



ANNALES
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MORPHISMS OF DOUBLE
(QUASI-)POISSON ALGEBRAS
AND ACTION-ANGLE DUALITY
OF INTEGRABLE SYSTEMS

MORPHISMES D'ALGÈBRES DE
(QUASI-)POISSON DOUBLES
ET DUALITÉ ACTION-ANGLE
DE SYSTÈMES INTÉGRABLES

ABSTRACT. — Double (quasi-)Poisson algebras were introduced by Van den Bergh as non-commutative analogues of algebras endowed with a (quasi-)Poisson bracket. In this work, we provide a study of morphisms of double (quasi-)Poisson algebras, which we relate to the H_0 -Poisson structures of Crawley–Boevey. We prove in particular that the double (quasi-)Poisson algebra structure defined by Van den Bergh for an arbitrary quiver only depends upon the quiver seen as an undirected graph, up to isomorphism. We derive from our results a representation theoretic description of action-angle duality for several classical integrable systems.

Keywords: Double bracket, Hamiltonian algebra, Quasi-Hamiltonian algebra, Non-commutative geometry, Integrable system.

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RÉSUMÉ. — Les algèbres de (quasi-)Poisson doubles ont été introduites par Van den Bergh comme des analogues non-commutatifs d’algèbres munies d’un crochet de (quasi-)Poisson. Dans ce travail, nous produisons une étude des morphismes d’algèbres de (quasi-)Poisson doubles, que nous relient aux structures H_0 -Poisson de Crawley–Boevey. Nous prouvons en particulier que la structure d’algèbre de (quasi-)Poisson double définie par Van den Bergh pour un carquois arbitraire dépend seulement du carquois vu comme un graphe non-dirigé, à isomorphisme près. Nous dérivons de nos résultats une description issue de la théorie des représentations de la dualité action-angle pour plusieurs systèmes intégrables classiques.

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1. Introduction

In the seminal paper [VdB08a], Van den Bergh laid the foundation of a noncommutative theory of Poisson geometry based on double brackets. Starting with an algebra A over a field of characteristic zero \mathbb{k} , he introduced the concept of a double bracket as an operation $A \times A \rightarrow A \otimes A$ satisfying noncommutative rules of derivation and antisymmetry; this is done in such a way that the corresponding affine scheme of representations $\text{Rep}(A, n)$ carries an antisymmetric biderivation for each $n \geq 1$. An interesting class of those structures are double Poisson brackets, which satisfy a version of Jacobi identity valued in $A^{\otimes 3}$ so that $\text{Rep}(A, n)$ is now endowed with a Poisson bracket, in agreement with the “non-commutative principle” of Kontsevich and Rosenberg [Kon93, KR00]. In fact, pushing this principle even further, Van den Bergh pointed out that the existence of a distinguished element on A permits to understand the process of Hamiltonian reduction of $\text{Rep}(A, n)$ with respect to the natural action of $\text{GL}_n(\mathbb{k})$ directly at the level of the algebra A . Due to this representation theoretic perspective, it seems interesting to examine the algebras endowed with a double Poisson bracket, which are called *double Poisson algebras*. They have been the object of several studies [Bie13, ORS13, ORS14, Pow16, PVdW08, VdW08], and our aim is to explore morphisms between double Poisson algebras. Noting that Van den Bergh also introduced the analogous notion of double quasi-Poisson brackets [VdB08a, VdB08b], it is natural to extend our investigation to morphisms between the corresponding algebras, called *double quasi-Poisson algebras*, which currently attract attention [Art18, CF17, CF20, Fai19a, Fai19b, Fai21, FH21, MT14].

A rich family of double (quasi-)Poisson algebras is associated with quivers [VdB08a] and encodes the Poisson geometry of quiver varieties [Nak94] as well as their multiplicative analogues [CBS06]. These varieties possess ramifications in numerous branches of mathematics (see e.g. [Gin12, Kir16, Sch08] for quiver varieties), in particular in the field of integrable systems. Indeed, in the complex setting it is known since 1998 and the work of Wilson [Wil98] that the phase space of the Calogero–Moser system can be realised as a quiver variety. Moreover, similar constructions of phase spaces as quiver varieties have been obtained for other related systems [BP11, CS17, Tac15], and they have been extended to Ruijsenaars–Schneider systems using multiplicative quiver varieties [CF17, CF20, Fai19a, Fai19b]. Some of these papers used the formalism of double brackets to grasp features of integrability directly on the path algebra of the relevant quivers. Thus, it raises the question of determining if additional properties of these integrable systems can be realised at the level of the quivers. This issue constitutes the main motivation behind our work, as we want to present a novel interpretation of action-angle duality, which we explain now.

Let M and M' be two manifolds of dimension $2n$ endowed with non-degenerate Poisson brackets such that the n functions H, H' on M, M' define (Liouville) integrable systems, i.e. H and H' form sets of n functionally independent Poisson-commuting elements. Let us furthermore assume that on a dense open subset of each manifold, there exist canonical Darboux coordinates (q, p) or (q', p') , and that there exists a Poisson diffeomorphism $\Psi : M \rightarrow M'$ such that (after restriction to dense subspaces) $H' \circ \Psi$ only depends on the coordinates q , while $H \circ \Psi^{-1}$ only depends on the coordinates q' . Due to the assumptions, we get action-angle variables⁽¹⁾ as follows. The coordinates q' become the action coordinates of H while the coordinates p' become the angle coordinates of H , and the same is true for q, p and H' . We thus say that the pairs (M, H) and (M', H') are *action-angle duals*. This construction has been largely investigated for systems of Calogero–Moser and Ruijsenaars–Schneider type following the pioneering work of Ruijsenaars [Rui88]. A widespread method used to unearth action-angle duality is the existence of two different slices in a suitable orbit space defined by Hamiltonian reduction, such that each slice provides one of the two sets of Darboux coordinates. This approach is inspired by the work of Kazhdan, Kostant and Sternberg [KKS78], and it has been considered both in the complex and the real settings, see [FM19, FGNR00, Pus12, Res16] and references therein.

One of the aims of this work is to obtain a different point of view on duality where representation theoretic considerations yield the action-angle map. Indeed, in view of the previous paragraphs, many classical integrable systems for which action-angle duality is known are defined on (multiplicative) quiver varieties. This leads us to the following natural question⁽²⁾:

⁽¹⁾By action-angle variables, we mean the existence of a set of $2n$ Darboux coordinates on a dense subspace of the manifold such that the integrable system only depends on the first n coordinates, the *action variables*. This ensures that the Hamiltonian flows are linearised with respect to the last n coordinates, the *angle variables*.

⁽²⁾This formulation is close to the original question posed to the author by V. Rubtsov at the conference *Geometric aspects of momentum maps and integrability* in Ascona, April 2018.

Is it possible to understand action-angle duality in terms of relevant quivers?

We end this work by answering this question positively in Section 6, as we note that action-angle duality can be realised at the level of quivers simply as a map “reversing arrows”. This is, for example, the case for the well-known duality between the hyperbolic Calogero–Moser system and the rational Ruijsenaars–Schneider system (see § 6.2), or the self-duality of the hyperbolic Ruijsenaars–Schneider system (see § 6.4). Note that this simple point of view also allows to derive new instances of action-angle duality for a generalisation of the rational Calogero–Moser system due to Chalykh and Silantyev [CS17] (see § 6.1), or a modification of the hyperbolic Ruijsenaars–Schneider system obtained jointly with Chalykh in [CF17] (see § 6.3).

The main tool needed to provide a precise construction of the action-angle duality maps is a study of morphisms of double (quasi-)Poisson algebras, which forms the core of this text. We will deduce from this formalism the following important result.

THEOREM 1.1 (See Theorems 3.12 and 4.12). — *Up to isomorphism, the double (quasi-)Poisson algebra associated with a quiver by Van den Bergh [VdB08a] only depends on the underlying quiver seen as an undirected graph.*

There exists a different notion of non-commutative Poisson structures due to Crawley–Boevey [CB11], called H_0 -Poisson structure, and we will show that morphisms of double (quasi-)Poisson algebras induce morphisms of H_0 -Poisson structures. If we work over an algebraically closed field of characteristic zero, we will note as part of Proposition 5.9 that a morphism of H_0 -Poisson structures induces a Poisson morphism between the rings of invariant functions on the corresponding moduli spaces of representations. Using the latter result for the H_0 -Poisson structures induced by the double Poisson algebra structures considered in Theorem 1.1, we obtain the well-known statement that a quiver variety only depends on the quiver seen as an undirected graph (for *fixed* dimension vector and parameter of the moment map). Using the double quasi-Poisson algebra structures of Van den Bergh instead, we get the same statement for multiplicative quiver varieties, which can be attributed to Yamakawa [Yam08]. Therefore, such Poisson isomorphisms of (multiplicative) quiver varieties have a non-commutative origin given by Theorem 1.1. This will be the key observation that we shall need to construct explicitly the Poisson isomorphisms that yield the action-angle duality maps that were mentioned earlier.

Let us finish by remarking a side result of the present work, which relates to the study of noncommutative algebras in their own right. We will establish that all the automorphisms of the first Weyl algebra A_1 and the quantum torus C_1^c (see § 5.2 for definitions) are induced by isomorphisms of double (quasi-)Hamiltonian algebras.

Outline of the paper

In Section 2, we introduce double brackets [VdB08a] and their morphisms. We also review the notion of fusion, which allows to identify idempotents in an algebra and which preserves double brackets. In Section 3, we study morphisms of double Poisson algebras, and their relation to fusion. This investigation ends up with Theorem 3.12 where we obtain that Van den Bergh’s double Poisson algebras associated

with quivers are independent of the orientation of the underlying quivers. The latter section sets the stage for Section 4 where we provide a non-trivial adaptation of the previous results to the case of double quasi-Poisson algebras, and to a subclass of these algebras associated with quivers, see Theorem 4.12. In Section 5, we introduce H_0 -Poisson structures following Crawley–Boevey [CB11], and we explain how morphisms of double (quasi-)Poisson algebras yield morphisms of H_0 -Poisson structures. We describe how they induce Poisson morphisms on moduli spaces of representations. Finally, we deal with the action-angle duality of various integrable systems in Section 6 using several results that are derived throughout the paper. There are three appendices containing ancillary statements and proofs.

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2. Basic definitions

Throughout the paper, \mathbb{k} is a field of characteristic 0, and we write $\otimes = \otimes_{\mathbb{k}}$. A \mathbb{k} -algebra is always assumed to be associative, unital and finitely generated. We use $d = d' \otimes d'' \in A \otimes A$ as a shorthand way for Sweedler’s notation $d = \sum_i d'_i \otimes d''_i$.

If A, B are \mathbb{k} -algebras, we say that A is a B -algebra if there is a morphism of \mathbb{k} -algebras $B \rightarrow A$. We only consider the following special case throughout the text: if the unit in A admits a decomposition in terms of a finite set of orthogonal idempotents $(e_s)_{s \in I}$, i.e.

$$(2.1) \quad 1 = \sum_{s \in I} e_s, \quad e_s e_t = \delta_{st} e_s, \quad |I| \in \mathbb{N}^\times,$$

we view A as a B -algebra for

$$(2.2) \quad B = \bigoplus_{s \in I} \mathbb{k} e_s.$$

2.1. Double brackets

We closely follow the exposition [Fai21] of the work of Van den Bergh [VdB08a]. Let A be a \mathbb{k} -algebra. A *double bracket* on A is a \mathbb{k} -bilinear map $\{\{-, -\}\}: A \times A \rightarrow A \otimes A$ satisfying

$$(2.3) \quad \{\{a, b\}\} = -\{\{b, a\}\}^\circ \quad \text{for all } a, b \in A, \quad (\text{cyclic antisymmetry})$$

where $(-)^\circ$ denotes the permutation of factors in $A \otimes A$, together with

$$(2.4) \quad \{\{a, bc\}\} = \{\{a, b\}\}c + b\{\{a, c\}\} \quad \text{for all } a, b, c \in A. \quad (\text{right derivation rule})$$

Here, the multiplication refers to the outer A -bimodule structure on $A \otimes A$, that is if $d \in A^{\otimes 2}$, then we have $adb = (ad') \otimes (d''b)$. Assuming that (2.3) holds, one can easily check that (2.4) is equivalent to

$$(2.5) \quad \{\{bc, a\}\} = \{\{b, a\}\} * c + b * \{\{c, a\}\} \quad \text{for all } a, b, c \in A, \quad (\text{left derivation rule})$$

where we use the inner A -bimodule structure on $A \otimes A$, i.e. $a * d * b = (d'b) \otimes (ad'')$. As a consequence of the derivation rules, it suffices to define double brackets on generators of A .

Given a double bracket, we can define an operation $A^{\times 3} \rightarrow A^{\otimes 3}$ by setting

$$(2.6) \quad \begin{aligned} \{\{a, b, c\}\} &= \{\{a, \{\{b, c\}\}'\}\} \otimes \{\{b, c\}\}'' \\ &+ \tau_{(123)} \{\{b, \{\{c, a\}\}'\}\} \otimes \{\{c, a\}\}'' \\ &+ \tau_{(123)}^2 \{\{c, \{\{a, b\}\}'\}\} \otimes \{\{a, b\}\}'' \quad \text{for all } a, b, c \in A. \end{aligned}$$

(Here, we define $\tau_{(123)} : A^{\otimes 3} \rightarrow A^{\otimes 3}$ by $\tau_{(123)}(a_1 \otimes a_2 \otimes a_3) = a_3 \otimes a_1 \otimes a_2$.) The map (2.6) is an instance of *triple bracket*, that is a \mathbb{k} -trilinear map, which satisfies a generalisation of the cyclic antisymmetry (2.3)

$$(2.7) \quad \tau_{(123)} \circ \{\{-, -, -\}\} \circ \tau_{(123)}^{-1} = \{\{-, -, -\}\},$$

and which is a derivation in its last argument for the outer A -bimodule structure of $A^{\otimes 3}$.

When we see A as a B -algebra for B given as in (2.2), it is convenient to work in the relative setting. In such a case, we assume that the double bracket $\{\{-, -\}\}$ vanishes whenever one of its entries is taken in B . We then say that the double bracket is B -linear. Equivalently, B -linearity can be stated as having a morphism of B -bimodules $\{\{a, -\}\} : A \rightarrow A \otimes A$ for any $a \in A$, where we use the outer B -bimodule structure on $A \otimes A$. Note that if $a = e_{r_a} a e_{s_a} \in A$ and $b = e_{r_b} b e_{s_b} \in A$ for $r_a, r_b, s_a, s_b \in I$, the derivation rules and B -linearity yield

$$(2.8) \quad \{\{a, b\}\} \in e_{r_b} A e_{s_a} \otimes e_{r_a} A e_{s_b}.$$

2.1.1. Morphisms of double brackets

Let A_1, A_2 be B -algebras endowed with B -linear double brackets $\{\{-, -\}\}_1, \{\{-, -\}\}_2$. We say that $\phi : A_1 \rightarrow A_2$ is a *morphism of double brackets* if it is a B -algebra homomorphism such that for any $a, b \in A_1$

$$(2.9) \quad \{\{\phi(a), \phi(b)\}\}_2 = (\phi \otimes \phi) \{\{a, b\}\}_1.$$

We say that ϕ is an *isomorphism (of double brackets)* if it is also an isomorphism of B -algebras. In that case, the inverse $\phi^{-1} : A_2 \rightarrow A_1$ is a morphism of double brackets.

Example 2.1. — The algebra $A = \mathbb{k}\langle x, y \rangle$ can be endowed with two \mathbb{k} -linear double brackets, denoted $\{\{-, -\}\}_1, \{\{-, -\}\}_2$, which are defined on generators by

$$\begin{aligned} \{\{x, x\}\}_1 &= 0 = \{\{x, x\}\}_2, & \{\{y, y\}\}_1 &= 0 = \{\{y, y\}\}_2, \\ \{\{x, y\}\}_1 &= 1 \otimes 1, & \{\{x, y\}\}_2 &= -1 \otimes 1. \end{aligned}$$

The automorphism $x \mapsto y, y \mapsto x$ defines an isomorphism of double brackets $(A, \{\{-, -\}_1\}) \rightarrow (A, \{\{-, -\}_2\})$. More generally, assume that A_1 is equipped with a B -linear double bracket $\{\{-, -\}$, and fix an isomorphism of B -algebras $\phi : A_1 \rightarrow A_2$. Then we get a B -linear double bracket on A_2 defined by

$$\{\{-, -\}_2 = (\phi \otimes \phi) \circ \{\{-, -\}_1 \circ (\phi^{-1} \times \phi^{-1}) ,$$

and ϕ naturally becomes an isomorphism of double brackets.

2.2. Morphisms and fusion

2.2.1. Fusion

Following Van den Bergh [VdB08a, § 2.5], we assume that A is a B -algebra such that there exist orthogonal idempotents $e_1, e_2 \in B$, and we construct the *fusion algebra* $A^f_{e_2 \rightarrow e_1}$ obtained by fusing e_2 onto e_1 as follows. (When the choice of idempotents is clear from the context, we simply write $A^f_{e_2 \rightarrow e_1}$ as A^f .)

First, we extend the algebra A along the pair (e_1, e_2) by adjoining generators e_{12}, e_{21} satisfying the usual matrix relations. Namely,

$$(2.10) \quad \bar{A} = A *_{\mathbb{k}e_1 \oplus \mathbb{k}e_2 \oplus \mathbb{k}\hat{e}} (\text{Mat}_2(\mathbb{k}) \oplus \mathbb{k}\hat{e}) ,$$

where $\hat{e} = 1 - e_1 - e_2$, and $\text{Mat}_2(\mathbb{k})$ is seen as the \mathbb{k} -algebra generated by $e_1 = e_{11}, e_{12}, e_{21}, e_2 = e_{22}$ with $e_{st}e_{uv} = \delta_{tu}e_{sv}$. Applying the same construction to B , we can introduce \bar{B} which is such that $\bar{A} = A *_B \bar{B}$.

Second, we consider the algebra A^f constructed from \bar{A} by dismissing elements of $e_2\bar{A} + \bar{A}e_2$. We call A^f the *fusion algebra of A* obtained by fusing e_2 onto e_1 . It is explicitly given by

$$(2.11) \quad A^f = \epsilon\bar{A}\epsilon, \quad \text{for } \epsilon = 1 - e_2 .$$

By construction, A^f is a B^f -algebra for $B^f = \epsilon\bar{B}\epsilon$. Using the map

$$(2.12) \quad A \rightarrow A^f : a \mapsto a^f := \epsilon a \epsilon + e_{12} a e_{21} + e_{12} a \epsilon + \epsilon a e_{21} ,$$

we can get a convenient set of generators in A^f , as observed in [VdB08a, § 5.3].

LEMMA 2.2. — *Generators of A^f can be chosen to be of the following four types:*

$$(2.13a) \quad (\text{first type}) \quad t^f = t, \quad t \in \epsilon A \epsilon ,$$

$$(2.13b) \quad (\text{second type}) \quad u^f = e_{12} u, \quad u \in e_2 A \epsilon ,$$

$$(2.13c) \quad (\text{third type}) \quad v^f = v e_{21}, \quad v \in \epsilon A e_2 ,$$

$$(2.13d) \quad (\text{fourth type}) \quad w^f = e_{12} w e_{21}, \quad w \in e_2 A e_2 .$$

2.2.2. Fusion and quivers

We fix some terminology related to quivers, i.e. directed graphs, which will be used throughout the text. A quiver Q consists of a vertex set I and an arrow set $\{a \in Q\}$. We define the head and tail maps $h, t : Q \rightarrow I$ which assign to any arrow a its head (ending vertex) $h(a)$ and its tail (starting vertex) $t(a)$. In other words, $a : t(a) \rightarrow h(a)$. The double \bar{Q} of Q is obtained by adding to Q the arrows $a^* : h(a) \rightarrow t(a)$ for each $a \in Q$. We then extend t, h to \bar{Q} and, if we extend the map $a \rightarrow a^*$ defined on $Q \subset \bar{Q}$ to \bar{Q} by setting $(a^*)^* = a$ for $a^* \in \bar{Q} \setminus Q$, we see that $h(a) = t(a^*)$ for each $a \in \bar{Q}$. The path algebra $\mathbb{k}Q$ of a quiver Q is generated by symbols e_s for $s \in I$, and $a \in Q$, subject to the relations $e_s e_t = \delta_{st} e_s$, $a = e_{t(a)} a e_{h(a)}$ (this implies that we read paths from left to right), where the multiplication is given by concatenation of paths.

Let us now see what is the meaning of fusion on a quiver. Fix a quiver Q with vertex set $I = \{1, \dots, k\}$ for some $k \geq 2$. The path algebra $\mathbb{k}Q$ of Q is an algebra over $B = \bigoplus_{s \in I} \mathbb{k}e_s$, where e_s denotes the idempotent attached to the s^{th} vertex. We can form the algebra $(\mathbb{k}Q)^f$ as in § 2.2.1 obtained by fusing e_2 onto e_1 . We can see that $(\mathbb{k}Q)^f$ is an algebra over $B^f = \bigoplus_{s \in I \setminus \{2\}} \mathbb{k}e_s$ generated by the following elements

$$\begin{array}{llll} a & \text{with} & a \in Q & \text{such that} & t(a), & h(a) \neq 2, \\ e_{12}a & \text{with} & a \in Q & \text{such that} & t(a) = 2, & h(a) \neq 2, \\ ae_{21} & \text{with} & a \in Q & \text{such that} & t(a) \neq 2, & h(a) = 2, \\ e_{12}ae_{21} & \text{with} & a \in Q & \text{such that} & t(a), & h(a) = 2, \end{array}$$

which are images of the arrows in Q , of the four types given in Lemma 2.2. At the same time, we can form the quiver Q^f with vertex set $I^f = I \setminus \{2\}$ by fusing together the vertices 1 and 2 in Q . The corresponding path algebra $\mathbb{k}(Q^f)$ is easily identified with $(\mathbb{k}Q)^f$, so that the fusion operation is the analogue at the level of the path algebras of identifying two vertices in a quiver.

Remark 2.3. — By repeatedly using the construction from § 2.2.1, we can obtain an algebra from A by identification of several idempotents. Note that there is a direct way to define this algebra using quivers (with relations), see [Brü01, § 2.1]. We will not follow this more general approach due to technicalities arising when performing fusion in double quasi-Poisson algebras, see Proposition 4.5.

2.2.3. Fusion and double brackets

We now assume that A is endowed with a B -linear double bracket. As noted in [VdB08a, § 2.5], the double bracket uniquely extends from A to \bar{A} by requiring it to be \bar{B} -linear, and it can then be restricted to A^f . If $a^f, b^f \in A^f$ are two generators as in Lemma 2.2, there exists $a, b \in A$ and $e_a, e_b \in \{\epsilon, e_{12}\}$, $f_a, f_b \in \{\epsilon, e_{21}\}$, such that $a^f = e_a a f_a$, $b^f = e_b b f_b$. We can then define the double bracket induced by A onto A^f using the following identity

$$(2.14) \quad \{\!\!\{ a^f, b^f \}\!\!\} = e_b \{\!\!\{ a, b \}\!\!\}' f_a \otimes e_a \{\!\!\{ a, b \}\!\!\}'' f_b.$$

In (2.14), the double bracket on the left-hand side is the one induced in A^f , while the double bracket on the right-hand side is the original one defined in A .

The next result shows that, up to isomorphism, the fusion algebra and its induced double bracket only depend on the unordered choice of idempotents $\{e_1, e_2\}$.

LEMMA 2.4. — *Let $A_1 = A_{e_2 \rightarrow e_1}^f$ be the fusion algebra obtained by fusing e_2 onto e_1 , and let $A_2 = A_{e_1 \rightarrow e_2}^f$ be the fusion algebra obtained by fusing e_1 onto e_2 . Then, the identity map on A induces an isomorphism of double brackets $A_1 \rightarrow A_2$.*

Proof. — We let $\widehat{e} = 1 - e_1 - e_2$ and note that $A_1 = (1 - e_2)\bar{A}(1 - e_2)$ and $A_2 = (1 - e_1)\bar{A}(1 - e_1)$. We can define a map $\phi : A_1 \rightarrow A_2$ which is given on generators of first type (2.13a) as

$$\begin{aligned} \phi(t) &= t, & \text{if } t \in \widehat{e}A\widehat{e}; & & \phi(t) &= e_{21}t, & \text{if } t \in e_1A\widehat{e}; \\ \phi(t) &= te_{12}, & \text{if } t \in \widehat{e}Ae_1; & & \phi(t) &= e_{21}te_{12}, & \text{if } t \in e_1Ae_1; \end{aligned}$$

on generators of second type (2.13b) as

$$\phi(e_{12}u) = u, \quad \text{if } u \in e_2A\widehat{e}; \quad \phi(e_{12}u) = ue_{12}, \quad \text{if } u \in e_2Ae_1;$$

on generators of third type (2.13c) as

$$\phi(ve_{21}) = v, \quad \text{if } v \in \widehat{e}Ae_2; \quad \phi(ve_{21}) = e_{21}v, \quad \text{if } v \in e_1Ae_2;$$

on generators of fourth type (2.13d) as

$$\phi(e_{12}we_{21}) = w, \quad \text{if } w \in e_2Ae_2.$$

The map ϕ is easily seen to be the image of the identity under the two projections $\pi_1 : A \rightarrow A_1, \pi_2 : A \rightarrow A_2$ given by (2.12), i.e. $\phi \circ \pi_1 = \pi_2 \circ \text{id}_A$. Moreover, swapping the labels 1, 2 provides the inverse of ϕ . Thus, we only need to check that it is a morphism of double brackets.

We can decompose A as $\bigoplus_{i,j=0,1,2} e'_i A e'_j$ if we set $e'_0 = \widehat{e}$ and $e'_i = e_i$ for $i = 1, 2$. Without loss of generality, take $a, b \in A$ belonging to one of these subsets. Then, there exists $e_a^1, e_b^1 \in \{e_{12}, e_1, \widehat{e}\}, f_a^1, f_b^1 \in \{e_{21}, e_1, \widehat{e}\}$, such that under the projection π_1 the elements $e_a^1 a f_a^1$ and $e_b^1 b f_b^1$ are generators of A_1 , see Lemma 2.2.

In the same way, there exists $e_a^2, e_b^2 \in \{e_{21}, e_2, \widehat{e}\}, f_a^2, f_b^2 \in \{e_{12}, e_2, \widehat{e}\}$, such that the elements $e_a^2 a f_a^2$ and $e_b^2 b f_b^2$ obtained from the projection π_2 are generators of A_2 . In particular, $\phi \circ \pi_1 = \pi_2$ implies that

$$\phi(e_a^1) = e_a^2, \quad \phi(e_b^1) = e_b^2, \quad \phi(f_a^1) = f_a^2, \quad \phi(f_b^1) = f_b^2.$$

Using the identity (2.14) for the double bracket $\{\!\{-, -\}\!\}_k$ induced by $\{\!\{-, -\}\!\}$ in A_k for $k = 1, 2$, we get

$$\{\!\{e_a^k a f_a^k, e_b^k b f_b^k\}\!\}_k = e_b^k \{\!\{a, b\}\!\}' f_a^k \otimes e_a^k \{\!\{a, b\}\!\}'' f_b^k,$$

where the double bracket on the right-hand side is taken in A . We can then directly see that (2.9) is satisfied. □

Fusion also preserves morphisms of double brackets.

LEMMA 2.5. — Let $\phi : A_1 \rightarrow A_2$ be a morphism of double brackets over B . Let $A_1^f = (A_1)_{e_2 \rightarrow e_1}^f$, $A_2^f = (A_2)_{e_2 \rightarrow e_1}^f$, be the fusion algebras with double brackets obtained by fusing e_2 onto e_1 . Then ϕ induces a morphism of double brackets $\phi^f : A_1^f \rightarrow A_2^f$.

Proof. — Recall that A_1^f is generated by elements $a^f = e_a a f_a$ for $a \in A_1$ and $e_a \in \{\epsilon, e_{12}\}$, $f_a \in \{\epsilon, e_{21}\}$. We can then set

$$(2.15) \quad \phi^f : A_1^f \rightarrow A_2^f : a^f = e_a a f_a \mapsto \phi^f(a^f) = e_a \phi(a) f_a.$$

For $k = 1, 2$, let $\{\{-, -\}_k$ and $\{\{-, -\}_k^f$ be respectively the double bracket on A_k and its induced double bracket on A_k^f . We have for any two generators $a^f = e_a a f_a$, $b^f = e_b b f_b \in A_1^f$ as described above that

$$\begin{aligned} \{\{\phi^f(a^f), \phi^f(b^f)\}_2^f &= e_b \{\{\phi(a), \phi(b)\}'_2 f_a \otimes e_a \{\{\phi(a), \phi(b)\}''_2 f_b \\ &= e_b \phi(\{\{a, b\}'_1) f_a \otimes e_a \phi(\{\{a, b\}''_1) f_b \\ &= (\phi^f \otimes \phi^f) e_b \{\{a, b\}'_1 f_a \otimes e_a \{\{a, b\}''_1 f_b \\ &= (\phi^f \otimes \phi^f) \{\{a^f, b^f\}_1. \end{aligned}$$

Here, we used formula (2.14) in A_2^f/A_1^f for the first/last equality, the morphism property (2.9) in the second equality, and the definition of ϕ^f for the remaining equality. \square

In Lemma 2.5, it is clear that if ϕ is an isomorphism, then so is ϕ^f . For the remainder of this text, we omit to mention such an extension to the case of an isomorphism in analogous results.

2.2.4. Identification of idempotents in different algebras

Let A_1 and A_2 be algebras endowed with double brackets respectively over B_1 and B_2 , and consider $A_1 \oplus A_2$ as a $(B_1 \oplus B_2)$ -algebra. It is easy to see that there exists a unique $(B_1 \oplus B_2)$ -linear double bracket $\{\{-, -\}^\oplus$ on $A_1 \oplus A_2$ extending $\{\{-, -\}_1$ and $\{\{-, -\}_2$, while it is such that $\{\{c_1, c_2\}^\oplus = 0$ whenever $c_1 = (a_1, 0)$, $c_2 = (0, a_2)$, with $a_1 \in A_1$, $a_2 \in A_2$.

For any idempotents $e_1 \in B_1$ and $e_2 \in B_2$, we can form the fusion algebra $(A_1 \oplus A_2)_{e_2 \rightarrow e_1}^f$ by fusing e_2 onto e_1 . This algebra inherits a $(B_1 \oplus B_2)_{e_2 \rightarrow e_1}^f$ -linear double bracket as noted in §2.2.3.

Example 2.6 (Extension by a central element). — Let A be a \mathbb{k} -algebra equipped with a double bracket $\{\{-, -\}$. Let $\mathbb{k}\langle y \rangle$ be equipped with the trivial double bracket $\{\{y, y\} = 0$. If we identify the units of these two algebras embedded as $(1, 0)$, $(0, 1)$ inside $A \oplus \mathbb{k}\langle y \rangle$, the fusion algebra $A_+ := A *_\mathbb{k} \mathbb{k}\langle y \rangle$ hence obtained is endowed with the double bracket $\{\{-, -\}_+$ given for any $a, b \in A$ by $\{\{a, b\}_+ = \{\{a, b\}$, $\{\{a, y\}_+ = 0$ and $\{\{y, y\}_+ = 0$. (Note that if $\{\{-, -\}$ is Poisson in the sense of Section 3, then $(A_+, \{\{-, -\}_+)$ is a double Poisson algebra.)

3. Morphisms of double Poisson algebras

As in Section 2, all algebras are over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. We will frequently identify I with $\{1, \dots, |I|\} \subset \mathbb{N}$.

3.1. Basic definitions and properties

3.1.1. Definitions

Following Van den Bergh [VdB08a], a *double Poisson bracket* $\{\{-, -\}$ on A is a double bracket for which the associated triple bracket $\{\{-, -, -\}$ defined by (2.6) identically vanishes. In such a case, we say that $(A, \{\{-, -\})$ (or simply A) is a *double Poisson algebra*. An element $\mu_s \in e_s A e_s$ is said to satisfy the *additive property* for e_s if, for any $a \in A$, we have

$$(3.1) \quad \{\{\mu_s, a\}\} = a e_s \otimes e_s - e_s \otimes e_s a.$$

Then, we define a *moment map* as an element $\mu = \sum_{s \in I} \mu_s$ with $\mu_s \in e_s A e_s$ such that for all $s \in I$ we have that μ_s satisfies the additive property for e_s (3.1). Note in particular that $e_s \mu e_t = \delta_{st} \mu_s$ for each $s, t \in I$. We call the triple $(A, \{\{-, -\}, \mu)$ (or simply A if no confusion can arise) a *Hamiltonian algebra*.

Let A_1, A_2 be B -algebras with B -linear double brackets $\{\{-, -\}_1, \{\{-, -\}_2$. If $\phi : A_1 \rightarrow A_2$ is a morphism of double brackets and A_1, A_2 are double Poisson algebras, we say that ϕ is a *morphism of double Poisson algebras*. If furthermore A_1, A_2 are Hamiltonian algebras with respective moment maps μ_1, μ_2 satisfying $\phi(\mu_1) = \mu_2$, we say that ϕ is a *morphism of Hamiltonian algebras*.

Example 3.1. — Consider the quiver Q with vertices $\{1, 2\}$ and one arrow $a : 1 \rightarrow 2$. Let \bar{Q} be its double obtained by adding the arrow $a^* : 2 \rightarrow 1$. Put $B = \mathbb{k}e_1 \oplus \mathbb{k}e_2$. Following Van den Bergh [VdB08a, § 6.3], the path algebra $\mathbb{k}\bar{Q}$ is a Hamiltonian algebra for the B -linear double bracket given on generators by

$$(3.2) \quad \{\{a, a\}\}_1 = 0, \quad \{\{a^*, a^*\}\}_1 = 0, \quad \{\{a, a^*\}\}_1 = e_2 \otimes e_1, \quad \{\{a^*, a\}\}_1 = -e_1 \otimes e_2,$$

and the moment map $\mu = [a, a^*]$. The moment map can be decomposed as $\mu_1 = aa^*$, $\mu_2 = -a^*a$. (See § 2.2.2 for the conventions that we follow with respect to quivers.)

Similarly, consider the quiver Q^{op} with vertices $\{1, 2\}$ and one arrow $b : 2 \rightarrow 1$, and let \bar{Q}^{op} be its double with new arrow $b^* : 1 \rightarrow 2$. We can also consider Van den Bergh’s Hamiltonian structure on $\mathbb{k}\bar{Q}^{op}$ which is given by the B -linear double bracket

$$(3.3) \quad \{\{b, b\}\}_2 = 0, \quad \{\{b^*, b^*\}\}_2 = 0, \quad \{\{b, b^*\}\}_2 = e_1 \otimes e_2, \quad \{\{b^*, b\}\}_2 = -e_2 \otimes e_1,$$

with moment map $\mu' = [b, b^*]$. We can see that $\phi : \mathbb{k}\bar{Q} \rightarrow \mathbb{k}\bar{Q}^{op}$ given by $\phi(a) = b^*$, $\phi(a^*) = -b$, is an isomorphism of B -algebras. One readily checks that

$$\begin{aligned} \phi^{\otimes 2} \{\{a, a\}\}_1 &= 0 = \{\{b^*, b^*\}\}_2, & \phi^{\otimes 2} \{\{a^*, a^*\}\}_1 &= 0 = \{\{b, b\}\}_2, \\ \phi^{\otimes 2} \{\{a, a^*\}\}_1 &= e_2 \otimes e_1 = \{\{b^*, -b\}\}_2, \end{aligned}$$

and $\phi(\mu) = \mu'$. Hence, ϕ is an isomorphism of Hamiltonian algebras.

3.1.2. First properties

The next result shows how injectivity or surjectivity of a morphism of double brackets can guarantee that it is a morphism of double Poisson (resp. Hamiltonian) algebras.

