\partial \Omega}(y)
$$

is bounded independently of $r$. Thus the result follows by Hölder's inequality.
Under a pointwise growth estimate on $g$ the convergence of $u_{\lambda \delta_{a}}$ is much more precise.

Lemma 3.5. Let the conditions of Theorem 3.1 be fulfilled. Assume also that there exists $\varepsilon_{0}>0$ such that the mapping $(k, x) \mapsto k^{N+1} g\left(k(x-a)+a, k^{1-N}\right)$ remains boundedfor $(k, x) \in\left(0, \varepsilon_{0}\right] \times\left\{x \in a+k^{-1}(\Omega-a): 1-\varepsilon_{0} \leq|x| \leq 1+\varepsilon_{0}\right\}$. Then for any $\eta_{1}>0$,

$$
\begin{equation*}
\lim _{\substack{x \rightarrow a \\(x-a) /|x-a| \rightarrow \eta_{1}}}|x-a|^{N-1} u_{\lambda \delta_{a}}(x)=\lambda C(N) \eta_{1} \tag{3.16}
\end{equation*}
$$

for some constant $C(N)>0$. Moreover, for any $\eta>0$, the convergence is uniform in the cone $\eta_{1} \geq \eta$.

Proof. We can assume $a=0$ and set $u_{k}(x)=k^{N-1} u(k x)$. Then

$$
\Delta u_{k}(x)=k^{N+1} g(k x, u(k x))
$$

Since $0 \leq u \leq \lambda P(x, 0) \leq \lambda C(N)|x|^{1-N}, u_{k}(x)$ and $(k, x) \mapsto k^{N+1} g(k x, u(k x))$ remains bounded for $(k, x) \in\left(0, \varepsilon_{0}\right] \times\left\{x \in k^{-1} \Omega: 1-\varepsilon_{0} \leq|x| \leq 1+\varepsilon_{0}\right\}$. Thus $\left\{u_{k}\right\}$ is relatively compact in $\left\{x \in k^{-1} \Omega: 1-\varepsilon_{0} / 2 \leq|x| \leq 1+\varepsilon_{0} / 2\right\}$, and there exist a sequence $\left\{k_{n}\right\}$ and some function $\zeta \in C^{1}\left(\mathbb{R}_{+}^{N} \cap\left(\bar{B}_{1+\varepsilon_{0}}(0) \backslash B_{1-\varepsilon_{0}}(0)\right)\right)$ such that $u_{k_{n}} \rightarrow \zeta$ and $\nabla u_{k_{n}} \rightarrow \nabla \zeta$ uniformly on $\mathbb{R}_{+}^{N} \cap\left(\bar{B}_{1+\varepsilon_{0} / 2}(0) \backslash B_{1-\varepsilon_{0} / 2}(0)\right)$. Putting $|x|=1$, it implies

$$
\lim _{k_{n} \rightarrow 0} k_{n}^{N-1} u\left(k_{n}, \sigma\right)=\zeta(\sigma)
$$

uniformly on any compact subset of $S_{+}^{N-1}$. Since $P(x, 0)=P(r, \sigma, 0)=$ $C(N) r^{1-N} \phi_{1}(\sigma)$, with $\phi_{1}(\sigma)=\left.x_{1}\right|_{S_{+}^{N-1}}$, the relation (3.15) yields to $\zeta(\sigma)=$ $C(N) \lambda \phi_{1}(\sigma)$, and finally

$$
\lim _{k \rightarrow 0} k_{n}^{N-1} u(k, .)=C(N) \lambda \phi_{1}(.)
$$

When $u_{\lambda \delta_{a}}$ is replaced by $w_{\lambda \delta_{a}}$, the convergence is comparable to the one of Lemma 3.4.

Lemma 3.6. Let the assumption of Theorem 3.1 be fulfilled. If $\lambda>0, a \in \partial \Omega$ and the coordinates are proper at a relatively to $\Omega$, there holds

$$
\begin{equation*}
\lim _{r \rightarrow 0} r^{N-2}\left(\int_{\Gamma_{r}^{\Omega}(a)}\left|w_{\lambda \delta_{a}(y)}-\gamma(a, \lambda) P(y, a)\right|^{q} \rho_{\partial \Omega}(y) d S\right)^{1 / q}=0 \tag{3.17}
\end{equation*}
$$

for any $1 \leq q<\infty$.

Proof. Since $\Delta w_{\lambda \delta_{a}}$ and $g\left(., w_{\lambda \delta_{a}}\right)$ belong to $L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$, there exists $\Phi \in L^{1}\left(\Omega ; \rho_{\partial \Omega} d x\right)$ such that

$$
\begin{aligned}
-\Delta w_{\lambda \delta_{a}} & =\Phi \quad \text { in } \Omega \\
w_{\lambda \delta_{a}} & =\gamma(a, \lambda) \delta_{0} \text { on } \partial \Omega
\end{aligned}
$$

Then $w_{\lambda \delta_{a}}=\mathbb{G}_{\Phi}+\gamma(a, \lambda) P(., a)$ and

$$
\left|w_{\lambda \delta_{a}}-\gamma(a, \lambda) P(., a)\right| \leq \mathbb{G}_{|\Phi|} .
$$

By Lemma 3.3

$$
\lim _{r \rightarrow 0} r^{N-2} \int_{\Gamma_{r}^{\Omega}(a)}\left|\mathbb{G}_{\Phi}(y)\right| d S=0
$$

thus

$$
\begin{equation*}
\lim _{r \rightarrow 0} r^{N-2} \int_{\Gamma_{r}^{\Omega}(a)}\left|w_{\lambda \delta_{a}}(y)-\gamma(a, \lambda) P(y, a)\right| \rho_{\partial \Omega}(y) d S=0 \tag{3.18}
\end{equation*}
$$

Since $0 \leq w_{\lambda \delta_{a}} \leq \lambda P(., a)$, and $r^{N-2} \rho_{\partial \Omega} P(., a)$ is bounded on $\Gamma_{r}^{\Omega}(a)$, (3.17) follows.

Proof of Theorem 3.1. Up to a translation, we can assume that $a=0$. We can assume that $S_{+}^{N-1}$ is the intersection of the unit sphere with the half space $\left\{x_{1}>0\right\}$ and $\partial S_{+}^{N-1}$ the intersection of $S^{N-1}$ with the hyperplane $\left\{x_{1}=0\right\}$. Thus $\phi_{1}$, the first eigenvalue of the Laplace-Beltrami operator $-\Delta_{\sigma}$ in $W_{0}^{1,2}\left(S_{+}^{N-1}\right)$ is the restriction to $S^{N-1}$ of the coordinate function $x \mapsto x_{1}$, and the corresponding eigenvalue is $N-1$. We normalize by $\max \phi_{1}=1$.
CASE 1. $\gamma(a)<\infty$. For $\lambda>\gamma(a)$ the following convergences hold in $L^{q}\left(S_{+}^{N-1}\right)$ for $1 \leq q<\infty$ :

$$
\lim _{r \rightarrow 0} r^{N-1} w_{\lambda \delta_{a}}(r, .)=C(N) \gamma(a, \lambda) \phi_{1}
$$

by Lemma 3.6, and

$$
\lim _{r \rightarrow 0} r^{N-1} u_{\lambda \delta_{a}}(r, \sigma)=C(N) \lambda \phi_{1}(\sigma)
$$

by Lemma 3.4. If $\left\{r_{n}\right\}$ is some sequence converging to 0 , there exists a subsequence $\left\{r_{n_{k}}\right\}$ such that

$$
\lim _{r_{n_{k}} \rightarrow 0} r_{n_{k}}^{N-1} w_{\lambda \delta_{a}}\left(r_{n_{k}}, \sigma\right)=C(N) \tilde{\gamma}_{u}(a, \lambda) \phi_{1}(\sigma),
$$

and

$$
\lim _{r_{n_{k}} \rightarrow 0} r_{n_{k}}^{N-1} u_{\lambda \delta_{a}}\left(r_{n_{k}}, \sigma\right)=C(N) \lambda \phi_{1}(\sigma),
$$

for almost all $\sigma \in S_{+}^{N-1}$. Therefore for almost all $\sigma \in S_{+}^{N-1}$, there exists $n_{k_{0}}$ such that for $n_{k} \geq n_{k_{O}}, w_{\lambda \delta_{a}}\left(r_{n_{k}}, \sigma\right)=u\left(r_{n_{k}}, \sigma\right)$. Consequently there holds

$$
\lim _{r \rightarrow 0} r^{N-1} u\left(r_{n_{k}}, \sigma\right)=C(N) \tilde{\gamma}_{u}(a, \lambda) \phi_{1}(\sigma)
$$

for almost all $\sigma \in S_{+}^{N-1}$. Let $\theta>\lambda$. It follows from Lemma 3.5 applied to the $w_{\theta \delta_{a}}\left(r_{n_{k}},.\right)$ and the previous argument, that, up to some subsequence $r_{n_{k}}$,

$$
\lim _{r \rightarrow 0} r^{N-1} u\left(r_{n_{k_{\ell}}}, \sigma\right)=C(N) \tilde{\gamma}_{u}(a, \theta) \phi_{1}(\sigma)
$$

almost everywhere. Therefore $\tilde{\gamma}_{u}(a, \theta)=\tilde{\gamma}_{u}(a, \lambda)=\tilde{\gamma}_{u}(a)$. This infers (i).
Case 2. $\tilde{\gamma}_{u}(a)=\infty$. From Lemma 2.1,

$$
\min \left(u, u_{\infty, a}\right) \geq u_{\infty, a} \Longrightarrow u \geq u_{\infty, a}>u_{\lambda \delta_{a}}, \quad \forall \lambda>0
$$

