

ANNALES DE L'I. H. P., SECTION A

M. A. PERELMUTER

Schrödinger operators with form-bounded potentials in L^p -spaces

Annales de l'I. H. P., section A, tome 52, n° 2 (1990), p. 151-161

http://www.numdam.org/item?id=AIHPA_1990__52_2_151_0

© Gauthier-Villars, 1990, tous droits réservés.

L'accès aux archives de la revue « Annales de l'I. H. P., section A » implique l'accord avec les conditions générales d'utilisation (<http://www.numdam.org/conditions>). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

NUMDAM

Article numérisé dans le cadre du programme
Numérisation de documents anciens mathématiques

<http://www.numdam.org/>

Schrödinger operators with form-bounded potentials in L^p -spaces

by

M. A. PERELMUTER

Urickij street, 11, apt. 153,
Kiev, 252035, U.S.S.R.

ABSTRACT. — Let

$$\begin{aligned} V &= V_+ - V_-, \quad 0 \leq V_{\pm} \in L^p_{\text{comp}}(\mathbb{R}^d), \\ \beta &< 1, \quad V_- \leq -\beta \Delta + C, \\ V_- &\in L^{k, \infty}_{\text{comp}}(\mathbb{R}^d), \quad k = (1 + \sqrt{1 - \beta}) \left/ \left(\frac{2}{p} - 1 + \sqrt{1 - \beta} \right) \right. \end{aligned}$$

($L^{p, \infty}$ is the weak- L^p -space). If $p \in \left(\frac{2}{1 + \sqrt{1 - \beta}}, \frac{2}{1 - \sqrt{1 - \beta}} \right)$, then the closure of the operator $(-\Delta + V + C) \upharpoonright C_0^\infty(\mathbb{R}^d)$ is generator of the contraction C_0 -semigroup in $L^p(\mathbb{R}^d)$. For $p=2$ we obtain a new theorem concerning the essential self-adjointness of the Schrödinger operator. The sharpness of the results is illustrated by the examples.

RÉSUMÉ. — $V = V_+ - V_-$, $0 \leq V_{\pm} \in L^p_{\text{comp}}(\mathbb{R}^d)$,
 $\beta < 1$, $V_- \leq -\beta \Delta + C$,

$$V_- \in L^{k, \infty}_{\text{comp}}(\mathbb{R}^d), \quad k = (1 + \sqrt{1 - \beta}) \left/ \left(\frac{2}{p} - 1 + \sqrt{1 - \beta} \right) \right.$$

($L^{p, \infty}$ est l'espace L^p faible). Si $p \in \left(\frac{2}{1 + \sqrt{1 - \beta}}, \frac{2}{1 - \sqrt{1 - \beta}} \right)$, la fermeture de l'opérateur $(-\Delta + V + C) \upharpoonright C_0^\infty(\mathbb{R}^d)$ est un générateur de semigrúpe C_0 de contraction dans $L^p(\mathbb{R}^d)$. Pour $p=2$ nous obtenons un

nouveau théorème concernant l'auto-adjonction essentielle de l'opérateur de Schrödinger. Ces résultats sont illustrés sur des exemples.

INTRODUCTION

This article deals with the Schrödinger operators $-\Delta + V$ acting in $L^p \equiv L^p(\mathbb{R}^l, dx)$. We want to investigate the first spectral problem, *i. e.* the problem of m -accretivity. Most of the results concerning this problem contain the condition of $-\Delta$ -boundedness of the operator $V_- \equiv \max\{-V, 0\}$. Our goal is to investigate the conditions of the following type:

$$\left. \begin{aligned} V_- &\leq -\beta\Delta + C(\beta), & \beta < 1, \\ V_- &\in F, \end{aligned} \right\} \quad (1)$$

where F is some function space.