LEMMA 3.2. — *Let $\phi : A_1 \rightarrow A_2$ be a morphism of double brackets.*

- (1) *Assume that ϕ is surjective as a B -algebra homomorphism.*
 - (a) *If A_1 is a double Poisson algebra, then A_2 is a double Poisson algebra.*
 - (b) *If A_1 is a Hamiltonian algebra, then A_2 admits a structure of Hamiltonian algebra such that ϕ is a morphism of Hamiltonian algebras.*
- (2) *Assume that ϕ is injective as a B -algebra homomorphism.*
 - (a) *If A_2 is a double Poisson algebra, then A_1 is a double Poisson algebra.*
 - (b) *If A_2 is a Hamiltonian algebra with moment map μ' such that $\mu' \in \text{Im}(\phi)$, then A_1 admits a structure of Hamiltonian algebra such that ϕ is a morphism of Hamiltonian algebras.*
- (3) *All the structures obtained in (1)–(2) are unique.*

Proof. — First, let us consider for $1 \leq i \leq 3$ the element $a_{1,i} \in A_1$, with image $a_{2,i} := \phi(a_{1,i}) \in A_2$. If we denote by σ the permutation (123), we get by (2.6) and (2.9)

$$\begin{aligned}
& \{\{a_{2,1}, a_{2,2}, a_{2,3}\}\}_2 \\
&= \sum_{i \in \mathbb{Z}_3} \tau_\sigma^i \left\{ \left\{ \phi(a_{1,\sigma^i(1)}), \left\{ \left\{ \phi(a_{1,\sigma^i(2)}), \phi(a_{1,\sigma^i(3)}) \right\}' \right\}_2 \right\}' \right\}_2 \otimes \left\{ \left\{ \phi(a_{1,\sigma^i(2)}), \phi(a_{1,\sigma^i(3)}) \right\}'' \right\}_2 \\
&= \sum_{i \in \mathbb{Z}_3} \tau_\sigma^i (\phi \otimes \phi \otimes \phi) \left\{ \left\{ a_{1,\sigma^i(1)}, \left\{ \left\{ a_{1,\sigma^i(2)}, a_{1,\sigma^i(3)} \right\}' \right\}' \right\}_1 \right\}_1 \otimes \left\{ \left\{ a_{1,\sigma^i(2)}, a_{1,\sigma^i(3)} \right\}'' \right\}_1 \\
&= (\phi \otimes \phi \otimes \phi) \{\{a_{1,1}, a_{1,2}, a_{1,3}\}\}_1 .
\end{aligned}$$

This yields in particular the identity of associated triple brackets

$$(3.4) \quad (\phi \otimes \phi \otimes \phi) \circ \{\{-, -, -\}\} = \{\{-, -, -\}\} \circ (\phi \times \phi \times \phi) .$$

(1.a) If $\{\{-, -\}\}$ is Poisson, $\{\{-, -, -\}\}$ vanishes and (3.4) yields that $\{\{-, -\}\}$ is Poisson.

(1.b) Since A_1 has a moment map $\mu = \sum_s \mu_s$, the element $\phi(\mu) = \sum_s \phi(\mu_s)$ is a moment map. Indeed, for any $s \in I$

$$\begin{aligned}
\{\{\phi(\mu_s), a_{2,1}\}\}_2 &= (\phi \otimes \phi) \{\{\mu_s, a_{1,1}\}\}_1 \\
&= (\phi \otimes \phi) (a_{1,1}e_s \otimes e_s - e_s \otimes e_s a_{1,1}) = a_{2,1}e_s \otimes e_s - e_s \otimes e_s a_{2,1} ,
\end{aligned}$$

where we have used that μ_s satisfies the additive property for e_s with respect to $\{\{-, -\}\}$.

(2.a-b) This is similar, and we need the fact that $\phi^{\otimes 2}, \phi^{\otimes 3}$ are injective.

(3) The uniqueness of the structures is a consequence of the constructions. \square

Given a double Poisson algebra $(A, \{\{-, -\}\})$, we introduce the set of *Casimir elements* of A by

$$(3.5) \quad \text{Cas}(A) = \{a \in A \mid \{\{a, b\}\} = 0 \text{ for all } b \in A\} .$$

In the case of a B -linear double bracket, we have that $B \subset \text{Cas}(A)$.

LEMMA 3.3. — *Let $\phi : A_1 \rightarrow A_2$ be an isomorphism of double Poisson algebras. If A_1, A_2 are Hamiltonian algebras with moment maps μ_1, μ_2 , then $\mu_2 - \phi(\mu_1) \in B$ and $\text{Cas}(A_1) = \text{Cas}(A_2) = B$. In particular, a moment map is unique up to the addition of an element of B .*

Proof. — Let us decompose $\mu_1 \in A_1$ with the idempotents as $\sum_{s \in I} \mu_{1,s}$ and do the same for $\mu_2 \in A_2$. We first note that for any $s \in I$, due to the additive property for e_s (3.1), we have for any $a \in A_2$ that

$$\{\{\mu_{2,s} - \phi(\mu_{1,s}), a\}\}_2 = \{\{\mu_{2,s}, a\}\}_2 - \phi^{\otimes 2} \{\{\mu_{1,s}, \phi^{-1}(a)\}\}_1 = 0.$$

Thus, $\mu_{2,s} - \phi(\mu_{1,s}) \in \text{Cas}(A_2)$ by definition.

Let us now show that $\text{Cas}(A_1) \subset B$, since we already know the opposite inclusion (the case of A_2 is obtained in the same way). If $c \in \text{Cas}(A_1)$, we have for any $s \in I$ that

$$0 = \{\{\mu_{1,s}, c\}\} = ce_s \otimes e_s - e_s \otimes ce_s,$$

where we used in the first equality that $\{\{-, c\}\}_1 = 0$ because c is a Casimir element, while in the second we use the property (3.1) of the moment map. This equality implies that $ce_s = \lambda_s e_s = e_s c$ for some $\lambda_s \in \mathbb{k}$, thus $c = \sum_{s \in I} \lambda_s e_s \in B$.

For the second part of the statement, it suffices to consider $\phi = \text{id}_{A_1} : A_1 \rightarrow A_1$. \square

Note that, in general $B \subsetneq \text{Cas}(A)$. For example, any B -algebra A can be endowed with the zero double bracket, which is Poisson and such that $\text{Cas}(A) = A$. This does not contradict Lemma 3.3 since the zero double Poisson bracket does not admit a moment map (except if $A = B$).

3.1.3. Sum of double Poisson algebras

Recall that we considered the sum of two double brackets in §2.2.4. The following result holds in the presence of double Poisson brackets.

LEMMA 3.4 ([VdB08a]). — *If A_1 and A_2 are endowed with double Poisson brackets, then the double bracket $\{\{-, -\}\}^\oplus$ on $A_1 \oplus A_2$ defined in §2.2.4 is Poisson. If μ_1 and μ_2 are moment maps in A_1 and A_2 , then $A_1 \oplus A_2$ is Hamiltonian with (μ_1, μ_2) as moment map.*

3.2. Double Poisson brackets and fusion

We assume that the index set I of B is such that $|I| > 1$.

3.2.1. Fusion in an algebra

Recall the fusion algebra $A^f = A_{e_2 \rightarrow e_1}^f$ defined in §2.2.1. We noted in §2.2.3 that if A is endowed with a double bracket, then A^f has an induced double bracket.

PROPOSITION 3.5 ([VdB08a, Corollary 2.5.6, Proposition 2.6.6]). — *If A is a double Poisson algebra over B , then A^f equipped with the induced double bracket is a double Poisson algebra over $B^f = \bigoplus_{s \in I \setminus \{2\}} \mathbb{k}e_s$. Furthermore, if μ is a moment map for A , then its projection μ^f under the map (2.12) is a moment map for A^f .*

Example 3.6. — Consider the path algebra $A = \mathbb{k}\bar{Q}$ analysed in Example 3.1 with its Hamiltonian algebra structure. Fusing e_2 onto e_1 , we get a Hamiltonian algebra structure on A^f which can be identified with the free algebra $\mathbb{k}\langle a, a^* \rangle$. The induced double bracket is determined by the first two identities of (3.2) and $\{\!\{a, a^*\}\!\} = 1 \otimes 1$. The moment map is $\mu = [a, a^*]$. The algebra is the path algebra of the double of the one-loop quiver, and the Hamiltonian structure obtained by fusion is just the one defined by Van den Bergh [VdB08a, § 6.3].

LEMMA 3.7. — *Let A be a double Poisson algebra over B . Consider the algebra $A_1 = A^f_{e_2 \rightarrow e_1}$ obtained by fusing e_2 onto e_1 , and the algebra $A_2 = A^f_{e_1 \rightarrow e_2}$ obtained by fusing e_1 onto e_2 . Then, the identity map on A induces an isomorphism of double Poisson algebras $\phi : A_1 \rightarrow A_2$. If A is a Hamiltonian algebra, then ϕ is an isomorphism of Hamiltonian algebras.*

Proof. — The first part directly follows from Lemma 2.4 and Proposition 3.5. In the Hamiltonian case, the morphism ϕ constructed in the proof of Lemma 2.4 satisfies

$$\begin{aligned} \phi(\mu^{(1)}) &= \mu^{(2)}, \quad \text{for } \mu^{(1)} = \mu_1 + e_{12}\mu_2e_{21} + \sum_{s \neq 1, 2} \mu_s, \\ \mu^{(2)} &= e_{21}\mu_1e_{12} + \mu_2 + \sum_{s \neq 1, 2} \mu_s. \end{aligned}$$

But $\mu^{(1)}, \mu^{(2)}$ are the moment maps of A_1 and A_2 by Proposition 3.5. □

In Lemma 3.7, the map $\phi : A_1 \rightarrow A_2$ hence obtained is a morphism of B' -algebras, where $B' = \mathbb{k}\tilde{e} \oplus \bigoplus_{s \in I \setminus \{1, 2\}} \mathbb{k}e_s$. Here, the base maps satisfy $\tilde{e} \mapsto e_1 \in A_1$, $\tilde{e} \mapsto e_2 \in A_2$ and are given in an obvious way on the other idempotents e_s , $s \in I \setminus \{1, 2\}$.

3.2.2. Fusion of several idempotents

LEMMA 3.8. — *Let A be a double Poisson algebra over B , and let $e_1, e_2, e_3 \in B$ be orthogonal idempotents. Let $A_1 := (A^f_{e_3 \rightarrow e_2})^f_{e_2 \rightarrow e_1}$ (resp. $A_2 := (A^f_{e_2 \rightarrow e_1})^f_{e_3 \rightarrow e_1}$) be the algebra obtained by fusing e_3 onto e_1 , then e_2 onto e_1 (resp. e_2 onto e_1 , then e_3 onto e_1). Then the identity map on A induces an isomorphism of double Poisson algebras $\psi : A_1 \rightarrow A_2$ over $B' = \bigoplus_{s \in I \setminus \{2, 3\}} \mathbb{k}e_s$. Furthermore, if A is a Hamiltonian algebra, then ψ is an isomorphism of Hamiltonian algebras.*

Proof. — Using Lemma 2.2 twice, we can write generators of A_1 and A_2 as

$$\begin{aligned} e_a a f_a, \quad & \text{for } a \in A, e_a \in \{1 - e_3 - e_2, e_{12}, e_{12}e_{23}\}, \\ & f_a \in \{1 - e_3 - e_2, e_{21}, e_{32}e_{21}\}, \quad (\text{in } A_1); \\ e'_a a f'_a, \quad & \text{for } a \in A, e'_a \in \{1 - e_3 - e_2, e_{12}, e_{13}\}, \\ & f'_a \in \{1 - e_3 - e_2, e_{21}, e_{31}\}, \quad (\text{in } A_2). \end{aligned}$$

We can then define $\phi : A_1 \rightarrow A_2$ on generators as $\phi(e_a a f_a) = e'_a a f'_a$ with $e'_a = e_a$ if $e_a \neq e_{12}e_{23}$ or $e'_a = e_{13}$ if $e_a = e_{12}e_{23}$, while $f'_a = f_a$ if $f_a \neq e_{32}e_{21}$ or $f'_a = e_{31}$ if $f_a = e_{32}e_{21}$. This is easily seen to be an isomorphism and, using (2.14) twice to induce the double bracket from A to A_1 or A_2 , we remark that it is also a morphism of double brackets. Hence, we have an isomorphism of double Poisson algebras.

In the Hamiltonian case, Proposition 3.5 yields that the moment maps are given by

$$\begin{aligned} \mu^{(1)} &= e_{12}\mu_2e_{21} + e_{12}e_{23}\mu_3e_{32}e_{21} + \sum_{s \neq 2,3} \mu_s \in A_1, \\ \mu^{(2)} &= e_{12}\mu_2e_{21} + e_{13}\mu_3e_{31} + \sum_{s \neq 2,3} \mu_s \in A_2, \end{aligned}$$

so that $\phi(\mu^{(1)}) = \mu^{(2)}$. □

Let us note the following interesting result. By gathering Lemmas 3.7 and 3.8, the double Poisson algebra structure obtained by fusion of three orthogonal idempotents does not depend on the order with respect to which we perform fusion, up to isomorphism. For example,

$$\begin{aligned} \left(A_{e_2 \rightarrow e_1}^f\right)_{e_3 \rightarrow e_1}^f &\xrightarrow{\sim} \left(A_{e_3 \rightarrow e_2}^f\right)_{e_2 \rightarrow e_1}^f \xrightarrow{\sim} \left(A_{e_3 \rightarrow e_2}^f\right)_{e_1 \rightarrow e_2}^f \\ &\xrightarrow{\sim} \left(A_{e_3 \rightarrow e_1}^f\right)_{e_1 \rightarrow e_2}^f \xrightarrow{\sim} \left(A_{e_3 \rightarrow e_1}^f\right)_{e_2 \rightarrow e_1}^f. \end{aligned}$$

Next, we establish that the order in which we perform fusion is not important if the idempotents are not related.

LEMMA 3.9. — *Let A be a double Poisson algebra over B , and let $e_1, e_2, e_3, e_4 \in B$ be orthogonal idempotents. Let $A_1 := (A_{e_4 \rightarrow e_3}^f)_{e_2 \rightarrow e_1}^f$ (resp. $A_2 := (A_{e_2 \rightarrow e_1}^f)_{e_4 \rightarrow e_3}^f$) be the algebra obtained by fusing e_4 onto e_3 , then e_2 onto e_1 (resp. e_2 onto e_1 , then e_4 onto e_3). Then the identity map on A induces an isomorphism of double Poisson algebras $\psi : A_1 \rightarrow A_2$ over $B' = \bigoplus_{s \in I \setminus \{2,4\}} \mathbb{k}e_s$. Furthermore, if A is a Hamiltonian algebra, then ψ is an isomorphism of Hamiltonian algebras.*

Proof. — Using Lemma 2.2 twice, we can write both generators of A_1 and A_2 as $e_a a f_a$, for $a \in A$, $e_a \in \{1 - e_4 - e_2, e_{12}, e_{34}\}$, $f_a \in \{1 - e_4 - e_2, e_{21}, e_{43}\}$.

In terms of these generators, the identity map provides the desired morphism. □

By gathering these results, we obtain that the double Poisson (or Hamiltonian) algebra structure of an algebra obtained by successive fusions is independent of the precise order in which we identify idempotents. To state the result, fix a double Poisson algebra A over $B = \bigoplus_{s \in I} \mathbb{k}e_s$, and consider a partition $I = \sqcup_{j \in J} I_j$. We say that A_1 is a *fusion algebra respecting the partition* $\sqcup_{j \in J} I_j$ if A_1 is obtained by a finite number of fusions starting from A , so that all the idempotents $(e_s)_{s \in I_j}$ end up being identified together for each $j \in J$. Equivalently, there exists $s_j \in I_j$ for each $j \in J$ and a map $\pi_1 : A \rightarrow A_1$ obtained by composing the morphisms (2.12) induced by a finite number of fusions such that

$$(3.6) \quad \pi_1(e_s) = e_{s_j} \quad \text{for each } s \in I_j \quad \text{and } j \in J.$$

Note that if A is a double Poisson algebra, then so too is A_1 by repeated use of Proposition 3.5.

THEOREM 3.10. — *Let A be a double Poisson algebra over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. Assume that A_1, A_2 are fusion algebras respecting a partition $I = \sqcup_{j \in J} I_j$. Then, there is a double Poisson algebra isomorphism $\phi : A_1 \rightarrow A_2$. If A is Hamiltonian, ϕ is an isomorphism of Hamiltonian algebras.*

Proof. — We prove the result by induction on $|J|$. The base case corresponds to the partition

$$(3.7) \quad |J| = 1, \quad I_1 = I,$$

which amounts to define A_1 and A_2 by fusing all the idempotents of B together. In that case, the result follows from Lemmas 3.7, 3.8 and 3.9 (see also the observation made right after Lemma 3.8).

Next, we prove that there exists an isomorphism between A_1 and a fusion algebra \widehat{A}_1 respecting the partition, such that \widehat{A}_1 is defined by first performing fusion of all the elements in $I_{|J|}$. If $|I_{|J|}| = 1$ there is nothing to prove. Assuming $|I_{|J|}| > 1$, we have that A_1 is obtained from A by a chain of fusions

$$(3.8) \quad e_{r_1} \rightarrow e_{t_1}, \quad \text{then } e_{r_2} \rightarrow e_{t_2}, \quad \text{then } \dots, \quad e_{r_\kappa} \rightarrow e_{t_\kappa}, \quad \kappa := |I| - |J|.$$

(Here $e_{r_k} \rightarrow e_{t_k}$ means that we fuse e_{r_k} onto e_{t_k} .) Note that r_1, \dots, r_κ are distinct, and there exists k such that $r_k \in I_{|J|}$ (or equivalently $t_k \in I_{|J|}$) by assumption. Let $k_1 \in \{1, \dots, \kappa\}$ be the smallest such integer for which $r_{k_1} \in I_{|J|}$. This means that all the fusions that are performed before are done by fusing idempotents corresponding to subsets distinct from $I_{|J|}$. we can thus use Lemma 3.9 to get that A_1 is isomorphic as a double Poisson algebra to the algebra which is obtained by the chain of fusions

$$\begin{aligned} e_{r_{k_1}} \rightarrow e_{t_{k_1}}, \quad \text{then } e_{r_1} \rightarrow e_{t_1}, \dots, e_{r_{k_1-1}} \rightarrow e_{t_{k_1-1}}, \\ \text{then } e_{r_{k_1+1}} \rightarrow e_{t_{k_1+1}}, \dots, e_{r_\kappa} \rightarrow e_{t_\kappa}. \end{aligned}$$

We then run this argument again on the last $\kappa - 1$ fusions of this chain. By a repeated use of this argument involving Lemma 3.9 only, we have that A_1 is isomorphic as a double Poisson algebra to \widehat{A}_1 obtained by a chain of fusion as (3.8), where this time $r_1, \dots, r_{|I_{|J|}|} \in I_{|J|}$. In particular, $r_k \notin I_{|J|}$ for $k > |I_{|J|}|$.

We now do the same with A_2 to get an isomorphism of double Poisson algebras with some \widehat{A}_2 , where the latter is obtained by a chain of fusion as (3.8), where $r_1, \dots, r_{|I_{|J|}|} \in I_{|J|}$. In particular, up to using the argument for $|J| = 1$, we can assume that the first $|I_{|J|}|$ fusions performed in \widehat{A}_1 and \widehat{A}_2 are precisely the same. Since $\widehat{A}_1, \widehat{A}_2$ only differ by fusions of the idempotents corresponding to the partition $I \setminus I_{|J|} = I_1 \sqcup \dots \sqcup I_{|J|-1}$, we have by induction that \widehat{A}_1 and \widehat{A}_2 are isomorphic. Thus A_1 and A_2 are also isomorphic as double Poisson algebras.

We can conclude since the isomorphisms involved are isomorphisms of Hamiltonian algebras if A admits a moment map. □

3.2.3. Fusion of morphisms

LEMMA 3.11. — *Let $\phi : A_1 \rightarrow A_2$ be a morphism of double Poisson algebras over B . Let $A_1^f = (A_1)_{e_2 \rightarrow e_1}^f, A_2^f = (A_2)_{e_2 \rightarrow e_1}^f$, be the fusion algebras with double brackets obtained by fusion of e_2 onto e_1 . Then ϕ induces a morphism of double*

Poisson algebras $\phi^f : A_1^f \rightarrow A_2^f$. Furthermore if ϕ is a morphism of Hamiltonian algebras, then so is ϕ^f .

Proof. — The map ϕ^f given by (2.15) is a morphism of double Poisson algebras by Lemma 2.5 and Lemma 3.2. In the Hamiltonian case, we can check from (2.15) that ϕ^f maps the moment map of A_1^f to that of A_2^f . Hence, it is a morphism of Hamiltonian algebras. \square

3.3. Application: Van den Bergh’s Hamiltonian structure for quivers

Following Van den Bergh [VdB08a, § 6.3], we can endow the path algebras of quivers with a Hamiltonian algebra structure. More precisely, given a quiver Q , we define on $\mathbb{k}\bar{Q}$ the double Poisson bracket given by

$$(3.9) \quad \{\{a, a^*\}\} = e_{h(a)} \otimes e_{t(a)} \quad \text{for } a \in Q, \quad \{\{a, a^*\}\} = -e_{h(a)} \otimes e_{t(a)} \quad \text{for } a \in \bar{Q} \setminus Q,$$

and such that $\{\{a, b\}\} = 0$ if $a \in \bar{Q}$ and $b \in \bar{Q} \setminus \{a, a^*\}$. The algebra admits the moment map

$$(3.10) \quad \mu = \sum_{a \in Q} (aa^* - a^*a) \quad \text{or} \quad \mu = \sum_{s \in I} \mu_s, \quad \text{for } \mu_s = \sum_{\substack{a \in Q \\ t(a)=s}} aa^* - \sum_{\substack{a \in Q \\ h(a)=s}} a^*a.$$

This was considered in a simple case in Example 3.1.

THEOREM 3.12. — *The Hamiltonian algebra $(\mathbb{k}\bar{Q}, \{\{-, -\}\}, \mu)$ only depends on Q seen as an undirected graph, up to isomorphism.*

Proof. — As in the proof of [VdB08a, Theorem 6.7.1], we begin with the “separated” quiver Q^{sep} which has vertex and arrow sets given by

$$(3.11) \quad I^{sep} = \{v_b, v_{b^*} \mid b \in Q\}, \quad Q^{sep} = \{b : v_b \rightarrow v_{b^*} \mid b \in Q\}.$$

We form the double \bar{Q}^{sep} of Q^{sep} , which amounts to add the arrows $\{b^* : v_{b^*} \rightarrow v_b \mid b \in Q\}$. We define on it the involution $*$ given by $b \mapsto b^*$ and $b^* \mapsto b$. By combining Example 3.1 (with $a = b$ for each $b \in Q^{sep}$) and Lemma 3.4, $A^{sep} = \mathbb{k}\bar{Q}^{sep}$ is Hamiltonian for the double Poisson bracket given by

$$(3.12) \quad \{\{b, b^*\}\} = e_{v_{b^*}} \otimes e_{v_b},$$

for all $b \in Q^{sep}$ and which is zero on each other pair of generators, while the moment map is defined as

$$(3.13) \quad \begin{aligned} \mu &= \sum_{b \in \bar{Q}^{sep}} \mu_{v_b}, \quad \mu_{v_b} = bb^* \quad \text{if } b \in Q^{sep}, \\ &\text{or } \mu_{v_b} = -bb^* \quad \text{if } b \in \bar{Q}^{sep} \setminus Q^{sep}. \end{aligned}$$

To get a Hamiltonian structure on $\mathbb{k}\bar{Q}$, it remains to fuse all these disjoint quivers of \bar{Q}^{sep} to form \bar{Q} . We can easily see that the Hamiltonian algebra structure induced by fusion coincides with (3.9)–(3.10). The independence of order in which we identify the vertices is obvious from the notations, and in fact it follows from Theorem 3.10.

It remains to see that this structure is independent of the directions of the arrows in Q , up to isomorphism. To do so, we first use the isomorphism from Example 3.1 to get that we can reverse the direction of any arrow $b \in Q^{sep}$ and obtain a Hamiltonian algebra isomorphic to A^{sep} . But Lemma 3.11 guarantees that fusion preserves isomorphisms of Hamiltonian algebras, so we are done. \square

4. Morphisms of double quasi-Poisson algebras

As in Section 2, all algebras are over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. With the exception of Remark 4.13, all the statements from this section have a counterpart in Section 3.

4.1. Basic definitions and properties

4.1.1. Definitions

Following Van den Bergh [VdB08a], we say that a *double quasi-Poisson bracket* $\{\!\{-, -\}\!\}$ on A is a double bracket for which the associated triple bracket $\{-, -, -\}$ defined by (2.6) satisfies

$$(4.1) \quad \{\!\{a, b, c\}\!\} = \frac{1}{4} \sum_{s \in I} \left(ce_s a \otimes e_s b \otimes e_s - ce_s a \otimes e_s \otimes be_s - ce_s \otimes ae_s b \otimes e_s + ce_s \otimes ae_s \otimes be_s - e_s a \otimes e_s b \otimes e_s c + e_s a \otimes e_s \otimes be_s c + e_s \otimes ae_s b \otimes e_s c - e_s \otimes ae_s \otimes be_s c \right),$$

on any $a, b, c \in A$. In such a case, we say that $(A, \{\!\{-, -\}\!\})$ (or simply A) is a *double quasi-Poisson algebra*. An element $\Phi_s \in e_s A e_s$ is said to satisfy the *multiplicative property for e_s* if, for any $a \in A$, we have

$$(4.2) \quad \{\!\{\Phi_s, a\}\!\} = \frac{1}{2} (ae_s \otimes \Phi_s - e_s \otimes \Phi_s a + a\Phi_s \otimes e_s - \Phi_s \otimes e_s a).$$

Then, we define a *multiplicative moment map* as an invertible element $\Phi = \sum_{s \in I} \Phi_s$ for $\Phi_s \in e_s A e_s$ which is such that for all $s \in I$, Φ_s satisfies the multiplicative property for e_s (4.2). The decomposition of Φ yields $e_s \Phi e_t = \delta_{st} \Phi_s$ for each $s, t \in I$, hence the invertibility condition implies that $\Phi_s^{-1} := e_s \Phi^{-1} e_s$ is an inverse for Φ_s in $e_s A e_s$. We call the triple $(A, \{\!\{-, -\}\!\}, \Phi)$ (or simply A if no confusion can arise) a *quasi-Hamiltonian algebra*.

Let A_1, A_2 be B -algebras with B -linear double brackets $\{\!\{-, -\}\!\}_1, \{\!\{-, -\}\!\}_2$. If $\psi : A_1 \rightarrow A_2$ is a morphism of double brackets and A_1, A_2 are double quasi-Poisson algebras, we say that ψ is a *morphism of double quasi-Poisson algebras*. If furthermore A_1, A_2 are quasi-Hamiltonian algebras with respective multiplicative moment maps Φ_1, Φ_2 satisfying $\psi(\Phi_1) = \Phi_2$, we say that ψ is a *morphism of quasi-Hamiltonian algebras*.

Example 4.1. — Consider the quiver Q with vertices $\{1, 2\}$ and one arrow $a : 1 \rightarrow 2$. Let \bar{Q} be its double obtained by adding the arrow $a^* : 2 \rightarrow 1$. Put $B = \mathbb{k}e_1 \oplus \mathbb{k}e_2$. Following Van den Bergh [VdB08a, § 6.5], the path algebra $\mathbb{k}\bar{Q}$ is a double quasi-Poisson algebra for the B -linear double bracket given on generators by

$$(4.3) \quad \begin{aligned} \{\!\!\{ a, a \}\!\!\}_1 &= 0, & \{\!\!\{ a^*, a^* \}\!\!\}_1 &= 0, \\ \{\!\!\{ a, a^* \}\!\!\}_1 &= e_2 \otimes e_1 + \frac{1}{2} (a^* a \otimes e_1 + e_2 \otimes a a^*). \end{aligned}$$

(We get $\{\!\!\{ a^*, a \}\!\!\}_1$ by cyclic antisymmetry (2.3).) To have a quasi-Hamiltonian algebra structure, we work in the localised algebra $A_1 = (\mathbb{k}\bar{Q})_{S_1}$ for $S_1 = \{1 + a a^*, 1 + a^* a\}$, where we introduce the multiplicative moment map $\Phi = (1 + a a^*)(1 + a^* a)^{-1}$. It can be decomposed as $\Phi = \Phi_1 + \Phi_2$ for $\Phi_1 = e_1 + a a^*$ and $\Phi_2 = (e_2 + a^* a)^{-1} := e_2(1 + a^* a)^{-1}e_2$.

Similarly, consider the quiver Q^{op} with vertices $\{1, 2\}$ and one arrow $b : 2 \rightarrow 1$, and let \bar{Q}^{op} be its double with new arrow $b^* : 1 \rightarrow 2$. We can also use Van den Bergh’s quasi-Hamiltonian structure on $A_2 = (\mathbb{k}\bar{Q}^{op})_{S_2}$, $S_2 = \{1 + b b^*, 1 + b^* b\}$, which is given by the B -linear double bracket

$$(4.4) \quad \begin{aligned} \{\!\!\{ b, b \}\!\!\}_2 &= 0, & \{\!\!\{ b^*, b^* \}\!\!\}_2 &= 0, \\ \{\!\!\{ b, b^* \}\!\!\}_2 &= e_1 \otimes e_2 + \frac{1}{2} (b^* b \otimes e_2 + e_1 \otimes b b^*). \end{aligned}$$

The multiplicative moment map $\Phi' = (1 + b b^*)(1 + b^* b)^{-1}$ can be decomposed as $\Phi' = (e_1 + b^* b)^{-1} + (e_2 + b b^*)$ as above.

Following Crawley–Boevey and Shaw [CBS06, Section 2], we introduce the isomorphism of B -algebras $\psi : A_1 \rightarrow A_2$ given by

$$(4.5) \quad \psi(a) = b^*, \quad \psi(a^*) = -(1 + b b^*)^{-1}b.$$

It is indeed a morphism since

$$(4.6) \quad \begin{aligned} \psi(1 + a^* a) &= 1 - (1 + b b^*)^{-1}b b^* = (1 + b b^*)^{-1}, \\ \psi(1 + a a^*) &= 1 - b^*(1 + b b^*)^{-1}b = (1 + b^* b)^{-1}, \end{aligned}$$

are both invertible. It is an isomorphism since the map $\theta : A_2 \rightarrow A_1$ given by

$$\theta(b) = -a^*(1 + a a^*)^{-1}, \quad \theta(b^*) = a,$$

is its inverse. We have, in fact, that $\psi : A_1 \rightarrow A_2$ is an isomorphism of quasi-Hamiltonian algebras. Using (4.6), we easily see that $\psi(\Phi_1) = \Phi'_1$ and $\psi(\Phi_2) = \Phi'_2$, so that we only need to show that ψ is a morphism of double brackets. By (2.9), this would follow from the identities

$$(4.7) \quad \begin{aligned} \psi^{\otimes 2} \{\!\!\{ a, a \}\!\!\}_1 &= \{\!\!\{ b^*, b^* \}\!\!\}_2, \\ \psi^{\otimes 2} \{\!\!\{ a^*, a^* \}\!\!\}_1 &= \{\!\!\{ (e_2 + b b^*)^{-1} b, (e_2 + b b^*)^{-1} b \}\!\!\}_2, \\ \psi^{\otimes 2} \{\!\!\{ a, a^* \}\!\!\}_1 &= -\{\!\!\{ b^*, (e_2 + b b^*)^{-1} b \}\!\!\}_2. \end{aligned}$$

The first equality in (4.7) is trivial as both sides vanish. We leave the proof of the second identity in (4.7) to the reader since it is similar to the third one which we

check now. Combining (4.3) and the definition of ψ (4.5), we have that

$$\begin{aligned}\psi^{\otimes 2} \{ \{ a, a^* \} \}_1 &= e_2 \otimes e_1 - \frac{1}{2} (e_2 + bb^*)^{-1} bb^* \otimes e_1 - \frac{1}{2} e_2 \otimes b^* (1 + bb^*)^{-1} b \\ &= (e_2 + bb^*)^{-1} \otimes e_1 + \frac{1}{2} (e_2 + bb^*)^{-1} bb^* \otimes e_1 - \frac{1}{2} e_2 \otimes b^* (e_2 + bb^*)^{-1} b.\end{aligned}$$

Since Φ' is a multiplicative moment map, we have from the multiplicative property (4.2) for e_2 that

$$\begin{aligned}\{ \{ b^*, (e_2 + bb^*)^{-1} \} \}_2 &= -(\Phi'_2)^{-1} \{ \{ b^*, \Phi'_2 \} \} (\Phi'_2)^{-1} = +(\Phi'_2)^{-1} (\{ \{ \Phi'_2, b^* \} \})^\circ (\Phi'_2)^{-1} \\ &= \frac{1}{2} (\Phi'_2)^{-1} (\Phi'_2 \otimes b^* e_2 - \Phi'_2 b^* \otimes e_2 + e_2 \otimes b^* \Phi'_2 - e_2 b^* \otimes \Phi'_2) (\Phi'_2)^{-1} \\ &= \frac{1}{2} (e_2 \otimes b^* (\Phi'_2)^{-1} + (\Phi'_2)^{-1} \otimes b^*),\end{aligned}$$

where $(\Phi'_2)^{-1} := e_2 (\Phi')^{-1} e_2 = (e_2 + bb^*)^{-1}$ and we used that $b^* \in e_1 A_2 e_2$. We thus get

$$\begin{aligned}- \{ \{ b^*, (e_2 + bb^*)^{-1} b \} \}_2 &= -(\Phi'_2)^{-1} \{ \{ b^*, b \} \}' - \{ \{ b^*, (\Phi'_2)^{-1} \} \}' b \\ &= (\Phi'_2)^{-1} \otimes e_1 + \frac{1}{2} (\Phi'_2)^{-1} bb^* \otimes e_1 - \frac{1}{2} e_2 \otimes b^* (\Phi'_2)^{-1} b,\end{aligned}$$

which coincides with $\psi^{\otimes 2} \{ \{ a, a^* \} \}_1$.

4.1.2. First properties

The following result can be proved in the same way as Lemma 3.2.

LEMMA 4.2. — *Let $\psi : A_1 \rightarrow A_2$ be a morphism of double brackets.*

- (1) *Assume that ψ is surjective as a B -algebra homomorphism.*
 - (a) *If A_1 is a double quasi-Poisson algebra, then A_2 is a double quasi-Poisson algebra.*
 - (b) *If A_1 is a quasi-Hamiltonian algebra, then A_2 admits a structure of quasi-Hamiltonian algebra such that ψ is a morphism of quasi-Hamiltonian algebras.*
- (2) *Assume that ψ is injective as a B -algebra homomorphism.*
 - (a) *If A_2 is a double quasi-Poisson algebra, then A_1 is a double quasi-Poisson algebra.*
 - (b) *If A_2 is a quasi-Hamiltonian algebra with multiplicative moment map Φ' such that $\Phi', (\Phi')^{-1} \in \text{Im}(\psi)$, then A_1 admits a structure of quasi-Hamiltonian algebra such that ψ is a morphism of quasi-Hamiltonian algebras.*
- (3) *All the structures obtained in (1)–(2) are unique.*

In analogy with the Hamiltonian case, a morphism of double brackets between two quasi-Hamiltonian algebras may fail to be a morphism of quasi-Hamiltonian algebras: we can rescale multiplicative moment maps. In fact, such a rescaling always exists if the morphism of double brackets is an isomorphism.

LEMMA 4.3. — *Let $\psi : A_1 \rightarrow A_2$ be an isomorphism of double quasi-Poisson algebras. If A_1, A_2 are quasi-Hamiltonian algebras with multiplicative moment maps Φ_1, Φ_2 , then $\Phi_2 = \nu \psi(\Phi_1)$ for some $\nu \in B^\times$. In particular, a multiplicative moment map is unique up to multiplication by an element of B^\times .*

Proof. — Let us decompose $\Phi_1 \in A_1$ with the idempotents as $\sum_{s \in I} \Phi_{1,s}$ and do the same for $\Phi_2 \in A_2$. We first note that, for any $\lambda_s \in \mathbb{k}$, the element $\Phi_{2,s} - \lambda_s \psi(\Phi_{1,s})$ satisfies the multiplicative property for e_s (4.2) because for any $a \in A_2$,

$$\{\{\Phi_{2,s} - \lambda_s \psi(\Phi_{1,s}), a\}\}_2 = \{\{\Phi_{2,s}, a\}\}_2 - \lambda_s \psi^{\otimes 2} \{\{\Phi_{1,s}, \psi^{-1}(a)\}\}_1,$$

and we can use (4.2) for the two double brackets on the right-hand side since Φ_1, Φ_2 are multiplicative moment maps. For any $\lambda = \sum_s \lambda_s e_s \in B$, introduce the element $\Phi^{(\lambda)} := \Phi_2 - \lambda \Phi_1$. Thus, for any $\lambda, \kappa \in B$, we get

$$\begin{aligned} (4.8) \quad & \{\{\Phi^{(\lambda)}, \Phi^{(\kappa)}\}\}_2 \\ &= \frac{1}{2} \sum_{s \in I} \left(\Phi_s^{(\kappa)} \otimes \Phi_s^{(\lambda)} - e_s \otimes \Phi_s^{(\lambda)} \Phi_s^{(\kappa)} + \Phi_s^{(\kappa)} \Phi_s^{(\lambda)} \otimes e_s - \Phi_s^{(\lambda)} \otimes \Phi_s^{(\kappa)} \right) \\ &= \frac{1}{2} \sum_{s \in I} \left(-\Phi_s^{(\kappa)} \otimes \Phi_s^{(\lambda)} + \Phi_s^{(\kappa)} \Phi_s^{(\lambda)} \otimes e_s - e_s \otimes \Phi_s^{(\lambda)} \Phi_s^{(\kappa)} + \Phi_s^{(\lambda)} \otimes \Phi_s^{(\kappa)} \right), \end{aligned}$$

where we used for the first (resp. second) equality that $\Phi_s^{(\lambda)}$ (resp. $\Phi_s^{(\kappa)}$) satisfies the multiplicative property for e_s (4.2). By decomposing (4.8) in terms of idempotents, we get that $\Phi_s^{(\kappa)} \otimes \Phi_s^{(\lambda)} = \Phi_s^{(\lambda)} \otimes \Phi_s^{(\kappa)}$ for all $s \in I$, which by definition of these elements is equivalent to

$$(\lambda_s - \kappa_s) \left(\psi(\Phi_{1,s}) \otimes \Phi_{2,s} - \Phi_{2,s} \otimes \psi(\Phi_{1,s}) \right) = 0.$$

Multiplying on both side with $\Phi_{2,s}^{-1} := e_s \Phi_{2,s}^{-1} e_s$, this is equivalent when $\lambda - \kappa \in B^\times$ to

$$\Phi_{2,s}^{-1} \psi(\Phi_{1,s}) \otimes e_s = e_s \otimes \psi(\Phi_{1,s}) \Phi_{2,s}^{-1} \in e_s A_2 e_s \otimes e_s A_2 e_s.$$

Thus, $\Phi_{2,s} = \nu_s \psi(\Phi_{1,s})$ for some $\nu_s \in \mathbb{k}^\times$.