By Lemma 3.4, for any $\varepsilon>0$,

$$
\begin{equation*}
\underset{\substack{x \rightarrow a \\\left(x_{1}-a_{1}\right) /|x-a| \rightarrow \eta_{1} \\ \eta_{1} \geq \varepsilon}}{\liminf }|x-a|^{N-1} u(x) \geq C(N) \lambda \varepsilon \tag{3.19}
\end{equation*}
$$

Since $\lambda$ is arbitrary, (ii) holds.

Remark 3.1. In the core of the proof in Case 1 we have seen that $\tilde{\gamma}_{u}(\lambda, a)=\tilde{\gamma}_{u}(a)$ for any $\lambda>\gamma(a)$. Actually the same proof gives also $\tilde{\gamma}_{u}(\lambda, a)=\tilde{\gamma}_{u}(a)$ for $\lambda=\tilde{\gamma}_{u}(a)$.

Remark 3.2. If it is supposed moreover that $(k, x) \mapsto k^{N+1} g(k(x-a)+$ $\left.a, k^{1-N}\right)$ remains bounded for $(k, x) \in\left(0, \varepsilon_{0}\right] \times\left\{x \in a+k^{-1}(\Omega-a): 1-\varepsilon_{0} \leq\right.$ $\left.|x| \leq 1+\varepsilon_{0}\right\}$, assertion (ii) can be replaced by:
or
(ii*) $\tilde{\gamma}_{u}(a)$ is infinite and

$$
\begin{equation*}
\lim _{\substack{x \rightarrow a \\\left(x_{1}-a_{1}\right) /|x-a| \rightarrow \eta_{1}}}|x-a|^{N-1} u(x)=\infty \tag{3.20}
\end{equation*}
$$

uniformly for $\eta_{1}>0$.

## 4. - Boundary trace of solutions with uniform absorption

In this section $\Omega$ is a $C^{2}$ bounded domain in $\mathbb{R}^{N}, g \in \mathcal{G}_{0}$ satisfies a uniform condition with respect to $x$ in the sense that

$$
\begin{align*}
& 0 \leq|g(x, r)| \leq f(r), \forall(x, r) \in \Omega \times \mathbb{R}_{+} \\
& \text {with } \quad \int_{0}^{1} f\left(\sigma s^{1-N}\right) s^{N} d s<\infty, \forall \sigma>0 \tag{4.1}
\end{align*}
$$

where $f$ is a continuous nondecreasing function defined on $\mathbb{R}_{+}$. The next result provides a precise characterisation of the boundary trace of solutions of inequalities with a uniform absorption in terms of outer regular Borel measures, without introducing the notion of coercivity and the strong barrier property as in [27].

Theorem 4.1. Assume $g \in \mathcal{G}_{0}$ satisfies (4.1) and $u$ is a nonnegative solution of (2.6) with boundary trace $\nu$. For any $a \in \partial \Omega$ the following dichotomy occurs. Either,
(i) $v(\mathcal{O})=\infty$ for any $\mathcal{O} \in \mathcal{N}_{a}$. In this case $a \in \mathcal{S}(u)$ and $u \geq u_{\infty, a}$. Consequently

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\mathcal{O}_{t}} u(y) d S_{t}=\infty, \forall \mathcal{O} \in \mathcal{N}_{a} \tag{4.2}
\end{equation*}
$$

Or
(ii) there exists $\mathcal{O} \in \mathcal{N}_{a}$ such that $\nu(\mathcal{O})<\infty$. In this case $a \in \mathcal{R}(u)$ and

$$
\begin{equation*}
\sup _{0<t \leq \beta_{0}} \int_{\mathcal{O}_{t}^{\prime}} u(y) d S_{t}<\infty \tag{4.3}
\end{equation*}
$$

for relatively every open subset $\mathcal{O}^{\prime} \subset \overline{\mathcal{O}^{\prime}} \subset \mathcal{O}$. Furthermore

$$
\begin{equation*}
\lim _{t \rightarrow 0} \int_{\Sigma_{t}} u(y) \phi(\sigma(y))(y) d S_{t}=\int_{\mathcal{R}(u)} \phi(y) d v(y), \quad \forall \phi \in C_{c}(\mathcal{R}(u)) \tag{4.4}
\end{equation*}
$$

A major point in the proof of the theorem is the following completion of Proposition 2.7 which gives a characterization of the regular part of the extended boundary trace of a solution $u$ based upon a local $L^{1}$ bound.

Proposition 4.1. Assume $g \in \mathcal{G}_{0}$ satisfies (4.1) and $u$ is a nonnegative solution of (2.6) with extended boundary trace $\nu$. Let $\mathcal{O}$ be a relatively open subset of $\partial \Omega$. If $\nu(\mathcal{O})<\infty$, then for any compact subset $K \subset \mathcal{O}, \int_{K_{t}} u(y) d S_{t}$ remains bounded independently of $t \in\left(0, \beta_{0}\right]$.

We recall some notations introduced in Section 2 : for $0<\beta \leq \beta_{0}$, we put $\Omega_{\beta}^{c}=\Omega \backslash \bar{\Omega}_{\beta}=\left\{x \in \Omega: \rho_{\partial \Omega}(x)>\beta\right\}$, and $\Sigma_{\beta}=\partial \Omega_{\beta}=\partial \bar{\Omega}_{\beta}^{c}$. The next result which extends Theorem 1.1 deals with the stability of the boundary value problem when the boundary is approximated from inside by a sequence of smooth domains.

Lemma 4.1. Let $\mu \in \mathfrak{M}_{+}(\partial \Omega),\left\{\varepsilon_{n}\right\}$ be a sequence of positive numbers converging to $0, \mu_{n} \in L_{+}^{1}\left(\Sigma_{\varepsilon_{n}}\right)$, with corresponding pull-back $\mu_{n}^{\varepsilon_{n}} \in L_{+}^{1}(\partial \Omega)$. If $\mu_{n}^{\varepsilon_{n}} \rightarrow \mu$ in the weak sense of measures, as $n \rightarrow \infty$, then the sequence of solutions $u_{n}=u_{\mu_{n}, \varepsilon_{n}}$ of

$$
\begin{align*}
-\Delta u_{n}+g\left(x, u_{n}\right) & =0 \quad \text { in } \Omega_{\varepsilon_{n}}^{c} \\
u_{n} & =\mu_{n} \text { on } \partial \Omega_{\varepsilon_{n}}^{c} \tag{4.5}
\end{align*}
$$

converges locally uniformly in $\Omega$ to the solution $u_{\mu}$ of (0.4).
Proof. Step 1. Since $g$ is continuous on $\bar{\Omega}_{\varepsilon_{n}}^{c} \times \mathbb{R}$, existence of a unique solution to (4.6) follows from Brezis' result. The shred of the proof of the convergence is parallel to Theorem 1.1 and Proposition 1.4. Set $P^{\varepsilon_{n}}$ the Poisson kernel in $\Omega_{\varepsilon_{n}}^{c}$ and $\mathbb{P}_{\nu}^{\varepsilon_{n}}$ the Poisson potential of any given Radon measure $v$ on $\Sigma_{\varepsilon_{n}}$. Since the mapping $\Pi$ is $C^{2}$ there exists $C_{2}>0$ independent $\varepsilon_{n}$ such that for any $(x, a) \in \Omega_{\varepsilon_{n}}^{c} \times \Sigma_{\varepsilon_{n}}$,

