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# The Stefan Problem with Kinetic Condition at the Free Boundary

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## 1. - Formulation of the problem

Consider the free boundary problem: find a function  $u(x, y, t)$  and a curve

$$(1.1) \quad \Gamma : y = g(x, t) \quad (g(x, t) > 0)$$

such that  $u$  satisfies the differential equation

$$(1.2) \quad \Delta u \equiv u_{xx} + u_{yy} = 0 \text{ in } \Omega_t \equiv \{(x, y); -\infty < x < \infty, 0 < y < g(x, t)\}$$

and the boundary conditions

$$(1.3) \quad u(x, 0, t) = b(x, t) \quad (-\infty < x < \infty),$$

$$(1.4) \quad \frac{\partial u}{\partial n} + u = 0 \text{ on } \Gamma,$$

$$(1.5) \quad u = V_n \text{ on } \Gamma,$$

and  $g$  satisfies the initial condition

$$(1.6) \quad g(x, 0) = g_0(x), \quad g_0(x) > 0 \quad (-\infty < x < \infty).$$

Here  $n$  is the outward normal to  $\Gamma$ ,

$$n = \frac{(-g_x, 1)}{\sqrt{1 + g_x^2}},$$

and  $V_n$  is the velocity of the free boundary,

$$V_n = \frac{gt}{\sqrt{1 + g_x^2}}.$$

In view of (1.5), equation (1.4) can also be written in the form

$$(1.4') \quad -\frac{\partial u}{\partial n} = V_n \text{ on } \Gamma.$$

The relations (1.4') and  $u = 0$  on  $\Gamma$  constitute the standard free boundary conditions for the Stefan problem. The modified condition  $u = V_n$  on  $\Gamma$  represents kinetic heating. For one space dimension with independent variable  $y$  and with  $\Delta u$  replaced by  $u_t = u_{yy}$ , this problem was studied by Dewynne, Howison, Ockendon and Xie [1] and by Xie [5].

Our interest in the present problem arises from the modeling of titanium silicide film growth. The problem actually involves three free boundaries; see [3; Chap. 8]. Here however we restrict ourselves to a subproblem whereby the lower part of  $\Omega_t$  is a fixed curve, which for simplicity is taken to be the  $x$ -axis. The function  $u$  represents the concentration of titanium silicide. Relation (1.4') is the conservation of mass, whereas (1.4) models the rate of conversion of titanium to titanium silicide.

Consider first the special one-dimension problem where  $g_0(x)$  and  $b(x, t)$  are independent of  $x$ . Given

$$(1.7) \quad g_0 = \bar{s}_0, \quad b = b(t), \quad (\bar{s} > 0, \quad b(t) > 0),$$

one easily finds a unique solution  $u_0(y, t)$ , with free boundary  $y = s_0(t)$ :

$$(1.8) \quad u_0(y, t) = -\frac{b(t)y}{1 + s_0(t)} + b(t), \quad 0 < y < s_0(t),$$

$$(1.9) \quad s_0'(t) = \frac{b(t)}{1 + s_0(t)}, \quad s_0(0) = \bar{s}_0.$$

From the last equation we get

$$(1.10) \quad \frac{1}{2} s_0^2(t) + s_0(t) = \frac{1}{2} \bar{s}_0^2 + \bar{s}_0 + \int_0^t b(\tau) d\tau.$$

Note that, if  $b(t) \sim C$  as  $t \rightarrow \infty$  ( $C > 0$ ), then  $s_0(t) \sim \sqrt{2Ct}$  as  $t \rightarrow \infty$ .

In this paper we shall prove the existence of a local classical solution of (1.1)–(1.6). We shall also prove that a global solution exists if the data  $g, b$  are “close” to the data (1.7):

$$(1.11) \quad g_0(x) = \bar{s}_0 + \varepsilon g_1(x),$$

$$(1.12) \quad b(x, t) = b(t) + \varepsilon b_1(x, t)$$

where  $\varepsilon$  is a sufficiently small positive constant. The global solution will have the form

$$(1.13) \quad u(x, t) = -\frac{b(t)y}{1 + s_0(t)} + b(t) + \varepsilon u_1(x, y, t),$$

$$(1.14) \quad g(x, t) = s_0(t) + \varepsilon g_1(x, t)$$

with suitable functions  $u_1, g_1$  (which depend on  $\varepsilon$ ).

In the modeling of titanium silicide film growth the free boundary is actually nearly flat, so that the assumptions (1.11), (1.12) possibly include practical cases.

In Section 2 we formulate the problem for  $u_1, g_1$ . In Sections 3–6 we derive a priori estimates on  $g_1$  and its  $C_x^{2,\alpha}$  norm. In Section 7 we establish local existence (for general data). Then, by combining local existence with the a priori estimates on  $g_1$ , global existence for the data (1.11), (1.12) immediately follows. Finally, uniqueness is proved in Section 8.

## 2. - The reduced problem

From (1.14) we get

$$n = \frac{(-\varepsilon g_{1x}, 1)}{\sqrt{1 + \varepsilon^2 g_{1x}^2}}.$$

Condition (1.4) written in terms of  $u_1$  is

$$(2.1) \quad \frac{\partial u_1}{\partial n} + u_1 - \frac{b(t)}{1 + s_0(t)} g_1(x, t) + \frac{\varepsilon b(t)}{1 + s_0(t)} F(\varepsilon g_{1x}) g_{1x}^2 = 0$$

where

$$F(\lambda) = \frac{1}{\sqrt{1 + \lambda^2} (1 + \sqrt{1 + \lambda^2})}.$$

Condition (1.5) can be reduced to

$$V_n = \frac{s'_0(t) + \varepsilon g_{1t}}{\sqrt{1 + \varepsilon^2 g_{1x}^2}} = \frac{-b(t)(s_0(t) + \varepsilon g_1)}{1 + s_0(t)} + b(t) + \varepsilon u_1$$

or, upon recalling (1.9),

$$(2.2) \quad g_{1t} + \frac{b(t)}{1 + s_0(t)} g_1 \sqrt{1 + \varepsilon^2 g_{1x}^2} - \frac{\varepsilon b(t)}{1 + s_0(t)} G(\varepsilon g_{1x}) g_{1x}^2 - \sqrt{1 + \varepsilon^2 g_{1x}^2} u_1(x, s_0(t) + \varepsilon g_1, t) = 0$$

where

$$G(\lambda) = \frac{1}{1 + \sqrt{1 + \lambda^2}}.$$

We also have

$$(2.3) \quad \Delta u_1 = 0 \text{ if } x \in \mathbb{R}, \quad 0 < y < s_0(t) + \varepsilon g_1(x),$$

and

$$(2.4) \quad u_1(x, 0, t) = b_1(x, t),$$

$$(2.5) \quad g_1(x_0) = g_1(x).$$

We shall henceforth assume:

$$(2.6) \quad b(t) \text{ is continuous, } 0 < \bar{b} \leq b(t) \leq \bar{\bar{b}} < \infty,$$

$$(2.7) \quad b_1(x, t) \text{ is continuous,}$$

$$(2.8) \quad |b_1(x, t)| \leq B(t), \quad B(t) \leq \frac{1}{1+t}, \quad \int_0^\infty B(t) dt < \infty.$$

We also assume that  $g_1(x)$  satisfies:

$$(2.9) \quad \|g_1\|_{L^\infty}, \|g_{1x}\|_{L^\infty}, \|g_{1xx}\|_{L^\infty}, [g_{1xx}]_{C_x^\alpha} \leq 1$$

where the norms are taken in  $\mathbb{R}$ .