For the second part of the statement, it suffices to consider $\psi = \text{id}_{A_1} : A_1 \rightarrow A_1$. \square

4.1.3. Sum of double quasi-Poisson algebras

LEMMA 4.4 ([VdB08a]). — *If A_1 and A_2 are endowed with double quasi-Poisson brackets, then the double bracket $\{\{-, -\}\}^\oplus$ on $A_1 \oplus A_2$ defined in §2.2.4 is quasi-Poisson. If Φ_1, Φ_2 are multiplicative moment maps in A_1 and A_2 , then $A_1 \oplus A_2$ is quasi-Hamiltonian with (Φ_1, Φ_2) as multiplicative moment map.*

4.2. Double quasi-Poisson brackets and fusion

We assume that the index set I of B is such that $|I| > 1$. Our aim is to adapt the results from §3.2 to the quasi-Hamiltonian setting. The main difficulty that we will encounter is that the double bracket given by (2.14) obtained by fusion from a double

quasi-Poisson algebra *is not* a double quasi-Poisson bracket in general, because the associated triple bracket does not satisfy (4.1). To overcome this issue, one needs to add a corrective term to the double bracket after fusion, see Proposition 4.5.

4.2.1. Fusion in an algebra

Recall the fusion algebra $A^f = A_{e_2 \rightarrow e_1}^f$ defined in § 2.2.1. The following Proposition was first observed by Van den Bergh under mild assumptions [VdB08a, Theorems 5.3.1, 5.3.2] and is proved in full generalities in [Fai21, Theorems 2.14, 2.15].

PROPOSITION 4.5. — *If A is a double quasi-Poisson algebra over B , then A^f is a double quasi-Poisson algebra over $B^f = \bigoplus_{s \in I \setminus \{2\}} \mathbb{k}e_s$. The double quasi-Poisson bracket in A^f is given by*

$$(4.9) \quad \{\{-, -\}^f := \{\{-, -\}_{ind} + \{\{-, -\}_{fus},$$

where the first double bracket on the right-hand side is induced in A^f by the one of A (see (2.14)), and the second double bracket $\{\{-, -\}_{fus}$ is defined in Appendix A. Furthermore, if Φ is a multiplicative moment map for A and Φ_s^f denotes the projection of the element $\Phi_s = e_s \Phi e_s$ under the map (2.12) for each $s \in I$, then $\Phi^{ff} = \Phi_1^f \Phi_2^f + \sum_{s \neq 1, 2} \Phi_s^f$ is a multiplicative moment map for A^f .

For the remainder of this section, if A is a double quasi-Poisson algebra and \tilde{A} is an algebra obtained from A by performing a finite number of fusions, we will see \tilde{A} as a double quasi-Poisson algebra using this last result. More precisely, if we define \tilde{A} using the following chain of algebras obtained by fusion

$$(4.10) \quad \begin{aligned} &A^{(0)} \longrightarrow A^{(1)} \longrightarrow \dots \longrightarrow A^{(k)} \longrightarrow \dots \longrightarrow A^{(n)}, \quad \text{where} \\ &A^{(0)} = A, \quad A^{(n)} = \tilde{A}, \quad A^{(k)} = \left(A^{(k-1)}\right)_{e_{j_k} \rightarrow e_{i_k}}^f \quad \text{for } k = 1, \dots, n, \end{aligned}$$

each algebra $A^{(k)}$ in the chain (4.10) is a double quasi-Poisson algebra whose double quasi-Poisson bracket is obtained from $A^{(k-1)}$ by using Proposition 4.5. If A is a quasi-Hamiltonian algebra, we see \tilde{A} as a quasi-Hamiltonian algebra using the same argument.

Example 4.6. — Consider the localised path algebra $A = \mathbb{k}\bar{Q}_S$ considered in Example 4.1 with its quasi-Hamiltonian algebra structure. Fusing e_2 onto e_1 , we get a quasi-Hamiltonian algebra structure on A^f which can be identified with the localised free algebra $\mathbb{k}\langle a, a^* \rangle_S$, $S = \{1 + aa^*, 1 + a^*a\}$. The double bracket is determined by

$$(4.11) \quad \begin{aligned} \{\{a, a\}^f &= \frac{1}{2} (a^2 \otimes 1 - 1 \otimes a^2), \quad \{\{a^*, a^*\}^f = -\frac{1}{2} ((a^*)^2 \otimes 1 - 1 \otimes (a^*)^2), \\ \{\{a, a^*\}^f &= e_2 \otimes e_1 + \frac{1}{2} (a^*a \otimes e_1 + e_2 \otimes aa^* + a^* \otimes a - a \otimes a^*), \end{aligned}$$

while the multiplicative moment map is $\Phi = (1 + aa^*)(1 + a^*a)^{-1}$. Indeed, to get the double brackets (4.11) we add to (4.3) the terms given by $\{\{-, -\}_{fus}$, so we

add (A.3c) to $\{\{a, a\}\}$, (A.2b) to $\{\{a^*, a^*\}\}$, and (A.3b) to $\{\{a, a^*\}\}$. The algebra A^f is the localised path algebra of the double of the one-loop quiver, and the quasi-Hamiltonian structure obtained by fusion is the one of Van den Bergh [VdB08a, § 6.5].

LEMMA 4.7. — *Let $(A, \{\{-, -\}\})$ be a double quasi-Poisson algebra over B . Consider the algebra $A_1 = A^f_{e_2 \rightarrow e_1}$ obtained by fusing e_2 onto e_1 and the algebra $A_2 = A^f_{e_1 \rightarrow e_2}$ obtained by fusing e_1 onto e_2 , which are endowed with a double quasi-Poisson bracket by Proposition 4.5.*

Assume that there exists an element $\Phi_2 \in e_2 A e_2$ invertible in $e_2 A e_2$ which satisfies the multiplicative property (4.2) for e_2 . Then there exists an isomorphism of double quasi-Poisson algebras $A_1 \rightarrow A_2$. If furthermore A is quasi-Hamiltonian, then the isomorphism $A_1 \rightarrow A_2$ is an isomorphism of quasi-Hamiltonian algebras.

The last part of this statement was announced in [Fai21, Proposition 3.8] as a non-commutative version of [AKSM02, Proposition 5.7].

Proof of Lemma 4.7. We first define an isomorphism $\psi : A_1 \rightarrow A_2$. (This is a morphism over the common base $B' = \mathbb{k}\tilde{e} \oplus \bigoplus_{s \in I \setminus \{1,2\}} \mathbb{k}e_s$, see the remark after Lemma 3.7.) As in Lemma 2.4, we define ψ on a specialisation of the set of generators of A_1 , i.e. on a suitably chosen subset of the generators given in Lemma 2.2. For $\hat{e} = 1 - e_1 - e_2$, ψ is given on generators of first type (2.13a) as

$$\begin{aligned} \psi(t) &= t, & \text{if } t \in \hat{e}A\hat{e}; & & \psi(t) &= \Phi_2 e_{21} t, & \text{if } t \in e_1 A \hat{e}; \\ \psi(t) &= t e_{12} \Phi_2^{-1}, & \text{if } t \in \hat{e}Ae_1; & & \psi(t) &= \Phi_2 e_{21} t e_{12} \Phi_2^{-1}, & \text{if } t \in e_1 Ae_1; \end{aligned}$$

on generators of second type (2.13b) as

$$\psi(e_{12}u) = u, \quad \text{if } u \in e_2 A \hat{e}; \quad \psi(e_{12}u) = u e_{12} \Phi_2^{-1}, \quad \text{if } u \in e_2 Ae_1;$$

on generators of third type (2.13c) as

$$\psi(v e_{21}) = v, \quad \text{if } v \in \hat{e}Ae_2; \quad \psi(v e_{21}) = \Phi_2 e_{21} v, \quad \text{if } v \in e_1 Ae_2;$$

on generators of fourth type (2.13d) as

$$\psi(e_{12}w e_{21}) = w, \quad \text{if } w \in e_2 Ae_2.$$

Similarly, we define a morphism $\theta : A_2 \rightarrow A_1$ on a specialisation of the set of generators of A_2 , which is given on generators of first type (2.13a) as

$$\begin{aligned} \theta(t') &= t', & \text{if } t' \in \hat{e}A\hat{e}; & & \theta(t') &= e_{12}t', & \text{if } t' \in e_2 A \hat{e}; \\ \theta(t') &= t' e_{21}, & \text{if } t' \in \hat{e}Ae_2; & & \theta(t') &= e_{12}t' e_{21}, & \text{if } t' \in e_2 Ae_2; \end{aligned}$$

on generators of second type (2.13b) as

$$\begin{aligned} \theta(e_{21}u') &= e_{12} \Phi_2^{-1} e_{21} u', & \text{if } u' \in e_1 A \hat{e}; \\ \theta(e_{21}u') &= e_{12} \Phi_2^{-1} e_{21} u' e_{21}, & \text{if } u' \in e_1 Ae_2; \end{aligned}$$

on generators of third type (2.13c) as

$$\begin{aligned} \theta(v' e_{12}) &= v' e_{12} \Phi_2 e_{21}, & \text{if } v' \in \hat{e}Ae_1; \\ \theta(v' e_{12}) &= e_{12} v' e_{12} \Phi_2 e_{21}, & \text{if } v' \in e_2 Ae_1; \end{aligned}$$

on generators of fourth type (2.13d) as

$$\theta(e_{21}w'e_{12}) = e_{12}\Phi_2^{-1}e_{21}w'e_{12}\Phi_2e_{21}, \quad \text{if } w' \in e_1Ae_1.$$

It is then a straightforward exercise to show that $\psi \circ \theta = \text{Id}_{A_2}$ and $\theta \circ \psi = \text{Id}_{A_1}$. For example, if $w' \in e_1Ae_1$,

$$\begin{aligned} \psi \circ \theta(e_{21}w'e_{12}) &= \psi(e_{12}\Phi_2^{-1}e_{21})\psi(w')\psi(e_{12}\Phi_2e_{21}) \\ &= \Phi_2^{-1}(\Phi_2e_{21}w'e_{12}\Phi_2^{-1})\Phi_2 = e_{21}w'e_{12}. \end{aligned}$$

We easily see that in the quasi-Hamiltonian case we get $\psi(\Phi_1e_{12}\Phi_2e_{21}) = \Phi_2e_{21}\Phi_1e_{12}$, so the only property that remains to be shown is that $\psi : A_1 \rightarrow A_2$ is a morphism of double quasi-Poisson algebras.

We note from Proposition 4.5 that the double quasi-Poisson bracket on A_1 is given by

$$(4.12) \quad \{\{-, -\}_1 := \{\{-, -\}_{ind}^{2 \rightarrow 1} + \{\{-, -\}_{fus}^{2 \rightarrow 1},$$

where $\{\{-, -\}_{ind}^{2 \rightarrow 1}$ is induced by the double bracket $\{\{-, -\}$ in A_1 using (2.14), while $\{\{-, -\}_{fus}^{2 \rightarrow 1}$ is the double bracket defined in Appendix A for $i = 1, j = 2$; the double quasi-Poisson bracket on A_2 is given by

$$(4.13) \quad \{\{-, -\}_2 := \{\{-, -\}_{ind}^{1 \rightarrow 2} + \{\{-, -\}_{fus}^{1 \rightarrow 2},$$

where $\{\{-, -\}_{ind}^{1 \rightarrow 2}$ is induced by the double bracket $\{\{-, -\}$ in A_2 using (2.14), while $\{\{-, -\}_{fus}^{1 \rightarrow 2}$ is the double bracket defined in Appendix A for $i = 2, j = 1$. Hence, we need to check that

$$(4.14) \quad (\psi \otimes \psi) \{\{c, d\}_1 = \{\{\psi(c), \psi(d)\}_2,$$

on each pair (c, d) of specialisations of generators of A_1 described above. There are 9 such specialisations, so making use of the cyclic antisymmetry we need to check (4.14) in 45 cases. Let us explain how to carry out the computations in one case; the remaining cases are treated in a similar way, and the corresponding double brackets that must be computed are gathered in Appendix B for the reader's convenience.

Consider $c = a, d = b$, where $a, b \in e_1A\hat{e}$. We have that c, d are generators of first type (2.13a) in A_1 , so that $\{\{c, d\}_{fus}^{2 \rightarrow 1} = 0$ by (A.1a) (for $i = 1, j = 2$). Also, we have by (2.14) that

$$\{\{c, d\}_{ind}^{2 \rightarrow 1} = e_1 * e_1 \{\{a, b\} \hat{e} * \hat{e} = \{\{a, b\},$$

since $\{\{a, b\}' = e_1 \{\{a, b\}' \hat{e}$ and the same holds for $\{\{a, b\}''$. Hence, by definition of (4.12) we can simply write the left-hand side of (4.14) as $\psi^{\otimes 2} \{\{a, b\}$. Meanwhile, we have that $\psi(c) = \Phi_2e_{21}a$ and $\psi(d) = \Phi_2e_{21}b$. Therefore, we get

$$(4.15) \quad \begin{aligned} \{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2} &= + \{\{\Phi_2, \Phi_2\}_{fus}^{1 \rightarrow 2} e_{21}b * e_{21}a + \Phi_2 * \{\{e_{21}a, \Phi_2\}_{fus}^{1 \rightarrow 2} e_{21}b \\ &+ \Phi_2 \{\{\Phi_2, e_{21}b\}_{fus}^{1 \rightarrow 2} * e_{21}a + \Phi_2 * \Phi_2 \{\{e_{21}a, e_{21}b\}_{fus}^{1 \rightarrow 2}, \end{aligned}$$

using the derivation rules of the double bracket $\{\{-, -\}_{fus}^{1 \rightarrow 2}$. We now use the explicit form of this double bracket given in Appendix A (for $i = 2, j = 1$). To do so, note that in A_2 , Φ_2 is a generator of first type (2.13a) while $e_{21}a, e_{21}b$ are generators of

second type (2.13b). Hence the double bracket appearing in the first term of (4.15) is given by (A.1a), the one in the second term by (A.2a), the one in the third term by (A.1b), and the one in the fourth term by (A.2b). Therefore,

$$\begin{aligned}
 \{\!\!\{ \psi(c), \psi(d) \}\!\!\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} \Phi_2 * (e_{21}a \otimes \Phi_2 - \Phi_2 e_{21}a \otimes e_2) e_{21}b \\
 (4.16) \qquad \qquad \qquad &+ \frac{1}{2} \Phi_2 (e_2 \otimes \Phi_2 e_{21}b - \Phi_2 \otimes e_{21}b) * e_{21}a \\
 &+ \frac{1}{2} \Phi_2 * \Phi_2 (e_2 \otimes e_{21}a e_{21}b - e_{21}b e_{21}a \otimes e_2) .
 \end{aligned}$$

Noting that $ae_2 = 0 = be_2$, we get that the last two terms of (4.16) vanish, and since the second and third terms cancel out this yields

$$(4.17) \qquad \{\!\!\{ \psi(c), \psi(d) \}\!\!\}_{fus}^{1 \rightarrow 2} = \frac{1}{2} (e_{21}a \otimes \Phi_2^2 e_{21}b - \Phi_2^2 e_{21}a \otimes e_{21}b) .$$

Finally, we compute that

$$\begin{aligned}
 \{\!\!\{ \psi(c), \psi(d) \}\!\!\}_{ind}^{1 \rightarrow 2} &= + \{\!\!\{ \Phi_2, \Phi_2 \}\!\!\}_{ind}^{1 \rightarrow 2} e_{21}b * e_{21}a + \Phi_2 * \{\!\!\{ e_{21}a, \Phi_2 \}\!\!\}_{ind}^{1 \rightarrow 2} e_{21}b \\
 (4.18) \qquad \qquad \qquad &+ \Phi_2 \{\!\!\{ \Phi_2, e_{21}b \}\!\!\}_{ind}^{1 \rightarrow 2} * e_{21}a + \Phi_2 * \Phi_2 \{\!\!\{ e_{21}a, e_{21}b \}\!\!\}_{ind}^{1 \rightarrow 2} \\
 &= + \{\!\!\{ \Phi_2, \Phi_2 \}\!\!\} e_{21}b * e_{21}a + \Phi_2 e_{21} * \{\!\!\{ a, \Phi_2 \}\!\!\} e_{21}b \\
 &+ \Phi_2 e_{21} \{\!\!\{ \Phi_2, b \}\!\!\} * e_{21}a + \Phi_2 e_{21} * \Phi_2 e_{21} \{\!\!\{ a, b \}\!\!\} ,
 \end{aligned}$$

where we have used the derivation rules of the double bracket $\{\!\!\{ -, - \}\!\!\}_{ind}^{1 \rightarrow 2}$ for the first equality, and its definition (2.14) using $\{\!\!\{ -, - \}\!\!\}$ for the second equality. By assumption, Φ_2 satisfies the multiplicative property (4.2) for e_2 in A , hence we can compute the double brackets in the first three terms of (4.18). We get that (since $e_2a = 0 = ae_2$ and the same holds for b , $\{\!\!\{ a, \Phi_2 \}\!\!\} = 0 = \{\!\!\{ \Phi_2, b \}\!\!\}$ in this case)

$$\begin{aligned}
 (4.19) \qquad \{\!\!\{ \psi(c), \psi(d) \}\!\!\}_{ind}^{1 \rightarrow 2} &= \frac{1}{2} (\Phi_2^2 e_{21}a \otimes e_{21}b - e_{21}a \otimes \Phi_2^2 e_{21}b) + \Phi_2 e_{21} \{\!\!\{ a, b \}\!\!\}' \otimes \Phi_2 e_{21} \{\!\!\{ a, b \}\!\!\}'' .
 \end{aligned}$$

Combining (4.17) and (4.19), we get that

$$(4.20) \qquad \{\!\!\{ \psi(c), \psi(d) \}\!\!\}_2 = \Phi_2 e_{21} \{\!\!\{ a, b \}\!\!\}' \otimes \Phi_2 e_{21} \{\!\!\{ a, b \}\!\!\}'' = \psi^{\otimes 2} \{\!\!\{ a, b \}\!\!\} ,$$

by definition of ψ , so this is precisely $\psi^{\otimes 2} \{\!\!\{ c, d \}\!\!\}_1$. □

4.2.2. Fusion of several idempotents

LEMMA 4.8. — *Let A be a double quasi-Poisson algebra over B , and let $e_1, e_2, e_3 \in B$ be orthogonal idempotents. Let $A_1 := (A_{e_3 \rightarrow e_2}^f)_{e_2 \rightarrow e_1}^f$ (resp. $A_2 := (A_{e_2 \rightarrow e_1}^f)_{e_3 \rightarrow e_1}^f$) be the algebra obtained by fusing e_3 onto e_1 , then e_2 onto e_1 (resp. e_2 onto e_1 , then e_3 onto e_1). Then the identity map on A induces an isomorphism of double quasi-Poisson algebras $\psi : A_1 \rightarrow A_2$ over $B' = \bigoplus_{s \in I \setminus \{2, 3\}} \mathbb{k}e_s$. If furthermore A is quasi-Hamiltonian, then ψ is an isomorphism of quasi-Hamiltonian algebras.*

Proof. — The map $\psi : A_1 \rightarrow A_2$ induced by the identity is explicitly spelled out in the proof of Lemma 3.8. In particular, if $\{\{-, -\}\}_{k, ind}$ denotes the double bracket from A induced in A_k , $k = 1, 2$, the above-mentioned proof also implies that

$$(4.21) \quad \{\{\psi(c), \psi(d)\}\}_{2, ind} = (\psi \otimes \psi) \{\{c, d\}\}_{1, ind}, \quad \text{for all } c, d \in A_1.$$

By construction A_1 and A_2 are obtained by the following composition of fusions

$$(4.22) \quad A \xrightarrow{e_3 \rightarrow e_2} A_{e_3 \rightarrow e_2}^f \xrightarrow{e_2 \rightarrow e_1} A_1, \quad A \xrightarrow{e_2 \rightarrow e_1} A_{e_2 \rightarrow e_1}^f \xrightarrow{e_3 \rightarrow e_1} A_2.$$

Thus, using Proposition 4.5 twice, the double quasi-Poisson bracket on A_1 is given by

$$(4.23) \quad \{\{-, -\}\}_1 := \{\{-, -\}\}_{1, ind} + \{\{-, -\}\}_{1, fus}^{3 \rightarrow 2} + \{\{-, -\}\}_{1, fus}^{2 \rightarrow 1}.$$

Here, $\{\{-, -\}\}_{1, fus}^{3 \rightarrow 2}$ denotes the double bracket from Appendix A which is added after the fusion $e_3 \rightarrow e_2$ to get a double quasi-Poisson bracket on $A_{e_3 \rightarrow e_2}^f$ and is then induced in A_1 ; $\{\{-, -\}\}_{1, fus}^{2 \rightarrow 1}$ denotes the double bracket from Appendix A which is added after the fusion $e_2 \rightarrow e_1$ to get a double quasi-Poisson bracket on A_1 . In the same way, the double quasi-Poisson bracket on A_2 is given by

$$(4.24) \quad \{\{-, -\}\}_2 := \{\{-, -\}\}_{2, ind} + \{\{-, -\}\}_{2, fus}^{2 \rightarrow 1} + \{\{-, -\}\}_{2, fus}^{3 \rightarrow 1},$$

where the double brackets are defined in analogy to the case of A_1 . In particular, we get from (4.21) that the map $\psi : A_1 \rightarrow A_2$ is an isomorphism of double quasi-Poisson algebras provided that for all $c, d \in A_1$,

$$(4.25) \quad \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} + \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} \\ = (\psi \otimes \psi) \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} + (\psi \otimes \psi) \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1}.$$

This is proved in Appendix C. If A is quasi-Hamiltonian with multiplicative moment map $\Phi = \sum_s \Phi_s$ where $\Phi_s = e_s \Phi_s e_s$, then using Proposition 4.5 the multiplicative moment map transforms under the morphisms (4.22) as

$$\sum_s \Phi_s \mapsto e_2 \Phi_2 e_{23} \Phi_3 e_{32} + \sum_{s \neq 2, 3} \Phi_s \mapsto e_1 \Phi_1 e_{12} \Phi_2 e_{23} \Phi_3 e_{32} e_{21} + \sum_{s \neq 1, 2, 3} \Phi_s \in A_1, \\ \sum_s \Phi_s \mapsto e_1 \Phi_1 e_{12} \Phi_2 e_{21} + \sum_{s \neq 1, 2} \Phi_s \mapsto e_1 \Phi_1 e_{12} \Phi_2 e_{21} e_{13} \Phi_3 e_{31} + \sum_{s \neq 1, 2, 3} \Phi_s \in A_2.$$

We can see that ψ sends the multiplicative moment map of A_1 to the one of A_2 . \square

LEMMA 4.9. — *Let A be a double Poisson algebra over B , and let $e_1, e_2, e_3, e_4 \in B$ be orthogonal idempotents. Let $A_1 := (A_{e_4 \rightarrow e_3}^f)_{e_2 \rightarrow e_1}^f$ (resp. $A_2 := (A_{e_2 \rightarrow e_1}^f)_{e_4 \rightarrow e_3}^f$) be the algebra obtained by fusing e_4 onto e_3 , then e_2 onto e_1 (resp. e_2 onto e_1 , then e_4 onto e_3). Then the identity map on A induces an isomorphism of double quasi-Poisson algebras $\psi : A_1 \rightarrow A_2$ over $B' = \bigoplus_{s \in I \setminus \{2, 4\}} \mathbb{k}e_s$. If furthermore A is quasi-Hamiltonian, then ψ is an isomorphism of quasi-Hamiltonian algebras.*

Proof. — As in Lemma 3.9, we note that we can write generators of A_1 and A_2 in the form

$$(4.26) \quad e_a a f_a, \quad \text{for } a \in A, \quad e_a \in \{1 - e_4 - e_2, e_{12}, e_{34}\}, \\ f_a \in \{1 - e_4 - e_2, e_{21}, e_{43}\}.$$

Let us prove that the morphism $\psi : A_1 \rightarrow A_2$ that we seek is simply the identity when written in terms of these elements. Using Proposition 4.5 twice, the double quasi-Poisson bracket on A_1 is given by

$$(4.27) \quad \{\!\!\{-, -\}\!\!\}_1 := \{\!\!\{-, -\}\!\!\}_{1, ind} + \{\!\!\{-, -\}\!\!\}_{1, fus}^{4 \rightarrow 3} + \{\!\!\{-, -\}\!\!\}_{1, fus}^{2 \rightarrow 1} .$$

Here, $\{\!\!\{-, -\}\!\!\}_{1, ind}$ denotes the double bracket from A induced in A_1 ; $\{\!\!\{-, -\}\!\!\}_{1, fus}^{4 \rightarrow 3}$ denotes the double bracket from Appendix A which is added after the fusion $e_4 \rightarrow e_3$ to get a double quasi-Poisson bracket on $A_{e_4 \rightarrow e_3}^f$ and is then induced in A_1 ; $\{\!\!\{-, -\}\!\!\}_{1, fus}^{2 \rightarrow 1}$ denotes the double bracket from Appendix A which is added after the fusion $e_2 \rightarrow e_1$ to get a double quasi-Poisson bracket on A_1 . In the same way, we can write

$$(4.28) \quad \{\!\!\{-, -\}\!\!\}_2 := \{\!\!\{-, -\}\!\!\}_{2, ind} + \{\!\!\{-, -\}\!\!\}_{2, fus}^{2 \rightarrow 1} + \{\!\!\{-, -\}\!\!\}_{2, fus}^{4 \rightarrow 3} .$$

By Lemma 3.9, we directly have that

$$(4.29) \quad \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, ind} = (\psi \otimes \psi) \{\!\!\{c, d\}\!\!\}_{1, ind} , \quad \text{for all } c, d \in A_1 ,$$

and we note that if we can show for all $c, d \in A_1$

$$(4.30) \quad \begin{aligned} \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{4 \rightarrow 3} &= (\psi \otimes \psi) \{\!\!\{c, d\}\!\!\}_{1, fus}^{4 \rightarrow 3} , \quad \text{and} \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= (\psi \otimes \psi) \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} , \end{aligned}$$

then $\psi : A_1 \rightarrow A_2$ will be a morphism of double quasi-Poisson algebras. It suffices to prove one of the two equalities in (4.30) since the other follows by symmetry. Checking that such an equality holds is easy on generators of the form (4.26). For example, if we take $c = e_{34}a\hat{e}$ and $d = e_{12}be_{43}$ where $\hat{e} = 1 - e_2 - e_4$, we have in A_1 that

$$\begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \{\!\!\{(e_{34}a), e_{12}(be_{43})\}\!\!\}_{1, fus}^{2 \rightarrow 1} \\ &= \frac{1}{2} \left(e_1 \otimes (e_{34}a) e_{12}(be_{43}) - e_1(e_{34}a) \otimes e_{12}(be_{43}) \right) = \frac{1}{2} e_1 \otimes cd , \end{aligned}$$

after using (A.1b) since $e_{34}a\hat{e} \in A_1$ is a generator of first type, while $e_{12}(be_{43}) \in A_1$ is a generator of second type. Meanwhile, we have in A_2 that

$$\begin{aligned} \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= e_{34} * \{\!\!\{a, e_{12}b\}\!\!\}_{2, fus}^{2 \rightarrow 1} e_{43} \\ &= \frac{1}{2} e_{34} * (e_1 \otimes ae_{12}b - e_1a \otimes e_{12}b) e_{43} = \frac{1}{2} e_1 \otimes \psi(c)\psi(d) , \end{aligned}$$

where we used that the double bracket $\{\!\!\{-, -\}\!\!\}_{2, fus}^{2 \rightarrow 1}$ is induced in A_2 from the double bracket in $(A_{e_2 \rightarrow e_1})^f$ given in Appendix A with $j = 2, i = 1$.

If A is quasi-Hamiltonian, the multiplicative moment map in A_1 and A_2 can be written as

$$\Phi = \Phi_1 e_{12} \Phi_2 e_{21} + \Phi_3 e_{34} \Phi_4 e_{43} + \sum_{s \neq 1, 2, 3, 4} \Phi_s ,$$

so ψ is obviously a morphism of quasi-Hamiltonian algebras. □

We can now derive the quasi-Poisson version of Theorem 3.10 by reproducing its proof with Lemmas 3.7, 3.8 and 3.9 replaced by Lemmas 4.7, 4.8 and 4.9 respectively. We use the notion of fusion algebra respecting a partition defined before

Theorem 3.10. Such an algebra inherits a double quasi-Poisson bracket by repeated use of Proposition 4.5.

THEOREM 4.10. — *Let A be a double quasi-Poisson algebra over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. Assume that A_1, A_2 are fusion algebras respecting a partition $I = \sqcup_{j \in J} I_j$.*

Let $J' = \{j \in J \mid |I_j| > 1\}$, and assume that for each $j' \in J'$ and $k \in I_{j'}$ there exists $\Phi_k \in e_k A e_k$ invertible in $e_k A e_k$ which satisfies the multiplicative property (4.2) for e_k . Then, there is a double quasi-Poisson algebra isomorphism $\psi : A_1 \rightarrow A_2$. If furthermore A is quasi-Hamiltonian, then the morphism $\psi : A_1 \rightarrow A_2$ is an isomorphism of quasi-Hamiltonian algebras.

In the statement, the set J' corresponds to the idempotents in $B \subset A$ which are involved in the fusions defining A_1 and A_2 . The extra assumption requiring the existence of Φ_k , $k \in J'$, is then sufficient to use Lemma 4.7. In particular, given specific A_1 and A_2 , it is possible to construct the isomorphism with the existence of such elements Φ_k in a proper subset of J' , which can be empty as is easily seen from the case of Lemma 4.8.

4.2.3. Fusion of morphisms

LEMMA 4.11. — *Let $\psi : A_1 \rightarrow A_2$ be a morphism of double quasi-Poisson algebras over B . Let $A_1^f = (A_1)_{e_2 \rightarrow e_1}^f$, $A_2^f = (A_2)_{e_2 \rightarrow e_1}^f$, be the fusion algebras with double brackets obtained by fusion of e_2 onto e_1 . Then ψ induces a morphism of double quasi-Poisson algebras $\psi^f : A_1^f \rightarrow A_2^f$. If furthermore ψ is a morphism of quasi-Hamiltonian algebras, then so is ψ^f .*

Proof. — By Proposition 4.5, the algebra A_k^f with $k = 1, 2$ has a double quasi-Poisson bracket given by

$$(4.31) \quad \{\!\!\{-, -\}\!\!\}_k^f := \{\!\!\{-, -\}\!\!\}_{k, ind} + \{\!\!\{-, -\}\!\!\}_{k, fus} .$$

We claim that the map $\psi^f : A_1^f \rightarrow A_2^f$ obtained from ψ by (2.15) is the morphism that we seek. By Lemma 2.5, we have that

$$(\psi^f \otimes \psi^f) \{\!\!\{c, d\}\!\!\}_{1, ind} = \{\!\!\{\psi^f(c), \psi^f(d)\}\!\!\}_{2, ind} , \quad \text{for all } c, d \in A_1^f .$$

Moreover, ψ^f preserves the type of generators in A_1^f, A_2^f as defined in Lemma 2.2, and the double brackets $\{\!\!\{-, -\}\!\!\}_{k, fus}$ are defined in terms of these types of generators. Therefore

$$(\psi^f \otimes \psi^f) \{\!\!\{c, d\}\!\!\}_{1, fus} = \{\!\!\{\psi^f(c), \psi^f(d)\}\!\!\}_{2, fus} , \quad \text{for all } c, d \in A_1^f .$$

Thus $\psi^f : A_1^f \rightarrow A_2^f$ is a morphism of double quasi-Poisson algebras.

In the quasi-Hamiltonian case, we can check that ψ^f maps the multiplicative moment map of A_1^f to that of A_2^f . Hence, it is a morphism of quasi-Hamiltonian algebras. \square

4.3. Application: Van den Bergh’s quasi-Hamiltonian structure for quivers

We can endow specific localisations of the path algebras of quivers with a quasi-Hamiltonian algebra structure defined by Van den Bergh [VdB08a, § 6.7]. The double bracket was explicitly spelled out first in [CF17, Proposition 2.6], then in full generalities in [Fai21, Theorem 3.3]. To define this structure, we fix a quiver Q , and consider the path algebra of its double $\mathbb{k}\bar{Q}$ following the conventions given in §2.2.2. We consider the algebra A_Q obtained by universal localisation of $\mathbb{k}\bar{Q}$ from the set $S = \{1 + aa^* \mid a \in \bar{Q}\}$. This localisation can be understood as adding local inverses $e_{t(a)} + aa^* \in e_{t(a)}A_Qe_{t(a)}$.

For each vertex $s \in I$, consider a total ordering $<_s$ on the set $T_s = \{a \in \bar{Q} \mid t(a) = s\}$. This induces an ordering function at the vertex s

$$o_s(-, -) : \bar{Q} \times \bar{Q} \rightarrow \{-1, 0, 1\},$$

which is defined on arrows $a, b \in \bar{Q}$ by $o_s(a, b) = +1$ if $a <_s b$, $o_s(a, b) = -1$ if $b <_s a$, while it is zero otherwise, i.e. if $a = b \in T_s$, if $a \notin T_s$ or if $b \notin T_s$. The algebra A_Q has a double quasi-Poisson bracket defined by

$$(4.32a) \quad \{\{a, a\}\} = \frac{1}{2}o_{t(a)}(a, a^*) \left(a^2 \otimes e_{t(a)} - e_{h(a)} \otimes a^2 \right), \quad \text{for } a \in \bar{Q},$$

$$(4.32b) \quad \begin{aligned} \{\{a, a^*\}\} &= e_{h(a)} \otimes e_{t(a)} + \frac{1}{2}a^*a \otimes e_{t(a)} + \frac{1}{2}e_{h(a)} \otimes aa^* \\ &+ \frac{1}{2}o_{t(a)}(a, a^*) (a^* \otimes a - a \otimes a^*), \quad \text{for } a \in Q, \end{aligned}$$

and for $b, c \in \bar{Q}$ such that $c \neq b, b^*$

$$(4.33) \quad \begin{aligned} \{\{b, c\}\} &= -\frac{1}{2}o_{t(b)}(b, c) (b \otimes c) - \frac{1}{2}o_{h(b)}(b^*, c^*) (c \otimes b) \\ &+ \frac{1}{2}o_{t(b)}(b, c^*) cb \otimes e_{t(b)} + \frac{1}{2}o_{h(b)}(b^*, c) e_{h(b)} \otimes bc. \end{aligned}$$

Moreover, A_Q is quasi-Hamiltonian for the multiplicative moment map

$$(4.34) \quad \Phi = \sum_{s \in I} \Phi_s, \quad \Phi_s = \prod_{a \in T_s}^{\rightarrow} (e_s + aa^*)^{\epsilon(a)}.$$

THEOREM 4.12. — *The quasi-Hamiltonian algebra $(A_Q, \{\{-, -\}\}, \Phi)$ only depends on Q seen as an undirected graph, up to isomorphism.*

Proof. — We first prove that the structure is independent of the direction of the arrow in a one-arrow quiver. If Q is the quiver $1 \rightarrow 2$, the quasi-Hamiltonian algebra structure on A_Q defined above is nothing else than the structure defined on A_1 in Example 4.1. The same is true for the opposite quiver: the quasi-Hamiltonian structure on $A_{Q^{op}}$ for Q^{op} given by $2 \rightarrow 1$ is the structure on A_2 given in Example 4.1. In that same example, it was proved that there exists an isomorphism $\psi : A_1 \rightarrow A_2$ of quasi-Hamiltonian algebras, thus the structures on A_Q and $A_{Q^{op}}$ are isomorphic, as claimed.

