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FRANK MERLE

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Asymptotics for L^2 minimal blow-up solutions of critical nonlinear Schrödinger equation

by

Frank MERLE

Université de Cergy-Pontoise, Centre de Mathématiques Avenue du Parc 8, Le Campus, 95033 Cergy-Pontoise, France

ABSTRACT. – In this note, we describe the behavior of a sequence $v_n: \mathbb{R}^N \to \mathbb{C}$ minimal in L^2 such that $\frac{1}{2} \int \left| \nabla v_n \right|^2 - \frac{1}{\frac{4}{N}+2} \int \left| v_n \right|^{\frac{4}{N}+2} \leq E_0$ and $|v_n|_{H^1} \to +\infty$.

Résumé. — Dans cette note, on explicite de façon optimale le comportement d'une suite $v_n:\mathbb{R}^N\to\mathbb{C}$ de norme L^2 minimale telle que $\frac{1}{2}\int \left|\nabla v_n\right|^2 - \frac{1}{\frac{4}{N}+2}\int \left|v_n\right|^{\frac{4}{N}+2} \leq E_0$ et $|v_n|_{H^1}\to +\infty$.

In the present note, we are interested in the behavior of a sequence $v_n:\mathbb{R}^N\to\mathbb{C}$ of H^1 functions such that

$$\int |v_n|^2 = \int Q^2,$$

(2)
$$E(v_n) = \frac{1}{2} \int |\nabla v_n|^2 - \frac{1}{\frac{4}{N} + 2} \int |v_n|^{\frac{4}{N} + 2} \le E_0,$$

(3)
$$\int |\nabla v_n|^2 \to +\infty,$$

where Q is the radial positive symetric solution of the equation

$$\Delta v + |v|^{\frac{4}{N}}v = v.$$

(See references [1], [4], [7] for existence and uniqueness of Q.)

This problem is related to the asymptotics of minimal blow-up solutions in \mathcal{H}^1 of the equation

(5)
$$iu_t = -\Delta u - k(x)|u|^{\frac{4}{N}}u \text{ and } u(0) = \varphi,$$

where

(6)
$$\max_{x \in \mathbf{R}^N} k(x) = 1.$$

Indeed, for all $\varphi \in H^1$, there is a unique solution in H^1 on [0,T] ([2], [4]) and

$$T = +\infty$$
 or $\lim_{t \to T} \int |\nabla u(t, x)|^2 = +\infty$.

In addition, $\forall t$

(7)
$$\int |u(t,x)|^2 dx = \int |\varphi(x)|^2 dx$$

(8)
$$E_k(u(t)) = E_k(\varphi)$$

where

$$E_k(v) = \frac{1}{2} \int |\nabla v|^2 - \frac{1}{\frac{4}{N} + 2} \int k(x)|v|^{\frac{4}{N} + 2}.$$

From [9]

(9)
$$\forall v \in H^1, \quad \frac{1}{\frac{4}{N} + 2} \int |v|^{\frac{4}{N} + 2} \le \frac{1}{2} \left(\frac{\int |v|^2}{\int |Q|^2} \right)^{\frac{2}{N}} \int |\nabla v|^2$$

and it follows from (6)-(9) ([6]) that

if
$$|\varphi|_{L^2} < |Q|_{L^2}$$
, then $T = +\infty$.

Moreover under some conditions on k(x), for any $\varepsilon > 0$ there are blow-up solutions $u_{\varepsilon}(t)$ such that

$$|u_{\varepsilon}(0)|_{L^{2}}^{2} = |\varphi|_{L^{2}}^{2} + \varepsilon \quad ([6]).$$

Thus the questions are about existence of minimal blow-up solution (that is such that u(t) blows up in finite time and $\int |\varphi|^2 = \int Q^2$ and on the

form of these solutions. In the case where $k(x) \equiv 1$, the question has been completely solved (see Merle [5]). The general case is still open. We remark that from (6)-(9), if u(t) is a blow up solution, the sequences $v_n = u(t_n)$ as $t_n \to T$ satisfies (1)-(3) and we ask about the constrains it imply on v_n .