We seek a solution  $u, g$  of the form (1.13), (1.14) such that  $g_1(x, t)$  satisfies:

$$(2.10) \quad \|g_1(\cdot, t)\|_{L^\infty} + \|g_{1x}(\cdot, t)\|_{L^\infty} + \|g_{1xx}(\cdot, t)\|_{L^\infty} \leq \frac{K}{\sqrt{1+t}},$$

$$[g_{1xx}(\cdot, t)]_{C_x^\alpha} \leq K \quad (K > 1).$$

In Sections 3–6 we assume that a classical solution exists for all  $0 \leq t \leq T$ , for some  $T > 0$  (all the derivatives  $g_{1t}, g_{1x}, g_{1xx}$  are continuous and  $g_{1xx}$  is Hölder continuous in  $x$ ), and that (2.10) holds for some  $K$  and all  $0 \leq t \leq T$ . We shall then prove that for all  $0 \leq t \leq T$

$$(2.11) \quad \|g_1(\cdot, t)\|_{L^\infty} \leq \frac{C}{\sqrt{1+t}} \quad (\text{in Section 3}),$$

$$(2.12) \quad \|u_1(\cdot, \cdot, t)\|_{C_{xy}^{2,\alpha}(D_t)} \leq \frac{C}{1+t} + \frac{C}{\sqrt{1+t}} \|g_1(\cdot, t)\|_{C_x^{1,\alpha}} \quad (\text{in Section 4}),$$

where  $D_t$  is defined in (4.1).

$$(2.13) \quad \|g_{1xx}(\cdot, t)\|_{L^\infty} \leq \frac{C}{\sqrt{1+t}} \quad (\text{in Section 5}),$$

and

$$(2.14) \quad [g_{1xx}(\cdot, t)]_{C_x^\alpha} \leq C \quad (\text{in Section 6}),$$

where all norms are taken in  $\mathbb{R}$ . The crucial facts here are that  $\varepsilon$  is assumed to be sufficiently small depending on  $K$  but not on  $T$ , and that  $C$  is a constant independent of  $K, T, \varepsilon$ .

In Section 7 we prove the existence of a classical solution of (1.1)–(1.6) for general data  $b(x, t), g_0(x)$ , for  $0 \leq t \leq \bar{t}$  where  $\bar{t}$  is sufficiently small. By combining this result with the estimates derived in Sections 3–6, we easily construct a global classical solution of (1.1)–(1.6) for the data (1.11), (1.12) provided  $\varepsilon$  is sufficiently small. In Section 8 we prove uniqueness of the solution.

### 3. - Proof of (2.11)

In this section we prove:

LEMMA 3.1. *If  $u, g$  is a solution satisfying (2.10), for  $0 \leq t \leq T$ , then (2.11) holds for  $0 \leq t \leq T$  provided  $0 < \varepsilon < 1/K$ ;  $C$  is a constant independent of  $K, T, \varepsilon$ .*

We shall need the following fact:

$$(3.1) \quad \begin{aligned} &\text{if } h(x) \geq 0 \ (x \in \mathbb{R}) \text{ and } \|h''\|_{L^\infty} \leq A \quad (0 < A < \infty), \text{ then} \\ &h(x) \geq \frac{(h'(x))^2}{2A} \text{ for all } x \in \mathbb{R}. \end{aligned}$$

To prove it we use Taylor's expansion and the assumption  $h(y) \geq 0$  to deduce that

$$h(x) \geq (x - y)h'(x) - \frac{1}{2}(x - y)^2 \|h''\|_{L^\infty} \quad \forall y \in \mathbb{R}.$$

Choosing  $y$  such that  $(x - y)h'(x) \geq 0$  and

$$(x - y)h'(x) = (x - y)^2 A,$$

(3.1) follows.

Denote by  $w_1(y, t)$ ,  $s_1(t)$  the solution corresponding to

$$\begin{aligned}\bar{s}_1 &= \bar{s}_0 - 2\varepsilon, \\ b_1(t) &= b(t) - \varepsilon B(t),\end{aligned}$$

i.e. (see Section 1),

$$(3.2) \quad \begin{aligned}w_1(y, t) &= -\frac{b_1(t)y}{1+s_1(t)} + b_1(t), \\ \frac{1}{2}s_1^2(t) + s_1(t) &= \frac{1}{2}\bar{s}_1^2 + \bar{s}_1 + \int_0^t b_1(\tau)d\tau.\end{aligned}$$

LEMMA 3.2. *If  $\varepsilon < 1/K$  then*

$$(3.3) \quad \begin{aligned}g(x, t) &> s_1(t), \\ u(x, y, t) &\geq w_1(y, t)\end{aligned}$$

for  $x \in \mathbb{R}$ ,  $0 \leq y \leq g(x, t)$ ,  $0 \leq t \leq T$ .

PROOF. We first show that

$$(3.4) \quad \begin{aligned}\text{if } g(x, t) &> s_1(t) \text{ for } x \in \mathbb{R}, 0 \leq t \leq T, \\ \text{then } u(x, y, t) &\geq w_1(y, t) \text{ for } x \in \mathbb{R}, 0 \leq y \leq g(x, t), 0 \leq t \leq T.\end{aligned}$$

Indeed,

$$u(x, 0, t) \geq b_1(t) = w_1(0, t).$$

Also, on  $\Gamma : y = g(x, t) = s_0(t) + \varepsilon g_1(x, t)$ ,

$$(3.5) \quad \begin{aligned}\frac{\partial w_1}{\partial n} + w_1 &= -\frac{b_1(t)}{1+s_1(t)} n_y \frac{b_1(t)}{1+s_1(t)} g(x, t) + b_1(t) \\ &= -\frac{b_1(t)}{1+s_1(t)} [n_y + g(x, t) - 1 - s_1(t)].\end{aligned}$$

Since  $g(x, t) \geq s_1(t)$ , we have by (3.1),

$$g(x, t) - s_1(t) = \frac{g_x^2}{2K\varepsilon} = \frac{1}{2K\varepsilon} \varepsilon^2 g_{1x}^2.$$

Also

$$n_y - 1 = \frac{1}{\sqrt{1 + \varepsilon^2 g_{1x}^2}} - 1 \geq -\frac{1}{2} \varepsilon^2 g_{1x}^2.$$

It follows that the right-hand side of (3.5) is  $\leq 0$  if  $\varepsilon < 1/K$ . Applying the

maximum principle to  $u - w_1$  in the region  $\{0 \leq y \leq g(x, t)\}$ , we conclude that  $u \geq w_1$ .

Next we establish that

$$(3.6) \quad \begin{aligned} &\text{if } u(x, g(x, t), t) \geq w_1(g(x, t), t) \text{ for } x \in \mathbb{R}, 0 \leq t \leq T, \\ &\text{then } g(x, t) \geq s_1(t) + \delta_T \text{ for } x \in \mathbb{R}, 0 \leq t \leq T, \end{aligned}$$

where  $\delta_T$  is some positive constant.

Indeed, by (1.5),

$$\begin{aligned} g_t &= \sqrt{1 + \varepsilon^2 g_{1x}^2} u(x, g(x, t), t) \\ &\geq w_1(g(x, t), t) = -\frac{b_1(t)}{1 + s_1(t)} g + b_1(t) \end{aligned}$$

and therefore

$$(3.7) \quad g(x, t) \geq \tilde{s}(t),$$

where  $\tilde{s}(t)$  is defined by

$$\begin{aligned} \tilde{s}'(t) &= -\frac{b_1(t)}{1 + s_1(t)} \tilde{s}(t) + b_1(t), \\ \tilde{s}(0) &= \bar{s}_0 - \varepsilon \end{aligned}$$

(recall that  $|g_1(x, 0)| \leq 1$ ). Since  $s_1(t)$  satisfies the same differential equation as  $\tilde{s}(t)$  with  $s_1(0) = \bar{s}_0 - 2\varepsilon < \tilde{s}(0)$ , we have, for some  $\delta_T > 0$ ,

$$\tilde{s}(t) \geq s_1(t) + \delta_T \quad (0 \leq t \leq T).$$

This together with (3.7) complete the proof of (3.6).