The first result in this direction is due to Povzner [1] [$p=2$, $\beta=1$, $F=C(\mathbb{R}^l)$]. Recently some results were also obtained. In particular the essential self-adjointness was proved provided that $F=L^k$, $k > \frac{2l}{l-\beta(l-2)}$ [2]; and $F=L^k$, $k=(1+\sqrt{1-\beta})/\sqrt{1-\beta}$ [3].

My purpose here is to extend the abovementioned results to the case of L^p -spaces and to choose the broader space F . Even for $p=2$ we obtain a new result which is very close to optimal.

It should be noted that our method is an extension of the Semenov's one [3].

Note that L^p - Δ -boundedness implies the condition (1) (*see* [4]) and in my opinion the condition (1) corresponds to the point of our problem. We want also to emphasize that the exact eigenfunction estimates for the operator $-\Delta + V$ and the exact results concerning the accretivity of the Schrödinger operators have been proved under the same condition (1) (*see* [5], [6]).

Some common symbols are:

$$\begin{aligned} \|f\|_p &= \left(\int_{\mathbb{R}^l} |f(x)|^p d^l x \right)^{1/p}; \\ \langle f, g \rangle &= \int_{\mathbb{R}^l} f(x) \overline{g(x)} d^l x; \\ (f * g)(x) &= \int_{\mathbb{R}^l} f(x-y) g(y) d^l y; \end{aligned}$$

$L^p_{loc} \equiv L^p_{loc}(\mathbb{R}^l, d^l x)$ = the complex functions f on \mathbb{R}^l such that $f|_{\varphi} \in L^p, \forall \varphi \in C^\infty_0$; $\text{meas}(\mathbf{B})$ = Lebesgue measure of $\mathbf{B} \subset \mathbb{R}^l$; $L^{p,\infty} \equiv L^{p,\infty}(\mathbb{R}^l, d^l x)$ = the complex functions f on \mathbb{R}^l such that

$$\sup_{t>0} t^p \text{meas} \{x: |f(x)| > t\} < \infty$$

(note that $L^p \subset L^{p,\infty}$ and $|x|^{-l/p} \in L^{p,\infty}$). $L^{p,\infty}_{comp}$ = the set of functions in $L^{p,\infty}$ with compact support. $\Delta = \sum_{j=1}^l \frac{\partial^2}{\partial x_j^2}$ the Laplacian acting on $\mathcal{D}'(\mathbb{R}^l)$.

MAIN RESULTS

LEMMA 1. - If $V \in L^{p,\infty}_{comp}$, then $\|V\|_{p-\varepsilon} = O(\varepsilon^{-1/p}), \varepsilon \downarrow 0$.

Proof. - Let $\mu(t) = \text{meas} \{x: |V(x)| > t\}$. The function $\mu(t)$ is nonincreasing and $\mu(0) = \text{meas} \{ \text{supp } V \} < \infty$. The definition of the space $L^{p,\infty}$ implies that $\mu(t) \leq C t^{-p}, \forall t > 0$, so

$$\begin{aligned} \|V\|_{p-\varepsilon}^{p-\varepsilon} &\equiv (p-\varepsilon) \int_0^\infty t^{p-\varepsilon-1} \mu(t) dt \\ &= (p-\varepsilon) \left(\int_0^1 t^{p-\varepsilon-1} \mu(t) dt + \int_1^\infty t^{p-\varepsilon-1} \mu(t) dt \right) \\ &\leq (p-\varepsilon) \left(\mu(0) \int_0^1 t^{p-\varepsilon-1} dt + C \int_1^\infty t^{-1-\varepsilon} dt \right) \\ &= \mu(0) - C + \frac{Cp}{\varepsilon}. \quad \blacksquare \end{aligned}$$

DEFINITION. - Let $0 \leq V, W \in L^1_{loc}$. Then $V \in PK_\beta(-\Delta + W)$ if and only if

$$\langle V\varphi, \varphi \rangle \leq \beta (\|\bar{V}\varphi\|_2^2 + \|W^{1/2}\varphi\|_2^2) + C(\beta) \|\varphi\|_2^2, \forall \varphi \in C^\infty_0.$$