$$
\begin{equation*}
C_{2}^{-1} \rho_{\partial \Omega_{\varepsilon_{n}}^{c}}(x)|x-a|^{-N} \leq P^{\varepsilon_{n}}(x, a) \leq C_{2} \rho_{\partial \Omega_{\varepsilon_{n}}^{c}}(x)|x-a|^{-N}, \tag{4.6}
\end{equation*}
$$

provided $0 \leq \varepsilon_{n} \leq \beta_{0}$, and $\rho_{\partial \Omega_{\varepsilon_{n}}^{c}}(x)=\rho_{\partial \Omega}(x)-\varepsilon_{n}$. Estimates (1.22), (1.23), (1.24) are valid under the form

$$
\begin{align*}
\left\|\mathbb{P}_{v}^{\varepsilon_{n}}\right\|_{M^{(N+1) /(N-1)}\left(\Omega_{\varepsilon_{n}}^{c} ; \rho_{\partial \Omega_{\varepsilon_{n}}^{c}} d x\right)} & \leq K\|\nu\|_{L^{1}\left(\Sigma_{\varepsilon_{n}}\right)}  \tag{4.7}\\
\left\|\mathbb{P}_{v}^{\varepsilon_{n}}\right\|_{M^{N /(N-1)}\left(\Omega_{\varepsilon_{n}}^{c}\right)} & \leq K\|\nu\|_{L^{1}\left(\Sigma_{\varepsilon_{n}}\right)}  \tag{4.8}\\
\left\|\mathbb{P}_{v}^{\varepsilon_{n}}\right\|_{L^{\infty}\left(\Omega_{r+\varepsilon_{n}}^{c}\right)} & \leq K r^{1-N}\|\nu\|_{L^{1}\left(\Sigma_{\varepsilon_{n}}\right)} \tag{4.9}
\end{align*}
$$

Since $0 \leq u_{n} \leq v_{n}=\mathbb{P}_{\mu_{n}}^{\varepsilon_{n}}$, estimates (4.9) and (1.16)-(1.18) and the classical regularity theory for elliptic equations imply that the set of $u_{n}$ is relatively compact in the $C_{\mathrm{loc}}^{1}(\Omega)$-topology, and any cluster point of the sequence $\left\{u_{n}\right\}$ is a solution of (0.1) in $\Omega$. If $\eta \in C_{c}^{1,1}\left(\bar{\Omega}_{\varepsilon_{n}}^{c}\right)$ there holds

$$
\begin{equation*}
\int_{\Omega_{\varepsilon_{n}}^{c}}\left(-u_{n} \Delta \eta+g\left(., u_{n}\right) \eta\right) d y=-\int_{\Sigma_{n}} \frac{\partial \eta}{\partial \mathbf{n}_{y}} \mu_{n} d S(y) \tag{4.10}
\end{equation*}
$$

If $\zeta \in C_{c}^{1,1}(\bar{\Omega})$, with support in $\bar{\Omega}_{\beta_{0}}$, we set

$$
\zeta_{n}(y)=\zeta\left(y+\varepsilon_{n} \mathbf{n}_{y}\right), \quad \forall y \in \bar{\Omega}_{\varepsilon_{n}}^{c} \Longleftrightarrow \zeta(x)=\zeta_{n}\left(x-\varepsilon_{n} \mathbf{n}_{x}\right), \quad \forall x \in \bar{\Omega}
$$

In relation (4.10) we take $\eta=\zeta_{n}$ and get

$$
\begin{equation*}
\int_{\Omega_{\varepsilon_{n}}^{c}}\left(-u_{n} \Delta \zeta_{n}+g\left(., u_{n}\right) \zeta_{n}\right) d y=-\int_{\Sigma_{n}} \frac{\partial \zeta_{n}}{\partial \mathbf{n}} \mu_{n} d S_{\varepsilon_{n}} \tag{4.11}
\end{equation*}
$$

But the pointing outward normal vector $\mathbf{n}_{y}$ at $y \in \Sigma_{n}$ is the same as the pointing outward normal vector at $y+\varepsilon_{n} \mathbf{n}_{y} y \in \partial \Omega$, therefore

$$
\begin{equation*}
\int_{\Sigma_{n}} \frac{\partial \zeta_{n}}{\partial \mathbf{n}} \mu_{n} d S_{\varepsilon_{n}}=\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} \mu_{n}^{\varepsilon_{n}} d S \tag{4.12}
\end{equation*}
$$

by (2.5). Moreover, if we perform the change of variable $x=y+\varepsilon_{n} \mathbf{n}_{y}$ in (4.11) and denote by $\mathbf{n}_{y}^{j}$ the coordinates of $\mathbf{n}_{y}$, we get

$$
\begin{aligned}
\frac{\partial \zeta_{n}}{\partial y_{i}}(y)= & \sum_{j} \frac{\partial \zeta}{\partial x_{j}}(x)\left(\delta_{i j}+\varepsilon_{n} \frac{\partial \mathbf{n}_{y}^{j}}{\partial y_{i}}\right) \\
\frac{\partial^{2} \zeta_{n}}{\partial y_{i}^{2}}(y)= & \sum_{k, j} \frac{\partial^{2} \zeta}{\partial x_{k} \partial x_{j}}(x)\left(\delta_{i j}+\varepsilon_{n} \frac{\partial \mathbf{n}_{y}^{j}}{\partial y_{i}}\right)\left(\delta_{i k}+\varepsilon_{n} \frac{\partial \mathbf{n}_{y}^{k}}{\partial y_{i}}\right) \\
& +\varepsilon_{n} \sum_{j} \frac{\partial \zeta}{\partial x_{j}}(x)\left(\frac{\partial^{2} \mathbf{n}_{y}^{j}}{\partial y_{i}^{2}}\right), \\
\Delta \zeta_{n}(y)= & \sum_{i} \frac{\partial^{2} \zeta}{\partial x_{i}}(x)\left(1+\varepsilon_{n} \frac{\partial \mathbf{n}_{y}^{i}}{\partial y_{i}}\right)^{2} \\
& +\sum_{\substack{i, j, k \\
k \neq i}} \frac{\partial^{2} \zeta}{\partial x_{k} \partial x_{j}}(x)\left(\delta_{i j}+\varepsilon_{n} \frac{\partial \mathbf{n}_{y}^{j}}{\partial y_{i}}\right)\left(\delta_{i k}+\varepsilon_{n} \frac{\partial \mathbf{n}_{y}^{k}}{\partial y_{i}}\right) \\
& +\varepsilon_{n} \sum_{j} \frac{\partial \zeta}{\partial x_{j}}(x) \Delta \mathbf{n}_{y}^{j} .
\end{aligned}
$$

Then

$$
\begin{equation*}
\Delta \zeta_{n}(y)=\Delta \zeta(x)+\varepsilon_{n} \mathcal{L}\left(D \zeta(x), D^{2} \zeta(x)\right) \tag{4.13}
\end{equation*}
$$

where $\mathcal{L}\left(D \zeta, D^{2} \zeta\right)$ is a linear second order operator with continuous coefficients. Plugging (4.12) and (4.13) into (4.11) yields

$$
\begin{equation*}
\int_{\Omega}\left(-\tilde{u}_{n}\left(\Delta \zeta+\varepsilon_{n} \mathcal{L}\left(D \zeta, D^{2} \zeta\right)\right)+g_{n}\left(., \tilde{u}_{n}\right) \zeta\right) J d x=-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} \mu_{n}^{\varepsilon_{n}} d S \tag{4.14}
\end{equation*}
$$

where $\tilde{u}_{n}(x)=u_{n}\left(x-\varepsilon_{n} \mathbf{n}_{x}\right), g_{n}(x, r)=g\left(x-\varepsilon_{n} \mathbf{n}_{x}, r\right)$ and $J(x)=\left|\operatorname{det}\left(I-\varepsilon_{n} D \mathbf{n}_{y}\right)\right|$.
Step 2 From (1.26) and (4.6) there exists $C_{3}>0$ independent of $\varepsilon_{n}$ and $(y, b) \in \Omega_{\varepsilon_{n}}^{c} \times \Sigma_{n}$ such that

$$
\begin{equation*}
C_{3}^{-1} P\left(y+\varepsilon_{n} \mathbf{n}_{y}, b+\varepsilon_{n} \mathbf{n}_{y}\right) \leq P^{\varepsilon_{n}}(y, b) \leq C_{3} P\left(y+\varepsilon_{n} \mathbf{n}_{y}, b+\varepsilon_{n} \mathbf{n}_{y}\right), \tag{4.15}
\end{equation*}
$$