Finally, by combining (3.4) with (3.6), the assertion of Lemma 3.2 follows.

We shall next compare  $u, g$  with the solution  $w_2(y, t), s_2(t)$  corresponding to

$$\begin{aligned} \bar{s}_2 &= \bar{s}_0 + 2\varepsilon, \\ b_2(t) &= b(t) + \varepsilon B(t). \end{aligned}$$

Clearly,

$$(3.8) \quad \begin{aligned} w_2(y, t) &= -\frac{b_2(t)y}{1 + s_2(t)} + b_2(t), \\ \frac{1}{2}s_2^2(t) + s_2(t) &= \frac{1}{2}\bar{s}_2^2 + \bar{s}_2 + \int_0^t b_2(\tau) d\tau. \end{aligned}$$



LEMMA 3.3. *There holds:*

$$(3.9) \quad \begin{aligned} g(x, t) &< s_2(t), \\ u(x, y, t) &\leq w_2(y, t), \end{aligned}$$

for  $x \in \mathbb{R}$ ,  $0 \leq y \leq g(x, t)$ ,  $0 \leq t \leq T$ .

PROOF. We first prove that

$$(3.10) \quad \begin{aligned} &\text{if } g(x, t) \leq s_2(t) \text{ for } x \in \mathbb{R}, 0 \leq t \leq T, \\ &\text{then } u(x, y, t) < w_1(y, t) \text{ for } x \in \mathbb{R}, 0 \leq y \leq g(x, t), \quad 0 \leq t \leq T. \end{aligned}$$

Indeed, the proof is similar to the proof of (3.4); it does not require that  $\varepsilon < 1/K$ , since

$$n_y + g(x, t) - 1 - s_2(t) \leq n_y - 1 \leq 0.$$

Next we prove:

$$(3.11) \quad \begin{aligned} &\text{if } u(x, g(x, t), t) \leq w_2(g(x, t), t) \text{ for } x \in \mathbb{R}, 0 \leq t \leq T, \\ &\text{then } g(x, t) \leq s_2(t) - \delta_T \text{ for } x \in \mathbb{R}, 0 \leq t \leq T, \end{aligned}$$

where  $\delta_T$  is some positive constant. To prove it we note that

$$g_t = \sqrt{1 + g_x^2} u(x, g, t) \leq \sqrt{1 + g_x^2} w_2(g, t) \equiv H(t, g, g_x).$$

Let  $\tilde{s}(t)$  be the solution to

$$\begin{aligned} \tilde{s}'(t) &= w_2(\tilde{s}(t), t), \\ \tilde{s}(0) &= \bar{s}_0 + \varepsilon. \end{aligned}$$

Then by comparison

$$(3.12) \quad \tilde{s}(t) \leq s_2(t) - \delta_T \quad (0 \leq t \leq T)$$

for some  $\delta_T > 0$ . Also

$$(3.13) \quad \tilde{s}' = H(t, \tilde{s}, 0).$$

Taking the difference of the inequality  $g_t \leq H(t, g, g_x)$  and (3.13) and noting that  $H(t, v, v_x)$  is Lipschitz in the variables  $v, v_x$ , we get, for  $z = g - \tilde{s}$ ,

$$z_t \leq az_x + bz$$

where  $a, b$  are bounded functions. Since also  $z(x, 0) \leq 0$ , it follows that  $z(x, t) \leq 0$ , i.e.,  $g(x, t) \leq \tilde{s}(t)$ . Upon recalling (3.12), the assertion (3.11) follows.

Finally, Lemma 3.3 follows by combining (3.10) and (3.11).

PROOF OF LEMMA 3.1. We shall estimate the function  $s_2(t) - s_1(t)$ . We have  $s_1(t) < s_2(t)$  and, from (3.2), (3.8) we easily find that

$$(s_2(t) - s_1(t)) \left( \frac{s_1(t) + s_2(t)}{2} + 1 \right) \leq \left( C + 2 \int_0^\infty B(s) ds \right) \varepsilon.$$

From (2.6), (2.8) and (3.2), (3.8) we also have

$$(3.14) \quad \frac{1}{C} \sqrt{1+t} \leq s_1(t) \leq s_0(t) \leq s_2(t) \leq C \sqrt{1+t}$$

and therefore

$$(3.15) \quad 0 \leq s_2(t) - s_1(t) \leq \frac{C\varepsilon}{\sqrt{1+t}}.$$

Recalling Lemmas 3.2, 3.3, (1.14) and using (3.14), (3.15), the assertion (2.11) follows.

#### 4. - Proof of (2.12)

Set

$$(4.1) \quad D_t = \left\{ x \in \mathbb{R}, g(x, t) - \frac{1}{2} \bar{s}_0 \leq y \leq g(x, t) \right\}.$$

In this section we prove:

LEMMA 4.1. *If  $u, g$  is a solution satisfying (2.10) for  $0 \leq t \leq T$ , then (2.12) holds for  $0 \leq t \leq T$  provided  $0 < \varepsilon < \varepsilon_K$ ;  $\varepsilon_K$  is a constant independent of  $T$  and  $C$  is a constant independent of  $K, T, \varepsilon$ .*

PROOF. From (1.13), Lemmas 3.2, 3.3, (3.15) and the assumption  $B(t) \leq 1/(1+t)$  we find that

$$(4.2) \quad |u_1(x, y, t)| \leq \frac{C}{1+t}.$$

Next, by (2.10),

$$(4.3) \quad \|g_x\|_{C_x^{1,\alpha}} = \|\varepsilon g_{1x}\|_{C_x^{1,\alpha}} \leq 1 \text{ if } \varepsilon K < 1,$$

so that  $\Gamma : y = g(x, t)$  is uniformly in  $C^{2,\alpha}$  (independently of  $K$  and  $t$ ).

Consider the expression

$$\tilde{F} = F(\varepsilon g_{1x}) g_{1x}^2$$

which appears in (2.1). Clearly

$$\tilde{F}_x = \varepsilon F'(\varepsilon g_{1x})g_{1x}^2g_{1xx} + F(\varepsilon g_{1x})2g_{1x}g_{1xx}.$$

Using (2.10) we easily find that

$$\|\tilde{F}\|_{L^\infty} \leq \frac{CK^2}{1+t},$$

$$\|\tilde{F}_x\|_{L^\infty} \leq \frac{CK^3}{1+t},$$

and

$$\|\tilde{F}_x\|_{C_x^\alpha} \leq \frac{C_K}{\sqrt{1+t}}.$$

It follows that

$$\|\varepsilon \tilde{F}\|_{C_x^{1,\alpha}} \leq \frac{\varepsilon C_K}{\sqrt{1+t}} < \frac{1}{\sqrt{1+t}}$$

if  $\varepsilon \leq \varepsilon_K$ , and consequently, by (2.1),

$$(4.4) \quad \left\| \left[ \frac{\partial u_1}{\partial n} + u_1 \right]_{y=g(x,t)} \right\|_{C_x^{1,\alpha}} \leq \frac{C}{1+t} + \frac{C}{\sqrt{1+t}} \|g_1\|_{C_x^{1,\alpha}}.$$

Using (4.2), (4.3) and (4.4), we can now apply the interior-boundary Schauder estimates to  $u_1$  to obtain the assertion (2.12).