THEOREM 2. - Let V be a real valued measurable function such that

$$\begin{aligned} V &= V_+ - V_-, \quad 0 \leq V_+ \in L^p_{\text{comp}}, \\ 0 &\leq V_- \in PK_\beta(-\Delta + V_+) \cap L^{\lambda, \infty}, \\ \beta &< 1, \quad p \in \left(1, \frac{2}{1 - \sqrt{1 - \beta}}\right), \\ k &= (1 + \sqrt{1 - \beta}) / \left(\frac{2}{p} - 1 + \sqrt{1 - \beta}\right). \end{aligned}$$

Then the range of the operator $(-\Delta + V + \lambda) \upharpoonright C_0^\infty$ is dense in L^p for any $\lambda > C(\beta)$.

Proof. - If $p \in \left(1, \frac{2}{1 - \sqrt{1 - \beta}}\right)$, then $1 < p < k < \infty$.

Suppose that the range is not dense, then by the Hahn-Banach theorem, there exists a function $0 \neq u \in L^p$ such that

$$\langle (-\Delta + V + \lambda)\varphi, u \rangle = 0, \quad \forall \varphi \in C_0^\infty. \quad (3)$$

Without loss we can assume $u = \text{Re } u$. Let $1/t = 1/p' + 1/k'$, $q = p'/k'$. Let $\varepsilon > 0$ be a small parameter and let us set $1/s = 1/t + \varepsilon/p'$, $1/s' = (q - 1 - \varepsilon)/p'$. Since $V_- \in L^{\lambda, \infty}_{\text{comp}}$ and $u \in L^p$. It follows that $V_- u \in L^1 \cap L^s$.

The equality (3) gives

$$(-\Delta + V_+ + \lambda)u = V_- u \in L^1 \cap L^s \quad (4)$$

in the distributional sense; hence

$$u \in L^1 \cap L^s \cap L^{p'}; \quad (5)$$

(cf. [7]).

Moreover, $u \in L^{p'}$ implies

$$|u|^{q-1} \in L^t. \quad (6)$$

Let us consider the operator H_z^+ with the domain

$$\mathcal{D}(H_z^+) = \{u \in L^2 : V_+ u \in L^1_{\text{loc}}, (-\Delta + V_+)u \in L^2\},$$

where $(-\Delta + V_+)$ is taken in the distributional sense and $H_z^+ u \equiv (-\Delta + V_+)u$, $\forall u \in \mathcal{D}(H_z^+)$.

H_z^+ is the generator of the strongly continuous contraction semigroup $T_z(t) \equiv \exp(-tH_z^+)$, $\forall t > 0$ in L^2 which has properties such as follows below

$$\left. \begin{aligned} \|T_z(t)\varphi\|_w &\leq C(t, w, y, l) \|\varphi\|_y, \\ \forall \varphi \in C_0^\infty, \quad \forall t > 0, \quad \forall w \geq y \geq l \end{aligned} \right\} \quad (7)$$

$$T_z(t) \upharpoonright [L^z \cap L^w] = T_w(t) \upharpoonright [L^z \cap L^w] \quad (8)$$

(cf. [7]).

Using the standard diagonal process we can choose the subsequence $u_r = \varphi_{d, m_r}$ such that

$$u_r \xrightarrow{L^y} u, \quad \forall y \in [1, p'], \quad (9)$$

$$(-\Delta + V_+) u_r \rightarrow (-\Delta + V_+) u, \quad (10)$$

$$u_r |u_r|^{q-\varepsilon-2} \xrightarrow{L^s} u |u|^{q-\varepsilon-2}, \quad (11)$$

$$|u_r|^{q-\varepsilon} \xrightarrow{L^y} |u|^{q-\varepsilon}, \quad \forall y \leq k'. \quad (12)$$