The first result in this direction was obtained by Weinstein in [9]. Using the concentration compactness method, he showed that there is a $\theta_n \in \mathbb{R}$, $x_n \in \mathbb{R}^N$ such that

(10)
$$v_n = \lambda_n^{\frac{N}{2}} e^{i\theta_n} Q\left(\lambda_n^{\frac{N}{2}} (x - x_n)\right) + \varepsilon_n,$$

where

(11)
$$\lambda_n = \frac{|\nabla v_n|_{L^2}}{|\nabla Q|_{L^2}},$$

(12)
$$|\varepsilon_n|_{L^2} \xrightarrow{n \to +\infty} 0 \text{ and } \frac{|\nabla \varepsilon_n|_{L^2}}{\lambda_n} \xrightarrow{n \to +\infty} 0.$$

In [5], Merle then showed that for all R > 0, there is a c > 0 such that

$$(13) \qquad \int_{|x-x_n| > R} |\nabla v_n|^2 \le c.$$

We now claim the following result

THEOREM. – Let (v_n) be a sequence of H^1 functions satisfying (1)-(3) and $\theta_n(x)$ be such that $v_n = |v_n|e^{i\theta_n}$.

i) Phase estimates. There is a c > 0 such that

$$\forall n, \quad \int |v_n|^2 |\nabla \theta_n|^2 \le c.$$

ii) asymptotics on the modulus.

There is a $\varepsilon_n(x)$, $x_n \in \mathbb{R}^N$, and c > 0 such that

$$\forall x, |v_n(x)| = \lambda_n^{\frac{N}{2}} Q(\lambda_n(x - x_n)) + \varepsilon_n(x)$$

where

$$|\nabla \varepsilon_n|_{L^2} \le c$$
, $|\varepsilon_n|_{L^2} \le \frac{c}{\lambda_n}$ and $\lambda_n \left(\frac{|\nabla Q|_{L^2}}{|\nabla v_n|_{L^2}}\right) \xrightarrow[n \to +\infty]{} 1$.

Remark. – This Theorem simplifies some proofs in [5], [6]. The case where v_n is real valued is also related to similar problems for the generalized Korteweg-de Vries equation with critical nonlinearity.

Remark. – The Theorem implies in particular for blow-up solution of equation (5) $u(t,x)=|u(t,x)|e^{i\theta(t,x)}$ and $\int |u(t)|^2=\int Q^2$ the phase gradient is uniformly bounded at the blow-up: there is a c>0 such that

$$\forall \ 0 < t < T, \quad \int |u(t,x)|^2 |\nabla \theta(t,x)|^2 dx \le c.$$

(Of course, we still have
$$\int |\nabla |u||^2(t,x) \xrightarrow{t \to T} + \infty$$
.)

Remark. – It is easy to check that the result is optimal. We remark that the residual term in the theorem $\varepsilon_n = O(1)$ (compared to $o(|\nabla v_n|_{L^2})$ in [9]).

Proof of the Theorem. – Let (v_n) a sequence of H^1 function satisfying (1)-(3) and $\theta_n(x)$ such that $v_n = |v_n|e^{i\theta_n}$. We have that

(14)
$$\frac{1}{2} \int |\nabla v_n|^2 = \frac{1}{2} \left\{ |\nabla |v_n||^2 + \int |v_n|^2 |\nabla \theta_n|^2 \right\}$$

and

(15)
$$E(v_n) = \frac{1}{2} \int |v_n|^2 |\nabla \theta_n|^2 + E(|v_n|).$$

The idea is to apply the variational identity (9) not with v_n but with $|v_n|$. Indeed, since $v_n \in H^1$ we have that $|v_n| \in H^1$. From (9) (applied with $|v_n|$)

$$(16) \quad \frac{1}{\frac{4}{n}+2} \int |v_n|^{\frac{4}{n}+2} \le \frac{1}{2} \left(\frac{\int |v_n|^2}{\int Q^2} \right)^{\frac{2}{N}} \int |\nabla |v_n||^2 \le \frac{1}{2} \int |\nabla |v_n||^2,$$

or equivalently

$$(17) E(|v_n|) \ge 0.$$

Thus (2), (15), (17) imply that

$$(18) \qquad \frac{1}{2} \int |v_n|^2 |\nabla \theta_n|^2 \le E_0$$

$$(19) E(|v_n|) \le E_0.$$

Part i). - It is implied by (18).

Part ii). – We claim that it is as a consequence of (18)-(19). We prove it in three steps:

- from Weinstein's results, we first obtain rough estimates on $|v_n|$,
- we then choose appropriate approximations parameters,
- we conclude the proof using a convexity property in certain directions of E near Q (and use in a crucial way that $|v_n|$ is a real-valued function).