### 5. - Proof of (2.13)

In this section we prove:

LEMMA 5.1. *If  $u, g$  is a solution satisfying (2.10) for  $0 \leq t \leq T$ , then (2.13) holds for  $0 \leq t \leq T$  provided  $0 < \varepsilon < \varepsilon_K$ ;  $\varepsilon_K$  is a constant independent of  $T$ , and  $C$  is a constant independent of  $K, T, \varepsilon$ .*

PROOF. We first assume that

$$(5.1) \quad g_{xxx}, g_{xxt} \text{ are continuous.}$$

We wish to differentiate (2.2) twice with respect to  $x$  so as to obtain an equation of the form

$$(g_{1xx})_t + A(g_{1xx})_x = B,$$

and then integrate along characteristics to derive (2.13). We begin with

$$\frac{\partial}{\partial x} \sqrt{1 + \varepsilon^2 g_{1x}^2} = \frac{\varepsilon^2 g_{1x} g_{1xx}}{\sqrt{1 + \varepsilon^2 g_{1x}^2}} \equiv J_1.$$

By (2.10)

$$\|J_1\|_{C_x^\alpha} \leq \frac{\varepsilon}{\sqrt{1+t}} \text{ if } \varepsilon < \varepsilon_K.$$

Next

$$(5.2) \quad \frac{\partial^2}{\partial x^2} \sqrt{1 + \varepsilon^2 g_{1x}^2} = \frac{\varepsilon^2 g_{1x}}{\sqrt{1 + \varepsilon^2 g_{1x}^2}} g_{1xxx} + J_2$$

where

$$J_2 = \varepsilon^2 \left( \frac{g_{1xx}}{\sqrt{1 + \varepsilon^2 g_{1x}^2}} - g_{1x} \frac{\varepsilon^2 g_{1x} g_{1xx}}{(1 + \varepsilon^2 g_{1x}^2)^{3/2}} \right) g_{1xx}.$$

Again using (2.10), we get

$$\|J_2\|_{C_x^\alpha} \leq \frac{\varepsilon}{\sqrt{1+t}} \text{ if } \varepsilon \leq \varepsilon_K.$$

We now turn to the expression  $\varepsilon G g_{1x}^2$  in (2.2). Clearly

$$\frac{\partial}{\partial x} (\varepsilon G(\varepsilon g_{1x}) g_{1x}^2) = [\varepsilon^2 G'(\varepsilon g_{1x}) g_{1x}^2 + 2\varepsilon G(\varepsilon g_{1x}) g_{1x}] g_{1xx}$$

so that

$$(5.3) \quad \begin{aligned} \frac{\partial^2}{\partial x^2} (\varepsilon G(\varepsilon g_{1x}) g_{1x}^2) &= [\varepsilon^2 G'(\varepsilon g_{1x}) g_{1x}^2 + 2\varepsilon G(\varepsilon g_{1x}) g_{1x}] g_{1xxx} + J_4 \\ &\equiv J_3 g_{1xxx} + J_4, \end{aligned}$$

where

$$J_4 = \frac{\partial J_3}{\partial x} g_{1xx} = \varepsilon Q(g_{1x}) g_{1xx}^2$$

and  $Q(s)$  is a smooth function. Using (2.10) we find that

$$\|J_3\|_{W_x^{1,\infty}} \leq \frac{\varepsilon^{2/3}}{\sqrt{1+t}} \text{ if } \varepsilon \leq \varepsilon_K,$$

and

$$\begin{aligned} \|J_4\|_{C_x^\alpha} &\leq \varepsilon \|Q(g_{1x})\|_{C_x^\alpha} \left\{ 2 \|g_{1xx}\|_{L^\infty} \|g_{1xx}\|_{C_x^\alpha} \right\} \\ &\leq \frac{1}{\sqrt{1+t}} \text{ if } \varepsilon \leq \varepsilon_K. \end{aligned}$$

We now differentiate (2.2) twice in  $x$  to obtain

$$(5.4) \quad \begin{aligned} (g_{1xx})_t + \frac{b(t)}{1+s_0(t)} \sqrt{1 + \varepsilon^2 g_{1x}^2} g_{1xx} + \frac{2b(t)}{1+s_0(t)} J_1 g_{1x} \\ + \frac{b(t)}{1+s_0(t)} g_{1x} \left( \frac{\varepsilon^2 g_{1x}}{\sqrt{1 + \varepsilon^2 g_{1x}^2}} g_{1xxx} + J_2 \right) \\ - \frac{b(t)}{1+s_0(t)} (J_3 g_{1xxx} + J_4) - \frac{\varepsilon^2 g_{1x} u_1}{\sqrt{1 + \varepsilon^2 g_{1x}^2}} g_{1xxx} - J_2 u_1 - J_5 = 0, \end{aligned}$$

where  $J_5$  involves the first and second derivatives of  $u_1$ . By (2.10), (2.12) and the estimates on  $J_1, J_2$  we find that

$$\|J_5\|_{C_x^\alpha} \leq \frac{C}{1+t} + \frac{C}{\sqrt{1+t}} \|g_1(\cdot, t)\|_{C_x^{1,\alpha}}$$

provided  $\varepsilon \leq \varepsilon_K$ . We can rewrite equation (5.4) in the form

$$(5.5) \quad (g_{1xx})_t + H(x, t)(g_{1xx})_x + \frac{b(t)}{1+s_0(t)} \sqrt{1+\varepsilon^2 g_{1x}^2} g_{1xx} = J_6$$

where, by (2.10) and the estimates on the  $J_i$  ( $i < 6$ ),

$$(5.6) \quad \|J_6\|_{C_x^\alpha} \leq \frac{C}{1+t} + \frac{C}{\sqrt{1+t}} \|g_1(\cdot, t)\|_{C_x^{1,\alpha}}$$

and

$$(5.7) \quad \|H(\cdot, t)\|_{W_x^{1,\infty}} \leq \frac{\sqrt{\varepsilon}}{1+t}$$

if  $\varepsilon \leq \varepsilon_K$ . Introduce the characteristics  $\xi = \xi(x, t)$  by

$$(5.8) \quad \frac{d\xi}{dt} = H(\xi, t), \quad \xi(x, t) = x;$$

in view of (5.7), these are uniquely defined and Lipschitz continuous in  $t$ . We can rewrite (5.5) as

$$(5.9) \quad \frac{d}{dt} g_{1xx}(\xi(x, t), t) + \frac{b(t)}{1+s_0(t)} \sqrt{1+\varepsilon^2 g_{1x}^2} g_{1xx}(\xi(x, t), t) = J_6(\xi(x, t), t).$$

By integration,

$$(5.10) \quad \begin{aligned} & g_{1xx}(\xi(x, t), t) - g_{1xx}(x, 0) \exp\left(-\int_0^t \frac{b(\tau)}{1+s_0(\tau)} \sqrt{1+\varepsilon^2 g_{1x}^2} d\tau\right) \\ &= \int_0^t \exp\left(-\int_s^t \frac{b(\tau)}{1+s_0(\tau)} \sqrt{1+\varepsilon^2 g_{1x}^2} d\tau\right) J_6(\xi(x, s), s) ds, \end{aligned}$$

so that

$$(5.11) \quad \begin{aligned} & \|g_{1xx}\|_{L^\infty(t)} \\ & \leq \|g_{1xx}\|_{L^\infty(0)} e^{-c_0\sqrt{1+t}} + \int_0^t e^{-c_0(\sqrt{1+t}-\sqrt{1+s})} \|J_6\|_{L^\infty(s)} ds \end{aligned}$$

where  $c_0$  is a positive constant such that

$$\frac{b(\tau)}{1+s_0(\tau)} \geq \frac{c_0}{2\sqrt{1+\tau}}.$$

Let

$$U_2(t) = \|g_{1xx}\|_{L^\infty}(t), \quad U_{1,\alpha}(t) = \|g_1\|_{C_x^{1,\alpha}}(t).$$