provided $\varepsilon_{n} \leq \beta_{0}$. Because $0 \leq u_{n}(y) \leq \mathbb{P}_{\mu_{n}}^{\varepsilon_{n}}$, the above inequality and the Poisson representation formula imply

$$
\begin{equation*}
0 \leq \tilde{u}_{n} \leq C_{3} \mathbb{P}_{\mu_{n}^{\varepsilon_{n}}} \tag{4.16}
\end{equation*}
$$

in $\bar{\Omega}_{\beta_{0}-\varepsilon_{n}}$. Jointly with (1.23) it implies that $\left\{\tilde{u}_{n}\right\}$ is uniformly integrable in $\Omega_{\beta_{0}-\varepsilon_{n}}$, and thus in $\Omega$.
Step 3 From (4.1), (4.16),

$$
\begin{equation*}
0 \leq g_{n}\left(x, \tilde{u}_{n}\right)(x) \leq f\left(C_{3} \mathbb{P}_{\mu_{n}^{\varepsilon_{n}}}(x)\right) \tag{4.17}
\end{equation*}
$$

for any $x \in \Omega_{\beta_{0}-\varepsilon_{n}}$. For $\lambda \geq 0$, put $\Gamma_{\lambda}=\{x \in \Omega: P(x, a) \geq \lambda\}$ and

$$
\beta_{\varepsilon_{n}}(\lambda)=\int_{\Gamma_{\lambda}} \rho_{\partial \Omega} d x
$$

Then (see Step 2 in the proof of Theorem 1.1),

$$
\begin{equation*}
\beta_{\varepsilon_{n}}(\lambda) \leq C_{2} \int_{0}^{\left(C_{1} / \lambda\right)^{1 /(N-1)}} s^{N} d s \leq \frac{C_{2}}{N+1}\left(\frac{C_{1}}{\lambda}\right)^{(N+1) /(N-1)} \tag{4.18}
\end{equation*}
$$

It follows from (4.17) that for any Borel set $G \subset \Omega$,

$$
\begin{equation*}
\int_{G} g_{n}\left(., \tilde{u}_{n}\right) \rho_{\partial \Omega} d x \leq \int_{G} f\left(C_{3} \mathbb{P}_{\mu_{n}^{\varepsilon_{n}}}\right) \rho_{\partial \Omega} d x \tag{4.19}
\end{equation*}
$$

In order to estimate the right-hand side of (4.19), we follow the proof of Theorem 1.1. Let $m>0$ and $a \in \partial \Omega$, then

$$
\begin{align*}
& \int_{G} f(m P(., a)) \rho_{\partial \Omega} d x  \tag{4.20}\\
& \quad \leq f(\lambda) \int_{G} \rho_{\partial \Omega} d x+C_{4} m^{(N+1) /(N-1)} \int_{\lambda}^{\infty} f(s) s^{-2 N /(N-1)} d s
\end{align*}
$$

We take $m=m_{n}=\int_{\partial \Omega} d \mu_{n}^{\varepsilon_{n}}$ and we deduce that the $\left\{g_{n}\left(., \tilde{u}_{n}\right)\right\}$ are uniformly integrable, as in the proof of Theorem 1.1-Step 7. If $\tilde{u}_{n_{k}}$ is a subsequence converging in the $C_{\text {loc }}^{1}$ topology to some function $u$, we can pass to the limit in (4.14) and get

$$
\begin{equation*}
\int_{\Omega}(-u \Delta \zeta+g(., u) \zeta) d y=-\int_{\partial \Omega} \frac{\partial \zeta}{\partial \mathbf{n}} d \mu \tag{4.21}
\end{equation*}
$$

Because of uniqueness, the whole sequence $\tilde{u}_{n}$ converges to $u$.

Remark 4.1. This above stability result of solutions with respect to approximations from inside is no longer valid if the absorption term is truly degenerate, for example if

$$
g(x, r)=\rho_{\partial \Omega}^{\alpha}(x)|u|^{q-1} u,
$$

for some $\alpha>0$ and $q>1$. In that case Problem (0.4) is solvable in $\Omega$ for any Radon measure $\mu$ when $1<q<(N+1+\alpha) /(N-1)$ and is not solvable when $\mu$ is a Dirac mass if $q \geq(N+1+\alpha) /(N-1)$ (see [27]). Therefore, even if the data $\mu_{n}$ are smooth functions on $\Sigma_{\varepsilon_{n}}$, if they concentrate too quickly to a Dirac mass on the boundary, the corresponding solutions $u_{n}$ of (4.6) converge to 0 .

Proof of Proposition 4.1. We proceed by contradiction in assuming that there exist a compact $K \subset \mathcal{O}$ and a sequence $\left\{\varepsilon_{n}\right\}$ converging to 0 such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{K_{\varepsilon_{n}}} u(y) d S_{\varepsilon_{n}}=\infty \tag{4.22}
\end{equation*}
$$

Since $K$ is compact, there exist a sequence $\left\{a_{n}\{\subset K\right.$ converging to some $a \in K$ and a sequence $\left\{t_{n}\right\}$ converging to 0 such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{K_{\varepsilon_{n} \cap B_{t_{n}}\left(a_{n}\right)}} u(y) d S_{\varepsilon_{n}}=\infty \tag{4.23}
\end{equation*}
$$

For $k>0$, there exists $\ell_{k}>0$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} \int_{K_{\varepsilon_{n}} \cap B_{t_{n}}\left(a_{n}\right)} \min \left\{\ell_{k}, u(y)\right\} d S_{\varepsilon_{n}}=k \tag{4.24}
\end{equation*}
$$

and $\ell_{k} \rightarrow 0$ as $n \rightarrow \infty$. We set $\mu_{n}=\min \left\{\ell_{k}, u(y)\right\} \chi_{K_{\varepsilon_{n} \cap B_{t_{n}}\left(a_{n}\right)}}$ and denote by $u_{n}$ the solution of (4.6) in $\Omega_{\varepsilon_{n}}^{c}$ with this specific boundary data. Then

$$
u \geq u_{n} \quad \text { in } \quad \Omega_{\varepsilon_{n}}^{c} .
$$

Since the corresponding measure $\gamma_{u}\left(\mu_{n}\right)=\mu_{n}^{\varepsilon_{n}}$ on $\partial \Omega$ converges to $k \delta_{a}$, and $u_{n}$ converges to $u_{k \delta_{a}}$ by Lemma 4.1, it leads to $u \geq u_{k \delta_{a}}$ in $\Omega$. Since $k$ is arbitrary,

$$
u \geq u_{\infty, a} \quad \text { in } \Omega
$$

Therefore

$$
v(\mathcal{O}) \geq v(a)=\tilde{\gamma}(a)=\infty
$$

by Proposition 2.7, a contradiction.

For any $a \in \partial \Omega$, we recall that $\mathcal{N}_{a}$ is the set of relatively open neighborhoods of $a$ in $\partial \Omega$.

Proof of Theorem 4.1. Let $a \in \partial \Omega$. If (i) holds, inequality $u \geq u_{\infty, a}$ follows from Proposition 2.7, and

$$
\lim _{t \rightarrow 0} \int_{\mathcal{O}_{t}} u(y) d S_{t} \geq \lim _{t \rightarrow 0} \int_{\mathcal{O}_{t}} u_{\infty, a}(y) d S_{t}=\infty
$$

which is equivalent to $v(\mathcal{O})=\infty$.
Next we assume that (i) does not hold, and there exists $\mathcal{O} \in \mathcal{N}_{a}$ such that $\nu(\mathcal{O})<\infty$. By Proposition 4.1, for any compact subset $\mathcal{K} \in \mathcal{N}_{a}$ such that $\overline{\mathcal{K}} \subset \mathcal{O}$, there exists a constant $M_{K}>0$ such that

$$
\int_{\mathcal{K}_{t}} u(y) d S_{t} \leq M_{K}, \quad \text { on }\left(0, \beta_{0}\right] .
$$