Step 1: First asymptotics. - Since

(20)
$$\int |\nabla v_n|^2 = \int |\nabla |v_n||^2 + \int |v_n|^2 |\nabla \theta_n|^2 \xrightarrow[n \to +\infty]{} + \infty,$$

and

$$\forall n, \int |v_n|^2 |\nabla \theta_n|^2 \le c,$$

we have

(21)
$$\int |\nabla |v_n||^2 \xrightarrow[n \to +\infty]{} + \infty.$$

Moreover,

(22)
$$\int |v_n|^2 = \int Q^2 \text{ and } E(|v_n|) \le E_0.$$

We conclude from Weinstein's result on the existence of \hat{x}_n , $\hat{\varepsilon}_n$ such that

(23)
$$|v_n|(x) = \hat{\lambda}_n^{N/2} Q(\hat{\lambda}_n x - \hat{x}_n) + \hat{\varepsilon}_n(x)$$

where

$$\hat{\lambda}_n = \frac{|\nabla |v_n||_{L^2}}{|\nabla Q|_{L^2}},$$

(25)
$$|\nabla \hat{\varepsilon}_n|_{L^2} = o(\hat{\lambda}_n), \quad |\hat{\varepsilon}_n|_{L^2} = o(1).$$

In order to obtain better estimates on the rest (that is $|\nabla \varepsilon_n|_{L^2} \leq c$), we have to choose appropriate parameters λ_n, x_n and use the structure of the functional $E(\cdot)$ near Q.

Step 2: Choice of the parameters of approximation. – Let us first renormalize the problem. We consider

(26)
$$w_{n,\lambda_1,x_1}(x) = \left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^{\frac{N}{2}} |v_n| \left((\lambda_1 x + \hat{x}_n + x_1) \frac{1}{\hat{\lambda}_n}\right)$$

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(27)
$$\tilde{\varepsilon}_{n,\lambda_1,x_1}(x) = \left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^{\frac{N}{2}} \hat{\varepsilon}_n \left((\lambda_1 x + \hat{x}_n + x_1) \frac{1}{\hat{\lambda}_n}\right).$$

We have from (23),

(28)
$$w_{n,\lambda_1,x_1}(x) = \lambda_1^{N/2} Q(\lambda_1 x + x_1) + \tilde{\varepsilon}_{n,\lambda_1,x_1}(x)$$

where

(29)
$$\frac{|\nabla \tilde{\varepsilon}_n|_{L^2}}{\lambda_1} + |\tilde{\varepsilon}_n|_{L^2} \xrightarrow[n \to +\infty]{} 0.$$

We write (28) as follows

$$w_{n,\lambda_1,x_1}(x) = Q(x) + \varepsilon_{n,\lambda_1,x_1}(x)$$

where

(30)
$$\varepsilon_{n,\lambda_1,x_1}(x) = \left[\lambda_1^{N/2} Q(\lambda_1 x + x_1) - Q(x)\right] + \tilde{\varepsilon}_{n,\lambda_1,x_1}(x).$$

From the implicit function Theorem, we derive easily for $|\tilde{\varepsilon}_n|_{H^1}$ small enough the existence of $\lambda_{1,n}, x_{1,n}$ such that

(31)
$$\forall i = 1, ..., N, \int \varepsilon_{n, \lambda_{1,n}, x_{1,n}} x_i Q = 0$$

(32)
$$\int \varepsilon_{n,\lambda_{1,n},x_{1,n}} |x|^2 Q = 0.$$

Moreover, from (29)

$$(33) \qquad (\lambda_{1,n}, x_{1,n}) \xrightarrow[n \to +\infty]{} (1,0).$$

Indeed, let us note

for
$$i = 1, ..., N$$
, $\rho_i(\lambda_1, x_1) = \int \varepsilon_{n, \lambda_1, x_1} x_i Q$,

$$\rho_{N+1}(\lambda_1, x_1) = \int \varepsilon_{n, \lambda_1, x_1} |x|^2 Q.$$

From (30), we have

$$\begin{split} &\frac{\partial \varepsilon_{n,1,0}}{\partial x_{1,i}} = \partial_i Q + \partial_i \ \tilde{\varepsilon}_{n,1,0} \\ &\frac{\partial \varepsilon_{n,1,0}}{\partial \lambda_1} = \frac{N}{2} Q + x. \nabla Q + \left(\frac{N}{2} \tilde{\varepsilon}_{n,1,0} + x. \nabla \tilde{\varepsilon}_{n,1,0}\right), \end{split}$$

where $x_1 = (x_{1,1}, ..., x_{1,N}).$

Therefore, from (29) and integration by parts, - for i = 1, ..., N, and j = 1, ..., N,