Using (5.6), we obtain from (5.11)

$$U_2(t) \leq U_2(0)e^{-c_0\sqrt{1+t}} + Ce^{-c_0\sqrt{1+t}} \int_0^t e^{c_0\sqrt{1+s}} \left( \frac{1}{1+s} + \frac{U_{1,\alpha}(s)}{\sqrt{1+s}} \right) ds.$$

By interpolation and Lemma 3.1,

$$U_{1,\alpha}(s) \leq \mu U_2(s) + \tilde{C}_\mu \|g_1\|_{L^\infty}(s) \leq \mu U_2 + \frac{C_\mu}{\sqrt{1+s}}$$

where  $\mu$  is any positive number and  $\tilde{C}_\mu, C_\mu$  depend only on  $\mu$  (but not on  $K, T$ ). It follows that

$$(5.12) \quad U_2(t) \leq U_2(0)e^{-c_0\sqrt{1+t}} + Ce^{-c_0\sqrt{1+t}} \int_0^t e^{c_0\sqrt{1+s}} \left( \frac{C_\mu^*}{1+s} + \frac{\mu U_2(s)}{\sqrt{1+s}} \right) ds$$

where  $C_\mu^* = C_\mu + 1$ .

We shall need the following inequality:

$$(5.13) \quad \int_0^t \frac{e^{\lambda\sqrt{1+s}}}{1+s} ds \leq \frac{\bar{C}_\lambda}{\sqrt{1+t}} e^{\lambda\sqrt{1+t}},$$

which holds for any  $\lambda > 0$  and suitable constant  $C_\lambda$ . To prove it denotes the difference of the right-hand side from the left-hand side by  $h(t)$ . If  $0 \leq t \leq t_0$ , where  $\sqrt{1+t_0} = 2/\lambda$ , then  $h(t) < 0$  provided  $\bar{C}_\lambda$  is sufficiently large. On the other hand, if  $t > t_0$

$$h'(t) = \frac{e^{\lambda\sqrt{1+t}}}{1+t} \left[ 1 - \frac{\bar{C}_\lambda}{2} \left( \lambda - \frac{1}{\sqrt{1+t}} \right) \right] < 0 \text{ if } \lambda \bar{C}_\lambda > 4,$$

and (5.13) follows.

We shall apply to (5.12) Gronwall's inequality which states:

$$(5.14) \quad \begin{aligned} & \text{if } \varphi(t) \leq \psi(t) + \int_0^t \chi(s)\varphi(s) \quad (\varphi > 0, \psi > 0, \chi > 0) \\ & \text{then } \varphi(t) \leq \psi(t) + \int_0^t \chi(s)\psi(s) \exp \left( \int_s^t \chi(\tau) d\tau \right) ds. \end{aligned}$$

We take

$$\begin{aligned}\varphi(t) &= e^{c_0\sqrt{1+t}}U_2(t), \\ \psi(t) &= U_2(0) + CC_\mu^* \int_0^t \frac{e^{c_0\sqrt{1+s}}}{1+s}, \\ \chi(t) &= \frac{C\mu}{\sqrt{1+t}}.\end{aligned}$$

By (5.13)

$$\psi(t) \leq \frac{\tilde{C}_\mu}{\sqrt{1+t}} e^{c_0\sqrt{1+t}}.$$

So

$$\int_0^t \chi(s)\psi(s) \exp\left(\int_s^t \chi(\tau)d\tau\right) ds \leq \int_0^t \frac{C\mu}{\sqrt{1+s}} \frac{\tilde{C}_\mu}{\sqrt{1+s}} e^{c_0\sqrt{1+s}} e^{2C\mu(\sqrt{1+t}-\sqrt{1+s})} ds.$$

Choosing  $\mu$  such that  $2C\mu < c_0$  and using (5.13), we find that the right-hand side is bounded by

$$e^{2C\mu\sqrt{1+t}} C\mu \hat{C}_\mu \int_0^t \frac{e^{(c_0-2C\mu)\sqrt{1+s}}}{1+s} ds \leq \bar{C}_\mu \frac{e^{c_0\sqrt{1+t}}}{\sqrt{1+t}}.$$

We now use Gronwall's inequality (5.14) to immediately deduce from (5.12) that

$$e^{c_0\sqrt{1+t}}U_2(t) \leq \frac{Ce^{c_0\sqrt{1+t}}}{\sqrt{1+t}}.$$

Thus Lemma 5.1 follows, provided (5.1) holds. The assumptions (5.1) can actually be avoided. Since we need only the integral equation (5.10) along characteristics, we may proceed as follows: we first differentiate (2.2) in  $x$  once, write the integral equation along characteristics, and then differentiate it once in  $x$ . In this way we avoid using  $g_{xxx}$ ,  $g_{xxt}$ .

## 6. - Proof of (2.14)

LEMMA 6.1. *If  $u, g$  is a solution satisfying (2.10) for  $0 \leq t \leq T$ , then (2.14) holds for  $0 \leq t \leq T$  provided  $0 < \varepsilon < \varepsilon_K$ ;  $\varepsilon_K$  is a constant independent of  $T$ , and  $C$  is a constant independent of  $K, T, \varepsilon$ .*

PROOF. For simplicity we again make the assumption (5.1). Then (5.7) is satisfied and, by (5.6) and Lemma 5.1,

$$(6.1) \quad \|J_6\|_{C_x^\alpha} \leq \frac{C}{1+t}.$$

For fixed  $x, \bar{x}$ , consider the function

$$w(t) = g_{1xx}(\xi(x, t), t) - g_{1xx}(\xi(\bar{x}, t), t).$$

Using (2.10) and (6.1) we find, from (5.9), that

$$(6.2) \quad \frac{d}{dt}w + \frac{b(t)}{1+s_0(t)}\sqrt{1+\varepsilon^2g_{1x}^2}w = R(x, t),$$

where

$$(6.3) \quad \|R(\cdot, t)\|_{L^\infty} \leq \frac{C}{1+t}|\xi(x, t) - \xi(\bar{x}, t)|^\alpha$$

and as before (cf. (5.11))

$$(6.4) \quad |w(t)| \leq |w(0)|e^{-c_0\sqrt{1+t}} + e^{-c_0\sqrt{1+t}} \int_0^t e^{c_0\sqrt{1+s}} \frac{C}{1+s} |\xi(x, s) - \xi(\bar{x}, s)|^\alpha ds.$$

From (5.8), (5.7) we have

$$\frac{d\xi_x(x, t)}{dt} = H_x(\xi, t)\xi_x, \quad \xi_x(x, 0) = 1$$

with

$$|H_x(x, t)| \leq \frac{\sqrt{\varepsilon}}{1+t}$$

and then, by comparison,

$$(6.5) \quad \frac{1}{(1+t)\sqrt{\varepsilon}} \leq \xi_x \leq (1+t)\sqrt{\varepsilon}.$$

This implies that the inverse function  $x = x(\xi, t)$  exists and satisfies:

$$(6.6) \quad \frac{1}{(1+t)\sqrt{\varepsilon}} \leq x_\xi \leq (1+t)\sqrt{\varepsilon}.$$

From (6.5) it follows that

$$|\xi(x, s) - \xi(\bar{x}, s)|^\alpha \leq C(1+s)^{\frac{1}{4}}|x - \bar{x}|^\alpha$$

if  $\sqrt{\varepsilon}\alpha < \frac{1}{4}$ . Using this in (6.4), we get

$$(6.7) \quad |w(t)| \leq |x - \bar{x}|^\alpha e^{-c_0\sqrt{1+t}} + e^{-c_0\sqrt{1+t}} \int_0^t \frac{C}{(1+s)^{3/4}} e^{c_0\sqrt{1+s}} |x - \bar{x}|^\alpha ds$$

$$\leq \frac{C}{(1+t)^{1/4}} |x - \bar{x}|^\alpha,$$



since

$$\int_0^t \frac{e^{\lambda\sqrt{1+s}}}{(1+s)^{3/4}} ds \leq \frac{\bar{C}_\lambda}{(1+t)^{1/4}} e^{\lambda\sqrt{1+t}}$$

(the proof is the same as for (5.13)).