Observe that $u_r \in C_0^\infty$, so that $u_r |u_r|^{q-\varepsilon-2}, |u_r|^{(q-\varepsilon)/2} \in \mathcal{D}(\vec{\nabla})$ (see, e. g., [10], Chapter 2), and the direct calculation yields the identity

$$\langle u_r |u_r|^{q-\varepsilon-2}, -\Delta u_r \rangle = \frac{4(q-\varepsilon-1)}{(q-\varepsilon)^2} \|\vec{\nabla} |u_r|^{(q-\varepsilon)/2}\|_2^2. \quad (13)$$

(4)-(6) imply the finiteness of the quantities

$$\begin{aligned} &\langle u |u|^{q-\varepsilon-2}, (-\Delta + V_+) u \rangle; \quad \langle u |u|^{q-\varepsilon-2}, V_- u \rangle; \\ &\langle u |u|^{q-\varepsilon-2}, \lambda u \rangle, \end{aligned}$$

so that the simple approximation arguments show that (3) is valid provided that $\varphi = u_r |u_r|^{q-\varepsilon-2}$ and (9)-(12) imply

$$\begin{aligned} &\lim_{r \rightarrow \infty} \langle u_r |u_r|^{q-\varepsilon-2}, (-\Delta + V + \lambda) u_r \rangle \\ &= \langle u |u|^{q-\varepsilon-2}, (-\Delta + V + \lambda) u \rangle = 0. \quad (14) \end{aligned}$$

Using (13), (14) and the condition $V_- \in \text{PK}_\beta(-\Delta + V_+)$ we have

$$\begin{aligned} &\lambda \|u\|_{q-\varepsilon}^{q-\varepsilon} + \langle u |u|^{q-\varepsilon-2}, (-\Delta + V_+) u \rangle \\ &= \langle u |u|^{q-\varepsilon-2}, V_- u \rangle = \lim_{r \rightarrow \infty} \langle u_r |u_r|^{q-\varepsilon-2}, V_- u_r \rangle \\ &= \lim_{r \rightarrow \infty} \langle V_- |u_r|^{(q-\varepsilon)/2}, |u_r|^{(q-\varepsilon)/2} \rangle \\ &\leq \lim_{r \rightarrow \infty} [\beta \|\vec{\nabla} |u_r|^{(q-\varepsilon)/2}\|_2^2 + \beta \\ &\quad \times \|V_+^{1/2} |u_r|^{(q-\varepsilon)/2}\|_2^2 + C(\beta) \| |u_r|^{(q-\varepsilon)/2} \|_2^2] \\ &= \lim_{r \rightarrow \infty} \left[\frac{\beta(q-\varepsilon)^2}{4(q-\varepsilon-1)} \langle u_r |u_r|^{q-\varepsilon-2}, -\Delta u_r \rangle \right. \\ &\quad \left. + \beta \langle u_r |u_r|^{q-\varepsilon-2}, V_+ u_r \rangle + C(\beta) \|u_r\|_{q-\varepsilon}^{q-\varepsilon} \right] \equiv \lim_{r \rightarrow \infty} I_r. \quad (15) \end{aligned}$$

Note that $\beta < \beta(q - \varepsilon)^2 / 4(q - \varepsilon - 1)$. Therefore,

$$\begin{aligned} \lim_{r \rightarrow \infty} I_r &\leq \frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} \lim_{r \rightarrow \infty} \langle u_r | u_r \rangle^{q - \varepsilon - 2}, (-\Delta + V_+) u_r \rangle \\ &\quad + C(\beta) \|u\|_{q - \varepsilon}^{q - \varepsilon} = \frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} \\ &\quad \times \langle u | u \rangle^{q - \varepsilon - 2}, (-\Delta + V_+) u \rangle + C(\beta) \|u\|_{q - \varepsilon}^{q - \varepsilon}. \end{aligned} \quad (16)$$