$$\begin{split} \frac{\partial \rho_i}{\partial x_{1,j}}(1,0) &= \int \partial_j Q \ x_i Q + o(1) = -2\delta_{i,j} \int Q^2 + o(1), \\ \frac{\partial \rho_i}{\partial \lambda_1}(1,0) &= \int \left(\frac{N}{2}Q + x.\nabla Q\right) x_i Q + o(1) = o(1), \\ \frac{\partial \rho_{N+1}}{\partial x_{1,j}}(1,0) &= \int \partial_j Q |x|^2 Q + o(1) = o(1), \\ \frac{\partial \rho_{N+1}}{\partial \lambda_1}(1,0) &= \int \left(\frac{N}{2}Q + x.\nabla Q\right) |x|^2 Q + o(1) \\ &= \frac{N}{2} \int |x|^2 Q^2 - \frac{N}{2} \int |x|^2 Q^2 - \frac{1}{2} \int x.x \ Q^2 + o(1) \\ &= -\frac{1}{2} \int |x|^2 Q^2 + o(1). \end{split}$$

Therefore, the implicit function theorem implies the existence of $(\lambda_{1,n}, x_{1,n})$ such that (31)-(33) hold.

In conclusion, we have proved the following. There exist $(\lambda_{1,n},x_{1,n}) \xrightarrow[n \to +\infty]{} (1,0)$ such that

(34)
$$w_{n,\lambda_{1,n},x_{1,n}}(x) = Q(x) + \varepsilon_{n,x_{1,n},x_{1,n}}(x)$$

where

(35)
$$\forall i = 1, ..., n, \quad \int \varepsilon_{n, \lambda_{1,n}, x_{1,n}} x_i Q = 0,$$

(36)
$$\int \varepsilon_{n,\lambda_{1,n},x_{1,n}} |x|^2 Q = 0,$$

(37)
$$\left|\varepsilon_{n,\lambda_{1,n},x_{1,n}}\right|_{H^1} \xrightarrow[n \to +\infty]{} 0.$$

We now note

$$w_n = w_{n,\lambda_{1,n},x_{1,n}},$$

$$\varepsilon_n = \varepsilon_{n,\lambda_{1,n},x_{1,n}}.$$

Step 3: Conclusion of the proof. – Geometry of energy functions at Q.

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We now use convexity properties of a functional (related to E) and the fact $\int w_n^2 = \int Q^2$ to conclude the proof. Let

$$H(v) = \frac{1}{2} \int |\nabla v|^2 - \frac{1}{\frac{4}{N} + 2} \int |v|^{\frac{4}{N} + 2} + \frac{1}{2} \int v^2 = E(v) + \frac{1}{2} \int v^2,$$

and

$$H_2(v) = \frac{1}{2} \int |\nabla v|^2 - \frac{4}{N} \int Q^{\frac{4}{N}} v^2 + \frac{1}{2} \int v^2.$$

We know that Q is a critical point of H, and H_2 is the quadratic part of H near Q (where v is real-valued). Moreover, it is classical that for $|\varepsilon|_{H^1} \leq 1$,

(38)
$$H(Q+\varepsilon) - H(Q) = H_2(\varepsilon) + \tilde{H}_2(\varepsilon)$$

where $\left| \tilde{H}_2(\varepsilon) \right| \, = \, o(|\varepsilon|_{H^1}^2).$

From a result of Weinstein [8] (see also Kwong [4]), we have the following convexity property of H_2 at Q.

Proposition 1. See [8]. – (Directions of convexity of H at Q in the set of real-valued functions.)

There is a constant $c_1 > 0$ such that $\forall \varepsilon \in H^1$ If

(i)
$$\forall i = 1, ..., N, \quad \int \varepsilon \ x_i Q = 0$$

(ii)
$$\int \varepsilon |x|^2 Q = 0,$$

(iii)
$$\int \varepsilon Q = 0,$$

then

$$H_2(\varepsilon) \ge c_1 \left(\int |\nabla \varepsilon|^2 + \varepsilon^2 \right) = c_1 |\varepsilon|_{H^1}^2.$$

Remark. – We have here a strict convexity property (up to the invariance of the equation) except in the direction Q which is not true for the quadratic part of H for complex valued functions (see [8]).

Remark. - This proposition is optimal. Other functions can be chosen also.