For any  $x_1, x_2, t$ , choose  $x$  and  $\bar{x}$  such that

$$x_1 = \xi(x, t), \quad x_2 = \xi(\bar{x}, t).$$

Then

$$|x_1 - x_2| = |\xi_x(\hat{x})| |x - \bar{x}| \geq \frac{1}{(1+t)\sqrt{\varepsilon}} |x - \bar{x}|$$

by (6.5). Using this in (6.7) we get

$$\begin{aligned} |g_{1xx}(x_1, t) - g_{1xx}(x_2, t)| &\leq \frac{C}{(1+t)^{1/4}} (1+t)^{\sqrt{\varepsilon}\alpha} |x_1 - x_2|^\alpha \\ &\leq C|x_1 - x_2|^\alpha \text{ if } \sqrt{\varepsilon}\alpha < \frac{1}{4}. \end{aligned}$$

This completes the proof of the lemma.

We summarize the results of Sections 3–6:

**THEOREM 6.2.** *Consider the problem (1.1)–(1.6), (1.11), (1.12) under the assumptions (2.6)–(2.9). If  $u, g$  is a solution for  $0 \leq t \leq T$  ( $0 < T < \infty$ ) satisfying (2.10), then it also satisfies (2.11)–(2.14), provided  $0 < \varepsilon < \varepsilon_K$ ;  $\varepsilon_K$  is a positive constant independent of  $T$ , and  $C$  is a positive constant independent of  $K, T, \varepsilon$ .*

**REMARK 6.1.** From (2.2) it follows that also

$$(6.8) \quad \|g_{1t}(\cdot, t)\|_{L^\infty} \leq \frac{C}{\sqrt{1+t}}, \quad \|g_{1xt}(\cdot, t)\|_{L^\infty} \leq \frac{C}{\sqrt{1+t}}.$$

**REMARK 6.2.** The proof of Theorem 6.2 breaks down if we relax the growth conditions on  $B(t)$  in (2.8). Indeed, suppose (instead of  $\int_0^\infty B(t)dt < \infty$ )

that

$$|B(t)| \leq \frac{C}{(1+t)^\kappa} \quad \left(\frac{1}{2} \leq \kappa < 1\right).$$

The comparison results of Section 2 are still valid, yielding the estimate

$$|g_1| \leq \frac{C}{(1+t)^{\kappa-1/2}}.$$

This suggests the extension of Theorem 6.2 with  $K/\sqrt{1+t}$  in (2.10) replaced by  $K/(1+t)^{\kappa-1/2}$ ; a term  $C/(1+t)^{\kappa-1/2}$  should then be added to the right-hand

sides of (2.11), (2.13), and  $C/(1+t)$  in (2.12) should be replaced by  $C/(1+t)^\kappa$ . Next, in (5.12),

$$\frac{C_\mu^*}{1+s} \text{ is replaced by } \frac{C_\mu^*}{(1+s)^\kappa},$$

and, analogously to (5.13),

$$\int_0^t \frac{e^{\lambda\sqrt{1+s}}}{(1+s)^\kappa} ds \leq \frac{\bar{C}_\mu}{(1+t)^{\kappa-1/2}} e^{\lambda\sqrt{1+t}}.$$

With these changes we can now proceed as before to derive the estimate

$$\|g_{1xx}(\cdot, t)\|_{L^\infty} \leq \frac{C}{(1+t)^{\kappa-1/2}}$$

which replaces (2.13). This estimate, however, is too weak for establishing the appropriate bound on the Hölder coefficient of  $g_{1xx}$ . Indeed, instead of (6.5) we only get

$$\exp\left\{-\sqrt{\varepsilon} \frac{(1+t)^{1-\kappa}}{1-\kappa}\right\} \leq \xi_x \leq \exp\left\{\sqrt{\varepsilon} \frac{(1+t)^{1-\kappa}}{1-\kappa}\right\},$$

which is insufficient for the proof of (2.14).

### 7. - Existence theorems

Consider (1.1)–(1.6) with

$$(7.1) \quad g_0(x) \geq c_0 > 0, \quad \|g_0\|_{C^{2,\alpha}} \leq K < \infty,$$

$$(7.2) \quad b(x, t) \text{ continuous and } |b(x, t)| \leq C < \infty \text{ for } x \in \mathbb{R}, t \geq 0.$$

We shall prove that for some small  $T > 0$  there exists a solution  $(u, g)$  with  $g$  in the class

$$B_{K,M} \equiv \left\{ g(x, t), 0 \leq t \leq T; g(x, t) \geq \frac{1}{2}c_0, \|g(\cdot, t)\|_{L^\infty} \leq 2K, \right. \\ \left. \|g_x(\cdot, t)\|_{L^\infty} \leq 2K, \|g_{xx}(\cdot, t)\|_{L^\infty} \leq 2K, [g_{xx}(\cdot, t)]_{C^\alpha} \leq 2K, \right. \\ \left. \|g_t(\cdot, t)\|_{L^\infty} \leq M, \text{ and } g(x, 0) = g_0(x) \right\},$$

where  $M$  is a positive constant to be determined.

For any  $g \in B_{K,M}$  define  $\Gamma : y = g(x, t)$  and let  $u$  be the solution of (1.2)–(1.4). By the maximum principle

$$(7.3) \quad |u| \leq C^* \quad (C^* \text{ independent of } K, M, T).$$

and by the Schauder estimates

$$(7.4) \quad \|u(\cdot, \cdot, t)\|_{C_{x,y}^{2,\alpha}(D_t)} \leq C_K \quad (C_K \text{ independent of } M, T),$$

where  $D_t$  is defined as in (4.1). Let

$$v(x, t) = u(x, g(x, t), t)$$

and let  $\varphi_\delta$  be mollifiers in  $x$ , and set

$$v_\delta(x, t) = (\varphi_\delta * v(\cdot, \cdot, t))(x).$$

Then

$$(7.5) \quad \|v_\delta(\cdot, t)\|_{C_x^{2,\alpha}} \leq C_K, \quad |v_\delta| \leq C^*,$$

We also introduce

$$g_\delta(x) = \varphi_\delta * g_0(x);$$

clearly

$$(7.6) \quad \|g_\delta\|_{C^{2,\alpha}} \leq K.$$

For any small  $\varepsilon > 0$ , let  $\tilde{g}(x, t)$  be the solution of

$$(7.7) \quad \tilde{g}_t = \sqrt{1 + \tilde{g}_x^2} v_\delta(x, t) + \varepsilon \tilde{g}_{xx},$$

$$(7.8) \quad \tilde{g}(x, 0) = g_\delta(x).$$

By comparison [2; p. 52]

$$(7.9) \quad \begin{aligned} \tilde{g}(x, t) &\leq C^* t + \max g_\delta(x), \\ \tilde{g}(x, t) &\geq -C^* t + \inf g_\delta(x), \end{aligned}$$

so that

$$(7.10) \quad \tilde{g}(x, t) > \frac{c_0}{2}, \quad \|\tilde{g}\|_{L^\infty} \leq 2K$$

provided  $T$  is small (depending only on  $C^*$ ,  $K$ ). Next differentiate (7.7) in  $x$  to obtain

$$(7.11) \quad \mathcal{L}\tilde{g}_x \equiv \frac{\partial}{\partial t}\tilde{g}_x - \frac{v_\delta(x,t)}{\sqrt{1+\tilde{g}_x^2}}\tilde{g}_x \frac{\partial}{\partial x}\tilde{g}_x - v_{\delta,x}(x,t)\sqrt{1+\tilde{g}_x^2} - \varepsilon \frac{\partial^2}{\partial x^2}\tilde{g}_x = 0.$$