Let $\mathcal{O}$ be any relatively open subset with compact closure in $\mathcal{R}(u), 0<\beta<\beta_{0}$ and

$$
\mathfrak{O}_{\beta}=\left\{x=\sigma(x)-t \mathbf{n}_{x}: \sigma(x) \in \mathcal{O}, \beta<t<\beta_{0}\right\}
$$

then $u \in L^{1}\left(\mathfrak{O}_{0}\right)$. As in the proof of [27, Lemma 1.8], if $\varphi \in C_{c}^{2}(\mathcal{O}), \varphi>0$, we define a test function which vanishes on $G s_{\beta}$ by

$$
\zeta(x)=\varphi(\sigma(x))\left(\rho_{\partial \Omega}(x)-\beta\right) \quad \forall x \in \mathfrak{O}_{\beta}
$$

and derive that the largest solution $u^{*}$ of (0.1) dominated by $u$ satisfies

$$
\begin{align*}
\int_{\mathfrak{O}_{\beta}}\left(-u^{*} \Delta \zeta+g\left(x, u^{*}\right) \zeta\right) d x= & \int_{\mathcal{O}_{\beta}} u^{*} \varphi^{\beta} d S_{\beta}-\int_{\mathcal{O}_{\beta_{0}}} u^{*} \varphi^{\beta_{0}} d S_{\beta_{0}} \\
& +\int_{\mathcal{O}_{\beta_{0}}} \frac{\partial u^{*}}{\partial \mathbf{n}} \zeta d S \tag{4.25}
\end{align*}
$$

Therefore $\int_{\mathfrak{O}_{\beta}} g\left(x, u^{*}\right) \zeta d x$ is bounded independently of $\beta$. Letting $\beta \rightarrow 0$ yields to

$$
\begin{equation*}
\int_{\mathfrak{O}_{0}} g\left(x, u^{*}\right) \varphi \rho_{\partial \Omega} d x<\infty \tag{4.26}
\end{equation*}
$$

Since $g\left(x, u^{*}\right) \in L^{1}\left(\mathfrak{O}_{0} ; \rho_{\partial \Omega} d x\right)$, [27, Corollary 1.3] implies that there exists a nonnegative Radon measure $\mu_{\mathcal{O}}$ on $\mathcal{O}$ such that

$$
\lim _{t \rightarrow 0} \int_{\mathcal{O}_{t}} \varphi_{t} u^{*}(y) d S_{t}=\int_{\mathcal{O}} \varphi d \mu
$$

for any $\varphi \in C_{c}(\mathcal{O})$. The measure $\nu$, which is equal to $\mu$ on $\mathcal{R}(u)$, is therefore a regular Borel measure.

Because of (4.3) from any sequence $\left\{\epsilon_{n}\right\}$ converging to 0 one can extract a subsequence, still denoted by $\left\{\epsilon_{n}\right\}$, such that $\left\{u\left(\epsilon_{n},\right) \chi_{\mathcal{R}(u)} d S\right\}$ converges in the weak sense of measures to some $\eta \in \mathfrak{M}_{+}(\mathcal{R}(u))$. We claim that

$$
\begin{equation*}
\eta=\left.\nu\right|_{\mathcal{R}(u)}=\chi_{\mathcal{R}(u)} \nu . \tag{4.27}
\end{equation*}
$$

Since $u \geq u^{*}, \eta \geq\left.\nu\right|_{\mathcal{R}(u)}$. If $\mathcal{O}$ is any relatively open subset of $\partial \Omega$ with compact closure in $\mathcal{R}(u)$, we put $\mu_{n}=u\left(\epsilon_{n},\right) \chi_{\mathcal{R}(u)} d S$ and denote by $u_{n}$ the solution of (4.6) in $\Omega_{\epsilon_{n}}^{c}$. By Lemma 4.1, $\left\{u_{n}\right\}$ converges locally uniformly in $\Omega$ to the solution $\tilde{u}$ of (0.4) with boundary data $\chi_{\mathcal{O}} \eta$. Since $u \geq u_{n}, u \geq \tilde{u}$ and thus $u^{*} \geq \tilde{u}$. Therefore $\chi_{\mathcal{O}} \eta \leq v_{\mathcal{O}}$. This implies (4.27). Finally, as $\eta$ is independent of the sequence $\left\{\epsilon_{n}\right\}$, it follows that $u(t,.) \chi_{\mathcal{R}(u)} d S$ converges to $\chi_{\mathcal{R}(u)} \nu$ in the sense of measures, as $t \rightarrow 0$. This ends the proof.

## 5. - Some examples

In this section $\Omega$ is a $C^{2}$ bounded domain and we consider absorption terms $g$ which are split under the form

$$
\begin{equation*}
g(x, r)=\tilde{h}(x) \tilde{g}(r) \tag{5.1}
\end{equation*}
$$

where both $\tilde{h}$ and $\tilde{g}$ are nonnegative continuous functions defined respectively on $\Omega$ and $\mathbb{R}_{+}$. We assume also that $\tilde{g}$ vanishes at 0 and is nondecreasing, and that $\tilde{h}(x)>0$ in $\Omega$. If the the Keller-Osserman condition is fulfilled, that is there exists some $c \geq 0$ such that

$$
\begin{equation*}
\int_{c}^{\infty} \frac{d s}{\sqrt{G(s)}}<\infty \tag{5.2}
\end{equation*}
$$

where $G(s)=\int_{0}^{s} \tilde{g}(t) d t$, then for any compact subset $K \subset \Omega$ there exists a constant $C(K)>0$ such that any nonnegative solution $u$ of

$$
\begin{equation*}
-\Delta u+g(x, u) \leq 0, \quad \text { in } \Omega \tag{5.3}
\end{equation*}
$$

satisfies

$$
u(x) \leq C(K), \quad \forall x \in K
$$

If the Keller-Osserman condition is not satisfied, and $\tilde{h}$ is a positive constant, no such a priori upper bound can exist [31]. If $\tilde{g}(r)=k r$ for some $k>0$ and $g$ is uniformly admissible, it is clear that

$$
\begin{equation*}
\lim _{\lambda \rightarrow \infty} u_{\lambda \delta_{a}}(x)=\infty, \quad \forall a \in \partial \Omega, \forall x \in \Omega \tag{5.4}
\end{equation*}
$$

but it appears difficult to find a general condition on $\tilde{g}$ which implies that (5.4) holds. However, it is proved in [13] that if

$$
g(x, r)=r(\ln (r+1))^{\gamma}
$$

with $0<\gamma \leq 2$, this property holds. As a consequence we have the following
Corollary 5.1. Let u be a nonnegative solution of

$$
\begin{equation*}
-\Delta u+u(\ln (u+1))^{\gamma}=0, \quad \text { in } \Omega \tag{5.5}
\end{equation*}
$$

Then the boundary trace of $u$ is a Radon measure $\mu, u(\ln (u+1))^{\gamma} \in L^{1}\left(\Omega ; \rho_{\partial \omega} d x\right)$ and

$$
u=\mathbb{P}_{\mu}-\mathbb{G}\left(u(\ln (u+1))^{\gamma}\right) .
$$

Proof. It is clear that $g(u)=u(\ln (u+1))^{\gamma}$ is uniformly admissible. By Theorem 4.1, $\operatorname{Tr}_{\partial \Omega}^{e}(u)$ is an outer regular Borel measure, which admits no singular part by [13], therefore it is a Radon measure. Thus the nonlinearity is integrable for the measure $\rho_{\partial \omega} d x$, and the representation follows.

Remark 5.1. If

$$
g(x, r)=r(\ln (r+1))^{\gamma}
$$

with $\gamma>2$, or if

$$
g(x, r)=\rho_{\partial \Omega}^{\alpha}(x)|r|^{q-1} r
$$

with $\alpha>-2$ and $1<q<(N+1+\alpha) /(N-1)$, it is proved respectively in [13] and [27] that for any $a \in \partial \Omega, u_{\infty, a}$ is a solution of (0.1) in $\Omega$ vanishing on $\partial \Omega \backslash\{a\}$, with a strong singularity at $a$. In those two cases the boundary trace of a nonnegative solution of (0.1) can be any outer regular Borel measure on $\partial \Omega$.

Another interesting type of problems deals with the situation in which the absorption term is strongly degenerate at the boundary. The model example is

$$
g(x, r)=\exp \left(-\kappa / \rho_{\partial \omega}(x)\right) u^{q}
$$

with $q>1$ and $\kappa>0$. In this case the function $g$ belongs to $\mathcal{H} \mathcal{G}_{0}$ for any $q>1$. Therefore fundamental solutions always exist, but a new phenomenon appears which is to be compared with what is called instantaneous or complete blow-up for parabolic equations ([2]) or elliptic equations ([7]), linear or nonlinear, with an inverse square potential.