Using now crucially estimates on the L^2 norm, we obtain the following

Proposition 2 (Control of the Q direction by the L^2 norm). Assume

(i)
$$\forall i = 1, ..., N, \qquad \int \varepsilon x_i Q = 0,$$

(ii)
$$\int \varepsilon |x|^2 Q = 0,$$

(iii)
$$\int (Q+\varepsilon)^2 = \int Q^2,$$

then there are $c_1 > 0$ and $c_2 > 0$ such that

$$|\nabla \varepsilon|_{L^2} + |\varepsilon|_{L^2} \le c_2 \text{ implies } H_2(\varepsilon) \ge c_1 (|\nabla \varepsilon|_{L^2}^2 + |\varepsilon|_{L^2}^2).$$

Remark. – We need control on 3 directions to obtain estimates on $|\varepsilon|_{H^1}$ with $H_2(Q+\varepsilon)$. Two directions can be controlled using the invariance of the equation. The last one is controlled by the condition of minimality on the L^2 norm (among sequence satisfying (2)).

Proof of Proposition 2. - Let us note

$$\tilde{H}_2(v_1, v_2) = \frac{1}{2} \int \nabla v_1 \nabla v_2 - \frac{4}{N} \int Q^{\frac{4}{N}} v_1 v_2 + \frac{1}{2} \int v_1 v_2.$$

We can write

$$\varepsilon = z + aQ + b|x|^2Q$$

with

$$\int zQ = \int zx_iQ = \int z|x|^2Q = 0$$
 for $i = 1, ..., N$.

Indeed a and b have to satisfy

$$\int \varepsilon Q = a \int Q^2 + b \int |x|^2 Q^2$$
$$o = a \int |x|^2 Q^2 + b \int |x|^4 Q^2$$

or equivalently

$$\begin{split} b &= -a \bigg(\frac{\int |x|^2 Q^2}{\int |x|^4 Q^2} \bigg) \\ a \bigg(\frac{\int Q^2 \int |x|^4 Q^2 - \left(\int |x|^2 Q^2\right)^2}{\int |x|^4 Q^2} \bigg) &= \int \varepsilon Q \end{split}$$

(which has always a solution since from the Schwarz inequality and the fact $|x|^2Q \neq Q$, $\int |x|^2Q^2 < \left(\int Q^2 \int |x|^4Q^4\right)^{1/2}$). On the other hand, we have from $\int (Q+\varepsilon)^2 = \int Q^2$

$$\begin{split} &2\int Q\varepsilon = -\int \varepsilon^2 \\ &2\bigg(a\int Q^2 + b\int |x|^2Q^2\bigg) = -\int \varepsilon^2 \\ &2a\bigg(\frac{\int Q^2\int |x|^4Q^2 - \left(\int |x|^2Q^2\right)^2}{\int |x|^4Q^2}\bigg) = -\int z^2 + O(a^2 + b^2) \end{split}$$

or equivalently,

$$ac_0 = -\int z^2 + O(a^2)$$
 where $c_0 \neq 0$

which implies that

$$a = O\left(\int z^2\right)$$
 and $b = O\left(\int z^2\right)$

and for $|\varepsilon|_{H^1}$ small enough

$$|\varepsilon|_{H^1}^2 \ge |z|_{H^1}^2 \ge \frac{1}{2} |\varepsilon|_{H^1}^2.$$

On the other hand, by bilinearity and Proposition 1, we have for $|\varepsilon|_{H^1}$ small enough

$$\begin{split} H_2(\varepsilon) &= H_2(z) + 2a\tilde{H}_2(z,Q) + 2b\tilde{H}_2\big(z,|x|^2Q\big) + 2ab\tilde{H}_2\big(Q,|x|^2Q\big) \\ &+ a^2H_2(Q) + b^2H_2\big(|x|^2Q\big) \\ &\geq H_2(z) - c\big(|z|_{H^1}(|a| + |b|) + a^2 + b^2\big) \\ &\geq H_2(z) - c\big(|z|_{H^1}^3 + |z|_{H^1}^4\big) \\ &\geq c_1\big(|z|_{H^1}^2\big) - c\big(|z|_{H^1} + |z|_{H^1}^4\big) \\ &\geq \frac{c_1}{2}|z|_{H^1}^2 \\ &\geq \frac{c_1}{4}|\varepsilon|_{H^1}^2. \end{split}$$

This concludes the proof of Proposition 2.