The function  $w = K + Ct$  satisfies

$$\mathcal{L}w = C - v_{\delta,x}(x,t)\sqrt{1+w^2} > 0$$

if  $C \geq C_K$  and  $T$  is small depending only on  $K$ . It follows, by comparison with  $\tilde{g}_x$ , that

$$(7.12) \quad \tilde{g}_x \leq w \leq 2K$$

if  $T$  is small enough, and similarly

$$(7.13) \quad \tilde{g}_x \geq -2K.$$

Differentiating (7.11) once more in  $x$  and using (7.10), (7.12) (7.13), we obtain by comparison, as before,

$$(7.14) \quad |\tilde{g}_{xx}| \leq Ct + K \leq 2K$$

where  $C = C_K$ , and  $T$  is small enough, depending only on  $K$ .

Finally, from (7.7),

$$(7.15) \quad |\tilde{g}_t| \leq M$$

where  $M$  depends only on  $K$ , provided  $T$  is small enough (depending only on  $K$ ).

Observe that  $u$  is continuous in  $t$  (by compactness and uniqueness). Therefore also

$$(7.16) \quad v_\delta \text{ is continuous in } (x, t).$$

We next observe that the problem

$$(7.17) \quad g_t = \sqrt{1+g_x^2}v_\delta(x, t),$$

$$(7.18) \quad g(x, 0) = g_\delta(x)$$

has at most one solution. Indeed, this follows by estimating the difference of two solutions, making use of the Lipschitz continuity of  $v_\delta(x, t)$  in  $x$  and its continuity in  $t$  (by (7.16)).

From the above observations and the estimates (7.10), (7.12), (7.13), (7.14), (7.15), it follows that the family  $\tilde{g} \equiv \tilde{g}_\varepsilon$  converges to a (unique) solution  $g^*$  of (7.17), (7.18) as  $\varepsilon \rightarrow 0$ .

By differentiating (7.17) formally twice in  $x$  we get

$$(7.19) \quad \begin{aligned} \frac{\partial}{\partial t} g_{xx}^* - \frac{g_x^*}{\sqrt{1+g_x^{*2}}} g_{xxx}^* &= \frac{1}{(1+g_x^{*2})^{3/2}} (g_{xx}^*)^2 v_\delta(x, t) \\ &+ \frac{2g_x^*}{\sqrt{1+g_x^{*2}}} g_{xx}^* v_{\delta,x}(x, t) + \sqrt{1+g_x^{*2}} v_{\delta,xx}(x, t). \end{aligned}$$

To justify this differentiation note that by differentiating (7.11) successively in  $x$  and comparing with functions of the form  $Ct + C_1$  we can estimate the derivatives  $\tilde{g}_{xxx}$ ,  $\tilde{g}_{xxxx}$ , etc. as we have done in (7.14). The constants depend on  $\delta$  but not on  $\varepsilon$ . Hence differentiating (7.7) twice in  $x$  and then letting  $\varepsilon \rightarrow 0$ , equation (7.19) follows.

Next we introduce the characteristics

$$(7.20) \quad \frac{d\xi}{dt} = - \left( \frac{g_x^*}{\sqrt{1+g_x^{*2}}} v_\delta \right) (\xi, t), \quad \xi(x, 0) = x$$

and note that

$$\frac{1}{2} \leq \frac{d\xi}{dx} \leq 2$$

if  $T$  is small. Writing (7.19) in integrated form along characteristics, we can derive the inequality

$$\begin{aligned} \left| g_{xx}^*(\xi(x_1, t), t) - g_{xx}^*(\xi(x_2, t), t) \right| &\leq \left| g_{0,xx}(x_1) - g_{0,xx}(x_2) \right| \\ &+ \int_0^t \left| A_1 \cdot [g_{xx}^*(\xi(x_1, s), s) - g_{xx}^*(\xi(x_2, s), s)] + A_2 \right| ds \end{aligned}$$

where  $|A_j| \leq C_K$ ,  $C_K$  independent of  $\delta$ . It easily follows that

$$[g_{xx}^*]_{C_x^\alpha} \leq 2K$$

if  $T$  is small.

Consider the mapping  $W$  defined by  $g \rightarrow Wg = g^*$ . We have proved that  $W$  maps  $B_{K,M}$  into itself provided  $T$  is sufficiently small (depending on  $K$ , but not on  $\delta$ ).

If we provide  $B_{K,M}$  with the uniform topology, then  $B_{K,M}$  is compact. From the uniqueness of solution to (7.17), (7.18) and compactness, it follows that  $W$  is continuous. Hence, by the Schauder fixed-point theorem,  $W$  has a fixed point  $g_\delta^*$ . Letting  $\delta \rightarrow 0$  through an appropriate subsequence, we obtain a

limiting function  $g$ , which together with the corresponding  $u$ , provide a solution to (1.1)–(1.6).

We summarize:

**THEOREM 7.1.** *If (7.1), (7.2) hold, then there exists a solution  $(u, g)$  of (1.1)–(1.6) with  $g$  in the set  $B_{K,M}$ , provided  $T$  is sufficiently small.*

The proof shows that  $T$  depends only on  $\|g_0\|_{C_x^{2,\alpha}}$  (and that all the norms in (2.10) are continuous in  $t$ ). Hence the solution can be extended step-by-step for all times as long as one can establish a priori estimate on

$$\|g(\cdot, t)\|_{C_x^{2,\alpha}}$$

independently of  $t$ . Such an estimate has already been derived in Theorem 6.2. We may therefore state:

**THEOREM 7.2.** *Consider the problem (1.1)–(1.6), (1.11), (1.12) under the assumptions (2.6)–(2.9). If  $\varepsilon$  is sufficiently small (depending on  $\|g_0\|_{C^{2,\alpha}}$ ), then there exists a global solution.*

The solution satisfies (2.10) and (6.8),  $g_{xx}$  is continuous in  $t$ , and the norms in (2.10) are continuous in  $t$ .

**REMARK 7.1.** The reason for introducing the mollifiers  $\varphi_\delta$  in the proof of Theorem 7.1 is to justify the calculations which involve third derivations of  $g$ . The diffusion term  $\varepsilon g_{xx}$  was introduced in (7.7) so that we can use a parabolic comparison theorem.

## 8. - Uniqueness

In Section 7 we proved the existence of solutions  $(u, g)$  such that

$$(8.1) \quad \text{all the norms in (2.10) and } g, g_x, g_{xx} \text{ are continuous in } t.$$

We shall now establish uniqueness of such solutions for general data.