Consider now an arbitrary quiver Q . As in the proof of Theorem 3.12, we note that \bar{Q} can be obtained by identifying vertices in the double \bar{Q}^{sep} of the separated quiver Q^{sep} . The algebra $A_{Q^{sep}}$ is quasi-Hamiltonian by combining Lemma 4.4 and the case of a one-arrow quiver described at the beginning of this proof. In particular, the double quasi-Poisson bracket is given by

$$(4.35) \quad \{\{b, b^*\}\} = e_{v_{b^*}} \otimes e_{v_b} + \frac{1}{2} (bb^* \otimes e_{v_b} + e_{v_{b^*}} \otimes b^*b) ,$$

for all $b \in Q^{sep}$ and it is zero on every other pair of generators. The multiplicative moment map is defined as

$$(4.36) \quad \begin{aligned} \Phi &= \sum_{b \in Q^{sep}} \Phi_{v_b}, \Phi_{v_b} = e_{v_b} + bb^* \quad \text{if } b \in Q^{sep}, \\ \text{or } \Phi_{v_b} &= (e_{v_b} + bb^*)^{-1} \quad \text{if } b \in \bar{Q}^{sep} \setminus Q^{sep}. \end{aligned}$$

Fusing idempotents in $A_{Q^{sep}}$ according to the identification of vertices from Q^{sep} to Q gives us A_Q . It is proved in [Fai21, Theorem 3.3] that the quasi-Hamiltonian algebra structure obtained on A_Q by fusion using Proposition 4.5 is precisely the one given above. In particular, it does not depend on the order in which we fuse the idempotents by Theorem 4.10.

It remains to see that this structure is independent of the directions of the arrows in Q , up to isomorphism. This follows from Lemma 4.11 and the isomorphism that exists for a one-arrow quiver which was given at the beginning of the proof. \square

Remark 4.13. — In [Fai21, § 3.2], modifications of the algebra A_Q and its quasi-Hamiltonian algebra structure were considered. Namely, fixing a choice of coefficients $\gamma_a \in \mathbb{k}$, $a \in \bar{Q}$, satisfying $\gamma_a = \gamma_{a^*}$, introduce the algebra $A_{Q,\gamma}$ obtained by universal localisation of $\mathbb{k}\bar{Q}$ from the set $S_\gamma = \{1 + (\gamma_a - 1)e_{t(a)} + aa^* \mid a \in \bar{Q}\}$. The double quasi-Poisson bracket is given by (4.32a), (4.33), while we consider (4.32b) with its first term multiplied by γ_a . The multiplicative moment map $\Phi = \sum_{s \in I} \Phi_s$ is such that

$$\Phi_s = \prod_{a \in T_s}^{\rightarrow} (\gamma_a e_s + aa^*)^{\epsilon(a)} .$$

The proof of Theorem 4.12 can then be adapted to $A_{Q,\gamma}$. If $\gamma_a \neq 0$ for all $a \in \bar{Q}$, we have furthermore an isomorphism of double quasi-Poisson algebras $\psi : A_Q \rightarrow A_{Q,\gamma}$ given on generators by

$$(4.37) \quad \psi(a) = \gamma_a^{-1}a, \quad \psi(a^*) = a^*, \quad \text{for all } a \in Q .$$

Since $\psi((e_{t(a)} + aa^*)^{\epsilon(a)}) = \gamma_a^{-\epsilon(a)} (\gamma_a e_{t(a)} + aa^*)^{\epsilon(a)}$ for all $a \in \bar{Q}$ due to the condition $\gamma_a = \gamma_{a^*}$, ψ is in fact an isomorphism of quasi-Hamiltonian algebras if a is a loop whenever $\gamma_a \neq 1$.

5. H_0 -Poisson structures

5.1. Definition and general results

Let A be a \mathbb{k} -algebra. Let $[A, A]$ be the vector space spanned by commutators in A , from which we can define $H_0(A) := A/[A, A]$, the zeroth Hochschild homology of A . Denote by $a \mapsto \bar{a}$ the map which sends an element of A to its image in $A/[A, A]$. Remark that any derivation $\partial \in \text{Der}(A)$ induces a linear map on $H_0(A)$ since $\partial([a, b]) \in [A, A]$ for any $a, b \in A$. Following Crawley–Boevey [CB11], we say that a \mathbb{k} -bilinear map $\langle -, - \rangle : H_0(A) \times H_0(A) \rightarrow H_0(A)$ is a H_0 -Poisson structure on A if it is a Lie bracket, i.e.

$$(5.1) \quad \langle \bar{a}, \bar{b} \rangle = - \langle \bar{b}, \bar{a} \rangle, \quad \langle \bar{a}, \langle \bar{b}, \bar{c} \rangle \rangle + \langle \bar{b}, \langle \bar{c}, \bar{a} \rangle \rangle + \langle \bar{c}, \langle \bar{a}, \bar{b} \rangle \rangle = 0,$$

and each linear map $\langle \bar{a}, - \rangle : H_0(A) \rightarrow H_0(A)$ is induced by a derivation $\partial_a \in \text{Der}(A)$. We will write $(A, \langle -, - \rangle)$ when we want to emphasise the H_0 -Poisson structure on A .

In the relative setting where A is a B -algebra (following the convention of Section 2), we require that the map $\langle \bar{e}, - \rangle$ is induced by the trivial derivation $\partial_e = 0_A$ for each $e \in B$. In that case, we say that the H_0 -Poisson structure $\langle -, - \rangle$ is B -linear.

Let $(A_1, \langle -, - \rangle_1)$ and $(A_2, \langle -, - \rangle_2)$ be two B -algebras with H_0 -Poisson structures. Note that if $\phi : A_1 \rightarrow A_2$ is a morphism of B -algebras, then $\phi([a, b]) = [\phi(a), \phi(b)]$ for any $a, b \in A$ so ϕ induces a morphism $\bar{\phi} : H_0(A_1) \rightarrow H_0(A_2)$. We say that $\phi : A_1 \rightarrow A_2$ is a H_0 -Poisson morphism if it is a morphism of B -algebras such that the induced map $\bar{\phi} : H_0(A_1) \rightarrow H_0(A_2)$ is a morphism of Lie algebras, i.e. for any $\bar{a}, \bar{b} \in H_0(A_1)$,

$$(5.2) \quad \langle \bar{\phi}(\bar{a}), \bar{\phi}(\bar{b}) \rangle = \bar{\phi}(\langle \bar{a}, \bar{b} \rangle).$$

We say that it is an H_0 -Poisson isomorphism if it is an isomorphism of B -algebras (hence $\bar{\phi}$ is an isomorphism of Lie algebras).

Let A be a B -algebra with double bracket $\{\!\{-, -\}\!\}$. Consider the bilinear map

$$(5.3) \quad \{-, -\} = m \circ \{\!\{-, -\}\!\} : A \times A \rightarrow A,$$

obtained by composing the double bracket with the multiplication m of A . It inherits the B -linearity of $\{\!\{-, -\}\!\}$ in the following form: for any $e \in B$, $\{e, -\}$ and $\{-, e\}$ are identically zero. Then, as noticed in [VdB08a, Lemma 2.4.1], we have that the operation (5.3) induces well-defined maps

$$(5.4) \quad \{-, -\} : H_0(A) \times A \rightarrow A, \quad \{-, -\} : H_0(A) \times H_0(A) \rightarrow H_0(A).$$

The first operation is such that for any $\bar{a} \in H_0(A)$ we have $\{\bar{a}, -\} \in \text{Der}(A)$, while the second operation is antisymmetric. By construction, this second linear map is given by

$$(5.5) \quad \{\bar{a}, \bar{b}\} = \overline{\{\bar{a}, b\}} = \overline{\{a, b\}},$$

for any $\bar{a}, \bar{b} \in H_0(A)$ and lifts $a, b \in A$. We have a B -linear H_0 -Poisson structure on A provided that the operation $\{-, -\}$ on $H_0(A)$ satisfies Jacobi identity.

LEMMA 5.1 ([VdB08a, Lemma 2.6.2]). — Let A be a B -algebra with double bracket $\{\{-, -\}\}$, and write $\{-, -, -\}_3 = m \circ \{\{-, -, -\}\}$ for the map obtained by composing the triple bracket $\{\{-, -, -\}\}$ given by (2.6) with the multiplication $m : A^{\times 3} \rightarrow A$. Then the induced map $\{-, -\} : H_0(A) \times H_0(A) \rightarrow H_0(A)$ is a H_0 -Poisson structure if $\{-, -, -\}_3$ vanishes identically.

In particular, we get a H_0 -Poisson structure if $\{\{-, -\}\}$ is a double (quasi-)Poisson bracket.

Proof. — From [VdB08a, Corollary 2.4.4], we have the following identity in A

$$(5.6) \quad \{a, \{b, c\}\} - \{b, \{a, c\}\} - \{\{a, b\}, c\} = \{a, b, c\}_3 - \{b, a, c\}_3.$$

Hence if the right-hand side vanishes, this equality induced in $H_0(A)$ is just Jacobi identity.

The second part of the statement is obvious for a double Poisson bracket, since $\{\{-, -, -\}\} = 0$ identically. For a double quasi-Poisson bracket, $\{-, -, -\}_3 = 0$ because applying the multiplication map to (4.1) gives zero. \square

5.1.1. Structures induced by Hamiltonian algebras

We assume that $(A, \{\{-, -\}\}, \mu)$ is a Hamiltonian algebra over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. Fix $(\lambda_s) \in \mathbb{k}^I$ or equivalently $\lambda := \sum_{s \in I} \lambda_s e_s \in B$, and note that by definition of the moment map, (3.1) implies that for any $a \in A$

$$\{\{a, \mu - \lambda\}\} = \sum_s (e_s a \otimes e_s - e_s \otimes a e_s).$$

In particular, $\{a, \mu - \lambda\} = 0$ after multiplication. This yields that, if $\overline{(\mu - \lambda)}$ denotes the vector subspace of $H_0(A)$ spanned by the image of the ideal $(\mu - \lambda)$ under the map $A \rightarrow A/[A, A]$, we have $\{\bar{a}, \overline{(\mu - \lambda)}\} \in \overline{(\mu - \lambda)}$ for any $\bar{a} \in H_0(A)$. In particular, the H_0 -Poisson structure descends from $H_0(A)$ to a Lie bracket $\{-, -\}^\lambda$ on $H_0(A)^\lambda := H_0(A)/\overline{(\mu - \lambda)}$ given by

$$(5.7) \quad \{\bar{a} + \overline{(\mu - \lambda)}, \bar{b} + \overline{(\mu - \lambda)}\}^\lambda = (\{a, b\} + [A, A]) + \overline{(\mu - \lambda)}, \quad \text{for any } a, b \in A.$$

Set $A^\lambda = A/\overline{(\mu - \lambda)}$, and remark that we can identify $H_0(A)^\lambda$ with $H_0(A^\lambda) := A^\lambda/[A^\lambda, A^\lambda]$. Under this identification, the Lie bracket $\{-, -\}^\lambda$ is given by

$$\{\bar{a} + \overline{(\mu - \lambda)}, \bar{b} + \overline{(\mu - \lambda)}\}^\lambda = \overline{\{a, b\} + (\mu - \lambda)}.$$

(Here, the bar in the right-hand side denotes the map $A^\lambda \rightarrow H_0(A^\lambda)$.) We get that A^λ is endowed with a H_0 -Poisson structure since the linear map $\{\bar{a} + \overline{(\mu - \lambda)}, -\}$ on $H_0(A^\lambda)$ is induced by $\partial_a \in \text{Der}(A^\lambda)$ given by

$$\partial_a(b + (\mu - \lambda)) = \{a, b\} + (\mu - \lambda),$$

for any lifts $a, b \in A$. Note that the induced linear map on $H_0(A^\lambda)$ is independent of the lift, though $\partial_a, \partial_{a+\mu-\lambda} \in \text{Der}(A^\lambda)$ are not the same in general. Indeed $\{\mu - \lambda, b\} = \sum_s (b e_s - e_s b)$ may be nonzero, but it vanishes modulo commutators. Combining this discussion with Lemma 5.1, we have obtained the following result.

PROPOSITION 5.2 ([VdB08a, Proposition 2.6.5]). — *Let $(A, \{\{-, -\}, \mu)$ be a Hamiltonian algebra. Then for any $\lambda \in B$, the H_0 -Poisson structure $\{-, -\}$ on A descends to a H_0 -Poisson structure $\{-, -\}^\lambda$ on A^λ .*

PROPOSITION 5.3. — *Let $\phi : (A_1, \{\{-, -\}_1\}) \rightarrow (A_2, \{\{-, -\}_2\})$ be a morphism of double Poisson algebras. Then $\phi : (A_1, \{-, -\}_1) \rightarrow (A_2, \{-, -\}_2)$ is a H_0 -Poisson morphism.*

If ϕ is a morphism of Hamiltonian algebras, then for any $\lambda \in B$, ϕ induces a H_0 -Poisson morphism $\phi^\lambda : (A_1/(\mu_1 - \lambda), \{-, -\}_1^\lambda) \rightarrow (A_2/(\mu_2 - \lambda), \{-, -\}_2^\lambda)$.

Proof. — In the first case, we just have to show that $\bar{\phi}$ is a morphism of Lie algebras. For any $a, b \in A_1$, (2.9) yields $\phi(\{a, b\}_1) = \{\phi(a), \phi(b)\}_2$. Hence, we get from (5.5) that for any $\bar{a}, \bar{b} \in H_0(A_1)$ with arbitrary lifts $a, b \in A_1$,

$$\begin{aligned} \{\bar{\phi}(\bar{a}), \bar{\phi}(\bar{b})\}_2 &= \{\overline{\phi(a)}, \overline{\phi(b)}\}_2 = \overline{\{\phi(a), \phi(b)\}_2} \\ &= \overline{\phi(\{a, b\}_1)} = \bar{\phi}(\overline{\{a, b\}_1}) = \bar{\phi}(\{\bar{a}, \bar{b}\}_1). \end{aligned}$$

In the second case, note that $\phi((\mu_1 - \lambda)) \subset (\mu_2 - \lambda)$ so ϕ induces a morphism $\phi^\lambda : A_1^\lambda \rightarrow A_2^\lambda$ hence a morphism $\bar{\phi}^\lambda : H_0(A_1^\lambda) \rightarrow H_0(A_2^\lambda)$ where $A_i^\lambda = A/(\mu_i - \lambda)$. In the same way, since $\bar{\phi}((\mu_1 - \lambda)) \subset (\mu_2 - \lambda)$, we have that $\bar{\phi}$ induces a map

$$\bar{\phi}^\lambda : H_0(A_1^\lambda) \simeq H_0(A_1)/\overline{(\mu_1 - \lambda)} \longrightarrow H_0(A_2)/\overline{(\mu_2 - \lambda)} \simeq H_0(A_2^\lambda),$$

which coincides with the map $\bar{\phi}^\lambda$ induced by ϕ^λ . We can thus use (5.7) to conclude. \square

Example 5.4. — Consider the Hamiltonian algebra structure on the path algebra of a double quiver \bar{Q} given in § 3.3. By Lemma 5.1, we get a H_0 -Poisson structure on $\mathbb{k}\bar{Q}$, and its associated Lie bracket on $H_0(\mathbb{k}\bar{Q})$ is the necklace Lie bracket [BLB02, Gin01]. By Proposition 5.2, the double bracket on $\mathbb{k}\bar{Q}$ descends to a H_0 -Poisson structure on $\Pi^\lambda(Q) := \mathbb{k}\bar{Q}/(\mu - \lambda)$ [CB11, VdB08a]. The algebra $\Pi^\lambda(Q)$ is called a *deformed preprojective algebra* [CBH98, Section 2]. It was proved by Crawley-Boevey and Holland in [CBH98, Lemma 2.2] that deformed preprojective algebras are independent of the orientation chosen on Q . Using Theorem 3.12, we obtain that the H_0 -Poisson structure hence defined is independent of the orientation chosen on Q up to isomorphism, and we can easily check that these isomorphisms are realised by the maps considered by Crawley-Boevey and Holland.

5.1.2. Structures induced by quasi-Hamiltonian algebras

We assume that $(A, \{\{-, -\}, \Phi)$ is a quasi-Hamiltonian algebra over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. Fix $(c_s) \in (\mathbb{k}^\times)^I$ or equivalently $c := \sum_{s \in I} c_s e_s \in B^\times$. By definition of the multiplicative moment map, for any $a \in A$ we have

$$\{\{a, \Phi - c\}\} = -\frac{1}{2} \sum_{s \in I} (\Phi_s \otimes a e_s - \Phi_s a \otimes e_s + e_s \otimes a \Phi_s - e_s a \otimes \Phi_s),$$

so that $\{a, \Phi - c\} = 0$ after multiplication. We can thus adapt the discussion from § 5.1.1 to the quasi-Hamiltonian setting with $A^c = A/(\Phi - c)$ and get the following results.

PROPOSITION 5.5 ([VdB08a, Proposition 5.1.5]). — *Let $(A, \{\{-, -\}, \Phi)$ be a quasi-Hamiltonian algebra. Then for any $c \in B^\times$, the H_0 -Poisson structure $\{-, -\}$ on A descends to a H_0 -Poisson structure $\{-, -\}^c$ on A^c .*

PROPOSITION 5.6. — *Let $\phi : (A_1, \{\{-, -\}_1\}) \rightarrow (A_2, \{\{-, -\}_2\})$ be a morphism of double quasi-Poisson algebras. Then $\phi : (A_1, \{-, -\}_1) \rightarrow (A_2, \{-, -\}_2)$ is a H_0 -Poisson morphism.*

If ϕ is a morphism of quasi-Hamiltonian algebras, then for any $c \in B^\times$, ϕ induces a H_0 -Poisson morphism $\phi^c : (A_1/(\Phi_1 - c), \{-, -\}_1^c) \rightarrow (A_2/(\Phi_2 - c), \{-, -\}_2^c)$.

Example 5.7. — Consider the quasi-Hamiltonian algebra structure on the localisation A_Q of the path algebra of a double quiver \bar{Q} given in § 4.3, which depends on an ordering of the arrows. The algebra $\Lambda^c(Q) := A_Q/(\Phi - c)$ is called a *multiplicative preprojective algebra* [CBS06]. As noticed by Van den Bergh [VdB08a, Proposition 6.8.1], the double bracket on A_Q descends to a H_0 -Poisson structure on $\Lambda^c(Q)$, see Proposition 5.5. It was proved by Crawley–Boevey and Shaw [CBH98, Theorem 1.4] that multiplicative preprojective algebras are independent of the orientation chosen on Q , and of the ordering of the arrows. Using Theorem 4.12, we obtain that the H_0 -Poisson structure on a multiplicative preprojective algebra is independent of the orientation and the ordering of the arrows up to isomorphism, and we can check that such isomorphisms are precisely realised by the maps considered by Crawley–Boevey and Shaw.

5.1.3. Relation to the affine moduli space of representations

We assume that the base field \mathbb{k} is an algebraically closed field of characteristic zero. Consider an algebra A over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. We denote the affine representation space (relative to B) of A with dimension vector $\alpha \in \mathbb{N}^I$ by $\text{Rep}(A, \alpha)$. Explicitly, $\text{Rep}(A, \alpha)$ parametrises representations ρ of A on \mathbb{k}^N , $N = \sum_s \alpha_s$, such that the idempotent matrix $\rho(e_s)$ is given by $\text{diag}(0_{\alpha_1}, \dots, 0_{\alpha_{s-1}}, \text{Id}_{\alpha_s}, 0_{\alpha_{s+1}}, \dots, 0_{\alpha_{|I|}})$ after identification of I with $\{1, \dots, |I|\}$. In other words, under the decomposition

$$(5.8) \quad \mathbb{k}^N = \mathbb{k}^{\alpha_1} \oplus \mathbb{k}^{\alpha_2} \oplus \dots \oplus \mathbb{k}^{\alpha_{|I|}},$$

$\rho(e_s)$ projects an element $(v_1, \dots, v_{|I|})$ onto $(0, \dots, 0, v_s, 0, \dots, 0)$, where for all $r \in I$ we have $v_r \in \mathbb{k}^{\alpha_r}$.

There is a natural action of $\text{GL}_\alpha := \prod_s \text{GL}_{\alpha_s}(\mathbb{k})$ by change of basis with respect to the decomposition (5.8). For any $a \in A$, we denote by $\mathcal{X}(a)$ the function on $\text{Rep}(A, \alpha)$ which returns the matrix representing a at each point, i.e. $\mathcal{X}(a)(\rho) = \rho(a)$. Following [CB11, VdB08a], we note that the map

$$(5.9) \quad \text{tr} : A \rightarrow \mathbb{k}[\text{Rep}(A, \alpha)], \quad \text{tr}(a) = \sum_{1 \leq i \leq N} \mathcal{X}(a)_{ii},$$

has its image which generates $\mathbb{k}[\text{Rep}(A, \alpha)]^{\text{GL}_\alpha}$. Furthermore, given $\bar{a} \in H_0(A)$ and two lifts $a_1, a_2 \in A$, we note that $\text{tr}(a_1) = \text{tr}(a_2)$ since the trace vanishes on commutators of matrices.

THEOREM 5.8 ([CB11, Theorem 4.5]). — *Let $\langle -, - \rangle$ be a H_0 -Poisson structure on A . Then, for any dimension vector α , there is a unique Poisson bracket $\{-, -\}$ on $\mathcal{O}[\text{Rep}(A, \alpha)]^{\text{GL}\alpha}$ such that for any $a, b \in A$*

$$(5.10) \quad \{\text{tr}(a), \text{tr}(b)\} = \text{tr} \left(\langle \bar{a}, \bar{b} \rangle^l \right),$$

where $\langle \bar{a}, \bar{b} \rangle^l \in A$ is an arbitrary lift of $\langle \bar{a}, \bar{b} \rangle \in H_0(A)$.

PROPOSITION 5.9. — *If $\phi : (A_1, \langle -, - \rangle_1) \rightarrow (A_2, \langle -, - \rangle_2)$ is a H_0 -Poisson morphism, then the morphism $\bar{\phi}_\alpha : \mathcal{O}[\text{Rep}(A_1, \alpha)]^{\text{GL}\alpha} \rightarrow \mathcal{O}[\text{Rep}(A_2, \alpha)]^{\text{GL}\alpha}$, uniquely defined by $\bar{\phi}_\alpha(\text{tr}(a)) = \text{tr}(\phi(a))$ for any $a \in A_1$, is a Poisson morphism.*

Proof. — Denote by $\{-, -\}_k$ the Poisson bracket on $\mathcal{O}[\text{Rep}(A_k, \alpha)]^{\text{GL}\alpha}$ induced by Theorem 5.8 for $k = 1, 2$. Then, on generators $\text{tr}(a), \text{tr}(b)$ of $\mathcal{O}[\text{Rep}(A_1, \alpha)]^{\text{GL}\alpha}$, we have that

$$(5.11a) \quad \begin{aligned} \{\bar{\phi}_\alpha(\text{tr}(a)), \bar{\phi}_\alpha(\text{tr}(b))\}_2 &= \{\text{tr}(\phi(a)), \text{tr}(\phi(b))\}_2 \\ &= \text{tr} \left(\langle \overline{\phi(a)}, \overline{\phi(b)} \rangle_2^l \right) = \text{tr} \left(\langle \bar{\phi}(\bar{a}), \bar{\phi}(\bar{b}) \rangle_2^l \right), \end{aligned}$$

$$(5.11b) \quad \begin{aligned} \bar{\phi}_\alpha(\{\text{tr}(a), \text{tr}(b)\}_1) &= \bar{\phi}_\alpha \left(\text{tr} \left(\langle \bar{a}, \bar{b} \rangle_1^l \right) \right) = \text{tr} \left(\phi \left(\langle \bar{a}, \bar{b} \rangle_1^l \right) \right) \\ &= \text{tr} \left(\left(\bar{\phi} \langle \bar{a}, \bar{b} \rangle_1 \right)^l \right). \end{aligned}$$

For (5.11a), we used the definition of $\bar{\phi}_\alpha$ in the first equality, the definition (5.10) of the Poisson bracket in the second, the fact that $\bar{\phi} : H_0(A_1) \rightarrow H_0(A_2)$ is induced by ϕ in the third; we used these results similarly in (5.11b). As the final terms in (5.11a) and (5.11b) are equal due to (5.2) since ϕ is a H_0 -Poisson morphism, we get that $\bar{\phi}_\alpha$ is a Poisson morphism. If ϕ is an isomorphism with inverse $\phi^{-1} : A_2 \rightarrow A_1$, then the inverse of $\bar{\phi}_\alpha$ is given by $\text{tr}(a) \mapsto \text{tr}(\phi^{-1}(a))$ for any $a \in A_2$. \square

5.2. Some applications

We have obtained in § 5.1 that morphisms of double (quasi-)Poisson algebras induce morphisms of quotients of these algebras obtained by fixing the (multiplicative) moment maps to a linear combination of idempotents. In particular, the induced morphisms define morphisms of coordinate rings of moduli spaces of representations. We investigate these constructions in the rest of this section. We will study several basic examples and their relation to the first Weyl algebra and quantum tori, as well as the morphisms that they induce on (q) -Calogero–Moser spaces which will appear in Section 6.

In this subsection, all algebras are B -algebras for the convention of Section 2. In that case, a B -linear morphism is a morphism satisfying $\phi(e) = e$ for all $e \in B$.

5.2.1. Using Hamiltonian algebras

Given a Hamiltonian algebra $(A, \{\{-, -\}, \mu)$ over B , let us denote by $\text{Aut}(A) := \text{Aut}_B(A)$ its group of B -linear automorphisms, and $\text{HAut}(A)$ the subgroup of automorphisms which are morphisms of Hamiltonian algebras. Following [BP11], let us also introduce $\text{Aut}(A; \mu)$ as the subgroup of automorphisms of A preserving μ . It is clear that we have the inclusions

$$(5.12) \quad \text{HAut}(A) \subset \text{Aut}(A; \mu) \subset \text{Aut}(A).$$

Remark 5.10. — The inclusions in (5.12) are not necessarily equalities. Consider $A = \mathbb{k}\langle x, y, z \rangle$ with double Poisson bracket given on generators by

$$\{\{x, y\}\} = 1 \otimes 1, \quad \{\{z, z\}\} = z \otimes 1 - 1 \otimes z, \quad \{\{x, z\}\} = 0 = \{\{y, z\}\},$$

and moment map $\mu = [x, y] + z$. Consider the automorphisms $\phi, \psi : A \rightarrow A$ defined on generators by

$$\phi(x) = \frac{1}{2}x, \quad \phi(y) = y, \quad \phi(z) = z + \frac{1}{2}[x, y], \quad \psi(x) = -x, \quad \psi(y) = y, \quad \psi(z) = z.$$

Then ϕ preserves μ but it is not a morphism of double brackets, while ψ does not preserve μ .

The first inclusion from (5.12) descends to an inclusion $\overline{\text{HAut}(A)} \subset \overline{\text{Aut}(A; \mu)}$ in the quotient $A^\lambda = A/(\mu - \lambda)$ for any $\lambda \in B$. Moreover, the automorphisms in $\overline{\text{HAut}(A)}$ preserve the induced H_0 -Poisson structure by Proposition 5.3.

Example 5.11. — Consider the Jordan quiver Q_\circ consisting of the vertex set $I = \{0\}$ and a single arrow which is a loop $a : 0 \rightarrow 0$. Its double \bar{Q}_\circ contains an additional arrow $a^* : 0 \rightarrow 0$. The path algebra of \bar{Q}_\circ has a Hamiltonian algebra structure constructed in §3.3. Under $a \mapsto x, a^* \mapsto y$ we can induce the Hamiltonian algebra structure on the free algebra $F_2 = \mathbb{k}\langle x, y \rangle$ by taking the double Poisson bracket

$$(5.13) \quad \{\{x, y\}\} = 1 \otimes 1, \quad \{\{x, x\}\} = 0 = \{\{y, y\}\},$$

with the moment map $\mu = [x, y]$. Meanwhile, we can start with the opposite quiver Q_\circ^{op} given by $b : 0 \rightarrow 0$, and the isomorphism $b \mapsto x, b^* \mapsto y$ induces the same Hamiltonian algebra structure on $F_2 \simeq \mathbb{k}\bar{Q}_\circ^{op}$. By Theorem 3.12, we have isomorphic Hamiltonian algebras if we start with Q_\circ or Q_\circ^{op} , and the isomorphism can be computed from Example 3.1 to be

$$(5.14) \quad \psi : \mathbb{k}\bar{Q}_\circ \rightarrow \mathbb{k}\bar{Q}_\circ^{op}, \quad \psi(a) = b^*, \quad \psi(a^*) = -b.$$

Under the identifications with F_2 , ψ induces an isomorphism of Hamiltonian algebras given by

$$(5.15) \quad \mathcal{F} : F_2 \rightarrow F_2, \quad \mathcal{F}(x) = y, \quad \mathcal{F}(y) = -x.$$

The automorphism \mathcal{F} satisfies $\mathcal{F}^4 = \text{id}$ and is sometimes called the formal Fourier transform. We can also easily check that for all $\gamma \in \mathbb{k}$ and $k \geq 0$ the automorphism

$$(5.16) \quad \phi_{k, \gamma} : F_2 \rightarrow F_2, \quad \phi_{k, \gamma}(x) = x + \gamma y^k, \quad \phi_{k, \gamma}(y) = y,$$

defines an isomorphism of Hamiltonian algebras as it preserves (5.13) and the moment map. Hence, we also get the isomorphism of Hamiltonian algebras

$$(5.17) \quad \phi'_{k,\gamma} = \mathcal{F}^{-1} \circ \phi_{k,-\gamma} \circ \mathcal{F} : F_2 \rightarrow F_2, \quad \phi'_{k,\gamma}(x) = x, \quad \phi'_{k,\gamma}(y) = y + \gamma x^k.$$

Imposing the relation $\mu = 1$, the automorphisms (5.16) and (5.17) descend to automorphisms of the first Weyl algebra $A_1 = \mathbb{k}\langle x, y \rangle / (xy - yx - 1)$.

It is a result of Dixmier [Dix68, Théorème 8.10] that the images in A_1 of the automorphisms (5.16) and (5.17) generate the whole group of automorphisms.

COROLLARY 5.12. — *Let F_2 be the free algebra on two generators and A_1 be the first Weyl algebra. Then $\text{HAut}(F_2)$ surjects onto $\text{Aut}(A_1)$.*

Remark 5.13. — Using Proposition 5.3, any automorphism of the first Weyl algebra preserves the H_0 -Poisson structure induced by (5.13). However, this result is not interesting as $H_0(A_1)$ is trivial. If we set instead the moment map μ to 0 in Example 5.11, we get that the H_0 -Poisson structure induced by (5.13) on $\mathbb{k}[x, y]$ is the canonical Poisson bracket defined by $\{x, y\} = 1$ and $\{x, x\} = 0 = \{y, y\}$.

Let us now assume that \mathbb{k} is of characteristic zero and algebraically closed. The elements of $H_0(A^\lambda)$ induce generators of representation spaces using the trace map as in § 5.1.3, and we can obtain the following diagram for any dimension vector $\alpha \in \mathbb{N}^I$

$$\begin{array}{ccccc} \text{HAut}(A) & \longrightarrow & \text{Aut}(A; \mu) & \longrightarrow & \text{Aut}(A) \\ \downarrow & & \downarrow & & \\ \overline{\text{HAut}(A)} & \longrightarrow & \overline{\text{Aut}(A; \mu)} & \longrightarrow & \text{Aut}(A^\lambda) \\ \downarrow & & \downarrow & & \\ \overline{\text{HAut}(A)}_\alpha & \longrightarrow & \overline{\text{Aut}(A; \mu)}_\alpha & \longrightarrow & \text{Aut}(\mathcal{A}^\lambda_\alpha) \end{array}$$

where $\mathcal{A}^\lambda_\alpha = \mathcal{O}[\text{Rep}(A^\lambda, \alpha)]^{\text{GL}_\alpha}$ denotes the coordinate ring of $\mathcal{X}(\mu)^{-1}(\lambda \text{Id}_\alpha) // \text{GL}_\alpha$ for $\lambda \text{Id}_\alpha = (\lambda_s \text{Id}_{\alpha_s})_s$. The automorphisms in $\overline{\text{HAut}(A)}_\alpha$ are Poisson by Proposition 5.9.

For a quiver Q , consider the Hamiltonian algebra structure on $\mathbb{k}\bar{Q}$ given in § 3.3. We get that $\text{HAut}(\mathbb{k}\bar{Q})$ induces Poisson automorphisms on the corresponding quiver varieties by the above argument. It is an interesting question to understand what are the properties of these morphisms.

Example 5.14. — We work over $\mathbb{k} = \mathbb{C}$. Consider the quiver Q_1 formed by the vertices $I = \{0, \infty\}$ and arrows $x : 0 \rightarrow \infty, v = 0 \rightarrow \infty$. The Hamiltonian automorphisms on $\mathbb{C}\langle x, y \rangle$ from Example 5.11 can be extended to $\mathbb{C}\bar{Q}_1$ using $x \mapsto x, y \mapsto x^*$ and acting as the identity on v, v^* . Fix $n \in \mathbb{N}^\times$. If we take $(\lambda_0, \lambda_\infty) = (1, -n)$ and $(\alpha_0, \alpha_\infty) = (n, 1)$ and denote the corresponding quiver variety as \mathcal{C}_n , the elements of $\text{HAut}(\mathbb{C}\bar{Q}_1)$ descend to Poisson automorphisms on \mathcal{C}_n . As the group of automorphisms generated by the images of $\phi_{k,\gamma}, \phi'_{k,\gamma}$ acts transitively on \mathcal{C}_n by [BW00, Theorems 1.2, 1.3], the same result holds for the image of $\text{HAut}(\mathbb{C}\bar{Q}_1)$. The variety \mathcal{C}_n is the n^{th} Calogero–Moser space [Wil98], see § 6.1.

If Q_2 is the quiver obtained from Q_1 by adding an arrow $\infty \rightarrow 0$, we can reproduce the same construction with the same parameters to get a quiver variety $\mathcal{C}_{n,2}$. It is proved in [BP11] that a subgroup of $\text{Aut}(\mathbb{C}\bar{Q}_2; \mu)$ acts transitively on $\mathcal{C}_{n,2}$ when induced onto this space, see also [MT13]. It is not known if the subgroup of Poisson automorphisms induced by $\text{HAut}(\mathbb{C}\bar{Q}_2)$ also acts transitively.

5.2.2. Using quasi-Hamiltonian algebras

Given a quasi-Hamiltonian algebra $(A, \{\{-, -\}, \Phi)$ over B , we can reproduce the construction from §5.2.1. Namely, we can define the group of B -linear automorphisms $\text{Aut}(A)$, the subgroup of quasi-Hamiltonian automorphisms $\text{qHAut}(A)$, and the subgroup $\text{Aut}(A; \Phi)$ of automorphisms of A preserving Φ . We get the inclusions

$$(5.18) \quad \text{qHAut}(A) \subset \text{Aut}(A; \Phi) \subset \text{Aut}(A).$$

The first inclusion descends to an inclusion of automorphisms of the quotient $A^c = A/(\Phi - c)$ for any $c \in B^\times$ such that the image of $\text{qHAut}(A)$ preserves the H_0 -Poisson structure obtained from Proposition 5.6. If \mathbb{k} is of characteristic zero and algebraically closed, they furthermore descend to automorphisms of affine moduli spaces of representations, for which the elements in the image of $\text{qHAut}(A)$ are Poisson isomorphisms by Proposition 5.9.