According to (15), (16),

$$(\lambda - C(\beta)) \|u\|_{q - \varepsilon}^{q - \varepsilon} \leq \left[\frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} - 1 \right] \langle u | u \rangle^{q - \varepsilon - 2}, (-\Delta + V_+) u \rangle.$$

From (14) we have that the right side of the last inequality equals

$$\begin{aligned} \left[\frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} - 1 \right] \langle u | u \rangle^{q - \varepsilon - 2}, (V_- - \lambda) u \rangle \\ \leq \left[\frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} - 1 \right] \langle u | u \rangle^{q - \varepsilon - 2}, V_- u \rangle. \end{aligned}$$

So, using Hölder inequality we obtain

$$\begin{aligned} (\lambda - C(\beta)) \|u\|_{q - \varepsilon}^{q - \varepsilon} &\leq \left[\frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} - 1 \right] \langle |u|^{q - \varepsilon}, V_- \rangle \\ &\leq \left[\frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} - 1 \right] \|u\|_{p'}^{q - \varepsilon} \|V_-\|_{z(\varepsilon)}, \end{aligned} \quad (17)$$

where $z(\varepsilon) = k / (1 + \varepsilon k / p')$.

According to the lemma $1 \|V_-\|_{z(\varepsilon)} = O(\varepsilon^{-1/k})$. On the other hand

$$\frac{\beta(q - \varepsilon)^2}{4(q - \varepsilon - 1)} - 1 = O(\varepsilon).$$

Now assuming that $\varepsilon \downarrow 0$ in (17) and using $k > 1$ we obtain

$$(\lambda - C(\beta)) \|u\|_q^q \leq 0.$$

Then $u \equiv 0$ as $\lambda > C(\beta)$. ■

Remarks. - 1. The condition $V_- \in PK_\beta(-\Delta + V_+) \cap L_{loc}^p$ is necessary in the theorem 2 but it is not sufficient as it may be seen in the example

$-\Delta - \frac{\kappa}{|x|^2}$ (see [11], Chapter X). The same example shows the sharpness of the theorem 2. Indeed, let $2p < m \leq l$,

$$V(x) = \kappa \chi(x) \left(\sum_{j=1}^m x_j^2 \right)^{-1},$$

where $\chi(\cdot)$ is the indicator of the unit ball in \mathbb{R}^l . Then the range of the operator $(-\Delta - V - \lambda) C_0^\infty$ is dense in L^p if and only if

$\kappa \leq m(m-2p)(p-1)/p^2$ [4]. On the other hand, $V \in L_{\text{comp}}^{m/2, \infty}$, i. e., $\kappa = m(m-2p)(p-1)/p^2$ implies $k = (1 + \sqrt{1-\beta}) \left(\frac{2}{p} - 1 + \sqrt{1-\beta} \right)$.

2. The condition $V_- \in \text{PK}_\beta(-\Delta + V_+)$ may be replaced by $V_- \in \text{PK}_\beta \left(-\Delta + \frac{p^2}{4(p-1)} V_+ \right)$.

3. The analysis of the proof shows that the restriction $V_- \in L_{\text{comp}}^{k, \infty}$ may be replaced by the weaker condition $\|V_-\|_{k-\varepsilon} = o(\varepsilon)$.

4. The main device in the proof of the theorem 2 is the inequality $\langle -\Delta u, u |u|^{p-2} \rangle \geq 4 \frac{p-1}{p^2} \|\bar{V} |u|^{p/2}\|_2^2$.

N. Th. Varopoulos [12] has proved the abstract version of this inequality for the generators of submarkovian semigroups. Thus our result may be generalized for these operators.