Proposition 5.1. For any $q>1, \kappa>0$ and $a \in \partial \Omega$

$$
u_{\infty, a}=u_{m}
$$

where $u_{m}$ is the minimal solution of

$$
\begin{align*}
&-\Delta u+\exp \left(-\kappa / \rho_{\partial \omega}(x)\right) u^{q}=0 \text { in } \Omega, \\
& \lim _{\partial \Omega}(x) \rightarrow 0 \tag{5.6}
\end{align*} u(x)=\infty .
$$

Proof. Without any loss of generality we assume $a=O$. Let $\left\{x=\left(x_{1}, x^{\prime}\right)\right\}$ be the coordinates in $\mathbb{R}^{N}$. We assume that the hyperplane $H_{0}=\left\{x=\left(0, x^{\prime}\right)\right\}$ is tangent to $\partial \Omega$ at $O$ and $S_{+}^{N-1}=B_{1}(O) \cap\left\{x=\left(x_{1}, x^{\prime}\right): x_{1}>0\right\}$. We write

$$
\exp \left(-\kappa / \rho_{\partial \omega}(x)\right)=h\left(\rho_{\partial \omega}(x)\right)
$$

Step 1 The case $1<q<(N+1) /(N-1)$. For $0<\varepsilon \leq \beta_{0}$, the function $u_{\infty, O}$ is minorized in $\Omega_{\varepsilon}^{m}=\left\{x \in \Omega:|x|<m, 0<\rho_{\partial \omega}(x)<\varepsilon\right\}$ ( $m>0$ small enough) by the function

$$
\ell(\varepsilon) U_{\varepsilon}
$$

where $\ell(\varepsilon)=h^{-1 /(q-1)}(\varepsilon)$ and $U_{\varepsilon}$ is the unique solution of

$$
\begin{align*}
-\Delta v=v^{q} & \text { in } \Omega_{\varepsilon}^{m} \\
v=\infty \delta_{O} & \text { on } \partial \Omega_{\varepsilon}^{m} \tag{5.7}
\end{align*}
$$

Moreover there holds (see [15])

$$
\begin{equation*}
\lim _{\substack{x \rightarrow O \\ x \in \Omega_{\varepsilon}}}|x|^{2 /(q-1)} U_{\varepsilon}(x)=\omega(x /|x|) \tag{5.8}
\end{equation*}
$$

where $\omega$ is the unique positive solution of

$$
\begin{align*}
-\Delta_{\sigma} \omega-\left(\frac{2}{q-1}\right)\left(\frac{2 q}{q-1}-N\right) \omega+\omega^{q} & =0
\end{aligned} \begin{aligned}
& \text { in } S_{+}^{N-1}  \tag{5.9}\\
\omega & =0
\end{align*} \begin{array}{r}
\text { on } \partial S_{+}^{N-1}
\end{array}
$$

If we write

$$
U_{\varepsilon}(x)=\varepsilon^{-2 /(q-1)} U_{1, \varepsilon}(x / \varepsilon)=\varepsilon^{-2 /(q-1)} U_{1, \varepsilon}(y), \quad y=x / \varepsilon
$$

the function $U_{1, \varepsilon}$ satisfies

$$
\Delta U_{1, \varepsilon}=U_{1, \varepsilon}^{q} \quad \text { in } \quad \mathfrak{D}_{\varepsilon}^{m}=\varepsilon^{-1} \Omega_{\varepsilon}^{m}
$$

When $\varepsilon \rightarrow 0, D_{\varepsilon}^{m}$ converges to $\mathfrak{D}_{0}=(0,1) \times \mathbb{R}^{N-1}$ in the sense of sets. Thus, for any $0<\theta_{1}<\theta_{2}<1$, there exists $\varepsilon_{0}$ such that if $0<\varepsilon \leq \varepsilon_{0}$, the following inclusion holds

$$
\mathfrak{G}_{\theta}^{m}=\left\{y=\left(y_{1}, y^{\prime}\right): \theta_{1}<y_{1}<\theta_{2}, 1<\left|y^{\prime}\right| \leq m / 2 \varepsilon\right\} \subset \mathfrak{D}_{\varepsilon}^{m}
$$

Because

$$
U_{1, \varepsilon}(y) \leq C|y|^{-2 /(q-1)}
$$

$U_{1, \varepsilon}$ converges uniformly, as $\varepsilon \rightarrow 0$, on any compact subset of $\mathfrak{D}_{0} \backslash\{O\}$ to the unique solution $U_{1}$ of

$$
\Delta U_{1}=U_{1}^{q} \quad \text { in } \quad \mathfrak{D}_{0}
$$

which vanishes on $\partial \mathfrak{D}_{0} \backslash\{O\}$ and satisfies

$$
\begin{equation*}
\lim _{\substack{y \rightarrow 0 \\ y \in \mathfrak{D}_{0}}}|y|^{2 /(q-1)} U_{1}(y)=\omega(y /|y|) \tag{5.10}
\end{equation*}
$$

Therefore there exists some $\eta \in(0,1)$ such that

$$
\begin{equation*}
U_{1, \varepsilon}\left(y_{1}, y^{\prime}\right) \geq \eta \sin \left(\frac{\pi\left(y_{1}-\theta_{1}\right)}{\theta_{2}-\theta_{1}}\right), \quad \text { for } \quad y_{1} \in\left[\theta_{1}, \theta_{2}\right] \text { and } \quad\left|y^{\prime}\right|=1 \tag{5.11}
\end{equation*}
$$

Notice that the function $y_{1} \mapsto \psi_{\theta}\left(y_{1}\right)=\sin \left(\pi\left(y_{1}-\theta_{1}\right) /\left(\theta_{2}-\theta_{1}\right)\right)$ vanishes for $y_{1}=\theta_{1}$ and for $y_{1}=\theta_{2}$.

We first suppose $N=2$. Then for any $\beta>0$ the function

$$
\begin{equation*}
y^{\prime} \mapsto \varphi_{\beta}\left(y^{\prime}\right)=\frac{\sinh \left(\beta\left(\frac{m}{2 \varepsilon}-\left|y^{\prime}\right|\right)\right)}{\sinh \left(\beta\left(\frac{m}{2 \varepsilon}-1\right)\right)}, \tag{5.12}
\end{equation*}
$$

is nonnegative takes the value 1 for $\left|y^{\prime}\right|=1$, and vanishes for $\left|y^{\prime}\right|=m /(2 \varepsilon)$. If we set

$$
\zeta_{\theta, \beta}\left(y_{1}, y^{\prime}\right)=\eta \psi_{\theta}\left(y_{1}\right) \varphi_{\beta}\left(y^{\prime}\right)
$$

there holds

$$
\begin{equation*}
\Delta \zeta_{\theta, \beta}=\left(\beta^{2}-\frac{\pi^{2}}{\left(\theta_{2}-\theta_{1}\right)^{2}}\right) \zeta_{\theta, \beta} \tag{5.13}
\end{equation*}
$$

Since $\zeta_{\theta, \beta} \leq \eta$, it follows

$$
\begin{equation*}
\Delta \zeta_{\theta, \beta} \geq\left(\beta^{2}-\frac{\pi^{2}}{\left(\theta_{2}-\theta_{1}\right)^{2}}\right) \eta^{1-q} \zeta_{\theta, \beta}^{q} \quad \text { in } \quad \mathfrak{G}_{\theta}^{m} \tag{5.14}
\end{equation*}
$$

furthermore

$$
\begin{aligned}
& \zeta_{\theta, \beta}\left(y_{1}, y^{\prime}\right)=0 \quad \text { for } \quad y_{1}=\theta_{i}, i=1,2 \\
& \zeta_{\theta, \beta}\left(y_{1}, y^{\prime}\right)=0 \quad \text { for } \quad\left|y^{\prime}\right|=m / 2 \varepsilon \\
& \zeta_{\theta, \beta}\left(y_{1}, y^{\prime}\right)=0 \quad \text { for } \quad\left|y^{\prime}\right|=1
\end{aligned}
$$

We chose $\beta$ such that $\frac{\pi^{2}}{\left(\theta_{2}-\theta_{1}\right)^{2}} \eta^{1-q}=1$. By (5.11) and the maximum principle one obtains

$$
\begin{equation*}
U_{1, \varepsilon}\left(y_{1}, y^{\prime}\right) \geq \zeta_{\theta, \beta}\left(y_{1}, y^{\prime}\right) \quad \text { in } \quad \mathfrak{G}_{\theta}^{m} \tag{5.15}
\end{equation*}
$$