Example 5.15. — Consider the Jordan quiver Q_\circ as in Example 5.11. The localisation A_{Q_\circ} of the path algebra of \bar{Q}_\circ has a quasi-Hamiltonian algebra structure constructed in §4.3. Let $F_{2,S} = \mathbb{k}\langle x, y \rangle_S$ be the universal localisation of $\mathbb{k}\langle x, y \rangle$ with respect to the set $S = \{1 + xy, 1 + yx\}$. Under $a \mapsto y, a^* \mapsto x$ we can induce the quasi-Hamiltonian algebra structure on $F_{2,S} \simeq A_{Q_\circ}$ with double quasi-Poisson bracket

$$(5.19) \quad \begin{aligned} \{\{x, x\}\}_1 &= -\frac{1}{2} (x^2 \otimes 1 - 1 \otimes x^2), & \{\{y, y\}\}_1 &= +\frac{1}{2} (y^2 \otimes 1 - 1 \otimes y^2), \\ \{\{y, x\}\}_1 &= 1 \otimes 1 + \frac{1}{2} (xy \otimes 1 + 1 \otimes yx + x \otimes y - y \otimes x), \end{aligned}$$

and the multiplicative moment map $\Phi = (1 + yx)(1 + xy)^{-1}$. This corresponds to taking the ordering $a < a^*$. If we consider the other ordering $a^* < a$, we get the double quasi-Poisson bracket

$$(5.20) \quad \begin{aligned} \{\{x, x\}\}_2 &= +\frac{1}{2} (x^2 \otimes 1 - 1 \otimes x^2), & \{\{y, y\}\}_2 &= -\frac{1}{2} (y^2 \otimes 1 - 1 \otimes y^2), \\ \{\{y, x\}\}_2 &= 1 \otimes 1 + \frac{1}{2} (xy \otimes 1 + 1 \otimes yx - x \otimes y + y \otimes x), \end{aligned}$$

and the multiplicative moment map $\Phi' = (1 + xy)^{-1}(1 + yx)$. We can read from the proof of Lemma 4.7 that the isomorphism of quasi-Hamiltonian algebras from $(F_{2,S}, \{\{-, -\}\}_1, \Phi)$ to $(F_{2,S}, \{\{-, -\}\}_2, \Phi')$ is given by

$$(5.21) \quad \psi : F_{2,S} \rightarrow F_{2,S}, \quad \psi(x) = x(1 + xy), \quad \psi(y) = (1 + xy)^{-1}y.$$

As in Example 5.11, we can start with the opposite quiver Q_\circ^{op} given by $b : 0 \rightarrow 0$, and considering the ordering $b^* < b$ we get a quasi-Hamiltonian algebra structure on $A_{Q_\circ^{op}}$. Under the identification $F_{2,S} \simeq A_{Q_\circ^{op}}$ given by $b \mapsto y, b^* \mapsto x$, we again get

a quasi-Hamiltonian algebra structure on $F_{2,S}$, which is given by $\{\{-, -\}\}_2$ and Φ_2 defined above. By Theorem 4.12, we have isomorphic quasi-Hamiltonian algebras if we start with Q_\circ or Q_\circ^{op} , and the isomorphism⁽³⁾ can be computed from Example 4.1 to be

$$a \mapsto b^*, \quad a^* \mapsto -(1 + bb^*)^{-1} b.$$

Hence it gives an isomorphism of quasi-Hamiltonian algebras from $(F_{2,S}, \{\{-, -\}\}_1, \Phi)$ to $(F_{2,S}, \{\{-, -\}\}_2, \Phi')$ as

$$(5.22) \quad \xi : F_{2,S} \rightarrow F_{2,S}, \quad \xi(x) = -(1 + yx)^{-1}y, \quad \xi(y) = x.$$

Note that $\xi \circ \psi^{-1} \in \text{qHAut}(F_{2,S})$ when $F_{2,S}$ is endowed with $\{\{-, -\}\}_2$ and Φ' . Finally, we note that the automorphism

$$(5.23) \quad \phi_\beta : F_{2,S} \rightarrow F_{2,S}, \quad \phi_\beta(x) = \beta^{-1}x, \quad \phi_\beta(y) = \beta y, \quad \beta \in \mathbb{k}^\times,$$

is such that $\phi_\beta \in \text{qHAut}(F_{2,S})$ for both quasi-Hamiltonian algebra structures. After imposing the relations $\Phi = c^{-1}, \Phi' = c^{-1}$ for $c \in \mathbb{k}^\times$, the automorphisms ψ, ξ, ϕ_β descend to automorphisms $\bar{\psi}, \bar{\xi}, \bar{\phi}_\beta$ of the localised first quantised Weyl algebra B_1^c defined as

$$(5.24) \quad B_1^c = (A_1^c)_{1+xy}, \quad A_1^c := \mathbb{k}\langle x, y \rangle / (1 + xy - c(1 + yx)).$$

They can be used to define the following elements of $\text{Aut}(B_1^c)$,

$$(5.25) \quad \begin{aligned} (x, y) &\mapsto ((1 + yx)x, y(1 + yx)^{-1}), & (x, y) &\mapsto (-(1 + yx)^{-1}y, x), \\ (x, y) &\mapsto (\beta^{-1}x, \beta y), & & \beta \in \mathbb{k}^\times. \end{aligned}$$

Alev and Dumas classified the automorphisms of B_1^c in [AD96, Théorème 1.7], and they obtained that for $c \neq \pm 1$, $\text{Aut}(B_1^c)$ is generated by the three automorphisms (5.25).

COROLLARY 5.16. — *Fix $c \in \mathbb{k} \setminus \{0, \pm 1\}$. Let $F_{2,S} = \mathbb{k}\langle x, y \rangle_S$ for $S = \{1 + xy, 1 + yx\}$ and B_1^c be the localisation of the first quantised Weyl algebra defined by (5.24). Then, the automorphisms of B_1^c are all induced by isomorphisms of quasi-Hamiltonian algebras on $F_{2,S}$ (possibly for different structures).*

Contrary to the case of the first Weyl algebra obtained in Corollary 5.12, we do not have a surjection $\text{qHAut}(F_{2,S}) \rightarrow \text{Aut}(B_1^c)$ since $\psi : F_{2,S} \rightarrow F_{2,S}$ given by (5.21) does not preserve the double bracket (5.19) or its multiplicative moment map Φ . This illustrates the key difference between Lemma 3.7 and Lemma 4.7 : in the quasi-Hamiltonian setting, performing fusion in the opposite order induces a non-trivial isomorphism.

Example 5.17. — Let $L_2 = \mathbb{k}\langle x^{\pm 1}, y^{\pm 1} \rangle$ denote the (non-commutative) algebra of Laurent polynomials in two variables. There is a quasi-Hamiltonian algebra structure on L_2 with double quasi-Poisson bracket

$$(5.26) \quad \begin{aligned} \{\{x, x\}\} &= +\frac{1}{2} (x^2 \otimes 1 - 1 \otimes x^2), & \{\{y, y\}\} &= -\frac{1}{2} (y^2 \otimes 1 - 1 \otimes y^2), \\ \{\{x, y\}\} &= \frac{1}{2} (yx \otimes 1 + 1 \otimes xy - x \otimes y + y \otimes x), \end{aligned}$$

⁽³⁾We take $a < a^*$ and $b^* < b$ to define the quasi-Hamiltonian structures.

and multiplicative moment map $\Phi = xyx^{-1}y^{-1}$. This quasi-Hamiltonian algebra was considered in [CF17] after localisation of A_{Q_\circ} . It can also be obtained from an intersection pairing on the fundamental group π of a punctured torus as we have

$$L_2 \simeq \mathbb{k} \langle x^{\pm 1}, y^{\pm 1}, \Phi^{\pm 1} \rangle / (\Phi - xyx^{-1}y^{-1}) =: \pi,$$

see the work of Massuyeau–Turaev [MT14]. Elements of $\text{Aut}(\pi)$ which are not acting by conjugation can be found in [Gol03, Appendix], and they include the Dehn twists

$$(5.27) \quad \begin{aligned} \tau : L_2 &\rightarrow L_2, & \tau(x) &= xy, & \tau(y) &= y, \\ \tilde{\tau} : L_2 &\rightarrow L_2, & \tilde{\tau}(x) &= x, & \tilde{\tau}(y) &= yx. \end{aligned}$$

We have $\tau, \tilde{\tau} \in \text{Aut}(L_2, \Phi)$, and we can prove that both automorphisms preserve the double bracket (5.26), hence these elements belong to $\text{qHAut}(L_2)$. We also get that

$$(5.28) \quad \sigma : L_2 \rightarrow L_2, \quad \sigma(x) = y^{-1}, \quad \sigma(y) = yxy^{-1},$$

belongs to $\text{qHAut}(L_2)$ in view of $\sigma = \tau^{-1} \circ \tilde{\tau} \circ \tau^{-1}$. Finally, we note that the automorphism

$$(5.29) \quad \phi_{\alpha, \beta} : L_2 \rightarrow L_2, \quad \phi_{\alpha, \beta}(x) = \alpha x, \quad \phi_{\alpha, \beta}(y) = \beta y, \quad \alpha, \beta \in \mathbb{k}^\times,$$

is such that $\phi_{\alpha, \beta} \in \text{qHAut}(L_2)$. Imposing the relation $\Phi = c$ for $c \in \mathbb{k}^\times$, all these automorphisms descend to automorphisms of the quantum torus

$$(5.30) \quad \mathbb{C}_1^c = \mathbb{k} \langle x^{\pm 1}, y^{\pm 1} \rangle / (xy - cyx).$$

They can be used to define the following elements of $\text{Aut}(\mathbb{C}_1^c)$,

$$(5.31) \quad (x, y) \mapsto (xy, y), \quad (x, y) \mapsto (y^{-1}, x), \quad (x, y) \mapsto (\alpha x, \beta y), \quad \alpha, \beta \in \mathbb{k}^\times.$$

We get from [KPS94, Theorem 1.5] or [AD96, Proposition 1.6] that for $c \neq \pm 1$, the automorphisms (5.31) generate $\text{Aut}(\mathbb{C}_1^c)$. The first two automorphisms in (5.31) generate a subgroup isomorphic to $\text{SL}_2(\mathbb{Z})$, and we have in particular $\text{Aut}(\mathbb{C}_1^c) \simeq \text{SL}_2(\mathbb{Z}) \times (\mathbb{k}^\times)^2$.

COROLLARY 5.18. — *Fix $c \in \mathbb{k} \setminus \{0, \pm 1\}$. Let $L_2 = \mathbb{k} \langle x^{\pm 1}, y^{\pm 1} \rangle$ and \mathbb{C}_1^c be the quantum torus defined by (5.30). Then $\text{qHAut}(L_2)$ surjects onto $\text{Aut}(\mathbb{C}_1^c)$.*

Remark 5.19. — Similarly to the case of the first Weyl algebra discussed in Remark 5.13, the H_0 -Poisson structure on \mathbb{C}_1^c is trivial for $c \neq +1$. For $c = +1$, we get that the H_0 -Poisson structure induced by (5.26) on $\mathbb{k}[x^{\pm 1}, y^{\pm 1}]$ is the Poisson bracket defined by $\{x, y\} = xy$ and $\{x, x\} = 0 = \{y, y\}$.

Example 5.20. — We work over $\mathbb{k} = \mathbb{C}$. Consider the quasi-Hamiltonian algebras L_2 from Example 5.17 and A_Q from §4.3, where Q consists of a unique arrow $v : 0 \rightarrow \infty$. After fusion of the idempotent $e_0 \in A_Q$ onto the unit of L_2 , we get an algebra A' which is quasi-Hamiltonian by Proposition 4.5. Fix $n \in \mathbb{N}^\times$ and $q \in \mathbb{C}^\times$ not a root of unity. Taking $(c_0, c_\infty) = (q, q^{-n})$ and $(\alpha_0, \alpha_\infty) = (n, 1)$, the corresponding affine moduli space of representations \mathcal{C}_n^q is the q -Calogero–Moser space [CF17, Obl04]. Since $\text{qHAut}(L_2) \subset \text{qHAut}(A')$ descends to an algebra acting by Poisson automorphisms on representation spaces, we get an action of $\text{SL}_2(\mathbb{Z})$ by Poisson automorphisms on \mathcal{C}_n^q .

6. Dual integrable systems from quivers

We have seen as part of Theorems 3.12 and 4.12 that the (quasi-)Hamiltonian algebra associated with a quiver only depends on the underlying undirected graph, up to isomorphism. In particular, we obtain from Proposition 5.9 an isomorphism of Poisson varieties after considering the orbit spaces obtained by (quasi-)Hamiltonian reduction of the representation spaces of these algebras. We will investigate this observation using a family of quivers denoted Q_m , and their opposites Q_m^{op} obtained by reversing all the arrows in each quiver Q_m . The motivation underlying this investigation is that the quiver should also dictate a particular choice of local coordinates on a subset of the associated Poisson variety, such that the isomorphism described above will lead us to dual integrable systems. When the quiver is a cyclic quiver extended by one arrow, the choice of local coordinates that we consider on the associated space will satisfy the principle **(P)** stated as follows

- the matrices representing the arrows in the cyclic quiver are diagonal matrices, which are either equal or related through the moment map equations;
- the matrix representing the additional arrow is a (co)vector with all entries equal to $+1$.

These two rules will serve us to fix the choice of representatives, up to a residual finite action. Furthermore, we will see that this choice naturally ensures that we obtain Lax matrices for integrable systems in the Calogero–Moser (CM) and Ruijsenaars–Schneider (RS) families [Cal71, Mos75, RS86, Sut71]. By following this general idea, we will see that the reversing of arrows $Q_m \rightsquigarrow Q_m^{op}$ yields

- (1) self-duality of the generalised rational CM systems from [CS17] in § 6.1;
- (2) the standard duality of the hyperbolic CM and rational RS systems in § 6.2;
- (3) the new dual integrable system of a modified hyperbolic RS system in § 6.3;
- (4) the standard self-duality of the hyperbolic RS system in § 6.4.

For the interested reader which has only a limited background on integrable systems, we recommend to consult [Eti07, Lectures 1-2] for an introduction to the subject of CM systems using the perspective of Hamiltonian reduction.

Remark 6.1. — Below, we work over $\mathbb{k} = \mathbb{C}$, we fix $m \geq 1$ and we let $I = \mathbb{Z}/m\mathbb{Z}$. All the computations provided are given under the assumption that $m \geq 2$ to simplify the presentation. The statements regarding the local coordinates and their Poisson brackets also hold in the case $m = 1$, and the computations in that case can be easily adapted by the reader.

6.1. Self-duality of rational CM systems

We consider the quiver Q_m with vertex set $\tilde{I} = I \cup \{\infty\}$, and $m + 1$ arrows given by $x_s : s \rightarrow s + 1$ for each $s \in I$, and $v : \infty \rightarrow 0$. In the double \bar{Q}_m , we denote the opposite arrows by $x_s^* : s + 1 \rightarrow s$, $v^* : 0 \rightarrow \infty$. We also consider the quiver Q_m^{op} with the same vertex set but with arrows $y_s : s + 1 \rightarrow s$, $w : 0 \rightarrow \infty$. The double \bar{Q}_m^{op} has additional arrows $y_s^* : s \rightarrow s + 1$, $w^* : \infty \rightarrow 0$. These quivers are depicted in Figure 6.1.

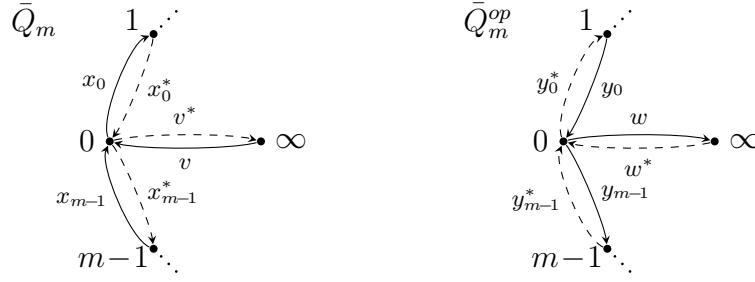


Figure 6.1. On the left : the double quiver \bar{Q}_m whose continuous arrows belong to the original quiver Q_m . On the right : the double quiver \bar{Q}_m^{op} whose continuous arrows belong to the original quiver Q_m^{op} .

Using § 3.3, the path algebra of each quiver has a Hamiltonian algebra structure, and by Theorem 3.12 and Example 3.1 we have an isomorphism of Hamiltonian algebras given by

$$(6.1) \quad \psi : \mathbb{C}\bar{Q}_m \rightarrow \mathbb{C}\bar{Q}_m^{op}, \quad \psi(x_s) = y_s^*, \quad \psi(x_s^*) = -y_s, \quad \psi(v) = w^*, \quad \psi(v^*) = -w.$$

Denote the moment maps by μ and μ^{op} . In view of Proposition 5.9, the map ψ induces a Poisson isomorphism

$$(6.2) \quad \Psi : \text{Rep} \left(\mathbb{C}\bar{Q}_m / (\mu - \tilde{\lambda}), \alpha \right) // \text{GL}_\alpha \rightarrow \text{Rep} \left(\mathbb{C}\bar{Q}_m^{op} / (\mu^{op} - \tilde{\lambda}), \alpha \right) // \text{GL}_\alpha,$$

for any $\alpha \in \mathbb{N}^{m+1}$ and $\tilde{\lambda} = \sum_{s \in I} \lambda_s e_s + \lambda_\infty e_\infty$ with $\lambda_s, \lambda_\infty \in \mathbb{C}$. We will consider the cases where the dimension vector α is such that for some $n \geq 1$ we put $\alpha_s = n$ for each $s \in I$ and $\alpha_\infty = 1$, while the parameters defining $\tilde{\lambda}$ are subject to $\lambda_\infty = -n|\lambda|$ for $|\lambda| := \sum_s \lambda_s$. Moreover the (λ_s) are assumed to satisfy the following regularity conditions

$$(6.3) \quad \sum_{s \in I} \lambda_s \neq 0, \quad k \sum_{s \in I} \lambda_s \neq \lambda_r + \dots + \lambda_{r'},$$

for all $k \in \mathbb{Z}$ and $1 \leq r \leq r' < m - 1$. Under these conditions, the spaces appearing in (6.2) are smooth and irreducible of dimension $2n$ [CS17].

6.1.1. Space associated with Q_m

Denote the reduced space $\text{Rep}(\mathbb{C}\bar{Q}_m / (\mu - \tilde{\lambda}), \alpha) // \text{GL}_\alpha$ by \mathcal{C}_n . By construction it can be described as the set of matrices

$$(6.4) \quad X_s, X_s^* \in \text{Mat}_{n \times n}(\mathbb{C}), \quad V \in \text{Mat}_{1 \times n}(\mathbb{C}), \quad V^* \in \text{Mat}_{n \times 1}(\mathbb{C}),$$

satisfying the m relations

$$(6.5) \quad X_s X_s^* - X_{s-1}^* X_{s-1} - \delta_{s0} V^* V = \lambda_s \text{Id}_n,$$

where we identify the elements in the same orbit of the action

$$(6.6) \quad g \cdot (X_s, X_s^*, V, V^*) = (g_s X_s g_{s+1}^{-1}, g_{s+1} X_s^* g_s^{-1}, V g_0^{-1}, g_0 V^*), \\ g = (g_s) \in \text{GL}_n(\mathbb{C})^m.$$

In view of the principle **(P)**, we consider the subspace $\mathcal{C}'_n \subset \mathcal{C}_n$ where there exists a representative with

$$(6.7) \quad X_s = \text{diag}(q_1, \dots, q_n), \quad s \in I, \quad V = (1, \dots, 1),$$

where $q_i \neq 0$ and $q_i^m \neq q_j^m$ for all $i \neq j$. Solving the constraint (6.5), we see that we can take

$$(6.8) \quad \begin{aligned} (X_s^*)_{ij} &= \delta_{ij} p_j + \delta_{ij} \frac{1}{q_j} (\lambda_1 + \dots + \lambda_s) - \delta_{(i \neq j)} |\lambda| \frac{q_i^{m-s-1} q_j^s}{q_i^m - q_j^m}, \quad s \in I, \\ V^* &= -|\lambda|(1, \dots, 1)^T, \end{aligned}$$

where the $(p_i) \in \mathbb{C}^n$ are free, and we recall $|\lambda| = \sum_s \lambda_s$. This choice is unique up to the action by the generalised symmetric group $\mathbb{Z}_m \wr S_n$, where S_n acts by simultaneous permutation of the (q_i, p_i) , while $(k_i) \in \mathbb{Z}_m^n$ acts by $(q_i, p_i) \mapsto (q_i \mu_m^{k_i}, p_i)$ where μ_m is a fixed primitive m^{th} root of unity.

LEMMA 6.2. — *The reduced Poisson bracket is such that $\{q_i, q_j\} = 0 = \{p_i, p_j\}$ and $\{q_i, p_j\} = \frac{1}{m} \delta_{ij}$.*

Proof. — By construction, the double bracket on $\mathbb{C}\bar{Q}_m$ is such that $\{\{x_r, x_s^*\}\} = \delta_{rs} e_{r+1} \otimes e_r$, $\{\{x_r, x_s\}\} = 0$ and $\{\{x_r^*, x_s^*\}\} = 0$. Denote by $\{-, -\}^\lambda$ the H_0 -Poisson structure on $A^\lambda := \mathbb{C}\bar{Q}_m / (\mu - \tilde{\lambda})$ and $\bar{a} \in H_0(A^\lambda)$ the image of an element $a \in A^\lambda$. We can compute from these results that for $x := x_0 \dots x_{m-1}$ and any $k, l \in \mathbb{N}$,

$$\begin{aligned} \{\overline{x^k}, \overline{x^l}\}^\lambda &= 0, \quad \{\overline{x^k}, \overline{x_{m-1}^* x_{m-1} x^l}\}^\lambda = k \overline{x^{k+l}}, \\ \{\overline{x_{m-1}^* x_{m-1} x^k}, \overline{x_{m-1}^* x_{m-1} x^l}\}^\lambda &= (k - l) \overline{x_{m-1}^* x_{m-1} x^{k+l}}. \end{aligned}$$

Hence by Theorem 5.8, if we let $X_{cyc} = X_0 \dots X_{m-1}$ we get on \mathcal{C}_n

$$(6.9) \quad \begin{aligned} \{\text{tr } X_{cyc}^k, \text{tr } X_{cyc}^l\} &= 0, \quad \{\text{tr } X_{cyc}^k, \text{tr } X_{m-1}^* X_{m-1} X_{cyc}^l\} = k \text{tr } X_{cyc}^{k+l}, \\ \{\text{tr } X_{m-1}^* X_{m-1} X_{cyc}^k, \text{tr } X_{m-1}^* X_{m-1} X_{cyc}^l\} &= (k - l) \text{tr } X_{m-1}^* X_{m-1} X_{cyc}^{k+l}. \end{aligned}$$

It remains to use these identities and the local coordinates as in the proof of [Eti07, Proposition 2.7]. □

The above description of \mathcal{C}_n was essentially given by Chalykh and Silantyev [CS17, Section V].

6.1.2. Space associated with Q_m^{op}

Denote the reduced space $\text{Rep}(\mathbb{C}\bar{Q}_m^{op} / (\mu^{op} - \tilde{\lambda}), \alpha) // \text{GL}_\alpha$ by \mathcal{C}_n^{op} . As in the previous case, it can be described as the set of matrices

$$(6.10) \quad Y_s, Y_s^* \in \text{Mat}_{n \times n}(\mathbb{C}), \quad W^* \in \text{Mat}_{1 \times n}(\mathbb{C}), \quad W \in \text{Mat}_{n \times 1}(\mathbb{C}),$$

satisfying

$$(6.11) \quad Y_{s-1} Y_{s-1}^* - Y_s^* Y_s + \delta_{s0} W W^* = \lambda_s \text{Id}_n,$$

under identifications of the elements of each orbit for the action

$$(6.12) \quad \begin{aligned} g \cdot (Y_s, Y_s^*, W, W^*) &= (g_{s+1} Y_s g_s^{-1}, g_s Y_s^* g_{s+1}^{-1}, g_0 W, W^* g_0^{-1}), \\ g &= (g_s) \in \mathrm{GL}_n(\mathbb{C})^m. \end{aligned}$$

Following the principle **(P)**, we consider the subspace $(\mathcal{C}_n^{op})' \subset \mathcal{C}_n^{op}$ where there exists a representative with

$$(6.13) \quad Y_s = \mathrm{diag}(\mathring{q}_1, \dots, \mathring{q}_n), \quad s \in I, \quad W = (1, \dots, 1)^T,$$

where $\mathring{q}_i \neq 0$ and $\mathring{q}_i^m \neq \mathring{q}_j^m$ for all $i \neq j$. We can then take

$$(6.14) \quad \begin{aligned} (Y_s^*)_{ij} &= \delta_{ij} \mathring{p}_j - \delta_{ij} \frac{1}{\mathring{q}_j} (\lambda_1 + \dots + \lambda_s) - \delta_{(i \neq j)} |\lambda| \frac{\mathring{q}_i^s \mathring{q}_j^{m-s-1}}{\mathring{q}_i^m - \mathring{q}_j^m}, \quad s \in I, \\ W^* &= |\lambda| (1, \dots, 1). \end{aligned}$$

This choice is unique up to the action of $\mathbb{Z}_m \wr S_n$.

LEMMA 6.3. — *The reduced Poisson bracket is such that $\{\mathring{q}_i, \mathring{q}_j\} = 0 = \{\mathring{p}_i, \mathring{p}_j\}$ and $\{\mathring{q}_i, \mathring{p}_j\} = \frac{1}{m} \delta_{ij}$.*

Proof. — By construction, the double bracket on $\mathbb{C}\bar{Q}_m^{op}$ is such that $\{\{y_r, y_s^*\}\} = \delta_{rs} e_r \otimes e_{r+1}$, $\{\{y_r, y_s\}\} = 0$ and $\{\{y_r^*, y_s^*\}\} = 0$. It then suffices to adapt the proof of Lemma 6.2. \square

6.1.3. Duality

After rescaling the Poisson brackets on $\mathcal{C}_n, \mathcal{C}_n^{op}$ by a factor m , we have a set of Darboux coordinates on a dense subset of each space. Let $X_{cyc} = X_0 \dots X_{m-1}$, $X_{cyc}^* = X_{m-1}^* \dots X_0^*$, and $Y_{cyc} = Y_{m-1} \dots Y_0$, $Y_{cyc}^* = Y_0^* \dots Y_{m-1}^*$. It is clear that the Poisson isomorphism (6.2) is such that

$$\mathrm{tr} Y_{cyc}^k \circ \Psi = (-1)^{km} \mathrm{tr} (X_{cyc}^*)^k, \quad \mathrm{tr} X_{cyc}^k \circ \Psi^{-1} = \mathrm{tr} (Y_{cyc}^*)^k.$$

In particular, these identities expressed in terms of Darboux coordinates yield that $(\mathrm{tr}(Y_{cyc}^*)^k)_{k=1}^n$ form an integrable system on \mathcal{C}_n^{op} , and the same holds for $(\mathrm{tr}(X_{cyc}^*)^k)_{k=1}^n$ on \mathcal{C}_n . Moreover, they are action-angle dual by definition. In the coordinates (q_i, p_i) , the functions $(\mathrm{tr}(X_{cyc}^*)^k)_{k=1}^n$ define a generalisation of the CM system which was introduced by Chalykh and Silantyev [CS17]. The same holds for the functions $(\mathrm{tr}(Y_{cyc}^*)^k)_{k=1}^n$ in the coordinates $(\mathring{q}_i, \mathring{p}_i)$, so that we get self-duality for CM systems having $\mathbb{Z}_m \wr S_n$ symmetry.

Remark 6.4. — In the case $m = 1$, we recover the well-known duality of the CM system of type A_{n-1} [KKS78, Rui88, Wil98]. For $m = 2$ we get the duality of type B_n in view of [CS17, Example 5.6].

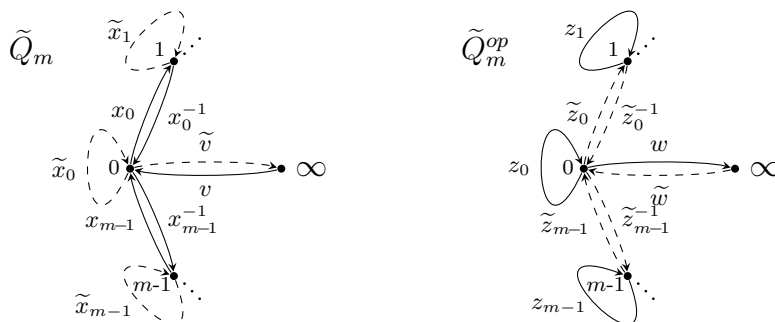


Figure 6.2. Representation as quivers with relations of the algebras obtained from $\mathbb{C}\bar{Q}_m$ and $\mathbb{C}\bar{Q}_m^{op}$ by localisation.

6.2. Duality of hyperbolic CM - rational RS systems

We modify §6.1 as follows. Consider the algebra \tilde{R}_m obtained by adding to $\mathbb{C}\bar{Q}_m$ local inverses $x_s^{-1} = e_{s+1}x_s^{-1}e_s$ satisfying

$$(6.15) \quad x_s x_s^{-1} = e_s, \quad x_s^{-1} x_s = e_{s+1}, \quad \text{for all } s \in I.$$

Introducing $\tilde{x}_s = x_s x_s^*$ and $\tilde{v} = v^*$, \tilde{R}_m is the path algebra of the quiver \tilde{Q}_m with relations (6.15) depicted on the left of Figure 6.2. The double bracket on $\mathbb{C}\bar{Q}_m$ uniquely extend to \tilde{R}_m by the derivation rule [VdB08a, Proposition 2.5.3], hence \tilde{R}_m is a Hamiltonian algebra.

Reproducing this construction, we can form the algebra \tilde{R}_m^{op} by adding local inverses $(y_s^*)^{-1} = e_{s+1}(y_s^*)^{-1}e_s$ to the path algebra $\mathbb{C}\bar{Q}_m^{op}$ of the double of the quiver Q_m^{op} . (Note that the superscript *op* indicates that we work with the opposite quiver, not that we take the opposite multiplication in \tilde{R}_m .) Taking $z_s = y_s^* y_s$, $\tilde{z}_s = y_s^*$, and $\tilde{w} = w^*$, \tilde{R}_m^{op} is the path algebra of the quiver \tilde{Q}_m^{op} with relation depicted on the right of Figure 6.2. We also obtain a Hamiltonian algebra structure on \tilde{R}_m^{op} such that the map ψ (6.1) extends to an isomorphism of Hamiltonian algebras

$$(6.16) \quad \tilde{\psi} : \tilde{R}_m \rightarrow \tilde{R}_m^{op}, \quad \tilde{\psi}(x_s) = \tilde{z}_s, \quad \tilde{\psi}(\tilde{x}_s) = -z_s, \quad \tilde{\psi}(v) = \tilde{w}, \quad \tilde{\psi}(\tilde{v}) = -w.$$

The map $\tilde{\psi}$ induces a Poisson isomorphism $\tilde{\Psi}$ on associated Poisson varieties as in §6.1. We consider the same dimension vector α and regular parameter $\tilde{\lambda}$ as in §6.1 for the rest of this subsection.

6.2.1. Space associated with \tilde{Q}_m

Denote the reduced space $\text{Rep}(\tilde{R}_m/(\mu - \tilde{\lambda}), \alpha) // \text{GL}_\alpha$ by $\tilde{\mathcal{C}}_n$. It is the subset of \mathcal{C}_n where each X_s is invertible, and where we use the elements $\tilde{X}_s = X_s X_s^*$, $\tilde{V} = V^*$. To

follow the principle **(P)** with the matrices (X_s, V) , we consider again the subspace $\mathcal{C}'_n \subset \tilde{\mathcal{C}}_n \subset \mathcal{C}_n$ where there exists a representative with

$$(6.17) \quad X_s = \text{diag}(x_1, \dots, x_n), \quad s \in I, \quad V = (1, \dots, 1),$$

where $x_i \neq 0$ and $x_i^m \neq x_j^m$ for all $i \neq j$, and where we take

$$(6.18) \quad (\tilde{X}_s)_{ij} = \delta_{ij}p_j + \delta_{ij}(\lambda_1 + \dots + \lambda_s) - \delta_{(i \neq j)}|\lambda| \frac{x_i^{m-s}x_j^s}{x_i^m - x_j^m}, \quad s \in I,$$

$$\tilde{V} = -|\lambda|(1, \dots, 1)^T,$$

for free parameters $(p_i) \in \mathbb{C}^n$. This choice is unique up to the action by $\mathbb{Z}_m \wr S_n$. The variables (x, p) correspond to (q, qp) in the choice associated with Q_m in §6.1, so that we get the following result from Lemma 6.2.

LEMMA 6.5. — *The reduced Poisson bracket is such that $\{x_i, x_j\} = 0 = \{p_i, p_j\}$ and $\{x_i, p_j\} = \frac{1}{m}x_i\delta_{ij}$.*

In particular, on a dense subspace of \mathcal{C}'_n we have Darboux coordinates (q_i, p_i) for $x_i = e^{\frac{1}{m}q_i}$.

6.2.2. Space associated with \tilde{Q}_m^{op}

Denote the reduced space $\text{Rep}(\tilde{R}_m^{op}/(\mu^{op} - \tilde{\lambda}), \alpha) // \text{GL}_\alpha$ by $\tilde{\mathcal{C}}_n^{op}$. The space $\tilde{\mathcal{C}}_n^{op} \subset \mathcal{C}_n^{op}$ can be described as the set of matrices

$$(6.19) \quad Z_s \in \text{Mat}_{n \times n}(\mathbb{C}), \quad \tilde{Z}_s \in \text{GL}_n(\mathbb{C}), \quad \tilde{W} \in \text{Mat}_{1 \times n}(\mathbb{C}), \quad W \in \text{Mat}_{n \times 1}(\mathbb{C}),$$

satisfying

$$(6.20) \quad \tilde{Z}_{s-1}^{-1}Z_{s-1}\tilde{Z}_{s-1} - Z_s + \delta_{s0}W\tilde{W} = \lambda_s \text{Id}_n,$$

under identification of the elements in each orbit for the action

$$(6.21) \quad g \cdot (Z_s, \tilde{Z}_s, W, \tilde{W}) = (g_s Z_s g_s^{-1}, g_s \tilde{Z}_s g_s^{-1}, g_0 W, \tilde{W} g_0^{-1}),$$

$$g = (g_s) \in \text{GL}_n(\mathbb{C})^m.$$

To apply the principle **(P)** to the matrices (Z_s, W) representing the continuous arrows of \tilde{Q}_m^{op} , we need the Z_s to be related diagonal matrices. In view of (6.20) with $s \neq 0$, Z_{s-1} and $Z_s + \lambda_s \text{Id}_n$ share the same spectrum. So we consider the subspace $(\tilde{\mathcal{C}}_n^{op})' \subset \tilde{\mathcal{C}}_n^{op}$ where there exists a representative with

$$(6.22) \quad Z_s = \text{diag}(\tilde{q}_1, \dots, \tilde{q}_n) - (\lambda_1 + \dots + \lambda_s) \text{Id}_n, \quad s \in I, \quad W = (1, \dots, 1)^T,$$

and we assume that

$$(6.23) \quad \tilde{q}_i \neq 0, \quad \tilde{q}_i \neq \tilde{q}_j, \quad \tilde{q}_i \neq \tilde{q}_j \pm |\lambda|, \quad \text{for all } i \neq j.$$

LEMMA 6.6. — *The subspace $(\tilde{\mathcal{C}}_n^{op})'$ can be parametrised by $(\tilde{q}_i) \in \mathbb{C}^n$ satisfying (6.23) and $(\omega_i) \in (\mathbb{C}^\times)^n$ where the matrices (Z_s, W) are given by (6.22), together with*

$$\begin{aligned} \tilde{Z}_s &= \text{Id}_n, \quad s \neq m-1, \quad (\tilde{Z}_{m-1})_{ij} = -|\lambda| \frac{\omega_i}{\tilde{q}_i - \tilde{q}_j - |\lambda|} \prod_{k \neq j} \frac{\tilde{q}_k - \tilde{q}_j - |\lambda|}{\tilde{q}_k - \tilde{q}_j}, \\ \tilde{W}_j &= |\lambda| \prod_{k \neq j} \frac{\tilde{q}_k - \tilde{q}_j - |\lambda|}{\tilde{q}_k - \tilde{q}_j}. \end{aligned}$$

Moreover, this choice is unique up to S_n action by simultaneous permutation of entries.