COROLLARY 3. — *Let*

$$\beta < 1, \quad V = V_+ - V_-, \quad 0 \leq V_\pm \in L_{\text{comp}}^2, \\ V_- \in \text{PK}_\beta(-\Delta + V_+) \cap L^{k, \infty}, \quad k = (1 + \sqrt{1-\beta}) / \sqrt{1-\beta}.$$

Then $(-\Delta + V) \upharpoonright C_0^\infty$ is essentially self-adjoint operator in L^2 .

Corollary 3 is proved under the restriction $V \in L_{\text{comp}}^2$, but using the results of C. Simader [13] or H. Brezis [14], we have

THEOREM 4. — *Let*

$$\beta < 1, \quad V = V_+ - V_-, \quad 0 \leq V_\pm \in L_{\text{loc}}^2, \\ V_- \in \text{PK}_\beta(-\Delta + V_+) \cap L_{\text{loc}}^{k, \infty}, \quad K = (1 + \sqrt{1-\beta}) / \sqrt{1-\beta}.$$

Then $(-\Delta + V) \upharpoonright C_0^\infty$ is essentially self-adjoint.

Remark. — In the case $l \geq 5$ theorem 4 is a generalization of the Kalf-Walter-Schmincke-Simon theorem (see [11], § X.4).

Example. — We consider the N-particle Hamiltonian

$$H = H_0 + V, \\ H_0 = -\frac{1}{2} \sum_{i=1}^N \Delta_i + \frac{1}{2N} \left(\sum_{i=1}^N \bar{V}_i \right)^2, \\ -\alpha \sum_{i < j} |x_i - x_j|^{-2} \leq V \in L_{\text{loc}}^2(\mathbb{R}^{l(N-1)}), \\ x_j \in \mathbb{R}^l, \quad l \geq 5.$$

Hardy inequality implies $V_- \in PK_\beta(H_0)$, $\beta = \frac{2N\alpha}{(l-2)^2}$. One sees easily that $V_- \in L_{loc}^{1/2, \infty}$. If $\alpha \leq \frac{(l-4)}{2N}$ it follows from theorem 4 that $H \upharpoonright C_0^\infty$ is essentially self-adjoint (cf. [15], [16]). Note that for the relativistic Hamiltonian

$$H_R + V_C \equiv \sum_{i=1}^N (-\Delta_i + m^2)^{1/2} - \alpha \sum_{i < j} |x_i - x_j|^{-1}, \quad l=3.$$

Lieb and Thirring [17] have proved that $V_C \in PK_\beta(H_R)$ with $\beta = O(N\alpha)$ and $H_R + V_C$ is unbounded from below for $\alpha = C/N$, $C < \infty$. This correlation between β and N appears to be true for our nonrelativistic Hamiltonian as well so that the dependence $\alpha = O(1/N)$ is optimal.

To prove m -accretivity we need.

THEOREM 5 [6]. — Let

$$\begin{aligned} V &= V_+ - V_-, \quad 0 \leq V_\pm \in L_{loc}^p, \\ V_- &\in PK_\beta(-\Delta + V_+), \quad \beta < 1. \end{aligned}$$

If $p \in \left[\frac{2}{1 + \sqrt{1-\beta}}, \frac{2}{1 - \sqrt{1-\beta}} \right]$, then the operator $(-\Delta + V + C(\beta)) \upharpoonright C_0^\infty$ is accretive operator in L^p .

Proof. — We will show that

$$\langle (-\Delta + V + C(\beta))u, u \rangle \geq 0, \quad \forall u \in C_0^\infty, \tag{18}$$

where $[v, w]$ is semi-inner product (see [11], § X.8). In the case of L^p spaces

$$[v, w] = \langle v, w | w|^{p-2} \rangle / \|w\|^{p-2},$$

so that (18) may be written in the following way

$$\langle V_- u, u | u|^{p-2} \rangle \leq \langle (-\Delta + V_+ + C(\beta))u, u | u|^{p-2} \rangle.$$