Therefore

$$
\begin{aligned}
u\left(x_{1}, x^{\prime}\right) & \geq \ell(\varepsilon) \varepsilon^{-2 /(q-1)} U_{1, \varepsilon}\left(x_{1} / \varepsilon, x^{\prime} / \varepsilon\right) \\
& \geq \ell(\varepsilon) \varepsilon^{-2 /(q-1)} \zeta_{\theta, \beta}\left(x_{1} / \varepsilon, x^{\prime} / \varepsilon\right)
\end{aligned}
$$

for $\theta_{1} \varepsilon \leq x_{1} \leq \theta_{1} \varepsilon$ and $\varepsilon \leq\left|x^{\prime}\right| \leq m / 2$. Take $x_{1}=\varepsilon\left(\theta_{1}+\theta_{2}\right) / 2=\theta \varepsilon$, then

$$
u\left(\theta \varepsilon, x^{\prime}\right) \geq \eta \ell(\varepsilon) \varepsilon^{-2 /(q-1)} \frac{\sinh \left(\beta\left(\frac{m-2\left|x^{\prime}\right|}{2 \varepsilon}\right)\right)}{\sinh \left(\beta\left(\frac{m}{2 \varepsilon}-1\right)\right)}
$$

If $|x|^{\prime} \leq m / 4$,

$$
\frac{\sinh \left(\beta\left(\frac{m-2\left|x^{\prime}\right|}{2 \varepsilon}\right)\right)}{\sinh \left(\beta\left(\frac{m}{2 \varepsilon}-1\right)\right)}=e^{\beta\left(1-|x|^{\prime} / \varepsilon\right)}(1+\circ(1)) \quad \text { as } \quad \varepsilon \rightarrow 0
$$

Thus

$$
u\left(\theta \varepsilon, x^{\prime}\right) \geq \eta \ell(\varepsilon) \varepsilon^{-2 /(q-1)} e^{\beta} e^{-\beta\left|x^{\prime}\right| / \varepsilon}(1+\circ(1))
$$

and

$$
\liminf _{\varepsilon \rightarrow 0} u\left(\theta \varepsilon, x^{\prime}\right) \geq \eta e^{\beta} \liminf _{\varepsilon \rightarrow 0} \ell(\varepsilon) \varepsilon^{-2 /(q-1)} e^{-\beta\left|x^{\prime}\right| / \varepsilon}
$$

Since $\ell(\varepsilon)=e^{\kappa /((q-1) \varepsilon}$,

$$
\ell(\varepsilon) \varepsilon^{-2 /(q-1)} e^{\beta\left(1-|x|^{\prime} / \varepsilon\right)} \geq \varepsilon^{-2 /(q-1)} \exp \left[\varepsilon^{-1}\left(\kappa /(q-1)-\beta|x|^{\prime}\right)\right]
$$

If we fix $|x|^{\prime}<\beta \kappa /(q-1)$, then

$$
\liminf _{\varepsilon \rightarrow 0} u\left(\theta \varepsilon, x^{\prime}\right)=\lim _{\rho_{\partial \Omega}(x) \rightarrow 0} u(x)=\infty
$$

and this limit is uniform on any compact subset of $\left\{x^{\prime}:|x|^{\prime}<\beta \kappa /(q-1)\right\}$, which is equivalent to any compact of $\{x:|\sigma(x)|<\beta \kappa /(q-1)\}$. Put $\tau=$ $\beta \kappa /(2(q-1))$. Because this blow-up holds in a fixed neighborhood $\partial \Omega \cap \bar{B}_{\tau}(O)$ of $O$, we can replace $O$ by any point $P$ in $\partial \Omega \cap \bar{B}_{\tau}(O)$ and conclude that

$$
\lim _{\rho_{\partial \Omega}(x) \rightarrow 0} u(x)=\infty,
$$

uniformly if $|\sigma(x)-P| \leq \tau$. Iterating this process infers that

$$
\lim _{\rho_{\partial \Omega}(x) \rightarrow 0} u(x)=\infty .
$$

Next we assume $N \geq 3$. Let $\beta>0$ to be fixed and

$$
\Gamma_{\varepsilon}=\left\{y^{\prime} \in \mathbb{R}^{N-1}: 1<\left|y^{\prime}\right|<m / 2 \varepsilon\right\}
$$

and let $B_{\beta}\left(y^{\prime}\right)$ be the solution of

$$
\begin{align*}
& \Delta_{y^{\prime}} B_{\beta}=\beta^{2} B_{\beta} \quad \text { in } \quad \Gamma_{\varepsilon} \\
& B_{\beta}\left(y^{\prime}\right)=1 \quad \text { if } \quad\left|y^{\prime}\right|=1  \tag{5.16}\\
& B_{\beta}\left(y^{\prime}\right)=0 \quad \text { if } \quad\left|y^{\prime}\right|=m / 2 \varepsilon
\end{align*}
$$

The function $\zeta_{\theta, \beta}\left(y_{1}, y^{\prime}\right)=\eta \psi_{\theta}\left(y_{1}\right) B_{\beta}\left(y^{\prime}\right)$ satisfies also (5.13) in $\mathfrak{G}_{\theta}^{m}$. Therefore, if we chose $\beta$ as in the case $N=2$, (5.15) is still valid. Since $B_{\beta}$ is a Bessel function, its behaviour at infinity is classical and there holds, for $|x|^{\prime} \leq m / 4$,

$$
B_{\beta}(y)=C_{\beta}\left(\frac{|x|^{\prime}}{\varepsilon}\right)^{1-N / 2} e^{-\beta|x|^{\prime} / \varepsilon}(1+\circ(1)), \quad \text { as } \quad \varepsilon \rightarrow 0
$$

We conclude as in the case $N=2$.
Step 2. The general case. If $q \geq(N+1) /(N-1)$ let $\alpha>0$ such that

$$
q<(N+1+\alpha) /(N-1)
$$

We write

$$
h\left(\rho_{\partial \Omega}(x)\right)=\rho_{\partial \Omega}^{\alpha}(x) \tilde{h}\left(\rho_{\partial \Omega}(x)\right)
$$

with

$$
\tilde{h}\left(\rho_{\partial \Omega}(x)\right)=\rho_{\partial \Omega}^{-\alpha}(x) h\left(\rho_{\partial \Omega}(x)\right)
$$

We can assume that $r \mapsto \tilde{h}(r)$ is nondecreasing near $r=0$ and we extend it by continuity at $r=0$ by putting $\tilde{h}(0)=0$. Thus there holds

$$
\Delta u \leq \tilde{h}(\varepsilon) \rho_{\partial \Omega}^{\alpha}(x) u^{q} \quad \text { in } \Omega_{\varepsilon}
$$

The equation

$$
-\Delta U+\rho_{\partial \Omega}^{\alpha}(x) U^{q}=0
$$

admits weak and strong isolated singularities on the boundary and any positive solution with a strong singularity at $x=O$ satisfies

$$
\begin{equation*}
\lim _{\substack{x \rightarrow O \\ x \in \Omega_{\varepsilon}}}|x|^{(2+\alpha) /(q-1)} U_{\varepsilon}(x)=\omega_{\alpha}(x /|x|) \tag{5.17}
\end{equation*}
$$

where $\omega_{\alpha}$ is the unique positive solution of

$$
\begin{align*}
-\Delta_{\sigma} \omega_{\alpha}-\left(\frac{2+\alpha}{q-1}\right)\left(\frac{2 q+\alpha}{q-1}-N\right) \omega_{\alpha}+\omega_{\alpha}^{q} & =0 \quad \text { in } S_{+}^{N-1}  \tag{5.18}\\
\omega & =0 \quad \text { on } \partial S_{+}^{N-1}
\end{align*}
$$

we proceed as in Step 1, with some minor changes of coefficients.
STEP 3. End of the proof. The minimal solution $u_{m}$ of (5.6) is constructed by considering the increasing sequence $u_{k}$ of solutions of

$$
\begin{align*}
-\Delta u_{k}+\exp \left(-\kappa / \rho_{\partial \omega}(x)\right) u_{k}^{q}=0 & \text { in } \Omega,  \tag{5.19}\\
u_{k}(x)=k & \text { on } \partial \Omega .
\end{align*}
$$