Proof. — Let us start with a representative such that (6.22) holds. By assumption on (\tilde{q}_i) and using (6.20) with $s \neq 0$, \tilde{Z}_s is diagonal, and we can fix the gauge (up to a residual permutation action) so that $\tilde{Z}_s = \text{Id}_n$. If we set $\omega_i = (\tilde{Z}_{m-1}W)_i$, we get from (6.20) with $s = 0$ that

$$(6.24) \quad (\tilde{Z}_{m-1})_{ij} = -\frac{\omega_i}{\tilde{q}_i - \tilde{q}_j - |\lambda|} \tilde{W}_j.$$

Using again the same equation in the form $\tilde{Z}_{m-1}^{-1}Z_{m-1}\tilde{Z}_{m-1} = Z_0 + \lambda_0 \text{Id}_n - W\tilde{W}$, we get that Z_{m-1} has the same spectrum as the right-hand side. Since the eigenvalues of Z_{m-1} are distinct, we must have the equality of characteristic polynomials in η

$$(6.25) \quad \begin{aligned} \prod_{k=1}^n (\tilde{q}_k - \lambda_1 - \dots - \lambda_{m-1} - \eta) &= \det(Z_{m-1} - \eta \text{Id}_n) \\ &= \det(Z_0 + (\lambda_0 - \eta) \text{Id}_n - W\tilde{W}). \end{aligned}$$

The right-hand side of (6.25) is a rank one deformation of a generically invertible matrix, so

$$\begin{aligned} \det(Z_0 + (\lambda_0 - \eta) \text{Id}_n - W\tilde{W}) &= \det(Z_0 + (\lambda_0 - \eta) \text{Id}_n) \left[1 - \tilde{W}(Z_0 + (\lambda_0 - \eta) \text{Id}_n)^{-1}W \right] \\ &= \prod_{k=1}^n (\tilde{q}_k + \lambda_0 - \eta) - \sum_{l=1}^n \tilde{W}_l \prod_{k \neq l} (\tilde{q}_k + \lambda_0 - \eta). \end{aligned}$$

Using this expression and evaluating (6.25) at $\eta = \tilde{q}_l + \lambda_0$, we get the claimed entries for \tilde{W} , hence for \tilde{Z}_{m-1} . It is then easy to check that $\omega_i = (\tilde{Z}_{m-1}W)_i$. By invertibility of \tilde{Z}_{m-1} we get $\omega_i \neq 0$ for all i . □

LEMMA 6.7. — *The reduced Poisson bracket is such that $\{\tilde{q}_i, \tilde{q}_j\} = 0 = \{\omega_i, \omega_j\}$ and $\{\tilde{q}_i, \omega_j\} = \delta_{ij}\omega_j$.*

Proof. — The double bracket on \tilde{R}_m^{op} is obtained by extending the one on $\mathbb{C}\tilde{Q}_m^{op}$ as explained at the beginning of this subsection. We can obtain in that way

$$\{\tilde{z}_r, \tilde{z}_s\} = 0, \quad \{\tilde{z}_0, \tilde{z}_s\} = \delta_{s0}e_0 \otimes \tilde{z}_0, \quad \{\tilde{z}_0, z_0\} = e_0 \otimes z_0 - z_0 \otimes e_0.$$

Denote by $\{-, -\}^\lambda$ the H_0 -Poisson structure on $A^\lambda := \tilde{R}_m^{op}/(\mu - \tilde{\lambda})$ and $\bar{a} \in H_0(A^\lambda)$ the image of an element $a \in A^\lambda$. We can compute from the above double brackets that for $\tilde{z} := \tilde{z}_0 \dots \tilde{z}_{m-1}$ and any $k, l \in \mathbb{N}$,

$$\{\overline{z_0^k}, \overline{z_0^l}\}^\lambda = 0, \quad \{\overline{z_0^k}, \overline{\tilde{z}z_0^l}\}^\lambda = k \overline{\tilde{z}z_0^{k+l-1}}, \quad \{\overline{\tilde{z}z_0^k}, \overline{\tilde{z}z_0^l}\}^\lambda = \left[\sum_{r=1}^{k-1} - \sum_{r=1}^{l-1} \right] \overline{\tilde{z}z_0^{k+l-r-1} \tilde{z}z_0^r}.$$

Hence by Theorem 5.8, if we let $\tilde{Z}_{cyc} = \tilde{Z}_0 \dots \tilde{Z}_{m-1}$ we get on $\tilde{\mathcal{C}}_n^{op}$

$$(6.26) \quad \begin{aligned} \{\mathrm{tr} Z_0^k, \mathrm{tr} Z_0^l\} &= 0, & \{\mathrm{tr} Z_0^k, \mathrm{tr} \tilde{Z}_{cyc} Z_0^l\} &= k \mathrm{tr} \tilde{Z} Z_0^{k+l-1}, \\ \{\mathrm{tr} \tilde{Z}_{cyc} Z_0^k, \mathrm{tr} \tilde{Z}_{cyc} Z_0^l\} &= \left[\sum_{r=1}^{k-1} - \sum_{r=1}^{l-1} \right] \mathrm{tr} \left(\tilde{Z}_{cyc} Z_0^{k+l-r-1} \tilde{Z}_{cyc} Z_0^r \right). \end{aligned}$$

It remains to use these identities and the local coordinates to get the reduced Poisson structure. □

As a corollary, on a dense subspace of $\tilde{\mathcal{C}}_n^{op}$ we have Darboux coordinates $(\tilde{q}_i, \tilde{p}_i)$ for $\omega_i = e^{\tilde{p}_i}$.

6.2.3. Duality

So far, we have obtained Darboux coordinates on dense subsets of $\tilde{\mathcal{C}}_n, \tilde{\mathcal{C}}_n^{op}$. Denoting $X_{cyc} = X_0 \dots X_{m-1}$ and $\tilde{Z}_{cyc} = \tilde{Z}_0 \dots \tilde{Z}_{m-1}$, we can use the Poisson diffeomorphism $\tilde{\Psi} : \tilde{\mathcal{C}}_n \rightarrow \tilde{\mathcal{C}}_n^{op}$ to obtain that

$$\mathrm{tr} Z_0^k \circ \tilde{\Psi} = (-1)^k \mathrm{tr} \tilde{X}_0^k, \quad \mathrm{tr} X_{cyc}^k \circ \tilde{\Psi}^{-1} = \mathrm{tr} (\tilde{Z}_{cyc})^k.$$

Hence, we can conclude that $(\mathrm{tr} \tilde{X}_0^k)_{k=1}^n$ and $(\mathrm{tr} (\tilde{Z}_{cyc})^k)_{k=1}^n$ form integrable systems on $\tilde{\mathcal{C}}_n$ and $\tilde{\mathcal{C}}_n^{op}$ respectively, which are in action-angle duality. In the coordinates (q_i, p_i) we can write the functions $(\mathrm{tr} \tilde{X}_0^k)_{k=1}^n$ as hyperbolic CM Hamiltonians of type A_{n-1} of order km in the momenta (p_i) . In the coordinates $(\tilde{q}_i, \tilde{p}_i)$, the functions $\mathrm{tr} (\tilde{Z}_{cyc})^k$ are all defining Hamiltonians of the rational RS system of type A_{n-1} . The dependence on m of this second family is only visible in the coupling λ_∞ .

The construction of this subsection using two quivers and their representation spaces is similar to the choice of two slices inside one phase space outlined by Gorsky and Rubtsov [GR01, § 4.5].

Remark 6.8. — The duality of the hyperbolic CM system and rational RS system in the real case goes back to Ruijsenaars [Rui88]. Since we work over \mathbb{C} , we do not distinguish the hyperbolic and trigonometric CM systems which are equivalent up to making the change of coordinates $q_i \mapsto \sqrt{-1}q_i$. Another consequence of the fact that we work in the complex setting is that the second integrable system is written in the coordinates $(\tilde{q}_i, \tilde{p}_i)$ as rational RS Hamiltonians in MacDonald form. In that case, the Lax matrix and the Hamiltonians do not involve square roots contrary to the original real case [RS86]. We also use the MacDonald form for the hyperbolic RS system and its variants presented in the next subsections.

6.3. Modified hyperbolic RS systems and their duals

We consider the quivers Q_m and Q_m^{op} as in § 6.1. We construct the path algebras over \mathbb{C} of their doubles, depicted in Figure 6.1, then we define the algebras $A_m := A_{Q_m}$, $A_m^{op} := A_{Q_m^{op}}$ by universal localisation as in § 4.3. To avoid confusion with the results in § 6.1, we will denote the elements of the double with a hat instead of a star, e.g. x_s^* will be denoted \hat{x}_s . We fix the following ordering $<_s$ on the elements a of the doubles such that $t(a) = s$,

$$(6.27) \quad \begin{aligned} x_s <_s \hat{x}_{s-1}, \quad s \in I \setminus \{0\}, \quad \text{and} \quad x_0 <_0 \hat{x}_{m-1} <_0 \hat{v}; \\ \hat{y}_s <_s y_{s-1}, \quad s \in I \setminus \{0\}, \quad \text{and} \quad \hat{y}_0 <_0 y_{m-1} <_0 w. \end{aligned}$$

(For $m = 1$, we only have $x_0 <_0 \hat{x}_0 <_0 \hat{v}$ or $\hat{y}_0 <_0 y_0 <_0 w$.) By § 4.3, this defines a quasi-Hamiltonian algebra structure on A_m and A_m^{op} . Using Theorem 4.12 and Example 4.1, the choice of orderings gives the following isomorphism of quasi-Hamiltonian algebras

$$(6.28) \quad \begin{aligned} \hat{\psi} : A_m \rightarrow A_m^{op}, \quad \hat{\psi}(x_s) = \hat{y}_s, \quad \hat{\psi}(\hat{x}_s) = -(e_{s+1} + y_s \hat{y}_s)^{-1} y_s, \\ \hat{\psi}(v) = \hat{w}, \quad \hat{\psi}(\hat{v}) = -(e_0 + w \hat{w})^{-1} w. \end{aligned}$$

Denote the multiplicative moment maps by Φ and Φ^{op} . In view of Proposition 5.9, the map $\hat{\psi}$ induces a Poisson isomorphism

$$(6.29) \quad \hat{\Psi} : \text{Rep}(A_m / (\Phi - \tilde{c}), \alpha) // \text{GL}_\alpha \rightarrow \text{Rep}(A_m^{op} / (\Phi^{op} - \tilde{c}), \alpha) // \text{GL}_\alpha,$$

for any $\alpha \in \mathbb{N}^{m+1}$ and $\tilde{c} = \sum_{s \in I} c_s e_s + c_\infty e_\infty$ with $c_s, c_\infty \in \mathbb{C}^\times$. We will consider the cases where $\alpha_s = n$ for each $s \in I$ with $n \geq 1$, $\alpha_\infty = 1$, while $c_\infty = (\prod_s c_s)^{-n}$ and the (c_s) are subject to the regularity conditions

$$(6.30) \quad \prod_{s \in I} c_s \neq 1, \quad \prod_{s \in I} c_s^k \neq \prod_{r \leq \rho \leq r'} c_\rho,$$

for all $k \in \mathbb{Z}$ and $1 \leq r \leq r' < m - 1$. Under these conditions, the spaces appearing in (6.29) are smooth and irreducible of dimension $2n$ [BEF20, CF17].

6.3.1. Space associated with Q_m

Denote the reduced space $\text{Rep}(A_m / (\Phi - \tilde{c}), \alpha) // \text{GL}_\alpha$ by $\hat{\mathcal{C}}_n$. By construction it can be described as the set of matrices

$$(6.31) \quad X_s, \hat{X}_s \in \text{Mat}_{n \times n}(\mathbb{C}), \quad V \in \text{Mat}_{1 \times n}(\mathbb{C}), \quad \hat{V} \in \text{Mat}_{n \times 1}(\mathbb{C}),$$

satisfying the m relations

$$(6.32) \quad (\text{Id}_n + X_s \hat{X}_s) (\text{Id}_n + \hat{X}_{s-1} X_{s-1})^{-1} (\text{Id}_n + \delta_{s0} \hat{V} V)^{-1} = c_s \text{Id}_n,$$

where all the factors appearing in (6.32) are invertible, and we identify the elements in the same orbit of the action

$$(6.33) \quad \begin{aligned} g \cdot (X_s, \hat{X}_s, V, \hat{V}) &= (g_s X_s g_{s+1}^{-1}, g_{s+1} \hat{X}_s g_s^{-1}, V g_0^{-1}, g_0 \hat{V}), \\ g &= (g_s) \in \text{GL}_n(\mathbb{C})^m. \end{aligned}$$

In view of the principle **(P)**, we consider the subspace $\widehat{\mathcal{C}}'_n \subset \widehat{\mathcal{C}}_n$ where there exists a representative with

$$(6.34) \quad X_s = \text{diag}(x_1, \dots, x_n), \quad s \in I, \quad V = (1, \dots, 1),$$

where the $(x_i) \in \mathbb{C}^n$ satisfy the following conditions with $t := \prod_{s \in I} c_s$,

$$(6.35) \quad x_i \neq 0, \quad x_i^m \neq x_j^m, \quad x_i^m \neq tx_j^m, \quad \text{for all } i \neq j.$$

LEMMA 6.9. — *The subspace $\widehat{\mathcal{C}}'_n$ can be parametrised by $(x_i) \in \mathbb{C}^n$ satisfying (6.35) and $(\rho_i) \in (\mathbb{C}^\times)^n$ where the matrices (X_s, V) are given by (6.34), together with⁽⁴⁾*

$$\begin{aligned} \widehat{X}_s &= -\delta_{ij} \frac{1}{x_i} + c_1 \dots c_s (1 - t^{-1}) \rho_j \frac{x_i^{m-s-1} x_j^s}{x_j^m - t^{-1} x_i^m} \prod_{l \neq i} \frac{x_l^m - t^{-1} x_i^m}{x_l^m - x_i^m}, \\ \widehat{V}_i &= - (1 - t^{-1}) \prod_{l \neq i} \frac{x_l^m - t^{-1} x_i^m}{x_l^m - x_i^m}. \end{aligned}$$

Moreover, this choice is unique up to $\mathbb{Z}_m \wr S_n$ action by simultaneous permutation of entries for S_n , and rescaling $x_i \mapsto \mu_m^{k_i} x_i$ for $(k_i) \in \mathbb{Z}_m^n$ with μ_m a fixed primitive m^{th} root of unity.

Proof. — In $\widehat{\mathcal{C}}'_n$, we can rewrite the moment map conditions (6.32) as

$$(6.36) \quad X_s (\widehat{X}_s + X_s^{-1}) X_{s-1}^{-1} (\widehat{X}_{s-1} + X_{s-1}^{-1})^{-1} (\text{Id}_n + \delta_{s0} \widehat{V} V)^{-1} = c_s \text{Id}_n.$$

Taking a representative such that (6.34) holds, we can use these identities to get

$$(\widehat{X}_s + X_s^{-1})_{ij} = c_1 \dots c_s x_i^{-s} x_j^s (\widehat{X}_0 + X_0^{-1})_{ij}, \quad (\widehat{X}_0 + X_0^{-1})_{ij} = t \frac{x_i^{-1} \widehat{V}_i \rho_j}{1 - tx_i^{-m} x_j^m},$$

where we have set

$$(6.37) \quad \rho_i = (VD^{-m+1} (\widehat{X}_0 + X_0^{-1}) D^m)_i, \quad \text{for } D = \text{diag}(x_1, \dots, x_n).$$

From this, we obtain that if \widehat{V} has the form claimed in the statement, then it will follow that it is true also for the matrices \widehat{X}_s . In particular, it can then be checked that (6.37) is satisfied for these particular matrices.

To determine the entries of \widehat{V} for a representative with (6.34), we note that the moment map condition implies

$$(6.38) \quad t^{-1} (\widehat{X}_0 + X_0^{-1}) D^{-m} (\widehat{X}_0 + X_0^{-1})^{-1} = D^{-m} + D^{-1} \widehat{V} V D^{-m+1}.$$

The matrices on both sides of this equality share the same spectrum hence, as in the proof of Lemma 6.6, equality of their characteristic polynomials in η yields

$$(6.39) \quad \prod_{k=1}^n (t^{-1} x_k^{-m} - \eta) = \prod_{k=1}^n (x_k^{-m} - \eta) + \sum_l \widehat{V}_l x_l^{-m} \prod_{k \neq l} (x_k^{-m} - \eta).$$

Evaluating this identity at $\eta = x_i^{-m}$, we get the desired \widehat{V}_i .

⁽⁴⁾In the local form of \widehat{X}_s , we consider that $c_1 \dots c_s$ for $s = 0$ is the empty product equal to $+1$. Hereafter, we follow this convention.

Finally, we note that $\rho_i \neq 0$ because the matrices $\widehat{X}_s + X_s^{-1}$ are invertible. The uniqueness of the representative up to $\mathbb{Z}_m \wr S_n$ action is easily obtained. \square

LEMMA 6.10. — *The reduced Poisson bracket is such that $\{x_i, x_j\} = 0 = \{\rho_i, \rho_j\}$ and $\{x_i, \rho_j\} = \frac{1}{m} \delta_{ij} x_i \rho_j$.*

Proof. — Due to the choice of ordering, we get from §4.3 that the double bracket on A_m satisfies

$$\begin{aligned} \{\{x_s, x_r\}\} &= \frac{1}{2} \delta_{s, r+1} x_r x_{r+1} \otimes e_{r+1} - \frac{1}{2} \delta_{s, r-1} e_r \otimes x_{r-1} x_r, \quad \{\{\widehat{x}_0, \widehat{x}_0\}\} = 0, \\ \{\{x_s, \widehat{x}_0\}\} &= \delta_{s, 0} \left[e_1 \otimes e_0 + \frac{1}{2} \widehat{x}_0 x_0 \otimes e_0 + \frac{1}{2} e_1 \otimes x_0 \widehat{x}_0 \right] \\ &\quad - \frac{1}{2} \delta_{s, 1} x_1 \otimes \widehat{x}_0 + \frac{1}{2} \delta_{s, m-1} \widehat{x}_0 \otimes x_{m-1}. \end{aligned}$$

Introducing $x := x_0 \dots x_{m-1}$ and $\check{x}_0 := e_0 + x_0 \widehat{x}_0$, we can then get

$$(6.40) \quad \begin{aligned} \{\{x, x\}\} &= \frac{1}{2} (x^2 \otimes e_0 - e_0 \otimes x^2), \quad \{\{\check{x}_0, \check{x}_0\}\} = \frac{1}{2} (\check{x}_0^2 \otimes e_0 - e_0 \otimes \check{x}_0^2), \\ \{\{x, \check{x}_0\}\} &= \frac{1}{2} (x \otimes \check{x}_0 + \check{x}_0 \otimes x + \check{x}_0 x \otimes e_0 - e_0 \otimes x \check{x}_0). \end{aligned}$$

Denote by $\{-, -\}^c$ the H_0 -Poisson structure on $A^c := A_m / (\Phi - \tilde{c})$ and $\bar{a} \in H_0(A^c)$ the image of an element $a \in A^c$. In a way similar to the proof of [CF17, Proposition 4.4], we can compute from (6.40) that for any $k, l \in \mathbb{N}$,

$$\begin{aligned} \{\overline{x^k}, \overline{x^l}\}^c &= 0, \quad \{\overline{x^k}, \overline{\check{x}_0 x^l}\}^c = k \overline{\check{x}_0 x^{k+l}}, \\ \{\overline{\check{x}_0 x^k}, \overline{\check{x}_0 x^l}\}^c &= \left[\sum_{r=1}^k - \sum_{r=1}^l \right] \overline{\check{x}_0 x^{k+l-r} \check{x}_0 x^r}. \end{aligned}$$

Hence by Theorem 5.8, if we let $X_{cyc} := X_0 \dots X_{m-1}$ and $\check{X}_0 := \text{Id}_n + X_0 \widehat{X}_0$, we get on $\widehat{\mathcal{C}}_n$

$$(6.41) \quad \begin{aligned} \{\text{tr } X_{cyc}^k, \text{tr } X_{cyc}^l\} &= 0, \quad \{\text{tr } X_{cyc}^k, \text{tr } \check{X}_0 X_{cyc}^l\} = k \text{tr } \check{X}_0 X_{cyc}^{k+l}, \\ \{\text{tr } \check{X}_0 X_{cyc}^k, \text{tr } \check{X}_0 X_{cyc}^l\} &= \left[\sum_{r=1}^k - \sum_{r=1}^l \right] \text{tr} \left(\check{X}_0 X_{cyc}^{k+l-r} \check{X}_0 X_{cyc}^r \right). \end{aligned}$$

We can then derive the reduced Poisson structure by writing these identities in local coordinates, which can be done e.g. by adapting [CF17, § 3.1]. \square

As an application of this lemma, we can get Darboux coordinates (q_i, p_i) on a dense subset of $\widehat{\mathcal{C}}_n$ by considering $x_i = e^{\frac{1}{m} q_i}$ and $\rho_i = e^{p_i}$.

Remark 6.11. — As part of the local computations needed to prove Lemma 6.10, we can get that

$$(6.42) \quad \{\nu_i, \nu_j\} = \nu_i \nu_j \frac{x_i^m + x_j^m}{x_i^m - x_j^m} \frac{(t^{-1} - 1)^2 x_i^m x_j^m}{(x_i^m - t^{-1} x_j^m)(x_j^m - t^{-1} x_i^m)}, \quad \nu_i := \rho_i \prod_{l \neq i} \frac{x_l^m - t^{-1} x_i^m}{x_l^m - x_i^m}.$$

The Poisson bracket of the elements (x_i, ν_i) was first obtained in this form by Fock and Rosly [FR99], see also [Obl04, § 2.4]. Note that the Poisson bracket (6.42) is invariant

under replacing t^{-1} by t . This explains the difference between our parametrisation in Lemma 6.9 and the one in [CF17, Sections 3,4]. There are another two ways to obtain Darboux coordinates, namely using the Ruijsenaars form as in [FR99, Appendix], or a form related to the q KP hierarchy [Ili00, Section 6]. We refer to [Fai19a, § 4.2.1] for a review of these different possibilities in the case $m = 1$, which are easily adapted to any $m \geq 1$.

6.3.2. Space associated with Q_m^{op}

Denote the reduced space $\text{Rep}(A_m^{op}/(\Phi^{op} - \tilde{c}), \alpha) // \text{GL}_\alpha$ by $\widehat{\mathcal{C}}_n^{op}$. It can be given as the set of matrices

$$(6.43) \quad Y_s, \widehat{Y}_s \in \text{Mat}_{n \times n}(\mathbb{C}), \quad \widehat{W} \in \text{Mat}_{1 \times n}(\mathbb{C}), \quad W \in \text{Mat}_{n \times 1}(\mathbb{C}),$$

satisfying the m relations

$$(6.44) \quad (\text{Id}_n + \widehat{Y}_s Y_s)^{-1} (\text{Id}_n + Y_{s-1} \widehat{Y}_{s-1}) (\text{Id}_n + \delta_{s0} W \widehat{W}) = c_s \text{Id}_n,$$

where all the factors appearing in (6.44) are invertible, and we identify the elements in the same orbit for the action

$$(6.45) \quad g \cdot (Y_s, \widehat{Y}_s, W, \widehat{W}) = (g_{s+1} Y_s g_s^{-1}, g_s \widehat{Y}_s g_{s+1}^{-1}, g_0 W, \widehat{W} g_0^{-1}), \\ g = (g_s) \in \text{GL}_n(\mathbb{C})^m.$$

In analogy with the space associated with Q_m , we consider the subspace $(\widehat{\mathcal{C}}_n^{op})' \subset \widehat{\mathcal{C}}_n^{op}$ where there exists a representative with

$$(6.46) \quad Y_s = \text{diag}(y_1, \dots, y_n), \quad s \in I, \quad W = (1, \dots, 1)^T,$$

where the $(y_i) \in \mathbb{C}^n$ satisfy the conditions (6.35), and we set again $t := \prod_{s \in I} c_s$.

LEMMA 6.12. — *The subspace $(\widehat{\mathcal{C}}_n^{op})'$ can be parametrised by $(y_i) \in \mathbb{C}^n$ satisfying (6.35) and $(\tau_i) \in (\mathbb{C}^\times)^n$ where the matrices (Y_s, W) are given by (6.46), together with*

$$\widehat{Y}_s = -\delta_{ij} \frac{1}{y_i} + c_1^{-1} \dots c_s^{-1} (1-t)^{\tau_i} \frac{y_i^s y_j^{m-s-1}}{y_i^m - t y_j^m} \prod_{k \neq j} \frac{y_k^m - t y_j^m}{y_k^m - y_j^m}, \\ \widehat{W}_j = -(1-t) \prod_{k \neq j} \frac{y_k^m - t y_j^m}{y_k^m - y_j^m}.$$

Moreover, this choice is unique up to $\mathbb{Z}_m \wr S_n$ action.

Proof. — The result can be derived in the same way as Lemma 6.9. To obtain the form of \widehat{W} , we remark that the moment map (6.44) implies the following identity :

$$(6.47) \quad (Y_0 + \widehat{Y}_0)^{-1} D^{-m} (Y_0 + \widehat{Y}_0) = t^{-1} D^{-m} + t^{-1} D^{1-m} W \widehat{W} D^{-1}, D \\ = \text{diag}(y_1, \dots, y_n).$$

It then suffices to equal both characteristic polynomials to determine \widehat{W} . \square

LEMMA 6.13. — *The reduced Poisson bracket is such that $\{y_i, y_j\} = 0 = \{\tau_i, \tau_j\}$ and $\{y_i, \tau_j\} = \frac{1}{m} \delta_{ij} y_i \tau_j$.*

Proof. — Due to the choice of ordering, the double bracket on A_m^{op} is such that

$$\begin{aligned} \{\{y_s, y_r\}\} &= \frac{1}{2} \delta_{s, r+1} e_{r+1} \otimes y_{r+1} y_r - \frac{1}{2} \delta_{s, r-1} y_r y_{r-1} \otimes e_r, \quad \{\{\hat{y}_0, \hat{y}_0\}\} = 0, \\ \{\{y_s, \hat{y}_0\}\} &= \delta_{s,0} \left[e_0 \otimes e_1 + \frac{1}{2} \hat{y}_0 y_0 \otimes e_1 + \frac{1}{2} e_0 \otimes y_0 \hat{y}_0 \right] \\ &\quad - \frac{1}{2} \delta_{s,1} \hat{y}_0 \otimes y_1 + \frac{1}{2} \delta_{s, m-1} y_{m-1} \otimes \hat{y}_0. \end{aligned}$$

Introducing $y := y_{m-1} \dots y_0$ and $\check{y}_0 := e_0 + \hat{y}_0 y_0$, we can then get

$$\begin{aligned} \{\{y, y\}\} &= \frac{1}{2} (e_0 \otimes y^2 - y^2 \otimes e_0), \quad \{\{\check{y}_0, \check{y}_0\}\} = \frac{1}{2} (e_0 \otimes \check{y}_0^2 - \check{y}_0^2 \otimes e_0), \\ \{\{y, \check{y}_0\}\} &= \frac{1}{2} (y \otimes \check{y}_0 + \check{y}_0 \otimes y + e_0 \otimes y \check{y}_0 - \check{y}_0 y \otimes e_0). \end{aligned}$$

If we let $Y_{cyc} := Y_{m-1} \dots Y_0$ and $\check{Y}_0 := \text{Id}_n + \hat{Y}_0 Y_0$, we get on $\hat{\mathcal{C}}_n^{op}$ the following identities

$$(6.48) \quad \begin{aligned} \{\{\text{tr } Y_{cyc}^k, \text{tr } Y_{cyc}^l\}\} &= 0, \quad \{\{\text{tr } Y_{cyc}^k, \text{tr } \check{Y}_0 Y_{cyc}^l\}\} = k \text{tr } \check{Y}_0 Y_{cyc}^{k+l}, \\ \{\{\text{tr } \check{Y}_0 Y_{cyc}^k, \text{tr } \check{Y}_0 Y_{cyc}^l\}\} &= \left[\sum_{r=1}^k - \sum_{r=1}^l \right] \text{tr} \left(\check{Y}_0 Y_{cyc}^{k+l-r} \check{Y}_0 Y_{cyc}^r \right), \end{aligned}$$

which can be derived as in Lemma 6.10. Notice that (6.41) and (6.48) are the same equations if we replace X_{cyc}, \check{X}_0 by Y_{cyc}, \check{Y}_0 , or vice-versa. If we also remark that the functions $\text{tr } X_{cyc}^k, \text{tr } \check{X}_0 X_{cyc}^l$ written in coordinates in $\hat{\mathcal{C}}_n'$, and the functions $\text{tr } Y_{cyc}^k, \text{tr } \check{Y}_0 Y_{cyc}^l$ written in coordinates in $(\hat{\mathcal{C}}_n^{op})'$ are exactly the same when we replace (x_j, ρ_j, t) by (y_j, τ_j, t^{-1}) , the statement follows from Lemma 6.10. \square

As a corollary, we can get Darboux coordinates (\hat{q}_i, \hat{p}_i) on a dense subset of $\hat{\mathcal{C}}_n^{op}$ by considering $y_i = e^{\frac{1}{m} \hat{q}_i}$ and $\tau_i = e^{\hat{p}_i}$.

6.3.3. Duality

We have obtained Darboux coordinates on dense subsets of $\hat{\mathcal{C}}_n, \hat{\mathcal{C}}_n^{op}$, which we now use to get integrable systems in action-angle duality. Introducing $X_{cyc} = X_0 \dots X_{m-1}, Y_{cyc} = Y_{m-1} \dots Y_0, \hat{Y}_{cyc} = \hat{Y}_0 \dots \hat{Y}_{m-1}$ and $L_{cyc} = L_{m-1} \dots L_0$ for $L_s = (\text{Id}_n + \hat{X}_s X_s)^{-1} \hat{X}_s$, we can use the Poisson diffeomorphism $\hat{\Psi} : \hat{\mathcal{C}}_n \rightarrow \hat{\mathcal{C}}_n^{op}$ to obtain that

$$\text{tr } \hat{Y}_{cyc}^k \circ \hat{\Psi} = \text{tr } X^k, \quad \text{tr } L_{cyc}^k \circ \hat{\Psi}^{-1} = (-1)^{km} \text{tr } Y_{cyc}^k.$$

Hence, we can conclude that

$$\left(\text{tr } L_{cyc}^k \right)_{k=1}^n \quad \text{and} \quad \left(\text{tr } \hat{Y}_{cyc}^k \right)_{k=1}^n$$

form integrable systems on $\hat{\mathcal{C}}_n$ and $\hat{\mathcal{C}}_n^{op}$ respectively, which are in action-angle duality. In the coordinates (\hat{q}_i, \hat{p}_i) we can write the functions $(\text{tr } \hat{Y}_{cyc}^k)_{k=1}^n$ as deformations of the hyperbolic (or trigonometric, see Remark 6.8) RS Hamiltonians of type A_{n-1} . They are written explicitly as the family $(\text{tr } Y^j)$ for $m = 1$, and the family $(H_{m,j})$ for $m \geq 2$ in [CF17], and their relation to other integrable systems is discussed.

In order to write the functions $(\text{tr } L_{cyc}^k)_{k=1}^n$ in the coordinates (q_i, p_i) , we note that on $\widehat{\mathcal{C}}'_n$

$$(6.49) \quad \text{Id}_n + \widehat{X}_s X_s = c_1 \dots c_s (1-t) T_s^L C T_s^R, \text{ where}$$

$$\left(T_s^L\right)_{ij} = \delta_{ij} x_i^{m-s-1} \prod_{k \neq i} \frac{x_k^m - t^{-1} x_i^m}{x_k^m - x_i^m}, \quad \left(T_s^R\right)_{ij} = \delta_{ij} \rho_j x_j^{s+1}, \quad C_{ij} = \frac{1}{x_i^m - t x_j^m}.$$

As C is a Cauchy matrix, its inverse can be computed to be

$$(6.50) \quad C_{ij}^{-1} = (1-t) (1-t^{-1}) \frac{x_i^m x_j^m}{x_i^m - t^{-1} x_j^m} \prod_{k \neq j} \frac{x_k^m - t^{-1} x_j^m}{x_k^m - x_j^m} \prod_{l \neq i} \frac{x_l^m - t x_i^m}{x_l^m - x_i^m}.$$

Therefore, the entries of $L_s = (\text{Id}_n + \widehat{X}_s X_s)^{-1} \widehat{X}_s$ are given by

$$(6.51) \quad (L_s)_{ij} = -\frac{1-t^{-1}}{c_1 \dots c_s} \rho_i^{-1} \frac{x_i^{m-s-1} x_j^s}{x_i^m - t^{-1} x_j^m} \prod_{l \neq i} \frac{x_l^m - t x_i^m}{x_l^m - x_i^m} \\ + (1-t) (1-t^{-1}) \sum_k \frac{x_i^{m-s-1} x_k^m x_j^s}{(x_i^m - t^{-1} x_k^m) (x_k^m - t x_j^m)} \frac{\rho_j}{\rho_i} \prod_{l \neq i} \frac{x_l^m - t x_i^m}{x_l^m - x_i^m} \prod_{a \neq k} \frac{x_a^m - t^{-1} x_k^m}{x_a^m - x_k^m}.$$

In the simplest case $m = 1$, the first element in the family $(\text{tr } L_{cyc}^k)_{k=1}^n$ can be written as

$$(6.52) \quad \text{tr } L_0 = -\sum_i \rho_i^{-1} \prod_{l \neq i} \frac{x_l - t x_i}{x_l - x_i} + (1-t^{-1})^2 \sum_{i,k} \frac{x_k}{(x_i - t^{-1} x_k)^2} \prod_{l \neq i} \frac{x_l - t x_i}{x_l - x_i} \prod_{a \neq k} \frac{x_a - t^{-1} x_k}{x_a - x_k}.$$

In terms of the Darboux coordinates (q_i, p_i) and $t = e^{2\gamma}$, $\gamma \in \mathbb{C}^\times$, this can be transformed into

$$(6.53) \quad \text{tr } L_0 = -\sum_i e^{-p_i} \prod_{l \neq i} \frac{\sinh\left(\frac{q_l - q_i}{2} - \gamma\right)}{\sinh\left(\frac{q_l - q_i}{2}\right)} \\ + \sum_{i,k} \frac{\sinh^2(\gamma) e^{-q_i}}{\sinh^2\left(\frac{q_i - q_k}{2} + \gamma\right)} \prod_{l \neq i} \frac{\sinh\left(\frac{q_l - q_i}{2} - \gamma\right)}{\sinh\left(\frac{q_l - q_i}{2}\right)} \prod_{a \neq k} \frac{\sinh\left(\frac{q_a - q_k}{2} + \gamma\right)}{\sinh\left(\frac{q_a - q_k}{2}\right)}.$$

The duality between the families $(\text{tr } L_{cyc}^k)_{k=1}^n$ and $(\text{tr } \widehat{Y}_{cyc}^k)_{k=1}^n$ seem to be new.

6.4. Self-duality of hyperbolic RS systems

We modify §6.3 as follows. Consider the algebra \check{A}_m obtained by adding to A_m local inverses $x_s^{-1} = e_{s+1} x_s^{-1} e_s$, i.e. these elements satisfy (6.15). We let $\check{x}_s = e_s + x_s \widehat{x}_s$ and $\check{v} = \widehat{v}$. Note that $x_s, v, \check{x}_s, \check{v}$, the idempotents and the inverses $x_s^{-1}, \check{x}_s^{-1}, (e_\infty + v\check{v})^{-1}, (e_0 + \check{v}v)^{-1}$ generate \check{A}_m . Moreover, \check{A}_m is a quasi-Hamiltonian algebra if we extend the double bracket from A_{Q_m} described in §6.3 by localisation.

In the same way, we introduce the algebra \check{A}_m^{op} obtained by adding to A_m^{op} defined in §6.3 local inverses $y_s^{-1} = e_s y_s^{-1} e_{s+1}$. We use $z_s = e_s + \widehat{y}_s y_s, \check{z}_s = \widehat{y}_s, \check{w} = \widehat{w}$ and

w together with the idempotents and the inverses belonging to \check{A}_m^{op} as generators of this algebra. It is also a quasi-Hamiltonian algebra by localisation. By construction, the map $\check{\psi}$ (6.28) extends to an isomorphism of quasi-Hamiltonian algebras which can be written as

$$(6.54) \quad \begin{aligned} \check{\psi} : \check{A}_m &\rightarrow \check{A}_m^{op}, & \check{\psi}(x_s) &= \check{z}_s, & \check{\psi}(\check{x}_s) &= z_s^{-1}, \\ & & \check{\psi}(v) &= \check{w}, & \check{\psi}(\check{v}) &= -(e_0 + w\check{w})^{-1}w. \end{aligned}$$

The map $\check{\psi}$ induces a Poisson isomorphism $\check{\Psi}$ on associated Poisson varieties as in § 6.3. We consider the same dimension vector α and regular parameter \tilde{c} as in § 6.3 for the rest of this subsection.

To use the principle **(P)** in order to find dual integrable systems, we will see \check{A}_m and \check{A}_m^{op} as algebras attached to the quivers with relations $\check{Q}_m, \check{Q}_m^{op}$ depicted in Figure 6.2. (We use \check{x}_s instead of \tilde{x}_s and do the same for $\check{v}, \check{z}_s, \check{w}$ to avoid confusion with the cases considered in § 6.2.)

6.4.1. Space associated with \check{Q}_m

Denote the reduced space $\text{Rep}(\check{A}_m/(\Phi - \tilde{c}), \alpha) // \text{GL}_\alpha$ by $\check{\mathcal{C}}_n$. It is the subset of $\hat{\mathcal{C}}_n$ described in § 6.3 where each X_s is invertible, and where we use the elements $\check{X}_s = \text{Id}_n + X_s \hat{X}_s, \check{V} = \hat{V}$. To follow the principle **(P)** with the matrices (X_s, V) , we consider again the subspace $\hat{\mathcal{C}}'_n \subset \check{\mathcal{C}}_n \subset \hat{\mathcal{C}}_n$ where there exists a representative with X_s, V satisfying (6.34). The diagonal entries (x_i) of the X_s satisfy (6.35), and we have

$$(6.55) \quad \begin{aligned} \check{X}_s &= c_1 \dots c_s (1 - t^{-1}) \rho_j \frac{x_i^{m-s} x_j^s}{x_j^m - t^{-1} x_i^m} \prod_{l \neq i} \frac{x_l^m - t^{-1} x_i^m}{x_l^m - x_i^m}, \\ \check{V}_i &= - (1 - t^{-1}) \prod_{l \neq i} \frac{x_l^m - t^{-1} x_i^m}{x_l^m - x_i^m}. \end{aligned}$$

for $t := \prod_s c_s$. The $(\rho_i) \in (\mathbb{C}^\times)^n$ are free, and this choice is unique up to the action by $\mathbb{Z}_m \wr S_n$. The Poisson bracket between the variables (x, ρ) is given in Lemma 6.10. In particular, on a dense subspace of $\check{\mathcal{C}}_n$ we have Darboux coordinates (q_i, p_i) for $x_i = e^{\frac{1}{m} q_i}$ and $\rho_i = e^{p_i}$.