As a corollary of Proposition 2 and (38), we have

COROLLARY. – There are $c_1 > 0$ and $c_2 > 0$ such that if

(i)
$$\forall i = 1, ..., N, \quad \int \varepsilon x_i Q = 0$$

(ii)
$$\int \varepsilon |x|^2 Q = 0$$

(iii)
$$\int (Q+\varepsilon)^2 = \int Q^2$$

(iv)
$$|\nabla \varepsilon|_{L^2} + |\varepsilon|_{L^2} \le c_2$$

then

$$H(Q + \varepsilon) - H(Q) \ge c_1 \left(\int \nabla \varepsilon^2 + \int \varepsilon^2 \right).$$

We now apply the corollary. If $w_n=Q+\varepsilon_n$, we have $-|\varepsilon_n|_{H^1}\xrightarrow[n\to+\infty]{}0$, and in particular there is n_0 such that

$$\forall n \geq n_0, |\varepsilon_n|_{H^1} \leq c_2,$$

- $\forall i = 1,...,N, \int \varepsilon_n x_i Q = 0.$ In addition,

$$\begin{split} H(Q) &= \frac{1}{2} \int |\nabla Q|^2 - \frac{1}{\frac{4}{N} + 2} \int |Q|^{\frac{4}{N} + 2} + \frac{1}{2} \int Q^2 \\ &= E(Q) + \frac{1}{2} \int Q^2 \\ &= \frac{1}{2} \int Q^2 \end{split}$$

(since the Pohozaev identity for equation (4) yields E(Q) = 0), and

$$\begin{split} H(Q+\varepsilon_n) &= H(w_n) = H\!\left(\!\left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^{\frac{N}{2}}\!|v_n|\!\left(\frac{x\lambda_1}{\hat{\lambda}_n} + \hat{x}_n + x_1\right)\right) \\ &= E\!\left(\!\left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^{\frac{N}{2}}\!|v_n|\!\left(\frac{x\lambda_1}{\hat{\lambda}_n}\right)\right) + \frac{1}{2}\int \left|v_n\right|^2 \\ &= \left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^2\!E(|v_n|) + \frac{1}{2}\int Q^2. \end{split}$$

Therefore $\forall n \geq n_0$

(40)
$$\left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^2 E(|v_n|) > c_3 \left(|\varepsilon_n|_{H^1}^2\right)$$

or equivalently from (19), (24) and the fact that $\lambda_1 \rightarrow 1$,

$$|\varepsilon_n|_{H^1}^2 \le \frac{c}{2} \frac{1}{\int |\nabla |v_n||^2} \le c \frac{1}{\int |\nabla v_n|^2},$$

where c is independent of n. Thus,

$$w_n = Q + \varepsilon_n$$

with

$$\left|\varepsilon_{n}\right|_{H^{1}}^{2} \leq \frac{c}{\int \left|\nabla v_{n}\right|^{2}}.$$

Therefore from (26), there is x_n such that

$$(43) |v_n|(x) = \left(\frac{\hat{\lambda}_n}{\lambda_1}\right)^{\frac{N}{2}} Q\left(x\left(\frac{\hat{\lambda}_n}{\lambda_1}\right) + x_n\right) + \left(\frac{\hat{\lambda}_n}{\lambda_1}\right)^{\frac{N}{2}} \varepsilon_n \left(\frac{\hat{\lambda}_n}{\lambda_1}x + x_n\right).$$

We remark that from (19), (42), the fact that $\lambda_1 \rightarrow 1$,

$$\frac{\hat{\lambda}_n}{\lambda_1} \frac{1}{\left(\int \nabla v_n^2\right)^{\frac{1}{2}}} = \frac{1}{\lambda_1} \left(\frac{\int |\nabla |v_n|^2}{\int |\nabla v_n|^2 \int |\nabla Q|^2} \right)^{\frac{1}{2}} \xrightarrow{n \to +\infty} 1,$$

$$\left| \left(\frac{\hat{\lambda}_n}{\lambda_1} \right)^{\frac{N}{2}} \varepsilon_n \left(\frac{\hat{\lambda}_n}{\lambda_1} x + x_n \right) \right|_{L^2}^2 = \left| \varepsilon_n \right|_{L^2}^2 \le \frac{c}{\int |\nabla v_n|^2},$$

$$\left| \nabla \left(\frac{\hat{\lambda}_n}{\lambda_1} \right)^{\frac{N}{2}} \varepsilon_n \left(\frac{\hat{\lambda}_n}{\lambda_1} x + x_n \right) \right|_{L^2}^2 \le c \left(\frac{\hat{\lambda}_n^2}{\int |\nabla v_n|^2} \right) \le c$$

conclude the proof of the Theorem.

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