When $k \rightarrow \infty, u_{k} \rightarrow u_{m}$, thus $u_{\infty, a} \geq u_{m}$. On the other hand, $u_{\lambda \delta_{a}}$ is constructed by approximating the Dirac mass on the boundary by bounded functions $g_{\lambda}$. Thus the corresponding solutions $u_{g_{\lambda}}$ of (0.4) are all dominated by $u_{m}$. Therefore

$$
u_{\lambda \delta_{a}} \leq u_{m} \Longrightarrow u_{\infty, a} \leq u_{m}
$$

REmARK 5.2. If the domain $\Omega$ is starshaped with respect to some point, say $O$, the Iscoe uniqueness method (see [17]) of scaling applies straightforwardly to prove the uniqueness of the solution of (5.6). We recall this method. Let $\ell>0$ and $u_{\ell}(x)=\ell^{2 /(q-1)} u(\ell x)$, then $u_{\ell}$ satisfies

$$
\begin{aligned}
-\Delta u_{\ell}+e^{-\kappa / \rho_{\partial \Omega}(\ell x)} u_{\ell}^{q} & =0 \\
& \text { in } \quad \Omega_{\ell}=\frac{1}{\ell} \Omega \\
u_{\ell} & =\infty
\end{aligned} \quad \begin{aligned}
& \text { on } \quad \partial \Omega_{\ell}
\end{aligned}
$$

But $\Omega_{\ell} \subset \Omega$ and $e^{-\kappa / \rho_{\partial \Omega}(\ell x)} \leq e^{-\kappa / \rho_{\partial \Omega}(x)}$ if $\ell>1$. Therefore $u_{\ell}$ satisfies

$$
-\Delta u_{\ell}+e^{-\kappa / \rho_{\partial \Omega}(x)} u_{\ell}^{q} \geq 0, \quad \text { in } \quad \Omega_{\ell}
$$

If $\hat{u}$ is another of (5.6) in $\Omega$, then $u_{\ell} \geq \hat{u}$. Letting $\ell \rightarrow 1$ infers $u \geq \hat{u}$. In the same way $\hat{u} \geq u$. In a much more elaborated manner, if $\Omega$ is locally a continuous graph, the method of local translations developped by the authors in [21] can be adapted and once again uniqueness of the solution of (5.6) holds.

Combining Proposition 5.1, the previous remark and Theorem 4.1 (in the case $1<q<(N+1) /(N-1)$ ) we derive,

Corollary 5.2. Let $q>1$ and $u$ be a nonnegative solution of

$$
\begin{equation*}
-\Delta u+\exp \left(-1 / \rho_{\partial \omega}(x)\right) u^{q}=0 \quad \text { in } \Omega \tag{5.20}
\end{equation*}
$$

Then either
(i) $\mathcal{S}(u)=\partial \Omega, \operatorname{Tr}_{\partial \Omega}^{e}(u)$ the Borel measure indentically equal to $\infty$, and

$$
u_{\infty, a}=u_{m}
$$

or
(ii) $\mathcal{R}(u)=\partial \Omega$ and $\operatorname{Tr}_{\partial \Omega}^{e}(u)=v$ is a bounded Borel measure. Moreover, if $1<q<(N+1) /(N-1)$, $v$ is a Radon measure and $u=u_{v}$.

## REFERENCES

[1] P. Baras - M. Pierre, Singularités éliminables pour des équations semi-linéaires, Ann. Inst. Fourier (Grenoble) 34 (1984), 185-206.
[2] P. Baras - J. Goldstein, The heat equation with a singular potential, Trans. Amer. Math. Soc. 284 (1984), 121-139.
[3] Ph. Benilan - H. Brezis, Nonlinear problems related to the Thomas-Fermi equation, unpublished paper, see [6].
[4] M. F. Bidaut-VÉron - L. Vivier, An elliptic semilinear equation with source term involving boundary measures: the subcritical case, Rev. Mat. Iberoamericana 16 (2000), 477-513.
[5] H. Brezis, Une équation semi-linéaire avec conditions aux limites dans $L^{1}$, unpublished paper. See also [32]-Chap. 4.
[6] H. Brezis, Some variational problems of the Thomas-Fermi type, in "Variational Inequalities", R. W. Cottle, F. Giannessi and J.-L. Lions (eds.), Wiley, Chichester (1980), 53-73.
[7] X. CABRE, Extremal solutions and instantaneous complete blow-up for elliptic and parabolic problems, preprint.
[8] R. Dautray - J. L. Lions, "Analyse Mathématique et Calcul Numérique", Masson, Paris, 1987.
[9] J. Doob, "Classical Potential Theory and its Probabilistic Counterpart", Springer-Verlag, Berlin-New York, 1984.
[10] E. B. Dynkin - S. E. Kuznetsov, Trace on the boundary for solutions of nonlinear differential equations, Trans. Amer. Math. Soc. 350 (1998), 4499-4519.
[11] E. B. Dynkin - S. E. Kuznetsov, Solutions of nonlinear differential equtions on a Riemannian manifold and their trace on the Martin boundary, Trans. Amer. Math. Soc. 350 (1998), 4521-4552.
[12] E. B. Dynkin - S. E. Kuznetsov, Fine topology and fine trace on the boundary associated with a class of quasilinear differential equations, Comm. Pure Appl. Math. 51 (1998), 897936.
[13] J. FabBRi - J. R. Licois, Behavior at boundary of solutions of a weakly superlinear elliptic equation, Adv. Nonlinear Stud. 2 (2002), 147-176.
[14] D. Gilbarg - N. S. Trudinger, "Partial Differential Equations of Second Order", 2nd Ed. Springer-Verlag, Berlin-New York, 1983.
[15] A. Gmira - L. Véron, Boundary singularities of solutions of some nonlinear elliptic equations, Duke Math. J. 64 (1991), 271-324.
[16] M. Grillot - L. Véron, Boundary trace of solutions of the Prescribed Gaussian curvature equation, Proc. Roy. Soc. Edinburgh 130 A (2000), 1-34.
[17] I. Iscoe, On the support of measure-valued critical branching Brownian motion, Ann. Prob. 16 (1988), 200-221.
[18] J. B. Keller, On solutions of $\Delta u=f(u)$, Comm. Pure Appl. Math. 10 (1957), 503-510.
[19] J. F. Le Gall, Les solutions positives de $\Delta u=u^{2}$ dans le disque unité, C.R. Acad. Sci. Paris 317 Ser. I (1993), 873-878.
[20] J. F. Le Gall, The brownian snake and solutions of $\Delta u=u^{2}$ in a domain, Probab. Theory Related Fields 102 (1995), 393-432.
[21] M. Marcus - L. Véron, Uniqueness and asymptotic behaviour of solutions with boundary blow-up for a class of nonlinear elliptic equations, Ann. Inst. H. Poincaré 14 (1997), 237274.
[22] M. Marcus - L. Véron, Traces au bord des solutions positives d'équations elliptiques non-linéaires, C.R. Acad. Sci. Paris 321 Ser. I (1995), 179-184.
[23] M. Marcus - L. Véron, Traces au bord des solutions positives d'équations elliptiques et paraboliques non-linéaires: résultats d'existence et d'unicité, C.R. Acad. Sci. Paris 323 Ser. I (1996), 603-608.
[24] M. Marcus - L. VÉron, The boundary trace of positive solutions of semilinear elliptic equations: the subcritical case, Arch. Ration. Mech. Anal. 144 (1998), 201-231.
[25] M. Marcus - L. Véron, The boundary trace of positive solutions of semilinear elliptic equations: the supercritical case, J. Math. Pures Appl. 77, 481-524 (1998).
[26] M. Marcus - L. Véron, Removable singularities and boundary traces, J. Math. Pures Appl. 80 (2001), 879-900.
[27] M. Marcus - L. VÉron, The boundary trace and generalyzed boundary value problem for semilinear elliptic equations with coercive absorption, Comm. Pure Appl. Math. 56 (2003), 0689-0731.
[28] M. MARCUS - L. VÉron, Initial trace of positive solutions to semilinear parabolic inequalities, Adv. Nonlinear Studies 2 (2002), 395-436.
[29] A. Ratto - M. Rigoli - L. Véron, Scalar curvature and conformal deformation of hyperbolic space, J. Funct. Anal. 121 (1994), 15-77.
[30] Y. Richard - L. Véron, Isotropic singularities of nonlinear elliptic inequalities, Ann. Inst. H. Poincaré 6 (1989), 37-72.
[31] J. L. VAZQUEZ, An a priori interior estimate for the solution of a nonlinear problem representing weak diffusion, Nonlinear Anal. 5 (1981), 119-135.
[32] L. Véron, "Singularities of Solutions of Second Order Quasilinear Equations", Pitman Research Notes in Math. 353, Addison Wesley Longman Inc., 1996.

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