6.4.2. Space associated with \check{Q}_m^{op}

Denote the reduced space $\text{Rep}(\check{A}_m^{op}/(\Phi - \tilde{c}), \alpha) // \text{GL}_\alpha$ by $\check{\mathcal{C}}_n^{op}$. It can be given as the set of matrices

$$(6.56) \quad Z_s, \check{Z}_s \in \text{GL}_n(\mathbb{C}), \quad \check{W} \in \text{Mat}_{1 \times n}(\mathbb{C}), \quad W \in \text{Mat}_{n \times 1}(\mathbb{C}),$$

satisfying the m relations

$$(6.57) \quad Z_s^{-1} \check{Z}_{s-1}^{-1} Z_{s-1} \check{Z}_{s-1} (\text{Id}_n + \delta_{s0} W \check{W}) = c_s \text{Id}_n,$$

where all the factors appearing in (6.57) are invertible, and we identify the elements in the same orbit of the action

$$(6.58) \quad \begin{aligned} g \cdot (Z_s, \check{Z}_s, W, \check{W}) &= (g_s Z_s g_s^{-1}, g_s \check{Z}_s g_s^{-1}, g_0 W, \widehat{W} g_0^{-1}), \\ g &= (g_s) \in \mathrm{GL}_n(\mathbb{C})^m. \end{aligned}$$

To apply the principle **(P)** to the matrices (Z_s, W) representing the continuous arrows of \widetilde{Q}_m^{op} , we need the Z_s to be related diagonal matrices. In view of (6.57) with $s \neq 0$, Z_{s-1} and $c_s Z_s$ share the same spectrum. So we consider the subspace $(\check{C}_n^{op})' \subset \check{C}_n^{op}$ where there exists a representative with

$$(6.59) \quad Z_s = c_1^{-1} \dots c_s^{-1} \mathrm{diag}(z_1, \dots, z_n), \quad s \in I, \quad W = (1, \dots, 1)^T,$$

and we assume that for $t := \prod_{s \in I} c_s$,

$$(6.60) \quad z_i \neq 0, \quad z_i \neq z_j, \quad z_i \neq tz_j, \quad \text{for all } i \neq j.$$

LEMMA 6.14. — *The subspace $(\check{C}_n^{op})'$ can be parametrised by $(z_i) \in \mathbb{C}^n$ satisfying (6.60) and $(\sigma_i) \in (\mathbb{C}^\times)^n$ where the matrices (Z_s, W) are given by (6.59), together with*

$$\begin{aligned} \check{Z}_s &= \mathrm{Id}_n, \quad s \neq m-1, \quad (\check{Z}_{m-1})_{ij} = (1-t)\sigma_i \frac{z_i}{z_i - tz_j} \prod_{k \neq j} \frac{z_k - tz_j}{z_k - z_j}, \\ \check{W}_j &= -(1-t) \prod_{k \neq j} \frac{z_k - tz_j}{z_k - z_j}. \end{aligned}$$

Moreover, this choice is unique up to S_n action by simultaneous permutation of entries.

Proof. — Let us start with a representative such that (6.59) holds. By assumption on (z_i) and using (6.57) with $s \neq 0$, \check{Z}_s is diagonal, and we can fix the gauge (up to a finite action) so that $\check{Z}_s = \mathrm{Id}_n$. Using the case $s = 0$ in (6.57), we get that

$$(6.61) \quad (\check{Z}_{m-1})_{ij} = -\frac{\sigma_i \check{W}_j}{1 - tz_i^{-1} z_j},$$

where we have set $\sigma_i = (\check{Z}_{m-1} W)_i$. The equation that we have just used can also be written as

$$(6.62) \quad t \check{Z}_{m-1}^{-1} D^{-1} \check{Z}_{m-1} = D^{-1} + W (\check{W} D^{-1}), \quad D = \mathrm{diag}(z_1, \dots, z_n).$$

Using that the two sides of (6.62) share the same spectrum, we get as in the proof of Lemma 6.6 the following equality of characteristic polynomials in η

$$(6.63) \quad \prod_{k=1}^n (tz_k^{-1} - \eta) = \prod_{k=1}^n (z_k^{-1} - \eta) + \sum_l \check{W}_l z_l^{-1} \prod_{k \neq l} (z_k^{-1} - \eta).$$

Evaluating this identity at $\eta = z_j^{-1}$, we get the desired \check{W}_j .

Finally, we can check that $\sigma_i = (\check{Z}_{m-1} W)_i$ holds, and we note that $\sigma_i \neq 0$ because the matrix \check{Z}_{m-1} is invertible. The residual S_n action is clear. \square

LEMMA 6.15. — *The reduced Poisson bracket is such that $\{z_i, z_j\} = 0 = \{\sigma_i, \sigma_j\}$ and $\{z_i, \sigma_j\} = \delta_{ij} z_i \sigma_j$.*

Proof. — We first work in A_m^{op} as in Lemma 6.13. Due to the choice of ordering, we have

$$\begin{aligned} \{\{\widehat{y}_s, \widehat{y}_r\}\} &= \frac{1}{2} \delta_{s, r+1} \widehat{y}_r \widehat{y}_{r+1} \otimes e_{r+1} - \frac{1}{2} \delta_{s, r-1} e_r \otimes \widehat{y}_{r-1} \widehat{y}_r, \quad \{\{y_0, y_0\}\} = 0, \\ \{\{y_0, \widehat{y}_s\}\} &= \delta_{s, 0} \left[e_0 \otimes e_1 + \frac{1}{2} \widehat{y}_0 y_0 \otimes e_1 + \frac{1}{2} e_0 \otimes y_0 \widehat{y}_0 \right] - \frac{1}{2} \delta_{s, m-1} \widehat{y}_{m-1} \otimes y_0 + \frac{1}{2} \delta_{s, 1} y_0 \otimes \widehat{y}_1. \end{aligned}$$

By localisation, these relations hold in \check{A}_m^{op} , where we can determine the double brackets between the generators $z_0 = e_0 + \widehat{y}_0 y_0$ and $\check{z}_s = \widehat{y}_s$. Introducing $\check{z} := \check{z}_0 \dots \check{z}_{m-1}$, we can then get

$$\begin{aligned} \{\{\check{z}, \check{z}\}\} &= \frac{1}{2} \left(\check{z}^2 \otimes e_0 - e_0 \otimes \check{z}^2 \right), \quad \{\{z_0, z_0\}\} = \frac{1}{2} \left(e_0 \otimes z_0^2 - z_0^2 \otimes e_0 \right), \\ \{\{\check{z}, z_0\}\} &= -\frac{1}{2} \left(e_0 \otimes \check{z} z_0 + z_0 \check{z} \otimes e_0 + \check{z} \otimes z_0 - z_0 \otimes \check{z} \right). \end{aligned}$$

If we let $\check{Z}_{cyc} := \check{Z}_0 \dots \check{Z}_{m-1}$, we get from these double brackets the following identities on $\check{\mathcal{C}}_n^{op}$

$$(6.64) \quad \begin{aligned} \{\{\text{tr } Z_0^k, \text{tr } Z_0^l\}\} &= 0, \quad \{\{\text{tr } Z_0^k, \text{tr } \check{Z}_{cyc} Z_0^l\}\} = k \text{tr} \left(\check{Z}_{cyc} Z_0^{k+l} \right), \\ \{\{\text{tr } \check{Z}_{cyc} Z_0^k, \text{tr } \check{Z}_{cyc} Z_0^l\}\} &= \left[\sum_{r=1}^k - \sum_{r=1}^l \right] \text{tr} \left(\check{Z}_{cyc} Z_0^{k+l-r} \check{Z}_{cyc} Z_0^r \right), \end{aligned}$$

which can be derived as in Lemma 6.10. Notice that (6.41) and (6.64) are the same equations if we replace X_{cyc}, \check{X}_0 by Z_0, \check{Z}_{cyc} (in that order), or vice-versa. If we also remark that the functions $\text{tr } X_{cyc}^k, \text{tr } \check{X}_0 X_{cyc}^l$ written in coordinates in $\check{\mathcal{C}}_n$, and the functions $\text{tr } Z_0^k, \text{tr } \check{Z}_{cyc} Z_0^l$ written in coordinates in $(\check{\mathcal{C}}_n^{op})'$ are exactly the same upon replacing (x_j^m, ρ_j, t) by (z_j, σ_j, t^{-1}) , then the statement follows from Lemma 6.10. \square

As a corollary, we can get Darboux coordinates $(\check{q}_i, \check{p}_i)$ on a dense subset of $\check{\mathcal{C}}_n^{op}$ by considering $z_i = e^{\check{q}_i}$ and $\sigma_i = e^{\check{p}_i}$.

6.4.3. Duality

So far, we have obtained Darboux coordinates on dense subsets of $\check{\mathcal{C}}_n, \check{\mathcal{C}}_n^{op}$. Denoting $X_{cyc} = X_0 \dots X_{m-1}$ and $\check{Z}_{cyc} = \check{Z}_0 \dots \check{Z}_{m-1}$, we can use the Poisson diffeomorphism $\check{\Psi} : \check{\mathcal{C}}_n \rightarrow \check{\mathcal{C}}_n^{op}$ to obtain that

$$\text{tr } Z_0^{-k} \circ \check{\Psi} = \text{tr } \check{X}_0^k, \quad \text{tr } X_{cyc}^k \circ \check{\Psi}^{-1} = \text{tr} \left(\check{Z}_{cyc} \right)^k.$$

Hence, we conclude that we have action-angle duality between the integrable systems $(\text{tr } \check{X}_0^k)_{k=1}^n$ and $(\text{tr}(\check{Z}_{cyc})^k)_{k=1}^n$ defined on $\check{\mathcal{C}}_n$ and $\check{\mathcal{C}}_n^{op}$ respectively. In the coordinates (q_i, p_i) we can write the functions $(\text{tr } \check{X}_0^k)_{k=1}^n$ as hyperbolic (or trigonometric, see Remark 6.8) RS Hamiltonians of type A_{n-1} , which are exponential of order km in the momenta (p_i) . In the coordinates $(\check{q}_i, \check{p}_i)$, the functions $\text{tr}(\check{Z}_{cyc})^k$ are also Hamiltonians of the hyperbolic RS system of type A_{n-1} , but they are simply exponential of order k in the momenta. The dependence on m of this second family is hidden in the coupling $t = \prod_{s \in I} c_s$.

Remark 6.16. — Fix $m = 1$. In that case, the duality map hence obtained is close to the one considered in [CF17, Proposition 3.8], though the duality map used in that paper relates the spaces with inverse parameters t and t^{-1} . Note also that, upon identifying the generators (x, y) from Example 5.17 with $(x_0, x_0^{-1}\check{x}_0)$, the quasi-Hamiltonian isomorphisms (5.31) induce Poisson automorphisms on $\check{\mathcal{C}}_n$ which act as the identity on V and \check{V} while X_0, \check{X}_0 transform as

$$(6.65) \quad \begin{aligned} (X_0, \check{X}_0) &\mapsto (X_0, \check{X}_0 X_0^{-1} \check{X}_0) , \\ (X_0, \check{X}_0) &\mapsto (\check{X}_0^{-1} X_0, X_0) , \\ (X_0, \check{X}_0) &\mapsto (\alpha X_0, \alpha \beta \check{X}_0) , \end{aligned}$$

with $\alpha, \beta \in \mathbb{C}^\times$. As noted in Example 5.20, we get in particular that $\mathrm{SL}_2(\mathbb{Z})$ acts on $\check{\mathcal{C}}_n$ by Poisson automorphism. This $\mathrm{SL}_2(\mathbb{Z})$ action was related to the self-duality of the hyperbolic RS system in [GR01, § 2.2]. The analogue of this result in the real setting can be found in [FK12, FK13]. The original discovery of self-duality of the hyperbolic RS system in the real case is due to Ruijsenaars [Rui88].

6.5. Additional remarks and further directions

6.5.1.

The flows of the Hamiltonian vector fields associated with the two integrable systems presented in § 6.1 can be computed explicitly by adapting [CS17]. In particular, they are complete in \mathcal{C}_n and \mathcal{C}_n^{op} . Similarly, the flows associated with the different integrable systems found in the other subsections can also be computed explicitly (see [CF17] to get the flows of $(\mathrm{tr} \widehat{Y}_{cyc}^k)$ from § 6.3 and the systems in § 6.4), and they are complete in the phase spaces that support them.

6.5.2.

The different integrable systems defined in the previous subsections admit spin extensions [GH84, KZ95], which can be described as adding internal degrees of freedom to the different particles. The phase space of such spin extensions can be obtained by adding multiple framing arrows from the vertex ∞ to the different vertices in the cyclic quiver, see e.g. [CF20, CS17, Fai19b, Sil18]. Motivated by the work of Reshetikhin on the duality of spin systems in the CM-RS family [Res03, Res16], it would be interesting to understand if the duality map can be realised at the level of quivers for systems with spins.

Appendix A. The fusion double bracket

Let A be an algebra over $B = \bigoplus_{s \in I} \mathbb{k}e_s$. Fix $i, j \in I, i \neq j$, and consider the fusion algebra $A^f = A^f_{e_j \rightarrow e_i}$ obtained by fusing e_j onto e_i as in §2.2.1. The fusion algebra A^f always admits a particular double bracket, denoted $\{\{-, -\}_{fus}$, which we define now. For $\epsilon = 1 - e_j$, recall from Lemma 2.2 that the generators of A^f are of the form $\epsilon t \epsilon$ for $t \in \epsilon A \epsilon$, $e_{ij} u \epsilon$ for $u \in e_j A \epsilon$, $\epsilon v e_{ji}$ for $v \in \epsilon A e_j$, and $e_{ij} w e_{ji}$ for $w \in e_j A e_j$. The double bracket $\{\{-, -\}_{fus}$ is given on generators as follows:

$$\begin{aligned} \text{(A.1a)} \quad & \{\{\epsilon t \epsilon, \tilde{\epsilon} t \epsilon\}\}_{fus} = 0, \\ \text{(A.1b)} \quad & \{\{\epsilon t \epsilon, e_{ij} u \epsilon\}\}_{fus} = \frac{1}{2} (e_i \otimes t e_{ij} u - e_i t \otimes e_{ij} u), \\ \text{(A.1c)} \quad & \{\{\epsilon t \epsilon, \epsilon v e_{ji}\}\}_{fus} = \frac{1}{2} (v e_{ji} t \otimes e_i - v e_{ji} \otimes t e_i), \\ \text{(A.1d)} \quad & \{\{\epsilon t \epsilon, e_{ij} w e_{ji}\}\}_{fus} = \frac{1}{2} (e_{ij} w e_{ji} t \otimes e_i + e_i \otimes t e_{ij} w e_{ji} \\ & \quad - e_{ij} w e_{ji} \otimes t e_i - e_i t \otimes e_{ij} w e_{ji}), \end{aligned}$$

when the first component $\epsilon t \epsilon$ is a generator of the first type (2.13a);

$$\begin{aligned} \text{(A.2a)} \quad & \{\{e_{ij} u \epsilon, \epsilon t \epsilon\}\}_{fus} = \frac{1}{2} (e_{ij} u \otimes e_i t - t e_{ij} u \otimes e_i), \\ \text{(A.2b)} \quad & \{\{e_{ij} u \epsilon, e_{ij} \tilde{u} \epsilon\}\}_{fus} = \frac{1}{2} (e_i \otimes e_{ij} u e_{ij} \tilde{u} - e_{ij} \tilde{u} e_{ij} u \otimes e_i), \\ \text{(A.2c)} \quad & \{\{e_{ij} u \epsilon, \epsilon v e_{ji}\}\}_{fus} = \frac{1}{2} (e_{ij} u \otimes e_i v e_{ji} - v e_{ji} \otimes e_{ij} u e_i), \\ \text{(A.2d)} \quad & \{\{e_{ij} u \epsilon, e_{ij} w e_{ji}\}\}_{fus} = \frac{1}{2} (e_i \otimes e_{ij} u e_{ij} w e_{ji} - e_{ij} w e_{ji} \otimes e_{ij} u e_i), \end{aligned}$$

when the first component $e_{ij} u \epsilon$ is a generator of the second type (2.13b);

$$\begin{aligned} \text{(A.3a)} \quad & \{\{\epsilon v e_{ji}, \epsilon t \epsilon\}\}_{fus} = \frac{1}{2} (t e_i \otimes v e_{ji} - e_i \otimes v e_{ji} t), \\ \text{(A.3b)} \quad & \{\{\epsilon v e_{ji}, e_{ij} u \epsilon\}\}_{fus} = \frac{1}{2} (e_{ij} u e_i \otimes v e_{ji} - e_i v e_{ji} \otimes e_{ij} u), \\ \text{(A.3c)} \quad & \{\{\epsilon v e_{ji}, \epsilon \tilde{v} e_{ji}\}\}_{fus} = \frac{1}{2} (\tilde{v} e_{ji} v e_{ji} \otimes e_i - e_i \otimes v e_{ji} \tilde{v} e_{ji}), \\ \text{(A.3d)} \quad & \{\{\epsilon v e_{ji}, e_{ij} w e_{ji}\}\}_{fus} = \frac{1}{2} (e_{ij} w e_{ji} v e_{ji} \otimes e_i - e_i v e_{ji} \otimes e_{ij} w e_{ji}), \end{aligned}$$

when the first component $\epsilon v e_{ji}$ is a generator of the third type (2.13c);

$$(A.4a) \quad \{\{e_{ij}we_{ji}, \epsilon t \epsilon\}\}_{fus} = \frac{1}{2} \left(te_i \otimes e_{ij}we_{ji} + e_{ij}we_{ji} \otimes e_i t - te_{ij}we_{ji} \otimes e_i - e_i \otimes e_{ij}we_{ji} t \right),$$

$$(A.4b) \quad \{\{e_{ij}we_{ji}, e_{ij}u\epsilon\}\}_{fus} = \frac{1}{2} (e_{ij}ue_i \otimes e_{ij}we_{ji} - e_{ij}ue_{ij}we_{ji} \otimes e_i),$$

$$(A.4c) \quad \{\{e_{ij}we_{ji}, \epsilon ve_{ji}\}\}_{fus} = \frac{1}{2} (e_{ij}we_{ji} \otimes e_i ve_{ji} - e_i \otimes e_{ij}we_{ji} ve_{ji}),$$

$$(A.4d) \quad \{\{e_{ij}we_{ji}, e_{ij}\tilde{w}e_{ji}\}\}_{fus} = 0,$$

when the first component $e_{ij}we_{ji}$ is a generator of the fourth type (2.13d). The case corresponding to $i = 1, j = 2$ gives the double bracket considered in Proposition 4.5 on $A_{e_2 \rightarrow e_1}^f$. The double bracket $\{\{-, -\}\}_{fus}$ was first introduced by Van den Bergh in [VdB08a, Theorem 5.3.1], and the explicit form given above was computed in [Fai21, Lemma 2.19].

Appendix B. Identities for the proof of Lemma 4.7

In this section, we collect the values of all the elements in $A_2 \otimes A_2$ given by

$$\{\{c, d\}\}_{ind}^{2 \rightarrow 1}, \quad \{\{c, d\}\}_{fus}^{2 \rightarrow 1}, \quad \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2}, \quad \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2},$$

for all the needed specialisations of the generators $c, d \in A_1$. It is then easy to check that (4.14) is satisfied in all such cases, i.e.

$$\psi^{\otimes 2} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} + \psi^{\otimes 2} \{\{c, d\}\}_{fus}^{2 \rightarrow 1} = \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} + \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2}.$$

As a result, the proof of Lemma 4.7 is completed.

Remark B.1. — In each case, we write $\{\{c, d\}\}_{ind}^{2 \rightarrow 1}$ as an element of $eAf \otimes e'Af'$ for some $e, e' \in \{\hat{e}, e_1, e_{12}\}$ and $f, f' \in \{\hat{e}, e_1, e_{21}\}$. In particular, this completely characterises how $\psi^{\otimes 2}$ acts on $\{\{c, d\}\}_{ind}^{2 \rightarrow 1}$.

B.1

We consider the first specialisation of a generator of first type, that is $c = a$ for $a \in \hat{e}A\hat{e}$.

Case 1.1: $d = b, b \in \hat{e}A\hat{e}$.

$$(B.1) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \hat{e} \{\{a, b\}\}' \hat{e} \otimes \hat{e} \{\{a, b\}\}'' \hat{e}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.2: $d = b, b \in e_1A\hat{e}$.

$$(B.2) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' \hat{e} \otimes \hat{e} \{\{a, b\}\}'' \hat{e}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.3: $d = b, b \in \widehat{e}Ae_1$.

$$(B.3) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \widehat{e} \{\{a, b\}\}' \widehat{e} \otimes \widehat{e} \{\{a, b\}\}'' e_1, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\} e_{12} \Phi_2^{-1}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.4: $d = b, b \in e_1Ae_1$.

$$(B.4) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' \widehat{e} \otimes \widehat{e} \{\{a, b\}\}'' e_1, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} \{\{a, b\}\} e_{12} \Phi_2^{-1}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.5: $d = e_{12}b, b \in e_2A\widehat{e}$.

$$(B.5) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \widehat{e} \otimes \widehat{e} \{\{a, b\}\}'' \widehat{e}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.6: $d = e_{12}b, b \in e_2Ae_1$.

$$(B.6) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \widehat{e} \otimes \widehat{e} \{\{a, b\}\}'' e_1, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\} e_{12} \Phi_2^{-1}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.7: $d = be_{21}, b \in \widehat{e}Ae_2$.

$$(B.7) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \widehat{e} \{\{a, b\}\}' \widehat{e} \otimes \widehat{e} \{\{a, b\}\}'' e_{21}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.8: $d = be_{21}, b \in e_1Ae_2$.

$$(B.8) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' \widehat{e} \otimes \widehat{e} \{\{a, b\}\}'' e_{21}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 1.9: $d = e_{12}be_{21}, b \in e_2Ae_2$.

$$(B.9) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \widehat{e} \otimes \widehat{e} \{\{a, b\}\}'' e_{21}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

B.2

We consider the second specialisation of a generator of first type, that is $c = a$ for $a \in e_1A\widehat{e}$.

Case 2.1: $d = b, b \in e_1A\widehat{e}$. This was done in the proof of Lemma 4.7.

Case 2.2: $d = b, b \in \widehat{e}Ae_1$.

$$(B.10) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \widehat{e} \{\{a, b\}\}' \widehat{e} \otimes e_1 \{\{a, b\}\}'' e_1, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} e_{12} \Phi_2^{-1} \\ &\quad - \frac{1}{2} (be_{12} \Phi_2 e_{21} a \otimes \Phi_2^{-1} - be_{12} \Phi_2^{-1} e_{21} a \otimes \Phi_2), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= \frac{1}{2} (be_{12} \Phi_2 e_{21} a \otimes \Phi_2^{-1} - be_{12} \Phi_2^{-1} e_{21} a \otimes \Phi_2). \end{aligned}$$

Case 2.3: $d = b$, $b \in e_1 A e_1$.

$$(B.11) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' \hat{e} \otimes e_1 \{\{a, b\}\}'' e_1, \quad \{\{c, d\}\}_{fus}^{2 \rightarrow 1} = 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \Phi_2 e_{21} \{\{a, b\}\} e_{12} \Phi_2^{-1} \\ &\quad - \frac{1}{2} \left(\Phi_2 e_{21} b e_{12} \Phi_2 e_{21} a \otimes \Phi_2^{-1} - \Phi_2^2 e_{21} a \otimes e_{21} b e_{12} \Phi_2^{-1} \right) \\ &\quad - \frac{1}{2} \left(e_{21} a \otimes \Phi_2^2 e_{21} b e_{12} \Phi_2^{-1} - \Phi_2 e_{21} b e_{12} \Phi_2^{-1} e_{21} a \otimes \Phi_2 \right), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} \left(\Phi_2 e_{21} b e_{12} \Phi_2 e_{21} a \otimes \Phi_2^{-1} - \Phi_2^2 e_{21} a \otimes e_{21} b e_{12} \Phi_2^{-1} \right) \\ &\quad + \frac{1}{2} \left(e_{21} a \otimes \Phi_2^2 e_{21} b e_{12} \Phi_2^{-1} - \Phi_2 e_{21} b e_{12} \Phi_2^{-1} e_{21} a \otimes \Phi_2 \right). \end{aligned}$$

Case 2.4: $d = e_{12} b$, $b \in e_2 A \hat{e}$.

$$(B.12) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \hat{e} \otimes e_1 \{\{a, b\}\}'' \hat{e}, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= -\frac{1}{2} a \otimes e_{12} b, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} - \frac{1}{2} (e_{21} a \otimes \Phi_2 b + \Phi_2 e_{21} a \otimes b), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= \frac{1}{2} e_{21} a \otimes \Phi_2 b. \end{aligned}$$

Case 2.5: $d = e_{12} b$, $b \in e_2 A e_1$.

$$(B.13) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \hat{e} \otimes e_1 \{\{a, b\}\}'' e_1, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= -\frac{1}{2} a \otimes e_{12} b, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} e_{12} \Phi_2^{-1} - \frac{1}{2} \Phi_2 e_{21} a \otimes b e_{12} \Phi_2^{-1} \\ &\quad - \frac{1}{2} \left(b e_{12} \Phi_2 e_{21} a \otimes \Phi_2^{-1} - b e_{12} \Phi_2^{-1} e_{21} a \otimes \Phi_2 + e_{21} a \otimes \Phi_2 b e_{12} \Phi_2^{-1} \right), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= \\ &\quad + \frac{1}{2} \left(b e_{12} \Phi_2 e_{21} a \otimes \Phi_2^{-1} - b e_{12} \Phi_2^{-1} e_{21} a \otimes \Phi_2 + e_{21} a \otimes \Phi_2 b e_{12} \Phi_2^{-1} \right). \end{aligned}$$

Case 2.6: $d = b e_{21}$, $b \in \hat{e} A e_2$.⁽⁵⁾

$$(B.14) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \hat{e} \{\{a, b\}\}' \hat{e} \otimes e_1 \{\{a, b\}\}'' e_{21}, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} b e_{21} a \otimes e_1, \end{aligned}$$

⁽⁵⁾ This is the first case where one needs to be careful when computing $\psi^{\otimes 2} \{\{c, d\}\}_{fus}^{2 \rightarrow 1}$ in order to verify (4.14). It is given in this case by $\frac{1}{2} \psi(b e_{21} a) \otimes \psi(e_1)$, and we have by definition of ψ that $\psi(e_1) = \Phi_2 e_{21} e_{12} \Phi_2^{-1} = e_2$. For the first factor, we need to remark that $b e_{21} a$ is *not* a specialisation of generators of A_1 as defined earlier, but it is a product of two. Hence, $\psi(b e_{21} a) = \psi(b e_{21}) \psi(a) = b \Phi_2 e_{21} a$.

$$\begin{aligned} \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} + \frac{1}{2} (be_{21}a \otimes \Phi_2 + b\Phi_2 e_{21}a \otimes e_2) , \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= -\frac{1}{2} be_{21}a \otimes \Phi_2 . \end{aligned}$$

Case 2.7: $d = be_{21}$, $b \in e_1 A e_2$.

$$\begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' \hat{e} \otimes e_1 \{\{a, b\}\}'' e_{21}, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} be_{21}a \otimes e_1, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \Phi_2 e_{21} \{\{a, b\}\} + \frac{1}{2} \Phi_2 e_{21} b \Phi_2 e_{21} a \otimes e_2 \\ &\quad - \frac{1}{2} (e_{21}a \otimes \Phi_2^2 e_{21}b - \Phi_2^2 e_{21}a \otimes e_{21}b - \Phi_2 e_{21} b e_{21} a \otimes \Phi_2) , \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= +\frac{1}{2} (e_{21}a \otimes \Phi_2^2 e_{21}b - \Phi_2^2 e_{21}a \otimes e_{21}b - \Phi_2 e_{21} b e_{21} a \otimes \Phi_2) . \end{aligned}$$

Case 2.8: $d = e_{12} b e_{21}$, $b \in e_2 A e_2$.

$$\begin{aligned} \text{(B.15)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \hat{e} \otimes e_1 \{\{a, b\}\}'' e_{21}, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} (e_{12} b e_{21} a \otimes e_1 - a \otimes e_{12} b e_{21}) , \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} + \frac{1}{2} (b \Phi_2 e_{21} a \otimes e_2 - \Phi_2 e_{21} a \otimes b) \\ &\quad - \frac{1}{2} (e_{21}a \otimes \Phi_2 b - b e_{21} a \otimes \Phi_2) , \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= +\frac{1}{2} (e_{21}a \otimes \Phi_2 b - b e_{21} a \otimes \Phi_2) . \end{aligned}$$

B.3

We consider the third specialisation of a generator of first type, that is $c = a$ for $a \in \hat{e} A e_1$.

Case 3.1: $d = b$, $b \in \hat{e} A e_1$.

$$\begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \hat{e} \{\{a, b\}\}' e_1 \otimes \hat{e} \{\{a, b\}\}'' e_1, \quad \{\{c, d\}\}_{fus}^{2 \rightarrow 1} = 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\} e_{12} \Phi_2^{-1} * e_{12} \Phi_2^{-1} \\ \text{(B.16)} \quad &\quad - \frac{1}{2} (b e_{12} \Phi_2^{-2} \otimes a e_{12} - b e_{12} \otimes a e_{12} \Phi_2^{-2}) , \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= +\frac{1}{2} (b e_{12} \Phi_2^{-2} \otimes a e_{12} - b e_{12} \otimes a e_{12} \Phi_2^{-2}) . \end{aligned}$$

Case 3.2: $d = b$, $b \in e_1 A e_1$.

$$\begin{aligned} \text{(B.17)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' e_1 \otimes \hat{e} \{\{a, b\}\}'' e_1, \quad \{\{c, d\}\}_{fus}^{2 \rightarrow 1} = 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} \{\{a, b\}\} e_{12} \Phi_2^{-1} * e_{12} \Phi_2^{-1} \end{aligned}$$

$$\begin{aligned}
& -\frac{1}{2} \left(\Phi_2 \otimes ae_{12}\Phi_2^{-1}e_{21}be_{12}\Phi_2^{-1} - \Phi_2^{-1} \otimes ae_{12}\Phi_2e_{21}be_{12}\Phi_2^{-1} \right) \\
& -\frac{1}{2} \left(\Phi_2e_{21}be_{12}\Phi_2^{-2} \otimes ae_{12} - \Phi_2e_{21}be_{12} \otimes ae_{12}\Phi_2^{-2} \right), \\
\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2} & = +\frac{1}{2} \left(\Phi_2 \otimes ae_{12}\Phi_2^{-1}e_{21}be_{12}\Phi_2^{-1} - \Phi_2^{-1} \otimes ae_{12}\Phi_2e_{21}be_{12}\Phi_2^{-1} \right) \\
& +\frac{1}{2} \left(\Phi_2e_{21}be_{12}\Phi_2^{-2} \otimes ae_{12} - \Phi_2e_{21}be_{12} \otimes ae_{12}\Phi_2^{-2} \right).
\end{aligned}$$

Case 3.3: $d = e_{12}b$, $b \in e_2A\hat{e}$.

$$\begin{aligned}
\text{(B.18)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} & = e_{12} \{\{a, b\}\}' e_1 \otimes \hat{e} \{\{a, b\}\}'' \hat{e}, \\
\{\{c, d\}\}_{fus}^{2 \rightarrow 1} & = \frac{1}{2} e_1 \otimes ae_{12}b, \\
\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2} & = \{\{a, b\}\} * e_{12}\Phi_2^{-1} + \frac{1}{2} \left(e_2 \otimes ae_{12}\Phi_2^{-1}b + \Phi_2^{-1} \otimes ae_{12}b \right), \\
\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2} & = -\frac{1}{2} \Phi_2^{-1} \otimes ae_{12}b.
\end{aligned}$$

Case 3.4: $d = e_{12}b$, $b \in e_2Ae_1$.

$$\begin{aligned}
\text{(B.19)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} & = e_{12} \{\{a, b\}\}' e_1 \otimes \hat{e} \{\{a, b\}\}'' e_1, \\
\{\{c, d\}\}_{fus}^{2 \rightarrow 1} & = \frac{1}{2} e_1 \otimes ae_{12}b, \\
\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2} & = \{\{a, b\}\} e_{12}\Phi_2^{-1} * e_{12}\Phi_2^{-1} + \frac{1}{2} e_2 \otimes ae_{12}\Phi_2^{-1}be_{12}\Phi_2^{-1} \\
& -\frac{1}{2} \left(be_{12}\Phi_2^{-2} \otimes ae_{12} - be_{12} \otimes ae_{12}\Phi_2^{-2} - \Phi_2^{-1} \otimes ae_{12}be_{12}\Phi_2^{-1} \right), \\
\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2} & = \\
& +\frac{1}{2} \left(be_{12}\Phi_2^{-2} \otimes ae_{12} - be_{12} \otimes ae_{12}\Phi_2^{-2} - \Phi_2^{-1} \otimes ae_{12}be_{12}\Phi_2^{-1} \right).
\end{aligned}$$

Case 3.5: $d = be_{21}$, $b \in \hat{e}Ae_2$.

$$\begin{aligned}
\text{(B.20)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} & = \hat{e} \{\{a, b\}\}' e_1 \otimes \hat{e} \{\{a, b\}\}'' e_{21}, \quad \{\{c, d\}\}_{fus}^{2 \rightarrow 1} = -\frac{1}{2} be_{21} \otimes a, \\
\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2} & = \{\{a, b\}\} * e_{12}\Phi_2^{-1} - \frac{1}{2} \left(b\Phi_2^{-1} \otimes ae_{12} + b \otimes ae_{12}\Phi_2^{-1} \right), \\
\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2} & = +\frac{1}{2} b\Phi_2^{-1} \otimes ae_{12}.
\end{aligned}$$

Case 3.6: $d = be_{21}$, $b \in e_1Ae_2$.

$$\begin{aligned}
\text{(B.21)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} & = e_1 \{\{a, b\}\}' e_1 \otimes \hat{e} \{\{a, b\}\}'' e_{21}, \\
\{\{c, d\}\}_{fus}^{2 \rightarrow 1} & = -\frac{1}{2} be_{21} \otimes a,
\end{aligned}$$

$$\begin{aligned} \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} \{\{a, b\}\} * e_{12} \Phi_2^{-1} - \frac{1}{2} \Phi_2 e_{21} b \otimes a e_{12} \Phi_2^{-1} \\ &\quad - \frac{1}{2} \left(\Phi_2 e_{21} b \Phi_2^{-1} \otimes a e_{12} + \Phi_2 \otimes a e_{12} \Phi_2^{-1} e_{21} b - \Phi_2^{-1} \otimes a e_{12} \Phi_2 e_{21} b \right), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= \\ &\quad + \frac{1}{2} \left(\Phi_2 e_{21} b \Phi_2^{-1} \otimes a e_{12} + \Phi_2 \otimes a e_{12} \Phi_2^{-1} e_{21} b - \Phi_2^{-1} \otimes a e_{12} \Phi_2 e_{21} b \right). \end{aligned}$$

Case 3.7: $d = e_{12} b e_{21}$, $b \in e_2 A e_2$.

$$\begin{aligned} \text{(B.22)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' e_1 \otimes \hat{e} \{\{a, b\}\}'' e_{21}, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes a e_{12} b e_{21} - e_{12} b e_{21} \otimes a), \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\} * e_{12} \Phi_2^{-1} \\ &\quad + \frac{1}{2} \left(e_2 \otimes a e_{12} \Phi_2^{-1} b - b \otimes a e_{12} \Phi_2^{-1} \right) - \frac{1}{2} \left(b \Phi_2^{-1} \otimes a e_{12} - \Phi_2^{-1} \otimes a e_{12} b \right), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} \left(b \Phi_2^{-1} \otimes a e_{12} - \Phi_2^{-1} \otimes a e_{12} b \right). \end{aligned}$$

B.4

We consider the fourth specialisation of a generator of first type, that is $c = a$ for $a \in e_1 A e_1$.