But $V_- \in PK_\beta(-\Delta + V_+)$ so that

$$\begin{aligned} \langle V_- u, u | u|^{p-2} \rangle &= \langle V_- |u|^{p/2}, |u|^{p/2} \rangle \\ &\leq \beta \| \vec{\nabla} |u|^{p/2} \|_2^2 + \beta \langle V_+ |u|^{p/2}, |u|^{p/2} \rangle + C(\beta) \| |u|^{p/2} \|_2^2 \\ &= \frac{\beta p^2}{4(p-1)} \langle -\Delta u, u | u|^{p-2} \rangle \\ &\quad + \beta \langle V_+ u, u | u|^{p-2} \rangle + C(\beta) \langle u, u | u|^{p-2} \rangle. \end{aligned}$$

Since $\frac{\beta p^2}{4(p-1)} \leq 1$ provided that $p \in \left[\frac{2}{1 + \sqrt{1-\beta}}, \frac{2}{1 - \sqrt{1-\beta}} \right]$ we have proved (18). ■

Remark. — 1. The sharpness of the result is shown in [6] by the example $-\Delta - \frac{\kappa}{|x|^2}$.

2. T. Kato [7] raised the problem of quasi-accretivity of the operator $-\Delta + V$ in the limiting case $p = 1$ for the potentials such that

$$\lim_{\alpha \downarrow 0} \sup_x \int_{|x-y| \leq \alpha} |x-y|^{-l+2} |V(y)| d^l y = 0, \quad l \geq 3.$$

We will show that the answer is negative, *i. e.*, if $0 \leq V \in L^1_{loc}$ and if $-\Delta - V$ is quasi-accretive in L^1 , then $V \in L^\infty$. Indeed, the semi-inner product in L^1 is $[v, w] = \left\langle v, \frac{w}{|w|} \right\rangle \|w\|_1$. The quasi-accretivity implies that

$$\left\langle (-\Delta - V)u, \frac{u}{|u|} \right\rangle \geq -C \|u\|_1, \\ C < \infty, \quad \forall u \in C^\infty_0.$$

Note that $\langle \Delta u, 1 \rangle = 0, \forall u \in C^\infty_0$. Therefore, $\langle V, u \rangle \leq C \|u\|_1, i. e. V \in L^\infty$.

3. If $0 \leq V \notin L^\infty, V \in L^p, p > l/2, l \geq 3$, then $-\Delta - V$ is the generator of the positivity preserving C_0 -semigroup in L^1 which is not quasi-contractive (other examples may be found in [18], [6]).

Theorems 2, 5 and Lumer-Phillips theorem (*see* [11], § X.8) imply

THEOREM 6. — *Let*

$$\beta < 1, \quad p \in \left(\frac{2}{1 + \sqrt{1-\beta}}, \frac{2}{1 - \sqrt{1-\beta}} \right), \\ V = V_+ - V_-, \quad 0 \leq V_\pm \in L^p_{comp}, \\ V_- \in PK_\beta(-\Delta + V_+) \cap L^{k, \infty}$$

where $k = (1 + \sqrt{1-\beta}) \left(\frac{2}{p} - 1 + \sqrt{1-\beta} \right)$. Then the closure of the operator $(-\Delta + V + C(\beta)) \upharpoonright C^\infty_0$ is the generator of C_0 -semigroup of contractions in L^p .

Remarks. — 1. Theorem 6 is generalization (when $\text{supp } V$ is compact) of the result stated in [4]. It is of interest to prove the L^p -version of the result of C. Simader [13] and H. Brezis [14].