**THEOREM 8.1.** *Assume that, for some  $T > 0$ ,  $(u, g), (\tilde{u}, \tilde{g})$  are two solutions of (1.1)–(1.6) satisfying (8.1). Then  $u \equiv \tilde{u}$  and  $g \equiv \tilde{g}$ .*

**PROOF.** By assumption

$$\|g(\cdot, t)\|_{C_x^{2,\alpha}}, \|\tilde{g}(\cdot, t)\|_{C_x^{2,\alpha}} \leq C$$

and therefore by Schauder's estimates, for any  $\delta > 0$ ,

$$(8.2) \quad \|u(\cdot, \cdot, t)\|_{C^{2,\alpha}(\Omega_t^\delta)}, \|\tilde{u}(\cdot, \cdot, t)\|_{C^{2,\alpha}(\tilde{\Omega}_t^\delta)} \leq C_\delta$$

where  $\Omega_t^\delta = \{(x, y); \delta < y < g(x, t)\}$  and  $\tilde{\Omega}_t^\delta = \{(x, y); \delta < y < \tilde{g}(x, t)\}$ . Set

$$(8.3) \quad V(t) = \sup_x |g(x, t) - \tilde{g}(x, t)|$$

and introduce the domain

$$G_t = \{(x, y); 0 < y < g(x, t) - V(t)\}.$$

Then  $\partial G_t = \{y = 0\} \cup S_t$ , where

$$S_t = \{y = g(x, t) - V(t)\}$$

is uniformly in  $C_x^{2,\alpha}$ . The outward normal along  $S_t$  is

$$n = \frac{(-g_x, 1)}{\sqrt{1 + g_x^2}}.$$

Set

$$J_1 = \left( \frac{\partial u}{\partial n} + u \right) \Big|_{y=g(x,t)-V(t)} - \left( \frac{\partial u}{\partial n} + u \right) \Big|_{y=g(x,t)},$$

$$J_2 = \left( \frac{\partial \tilde{u}}{\partial n} + \tilde{u} \right) \Big|_{y=g(x,t)-V(t)} - \left( \frac{\partial \tilde{u}}{\partial n} + \tilde{u} \right) \Big|_{y=\tilde{g}(x,t)}.$$

Then, by (8.2),

$$(8.4) \quad \|J_1\|_{C_x^\alpha} \leq CV(t)$$

and

$$(8.5) \quad \|J_2\|_{C_x^\alpha} \leq CV(t) + C \|g(\cdot, t) - \tilde{g}(\cdot, t)\|_{C_x^\alpha}.$$

Introducing the normal to  $\{y = \tilde{g}(x, t)\}$ ,

$$\tilde{n} = \frac{(-\tilde{g}_x, 1)}{\sqrt{1 + \tilde{g}_x^2}},$$

we also have

$$(8.6) \quad \left\| \left( \frac{\partial \tilde{u}}{\partial n} - \frac{\partial \tilde{u}}{\partial \tilde{n}} \right) \Big|_{y=\tilde{g}(x,t)} \right\|_{C_x^\alpha} \leq C \|g_x - \tilde{g}_x\|_{C_x^\alpha}(t).$$

Using the free boundary condition (1.4) for both  $u$  and  $\tilde{u}$  and the estimates (8.4)–(8.6), we easily obtain

$$(8.7) \quad \left\| \left( \frac{\partial(u - \tilde{u})}{\partial n} + (u - \tilde{u}) \right) \Big|_{y=g(x,t)-V(t)} \right\|_{C_x^\alpha} \leq C \|g - \tilde{g}\|_{C_x^{1,\alpha}}(t).$$

By (8.7) and the maximum principle

$$(8.8) \quad \|(u - \tilde{u})(\cdot, t)\|_{L^\infty(G_t)} \leq C\|g_x - \tilde{g}_x\|_{C_x^\alpha(t)}.$$

Denote by  $w$  the harmonic conjugate of  $u - \tilde{u}$ . Then, by (8.7),

$$\left\| \frac{\partial w}{\partial s} \right\|_{C_x^\alpha} \leq C\|u - \tilde{u}\|_{C_x^\alpha} + C\|g - \tilde{g}\|_{C_x^{1,\alpha}},$$

where  $w$  and  $u - \tilde{u}$  are evaluated on  $S_t$ . Applying to  $w$  elliptic  $C^{1,\alpha}$  estimates [4, Theorem 2.4], we then easily get

$$\|w\|_{C_{x,y}^{1,\alpha}(G_t)} \leq C\|u - \tilde{u}\|_{C_x^\alpha(S_t)} + C\|g - \tilde{g}\|_{C_x^{1,\alpha}}.$$

Since, for any  $\delta > 0$ ,

$$\|u - \tilde{u}\|_{C_x^\alpha} \leq \delta\|u - \tilde{u}\|_{C_x^{1,\alpha}} + C_\delta\|u - \tilde{u}\|_{L^\infty} \text{ on } S_t$$

and

$$\|u - \tilde{u}\|_{C_x^{1,\alpha}(S_t)} \leq C\|w\|_{C_{x,y}^{1,\alpha}(G_t)},$$

it follows (by choosing  $\delta$  small enough) that

$$(8.9) \quad \|(u - \tilde{u})\|_{C_{x,y}^{1,\alpha}(G_t)} \leq C\|g - \tilde{g}\|_{C_x^{1,\alpha}(t)}.$$

Next, differentiating in  $x$  the free boundary condition

$$g_t = \sqrt{1 + g_x^2}u(x, g, t),$$

we get

$$(8.10) \quad g_{xt} = H(x, t)g_{xx} + K(x, t),$$

where

$$H(x, t) = \frac{g_x}{\sqrt{1 + g_x^2}}u(x, g(x, t), t),$$

$$K(x, t) = \sqrt{1 + g_x^2} \left( u_x(x, g(x, t), t) + u_y(x, g(x, t), t)g_x \right).$$

A similar formula holds for  $\tilde{g}$ . Using (8.2) and (8.9) we can estimate

$$\|K - \tilde{K}\|_{C_x^\alpha} + \|(H - \tilde{H})\tilde{g}_{xx}\|_{C_x^\alpha} \leq C\|g(\cdot, t) - \tilde{g}(\cdot, t)\|_{C_x^{1,\alpha}}.$$

Consequently, the function  $g^* = g - \tilde{g}$  satisfies

$$(8.11) \quad \|g_{xt}^* - H(x, t)g_{xx}^*\|_{C_x^\alpha} \leq C\|g^*\|_{C_x^{1,\alpha}(t)};$$



also,

$$(8.12) \quad \|H(\cdot, t)\|_{C_x^{1,\alpha}} \leq C.$$

Introduce the characteristics

$$(8.13) \quad \frac{d\xi}{dt} = -H(\xi, t), \quad \xi(x, 0) = 0.$$

By (8.12),

$$(8.14) \quad \frac{1}{2} \leq \frac{d\xi}{dx} \leq 2, \quad \text{for } 0 < t < t_0,$$

if  $t_0$  is small enough. Integrating (8.11) along characteristics, we obtain

$$(8.15) \quad \|g_x^*\|_{L_x^\infty}(t) \leq Ct \max_{0 \leq \tau \leq t} \|g^*(\cdot, \tau)\|_{C_x^{1,\alpha}}.$$

Using (8.14) and proceeding as in Section 6, we can also get (much more simply) the estimate

$$(8.16) \quad [g_x^*]_{C_x^\alpha}(t) \leq Ct \max_{0 \leq \tau \leq t} \|g^*(\cdot, \tau)\|_{C_x^{1,\alpha}}.$$

From (8.15), (8.16), it follows that

$$g(x, t) \equiv \tilde{g}(x, t) \quad \text{for } 0 \leq t \leq \tau,$$

if  $\tau$  is small enough, and then also  $u(x, y, t) \equiv \tilde{u}(x, y, t)$  for  $0 \leq t < \tau$ .

We can now proceed step-by-step to prove that  $\tilde{g} = g$ ,  $u = \tilde{u}$ , for all  $0 \leq t \leq T$ .

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