Case 4.1: $d = b$, $b \in e_1 A e_1$.

$$\begin{aligned} \text{(B.23)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' e_1 \otimes e_1 \{\{a, b\}\}'' e_1, \quad \{\{c, d\}\}_{fus}^{2 \rightarrow 1} = 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \Phi_2 e_{21} \{\{a, b\}\} e_{12} \Phi_2^{-1} * e_{12} \Phi_2^{-1} \\ &\quad - \frac{1}{2} \left(\Phi_2^2 e_{21} a e_{12} \Phi_2^{-1} \otimes e_{21} b e_{12} \Phi_2^{-1} - e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^2 e_{21} b e_{12} \Phi_2^{-1} \right) \\ &\quad - \frac{1}{2} \left(\Phi_2 e_{21} b e_{12} \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-2} - \Phi_2 e_{21} b e_{12} \Phi_2^{-2} \otimes \Phi_2 e_{21} a e_{12} \right) \\ &\quad - \frac{1}{2} \left(\Phi_2 e_{21} b e_{12} \Phi_2^{-1} e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 - \Phi_2 e_{21} b e_{12} \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^{-1} \right) \\ &\quad - \frac{1}{2} \left(\Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} \Phi_2 e_{21} b e_{12} \Phi_2^{-1} - \Phi_2 \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} e_{21} b e_{12} \Phi_2^{-1} \right), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= \\ &\quad + \frac{1}{2} \left(\Phi_2^2 e_{21} a e_{12} \Phi_2^{-1} \otimes e_{21} b e_{12} \Phi_2^{-1} - e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^2 e_{21} b e_{12} \Phi_2^{-1} \right) \\ &\quad + \frac{1}{2} \left(\Phi_2 e_{21} b e_{12} \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-2} - \Phi_2 e_{21} b e_{12} \Phi_2^{-2} \otimes \Phi_2 e_{21} a e_{12} \right) \\ &\quad + \frac{1}{2} \left(\Phi_2 e_{21} b e_{12} \Phi_2^{-1} e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 - \Phi_2 e_{21} b e_{12} \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^{-1} \right) \\ &\quad + \frac{1}{2} \left(\Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} \Phi_2 e_{21} b e_{12} \Phi_2^{-1} - \Phi_2 \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} e_{21} b e_{12} \Phi_2^{-1} \right). \end{aligned}$$

Case 4.2: $d = e_{12}b$, $b \in e_2A\hat{e}$.

$$\begin{aligned}
(B.24) \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' e_1 \otimes e_1 \{\{a, b\}\}'' \hat{e}, \\
\{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes ae_{12}b - a \otimes e_{12}b), \\
\{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} * e_{12} \Phi_2^{-1} \\
&\quad + \frac{1}{2} \left(e_2 \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} b - \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes b \right) \\
&\quad - \frac{1}{2} \left(e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 b - \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} b \right), \\
\{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} \left(e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 b - \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} b \right).
\end{aligned}$$

Case 4.3: $d = e_{12}b$, $b \in e_2Ae_1$.

$$\begin{aligned}
(B.25) \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' e_1 \otimes e_1 \{\{a, b\}\}'' e_1, \\
\{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes ae_{12}b - a \otimes e_{12}b), \\
\{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} e_{12} \Phi_2^{-1} * e_{12} \Phi_2^{-1} \\
&\quad + \frac{1}{2} \left(e_2 \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} b e_{12} \Phi_2^{-1} - \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes b e_{12} \Phi_2^{-1} \right) \\
&\quad - \frac{1}{2} \left(e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 b e_{12} \Phi_2^{-1} - \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} b e_{12} \Phi_2^{-1} \right) \\
&\quad - \frac{1}{2} \left(b e_{12} \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^{-1} - b e_{12} \Phi_2^{-1} e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 \right) \\
&\quad - \frac{1}{2} \left(b e_{12} \Phi_2^{-2} \otimes \Phi_2 e_{21} a e_{12} - b e_{12} \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-2} \right), \\
\{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} \left(e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 b e_{12} \Phi_2^{-1} - \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} b e_{12} \Phi_2^{-1} \right) \\
&\quad + \frac{1}{2} \left(b e_{12} \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^{-1} - b e_{12} \Phi_2^{-1} e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 \right) \\
&\quad + \frac{1}{2} \left(b e_{12} \Phi_2^{-2} \otimes \Phi_2 e_{21} a e_{12} - b e_{12} \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-2} \right).
\end{aligned}$$

Case 4.4: $d = be_{21}$, $b \in \hat{e}Ae_2$.

$$\begin{aligned}
(B.26) \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \hat{e} \{\{a, b\}\}' e_1 \otimes e_1 \{\{a, b\}\}'' e_{21}, \\
\{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} (be_{21}a \otimes e_1 - be_{21} \otimes ae_1), \\
\{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} * e_{12} \Phi_2^{-1} \\
&\quad + \frac{1}{2} \left(b \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes e_2 - b \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \right) \\
&\quad - \frac{1}{2} \left(b \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} - b e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 \right), \\
\{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} \left(b \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} - b e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 \right).
\end{aligned}$$

Case 4.5: $d = be_{21}$, $b \in e_1 A e_2$.

$$\begin{aligned}
 \text{(B.27)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}'\} e_1 \otimes e_1 \{\{a, b\}''\} e_{21}, \\
 \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} (be_{21}a \otimes e_1 - be_{21} \otimes ae_1), \\
 \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \Phi_2 e_{21} \{\{a, b\}\} * e_{12} \Phi_2^{-1} \\
 &\quad + \frac{1}{2} (\Phi_2 e_{21} b \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes e_2 - \Phi_2 e_{21} b \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1}) \\
 &\quad - \frac{1}{2} (e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^2 e_{21} b - \Phi_2^2 e_{21} a e_{12} \Phi_2^{-1} \otimes e_{21} b) \\
 &\quad - \frac{1}{2} (\Phi_2 \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} e_{21} b - \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} \Phi_2 e_{21} b) \\
 &\quad - \frac{1}{2} (\Phi_2 e_{21} b \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} - \Phi_2 e_{21} b e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2), \\
 \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} (e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2^2 e_{21} b - \Phi_2^2 e_{21} a e_{12} \Phi_2^{-1} \otimes e_{21} b) \\
 &\quad + \frac{1}{2} (\Phi_2 \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} e_{21} b - \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} \Phi_2 e_{21} b) \\
 &\quad + \frac{1}{2} (\Phi_2 e_{21} b \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} - \Phi_2 e_{21} b e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2).
 \end{aligned}$$

Case 4.6: $d = e_{12} b e_{21}$, $b \in e_2 A e_2$.

$$\begin{aligned}
 \text{(B.28)} \quad \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}'\} e_1 \otimes e_1 \{\{a, b\}''\} e_{21}, \\
 \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} (e_{12} b e_{21} a \otimes e_1 + e_1 \otimes a e_{12} b e_{21} - e_{12} b e_{21} \otimes a - a \otimes e_{12} b e_{21}), \\
 \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} * e_{12} \Phi_2^{-1} \\
 &\quad + \frac{1}{2} (b \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes e_2 + e_2 \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} b) \\
 &\quad - \frac{1}{2} (b \otimes \Phi_2 e_{21} a e_{12} \Phi_2^{-1} + \Phi_2 e_{21} a e_{12} \Phi_2^{-1} \otimes b) \\
 &\quad - \frac{1}{2} (b \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} + e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 b) \\
 &\quad + \frac{1}{2} (b e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 + \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} b), \\
 \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} (b \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} + e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 b) \\
 &\quad - \frac{1}{2} (b e_{21} a e_{12} \Phi_2^{-1} \otimes \Phi_2 + \Phi_2^{-1} \otimes \Phi_2 e_{21} a e_{12} b).
 \end{aligned}$$

B.5

We consider the first specialisation of a generator of second type, that is $c = e_{12} a$ for $a \in e_2 A \hat{e}$.

Case 5.1: $d = e_{12}b$, $b \in e_2A\hat{e}$.

$$(B.29) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \hat{e} \otimes e_{12} \{\{a, b\}\}'' \hat{e}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 5.2: $d = e_{12}b$, $b \in e_2Ae_1$.

$$(B.30) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \hat{e} \otimes e_{12} \{\{a, b\}\}'' e_1, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= -\frac{1}{2} e_{12} b e_{12} a \otimes e_1, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\} e_{12} \Phi_2^{-1} - \frac{1}{2} (b e_{12} a \otimes \Phi_2^{-1} + b e_{12} \Phi_2^{-1} a \otimes e_2), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= +\frac{1}{2} b e_{12} a \otimes \Phi_2^{-1}. \end{aligned}$$

Case 5.3: $d = b e_{21}$, $b \in \hat{e}Ae_2$.

$$(B.31) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= \hat{e} \{\{a, b\}\}' \hat{e} \otimes e_{12} \{\{a, b\}\}'' e_{21}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

Case 5.4: $d = b e_{21}$, $b \in e_1Ae_2$.

$$(B.32) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' \hat{e} \otimes e_{12} \{\{a, b\}\}'' e_{21}, \\ \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} e_{12} a \otimes b e_{21}, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} \{\{a, b\}\} + \frac{1}{2} (\Phi_2 a \otimes e_{21} b + a \otimes \Phi_2 e_{21} b), \\ \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= -\frac{1}{2} \Phi_2 a \otimes e_{21} b. \end{aligned}$$

Case 5.5: $d = e_{12}b e_{21}$, $b \in e_2Ae_2$.

$$(B.33) \quad \begin{aligned} \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' \hat{e} \otimes e_{12} \{\{a, b\}\}'' e_{21}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0. \end{aligned}$$

B.6

We consider the second specialisation of a generator of second type, that is $c = e_{12}a$ for $a \in e_2Ae_1$.

Case 6.1: $d = e_{12}b$, $b \in e_2Ae_1$.

$$\begin{aligned}
 \{\{c, d\}_{ind}^{2 \rightarrow 1}\} &= e_{12} \{\{a, b\}'\} e_1 \otimes e_{12} \{\{a, b\}''\} e_1, \\
 \{\{c, d\}_{fus}^{2 \rightarrow 1}\} &= \frac{1}{2} (e_1 \otimes e_{12}ae_{12}b - e_{12}be_{12}a \otimes e_1), \\
 \{\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2}\} &= \{\{a, b\}\} e_{12}\Phi_2^{-1} * e_{12}\Phi_2^{-1} \\
 &\quad + \frac{1}{2} (e_2 \otimes ae_{12}\Phi_2^{-1}be_{12}\Phi_2^{-1} - be_{12}\Phi_2^{-1}ae_{12}\Phi_2^{-1} \otimes e_2) \\
 &\quad - \frac{1}{2} (be_{12}ae_{12}\Phi_2^{-1} \otimes \Phi_2^{-1} - \Phi_2^{-1} \otimes ae_{12}be_{12}\Phi_2^{-1}) \\
 &\quad - \frac{1}{2} (be_{12}\Phi_2^{-2} \otimes ae_{12} - be_{12} \otimes ae_{12}\Phi_2^{-2}), \\
 \{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2}\} &= + \frac{1}{2} (be_{12}ae_{12}\Phi_2^{-1} \otimes \Phi_2^{-1} - \Phi_2^{-1} \otimes ae_{12}be_{12}\Phi_2^{-1}) \\
 &\quad + \frac{1}{2} (be_{12}\Phi_2^{-2} \otimes ae_{12} - be_{12} \otimes ae_{12}\Phi_2^{-2}).
 \end{aligned}
 \tag{B.34}$$

Case 6.2: $d = be_{21}$, $b \in \hat{e}Ae_2$.

$$\begin{aligned}
 \{\{c, d\}_{ind}^{2 \rightarrow 1}\} &= \hat{e} \{\{a, b\}'\} e_1 \otimes e_{12} \{\{a, b\}''\} e_{21}, \\
 \{\{c, d\}_{fus}^{2 \rightarrow 1}\} &= -\frac{1}{2} be_{21} \otimes e_{12}a, \\
 \{\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2}\} &= \{\{a, b\}\} * e_{12}\Phi_2^{-1} \\
 &\quad - \frac{1}{2} (b \otimes ae_{12}\Phi_2^{-1} + b\Phi_2^{-1} \otimes ae_{12}), \\
 \{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2}\} &= +\frac{1}{2} b\Phi_2^{-1} \otimes ae_{12}.
 \end{aligned}
 \tag{B.35}$$

Case 6.3: $d = be_{21}$, $b \in e_1Ae_2$.

$$\begin{aligned}
 \{\{c, d\}_{ind}^{2 \rightarrow 1}\} &= e_1 \{\{a, b\}'\} e_1 \otimes e_{12} \{\{a, b\}''\} e_{21}, \\
 \{\{c, d\}_{fus}^{2 \rightarrow 1}\} &= \frac{1}{2} (e_{12}a \otimes be_{21} - be_{21} \otimes e_{12}a), \\
 \{\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2}\} &= \Phi_2 e_{21} \{\{a, b\}\} * e_{12}\Phi_2^{-1} \\
 &\quad + \frac{1}{2} (ae_{12}\Phi_2^{-1} \otimes \Phi_2 e_{21}b - \Phi_2 e_{21}b \otimes ae_{12}\Phi_2^{-1}) \\
 &\quad - \frac{1}{2} (\Phi_2 \otimes ae_{12}\Phi_2^{-1}e_{21}b - \Phi_2^{-1} \otimes ae_{12}\Phi_2 e_{21}b) \\
 &\quad - \frac{1}{2} (\Phi_2 e_{21}b\Phi_2^{-1} \otimes ae_{12} - \Phi_2 ae_{12}\Phi_2^{-1} \otimes e_{21}b), \\
 \{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2}\} &= +\frac{1}{2} (\Phi_2 \otimes ae_{12}\Phi_2^{-1}e_{21}b - \Phi_2^{-1} \otimes ae_{12}\Phi_2 e_{21}b) \\
 &\quad + \frac{1}{2} (\Phi_2 e_{21}b\Phi_2^{-1} \otimes ae_{12} - \Phi_2 ae_{12}\Phi_2^{-1} \otimes e_{21}b).
 \end{aligned}
 \tag{B.36}$$

Case 6.4: $d = e_{12}be_{21}$, $b \in e_2Ae_2$.

$$\begin{aligned}
\{\{c, d\}_{ind}^{2 \rightarrow 1}\} &= e_{12} \{\{a, b\}'\} e_1 \otimes e_{12} \{\{a, b\}''\} e_{21}, \\
\{\{c, d\}_{fus}^{2 \rightarrow 1}\} &= \frac{1}{2} (e_1 \otimes e_{12}ae_{12}be_{21} - e_{12}be_{21} \otimes e_{12}a), \\
\{\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2}\} &= \{\{a, b\}\} * e_{12}\Phi_2^{-1} \\
&+ \frac{1}{2} (e_2 \otimes ae_{12}\Phi_2^{-1}b - b \otimes ae_{12}\Phi_2^{-1}) \\
&- \frac{1}{2} (b\Phi_2^{-1} \otimes ae_{12} - \Phi_2^{-1} \otimes ae_{12}b), \\
\{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2}\} &= +\frac{1}{2} (b\Phi_2^{-1} \otimes ae_{12} - \Phi_2^{-1} \otimes ae_{12}b).
\end{aligned}
\tag{B.37}$$

B.7

We consider the first specialisation of a generator of third type, that is $c = ae_{21}$ for $a \in \hat{e}Ae_2$.

Case 7.1: $d = be_{21}$, $b \in \hat{e}Ae_2$.

$$\begin{aligned}
\{\{c, d\}_{ind}^{2 \rightarrow 1}\} &= \hat{e} \{\{a, b\}'\} e_{21} \otimes \hat{e} \{\{a, b\}''\} e_{21}, \quad \{\{c, d\}_{fus}^{2 \rightarrow 1}\} = 0, \\
\{\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2}\} &= \{\{a, b\}\}, \quad \{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2}\} = 0.
\end{aligned}
\tag{B.38}$$

Case 7.2: $d = be_{21}$, $b \in e_1Ae_2$.

$$\begin{aligned}
\{\{c, d\}_{ind}^{2 \rightarrow 1}\} &= e_1 \{\{a, b\}'\} e_{21} \otimes \hat{e} \{\{a, b\}''\} e_{21}, \\
\{\{c, d\}_{fus}^{2 \rightarrow 1}\} &= -\frac{1}{2} e_1 \otimes ae_{21}be_{21}, \\
\{\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2}\} &= \Phi_2 e_{21} \{\{a, b\}\} - \frac{1}{2} (e_2 \otimes a\Phi_2 e_{21}b + \Phi_2 \otimes ae_{21}b), \\
\{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2}\} &= +\frac{1}{2} \Phi_2 \otimes ae_{21}b.
\end{aligned}
\tag{B.39}$$

Case 7.3: $d = e_{12}be_{21}$, $b \in e_2Ae_2$.

$$\begin{aligned}
\{\{c, d\}_{ind}^{2 \rightarrow 1}\} &= e_{12} \{\{a, b\}'\} e_{21} \otimes \hat{e} \{\{a, b\}''\} e_{21}, \quad \{\{c, d\}_{fus}^{2 \rightarrow 1}\} = 0, \\
\{\{\psi(c), \psi(d)\}_{ind}^{1 \rightarrow 2}\} &= \{\{a, b\}\}, \quad \{\{\psi(c), \psi(d)\}_{fus}^{1 \rightarrow 2}\} = 0.
\end{aligned}
\tag{B.40}$$

B.8

We consider the second specialisation of a generator of third type, that is $c = ae_{21}$ for $a \in e_1Ae_2$.

Case 8.1: $d = be_{21}$, $b \in e_1Ae_2$.

$$\begin{aligned}
 \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_1 \{\{a, b\}\}' e_{21} \otimes e_1 \{\{a, b\}\}'' e_{21}, \\
 \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2}(be_{21}ae_{21} \otimes e_1 - e_1 \otimes ae_{21}be_{21}), \\
 \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \Phi_2 e_{21} \{\{a, b\}\} \\
 &\quad + \frac{1}{2} (\Phi_2 e_{21} b \Phi_2 e_{21} a \otimes e_2 - e_2 \otimes \Phi_2 e_{21} a \Phi_2 e_{21} b) \\
 &\quad - \frac{1}{2} (\Phi_2 \otimes \Phi_2 e_{21} a e_{21} b - \Phi_2 e_{21} b e_{21} a \otimes \Phi_2) \\
 &\quad - \frac{1}{2} (e_{21} a \otimes \Phi_2^2 e_{21} b - \Phi_2^2 e_{21} a \otimes e_{21} b), \\
 \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} (\Phi_2 \otimes \Phi_2 e_{21} a e_{21} b - \Phi_2 e_{21} b e_{21} a \otimes \Phi_2) \\
 &\quad + \frac{1}{2} (e_{21} a \otimes \Phi_2^2 e_{21} b - \Phi_2^2 e_{21} a \otimes e_{21} b).
 \end{aligned}
 \tag{B.41}$$

Case 8.2: $d = e_{12}be_{21}$, $b \in e_2Ae_2$.

$$\begin{aligned}
 \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' e_{21} \otimes e_1 \{\{a, b\}\}'' e_{21}, \\
 \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= \frac{1}{2} (e_{12}be_{21}ae_{21} \otimes e_1 - ae_{21} \otimes e_{12}be_{21}), \\
 \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \Phi_2 e_{21} * \{\{a, b\}\} + \frac{1}{2} (b\Phi_2 e_{21} a \otimes e_2 - \Phi_2 e_{21} a \otimes b) \\
 &\quad - \frac{1}{2} (e_{21} a \otimes \Phi_2 b - be_{21} a \otimes \Phi_2), \\
 \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= + \frac{1}{2} (e_{21} a \otimes \Phi_2 b - be_{21} a \otimes \Phi_2).
 \end{aligned}
 \tag{B.42}$$

B.9

We consider a generator of fourth type, that is $c = e_{12}ae_{21}$ for $a \in e_2Ae_2$.

Case 9.1: $d = e_{12}be_{21}$, $b \in e_2Ae_2$.

$$\begin{aligned}
 \{\{c, d\}\}_{ind}^{2 \rightarrow 1} &= e_{12} \{\{a, b\}\}' e_{21} \otimes e_{12} \{\{a, b\}\}'' e_{21}, & \{\{c, d\}\}_{fus}^{2 \rightarrow 1} &= 0, \\
 \{\{\psi(c), \psi(d)\}\}_{ind}^{1 \rightarrow 2} &= \{\{a, b\}\}, & \{\{\psi(c), \psi(d)\}\}_{fus}^{1 \rightarrow 2} &= 0.
 \end{aligned}
 \tag{B.43}$$

Appendix C. Identities for the proof of Lemma 4.8

In this appendix, we show that (4.25) holds. It suffices to verify this identity on a choice of generators of A_1 . We will consider nine cases, which are given below with their images in A_2 under ψ . For $\hat{e} = 1 - e_2 - e_3$, we have

$$\psi(t) = t, \quad \text{if } t \in \hat{e}A\hat{e};$$

on generators of first type (2.13b) of A_1 ,

$$\psi(e_{12}u) = e_{12}u, \quad \text{if } u \in e_2A\widehat{e}; \quad \psi(e_{12}e_{23}u) = e_{13}u, \quad \text{if } u \in e_3A\widehat{e};$$

on generators of second type (2.13b) of A_1 ,

$$\psi(ve_{21}) = ve_{21}, \quad \text{if } v \in \widehat{e}Ae_2; \quad \psi(ve_{32}e_{21}) = ve_{31}, \quad \text{if } v \in \widehat{e}Ae_3;$$

on generators of third type (2.13c) of A_1 ,

$$\begin{aligned} \psi(e_{12}we_{21}) &= e_{12}we_{21}, & \text{if } w \in e_2Ae_2; \\ \psi(e_{12}we_{32}e_{21}) &= e_{12}we_{31}, & \text{if } w \in e_2Ae_3; \\ \psi(e_{12}e_{23}we_{21}) &= e_{13}we_{21}, & \text{if } w \in e_3Ae_2; \\ \psi(e_{12}e_{23}we_{32}e_{21}) &= e_{13}we_{31}, & \text{if } w \in e_3Ae_3; \end{aligned}$$

on generators of fourth type (2.13d) of A_1 .

Thanks to the cyclic antisymmetry, we have 45 cases where we have to check (4.25). If both $c, d \in A_1$ are of the form t for $t \in \widehat{e}A\widehat{e}$, $e_{12}u$ for $u \in e_2A\widehat{e}$, ve_{21} for $v \in \widehat{e}Ae_2$, or $e_{12}we_{21}$ for $w \in e_2Ae_2$, it is easy to derive that

$$(C.1) \quad \begin{aligned} \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= (\psi \otimes \psi) \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1}, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= 0, \quad (\psi \otimes \psi) \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} = 0, \end{aligned}$$

so that (4.25) is satisfied. So we are left with 35 cases to check. We will compute one case explicitly, and in the remaining cases we will only collect the four double brackets involved in (4.25) for the reader's convenience.

C.1

We consider that c is a generator of first type, that is $c = a$ for $a \in \widehat{e}A\widehat{e}$.

Case 1.1: $d = e_{12}be_{32}e_{21}$, $b \in e_2Ae_3$.

By definition, the double bracket $\{\{c, d\}\}_{1, fus}^{3 \rightarrow 2}$ is obtained from the fusion bracket in $A_{e_3 \rightarrow e_2}^f$ by inducing it in A_1 , so that we can write

$$\{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} = e_{12} \{\{a, be_{32}\}\}_{fus}^{3 \rightarrow 2} e_{21},$$

with the double bracket on the right-hand side given by the fusion bracket from Appendix A with $j = 3$, $i = 2$. In $A_{e_3 \rightarrow e_2}^f$, a is a generator of first type (2.13a) while be_{32} is a generator of third type (2.13c), so that the double bracket between them is given by (A.1c), and we get

$$\{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} = \frac{1}{2} e_{12} [be_{32}a \otimes e_2 - be_{32} \otimes ae_2] e_{21} = 0,$$

as $ae_2 = 0 = e_2a$. Next, we have

$$\{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} = \{\{a, e_{12}(be_{32})e_{21}\}\}_{fus}^{2 \rightarrow 1},$$

with the double bracket on the right-hand side given by the fusion bracket from Appendix A with $j = 2$, $i = 1$. In A_1 , a is a generator of first type (2.13a) while $e_{12}(be_{32})e_{21}$ is a generator of fourth type (2.13d) so that we get from (A.1d)

$$\begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2} \left(e_{12}(be_{32})e_{21}a \otimes e_1 + e_1 \otimes ae_{12}(be_{32})e_{21} \right. \\ &\quad \left. - e_{12}(be_{32})e_{21} \otimes ae_1 - e_1a \otimes e_{12}(be_{32})e_{21} \right) \\ &= \frac{1}{2} (dc \otimes e_1 + e_1 \otimes cd - d \otimes ce_1 - e_1c \otimes d). \end{aligned}$$

We now get the right-hand side of (4.25) by applying $\psi \otimes \psi$ to these two double brackets.

Meanwhile, we have in A_2 that $\psi(c) = a$ and $\psi(d) = e_{12}be_{31}$. We note that by inducing the double bracket from $A_{e_2 \rightarrow e_1}^f$ to A_2 ,

$$\{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} = \{\!\!\{a, e_{12}b\}\!\!\}_{2, fus}^{2 \rightarrow 1} e_{31},$$

where the double on the right-hand side given by the fusion bracket from Appendix A with $j = 2$, $i = 1$. We have that a is a generator of first type (2.13a) while $e_{12}b$ is a generator of second type (2.13b), so that using (A.1b)

$$\begin{aligned} \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes ae_{12}b - e_1a \otimes e_{12}b) e_{31} \\ &= \frac{1}{2} (e_1 \otimes \psi(c)\psi(d) - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

Finally, we have

$$\{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} = \{\!\!\{a, (e_{12}b)e_{31}\}\!\!\}_{2, fus}^{2 \rightarrow 1},$$

where the double on the right-hand side given by the fusion bracket from Appendix A with $j = 3$, $i = 1$. We have that a is a generator of first type (2.13a) while $(e_{12}b)e_{31}$ is a generator of third type (2.13c), so that using (A.1c)

$$\begin{aligned} \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (e_{12}be_{31}a \otimes e_1 - e_{12}be_{31} \otimes ae_1) \\ &= \frac{1}{2} (\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)e_1). \end{aligned}$$

Gathering the four expressions obtained for the double brackets, we can see that (4.25) is satisfied.

Case 1.2: $d = e_{12}e_{23}b$, $b \in e_3A\hat{e}$.

$$\begin{aligned} \text{(C.2)} \quad \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= 0, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes cd - e_1c \otimes d), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, & \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes \psi(c)\psi(d) - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

Case 1.3: $d = be_{32}e_{21}$, $b \in \widehat{e}Ae_3$.

$$(C.3) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= 0, & \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 - d \otimes ce_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, & \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)e_1). \end{aligned}$$

Case 1.4: $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.4) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= 0, \\ \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 + e_1 \otimes cd - d \otimes ce_1 - e_1c \otimes d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)e_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

Case 1.5: $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.5) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= 0, & \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 + e_1 \otimes cd - d \otimes ce_1 - e_1c \otimes d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \\ &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 + e_1 \otimes \psi(c)\psi(d) - \psi(d) \otimes \psi(c)e_1 - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

C.2

We consider that c is a generator of second type of the form $c = e_{12}a$ for $a \in e_2A\widehat{e}$.

Case 2.1: $d = e_{12}e_{23}b$, $b \in e_3A\widehat{e}$.

$$(C.6) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= -\frac{1}{2}c \otimes d, & \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes cd - dc \otimes e_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= -\frac{1}{2}\psi(d)\psi(c) \otimes e_1, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - \psi(c) \otimes \psi(d)). \end{aligned}$$

Case 2.2: $d = be_{32}e_{21}$, $b \in \widehat{e}Ae_3$.

$$(C.7) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}dc \otimes e_1, & \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(c \otimes e_1d - d \otimes ce_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}\psi(c) \otimes e_1\psi(d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)e_1). \end{aligned}$$

Case 2.3: $d = e_{12}be_{32}e_{21}$, $b \in e_2Ae_3$.

$$(C.8) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}dc \otimes e_1, & \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes cd - d \otimes ce_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}e_1 \otimes \psi(c)\psi(d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)e_1). \end{aligned}$$

Case 2.4: $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.9) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= -\frac{1}{2}c \otimes d, & \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes cd - d \otimes ce_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= -\frac{1}{2}\psi(d) \otimes \psi(c)e_1, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - \psi(c) \otimes \psi(d)). \end{aligned}$$

Case 2.5: $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.10) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}(dc \otimes e_1 - c \otimes d), \\ \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes cd - d \otimes ce_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 + e_1 \otimes \psi(c)\psi(d) - \psi(d) \otimes \psi(c)e_1 - \psi(c) \otimes \psi(d)). \end{aligned}$$

C.3

We consider that c is a generator of third type of the form $c = ae_{21}$ for $a \in \hat{e}Ae_2$.

Case 3.1: $d = e_{12}e_{23}b$, $b \in e_3A\hat{e}$.

$$(C.11) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}e_1 \otimes cd, & \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(de_1 \otimes c - e_1c \otimes d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}\psi(d)e_1 \otimes \psi(c), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

Case 3.2: $d = be_{32}e_{21}$, $b \in \hat{e}Ae_3$.

$$(C.12) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= -\frac{1}{2}d \otimes c, & \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 - e_1 \otimes cd), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= -\frac{1}{2}e_1 \otimes \psi(c)\psi(d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)). \end{aligned}$$

Case 3.3: $d = e_{12}be_{32}e_{21}$, $b \in e_2Ae_3$.

$$(C.13) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= -\frac{1}{2}d \otimes c, \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = \frac{1}{2}(dc \otimes e_1 - e_1c \otimes d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= -\frac{1}{2}e_1\psi(c) \otimes \psi(d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)). \end{aligned}$$

Case 3.4: $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.14) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}e_1 \otimes cd, \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = \frac{1}{2}(dc \otimes e_1 - e_1c \otimes d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}\psi(d)\psi(c) \otimes e_1, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

Case 3.5: $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.15) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}(e_1 \otimes cd - d \otimes c), \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = \frac{1}{2}(dc \otimes e_1 - e_1c \otimes d), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 + e_1 \otimes \psi(c)\psi(d) - \psi(d) \otimes \psi(c) - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

C.4

We consider that c is a generator of fourth type of the form $c = e_{12}ae_{21}$ for $a \in e_2Ae_2$.

Case 4.1: $d = e_{12}e_{23}b$, $b \in e_3A\hat{e}$.

$$(C.16) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}(e_1 \otimes cd - c \otimes d), \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = \frac{1}{2}(de_1 \otimes c - dc \otimes e_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}(\psi(d)e_1 \otimes \psi(c) - \psi(d)\psi(c) \otimes e_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - \psi(c) \otimes \psi(d)). \end{aligned}$$

Case 4.2: $d = be_{32}e_{21}$, $b \in \hat{e}Ae_3$.

$$(C.17) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}(dc \otimes e_1 - d \otimes c), \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = \frac{1}{2}(c \otimes e_1d - e_1 \otimes cd), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}(\psi(c) \otimes e_1\psi(d) - e_1 \otimes \psi(c)\psi(d)), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)). \end{aligned}$$

Case 4.3: $d = e_{12}be_{32}e_{21}$, $b \in e_2Ae_3$.

$$(C.18) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2} (dc \otimes e_1 - d \otimes c), \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (\psi(d)\psi(c) \otimes e_1 - \psi(d) \otimes \psi(c)). \end{aligned}$$

Case 4.4: $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.19) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2} (e_1 \otimes cd - c \otimes d), \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes \psi(c)\psi(d) - \psi(c) \otimes \psi(d)). \end{aligned}$$

Case 4.5: $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.20) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2} (dc \otimes e_1 + e_1 \otimes cd - d \otimes c - c \otimes d), \\ \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} &= 0, \quad \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} = 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \\ &= \frac{1}{2} (\psi(d)\psi(c) \otimes e_1 + e_1 \otimes \psi(c)\psi(d) - \psi(d) \otimes \psi(c) - \psi(c) \otimes \psi(d)). \end{aligned}$$

C.5

We consider that c is a generator of second type of the form $c = e_{12}e_{23}a$ for $a \in e_3A\hat{e}$.

Case 5.1: $d = e_{12}e_{23}b$, $b \in e_3A\hat{e}$.

$$(C.21) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= 0, \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = \frac{1}{2} (e_1 \otimes cd - dc \otimes e_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes \psi(c)\psi(d) - \psi(d)\psi(c) \otimes e_1). \end{aligned}$$

Case 5.2: $d = be_{32}e_{21}$, $b \in \hat{e}Ae_3$.

$$(C.22) \quad \begin{aligned} \{\{c, d\}\}_{1, fus}^{3 \rightarrow 2} &= 0, \quad \{\{c, d\}\}_{1, fus}^{2 \rightarrow 1} = \frac{1}{2} (c \otimes e_1d - d \otimes ce_1), \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\{\psi(c), \psi(d)\}\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (\psi(c) \otimes e_1\psi(d) - \psi(d) \otimes \psi(c)e_1). \end{aligned}$$

Case 5.3: $d = e_{12}be_{32}e_{21}$, $b \in e_2Ae_3$.

$$(C.23) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}c \otimes d, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes cd - d \otimes ce_1), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}e_1 \otimes \psi(c)\psi(d), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(c) \otimes \psi(d) - \psi(d) \otimes \psi(c)e_1). \end{aligned}$$

Case 5.4: $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.24) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= -\frac{1}{2}dc \otimes e_1, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes cd - d \otimes ce_1), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= -\frac{1}{2}\psi(d) \otimes \psi(c)e_1, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - \psi(d)\psi(c) \otimes e_1). \end{aligned}$$

Case 5.5: $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.25) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= 0, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes cd - d \otimes ce_1), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(e_1 \otimes \psi(c)\psi(d) - \psi(d) \otimes \psi(c)e_1). \end{aligned}$$

C.6

We consider that c is a generator of third type of the form $c = ae_{32}e_{21}$ for $a \in \widehat{e}Ae_3$.

Case 6.1: $d = be_{32}e_{21}$, $b \in \widehat{e}Ae_3$.

$$(C.26) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= 0, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 - e_1 \otimes cd), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - e_1 \otimes \psi(c)\psi(d)). \end{aligned}$$

Case 6.2: $d = e_{12}be_{32}e_{21}$, $b \in e_2Ae_3$.

$$(C.27) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= -\frac{1}{2}e_1 \otimes cd, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 - e_1c \otimes d), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= -\frac{1}{2}e_1\psi(c) \otimes \psi(d), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - e_1 \otimes \psi(c)\psi(d)). \end{aligned}$$

Case 6.3: $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.28) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}d \otimes c, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 - e_1c \otimes d), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= \frac{1}{2}\psi(d)\psi(c) \otimes e_1, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d) \otimes \psi(c) - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

Case 6.4: $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.29) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= 0, & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= \frac{1}{2}(dc \otimes e_1 - e_1c \otimes d), \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - e_1\psi(c) \otimes \psi(d)). \end{aligned}$$

C.7

We consider the remaining cases. The elements c, d are generators of fourth type in each of them.

Case 7.1: $c = e_{12}ae_{32}e_{21}$, $a \in e_2Ae_3$, and $d = e_{12}be_{32}e_{21}$, $b \in e_2Ae_3$.

$$(C.30) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}(dc \otimes e_1 - e_1 \otimes cd), & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - e_1 \otimes \psi(c)\psi(d)). \end{aligned}$$

Case 7.2: $c = e_{12}ae_{32}e_{21}$, $a \in e_2Ae_3$, and $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.31) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}(d \otimes c - c \otimes d), & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d) \otimes \psi(c) - \psi(c) \otimes \psi(d)). \end{aligned}$$

Case 7.3: $c = e_{12}ae_{32}e_{21}$, $a \in e_2Ae_3$, and $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.32) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2}(dc \otimes e_1 - c \otimes d), & \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2}(\psi(d)\psi(c) \otimes e_1 - \psi(c) \otimes \psi(d)). \end{aligned}$$

Case 7.4: $c = e_{12}e_{23}ae_{21}$, $a \in e_3Ae_2$, and $d = e_{12}e_{23}be_{21}$, $b \in e_3Ae_2$.

$$(C.33) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2} (e_1 \otimes cd - dc \otimes e_1), \quad \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} = 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes \psi(c)\psi(d) - \psi(d)\psi(c) \otimes e_1). \end{aligned}$$

Case 7.5: $c = e_{12}e_{23}ae_{21}$, $a \in e_3Ae_2$, and $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.34) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= \frac{1}{2} (e_1 \otimes cd - d \otimes c), \quad \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} = 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} &= \frac{1}{2} (e_1 \otimes \psi(c)\psi(d) - \psi(d) \otimes \psi(c)). \end{aligned}$$

Case 7.6: $c = e_{12}e_{23}ae_{32}e_{21}$, $a \in e_3Ae_3$, and $d = e_{12}e_{23}be_{32}e_{21}$, $b \in e_3Ae_3$.

$$(C.35) \quad \begin{aligned} \{\!\!\{c, d\}\!\!\}_{1, fus}^{3 \rightarrow 2} &= 0, \quad \{\!\!\{c, d\}\!\!\}_{1, fus}^{2 \rightarrow 1} = 0, \\ \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{2 \rightarrow 1} &= 0, \quad \{\!\!\{\psi(c), \psi(d)\}\!\!\}_{2, fus}^{3 \rightarrow 1} = 0. \end{aligned}$$

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