2. Having completed this paper the author was informed by Yu. A. Semenov about the proof of the theorem 6 provided $0 \leq V_\pm \in L^p_{loc}$,

$$p \in \left(\frac{2}{1 + \sqrt{1-\beta}}, \frac{2}{1 - \sqrt{1-\beta}} \right).$$

REFERENCES

- [1] A. Ya. POVZNER, On the Expansions of Arbitrary Functions in Terms of Eigenfunctions of the Operator $-\Delta u + Cu$, *Mat. Sbornik*, Vol. **32**, No. 1, 1953, pp. 109-156; *A.M.S. Trans. Ser. 2*, Vol. **60**, 1967, pp. 1-49.
- [2] M. A. PERELMUTER, Positivity Preserving Operators and One Criterion of Essential Self-Adjointness, *J. Math. Anal. Appl.*, Vol. **82**, No. 2, 1981, pp. 406-419.
- [3] Yu. A. SEMENOV, *One Criterion of Essential Self-Adjointness of the Schrödinger Operator with Negative Potential*, Kiev, preprint, 1987.
- [4] V. F. KOVALENKO, M. A. PERELMUTER, Yu. A. SEMENOV, Schrödinger Operators with $L_w^{1/2}(\mathbb{R}^d)$ -Potentials, *J. Math. Phys.*, Vol. **22**, No. 5, 1981, pp. 1033-1044.
- [5] Yu. A. SEMENOV, On the Spectral Theory of Second-Order Elliptic Differential Operators, *Math. U.S.S.R. Sbornik*, Vol. **56**, No. 1, 1987, pp. 221-247.
- [6] V. F. KOVALENKO and Yu. A. SEMENOV, L^p -Contractivity of the Semigroup Generated by the Schrödinger Operator with Singular Negative Potential, Kiev, preprint, 1985.
- [7] T. KATO, L^p -theory of Schrödinger Operators with a Singular Potential, in *Aspects of Positivity Funct. Anal. Proc. Conf.*, Tübingen, 24-28 June 1985, Amsterdam e. a., 1986, pp. 41-48.
- [8] W. P. NOVINGER, Mean Convergence in L^p Spaces, *Proc. Am. Math. Soc.*, Vol. **34**, No. 2, 1972, pp. 627-628.
- [9] R. A. ADAMS, *Sobolev Spaces*, New York e. a., Academic Press, 1975.
- [10] D. KINDERLEHRER and G. STAMPACCHIA, *An Introduction to Variational Inequalities and Their Applications*, New York e. a., Academic Press, 1980.
- [11] M. REED and B. SIMON, *Methods of Modern Mathematical Physics*, New York, San Francisco, London, Academic Press, 1978.
- [12] N. Th. VAROPOULOS, Hardy-Littlewood Theory for Semi-Groups, *J. Funct. Anal.*, Vol. **63**, No. 2, 1985, pp. 240-260.
- [13] C. G. SIMADER, Essential Self-Adjointness of Schrödinger Operators Bounded from Below, *Math. Z.*, Vol. **159**, No. 1, 1978, pp. 47-50.
- [14] H. BREZIS, "Localized" Self-Adjointness of Schrödinger Operators, *J. Oper. Theor.*, Vol. **1**, No. 2, 1979, pp. 287-290.
- [15] M. COMBESURE-MOULIN and J. GINIBRE, Essential Self-Adjointness of Many-Particle Schrödinger Hamiltonians with Singular Twobody Potentials, *Ann. Inst. H. Poincaré*, Vol. **A23**, No. 3, 1975, p. 211-234.
- [16] V. F. KOVALENKO and Yu. A. SEMENOV, Essential Self-Adjointness of Many-Particle Hamiltonian Operators of Schrödinger Type with Singular Two-Particle Potentials, *Ann. Inst. H. Poincaré*, Vol. **A24**, No. 4, 1977, pp. 325-334.
- [17] E. H. LIEB and W. E. THIRRING, Gravitational Collapse in Quantum Mechanics with Relativistic Kinetic Energy, *Ann. Phys. (U.S.A.)*, Vol. **155**, No. 2, 1984, pp. 494-512.
- [18] C. J. K. BATTY and E. B. DAVIES, Positive Semi-Groups and Resolvents, *J. Oper. Theor.*, Vol. **10**, No. 2, 1983, pp. 357-363.

(Manuscript received March 23, 1989.)