

CORRECTORS AND FIELD FLUCTUATIONS FOR THE $p_\epsilon(x)$ -LAPLACIAN WITH ROUGH EXPONENTS: THE SUBLINEAR GROWTH CASE

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Abstract. A corrector theory for the strong approximation of gradient fields inside periodic composites made from two materials with different power law behavior is provided. Each material component has a distinctly different exponent appearing in the constitutive law relating gradient to flux. The correctors are used to develop bounds on the local singularity strength for gradient fields inside micro-structured media. The bounds are multi-scale in nature and can be used to measure the amplification of applied macroscopic fields by the microstructure. The results in this paper are developed for materials having power law exponents strictly between -1 and zero.

Mathematics Subject Classification. 35J66, 35A15, 35B40, 74Q05.

Received September 11, 2011. Revised April 13, 2012.

Published online January 11, 2013.

1. INTRODUCTION

In this paper, we develop a corrector theory for the strong approximation of gradient fields inside periodic composites made from two materials with different power law behavior. Here the flux is related to the gradient ∇u by the power law $|\nabla u|^r \nabla u$. Each material component has a distinctly different exponent r appearing in the constitutive law relating gradient to flux. The correctors are used to develop bounds on the local singularity strength for gradient fields inside micro-structured media. The bounds are multi-scale in nature and can be used to measure the amplification of applied macroscopic fields by the microstructure. The novelty of the work presented in this paper is that it is carried out for materials having power law exponents r strictly between -1 and zero. Previous work [15], addresses power law materials with power law exponents r greater than or equal to zero. For that case strong approximations to the gradient fields are developed and lower bounds on the L^q norms ($q \geq 2$) of the gradient fields inside each material are provided. The bounds and strong approximations are expressed in terms of the correctors presented in Theorem 2.6 of [15].

The corrector theory for the linear case can be found in [20]. The earlier work of [9] provides the corrector theory for homogenization of monotone operators that in our case applies to composite materials made from constituents having the same power-law growth but with rough coefficients $\sigma(x)$. More recently, the homogenization of $p_\epsilon(x)$ -Laplacian boundary value problems for smooth exponential functions $p_\epsilon(x)$ uniformly converging to a

Keywords and phrases. Correctors, field concentrations, dispersed media, homogenization, layered media, p -Laplacian, periodic domain, power-law materials, young measures.

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limit function $p_0(x)$ has been studied in [1]. The convergence of the family of solutions for these homogenization problems is given in the topology of $L^{p_0(\cdot)}(\Omega)$.

Here we assume that the geometry of the composite is periodic and is specified by the indicator function of the sets occupied by each of the materials. The indicator function of material 1 and 2 are denoted by χ_1 and χ_2 , where $\chi_1(y) = 1$ in material 1 and is zero outside and $\chi_2(y) = 1 - \chi_1(y)$. The constitutive law for the heterogeneous medium is described by $A : \mathbb{R}^n \times \mathbb{R}^n \rightarrow \mathbb{R}^n$,

$$A(y, \xi) = \sigma(y) |\xi|^{p(y)-2} \xi, \quad (1.1)$$

with $\sigma(y) = \chi_1(y) \sigma_1 + \chi_2(y) \sigma_2$ and $0 < \sigma_1, \sigma_2 < \infty$; with $p(y) = \chi_1(y) p_1 + \chi_2(y) p_2$ and $1 < p_1 \leq p_2 < 2$ or $1 < p_1 \leq 2 \leq p_2$; and with both, $\sigma(y)$ and $p(y)$, periodic in y , with unit period cell $Y = (0, 1)^n$. This constitutive model occurs in several mathematical models of physical processes including nonlinear dielectrics [11, 12, 18, 27, 28], fluid flow (electrorheological fluids) [2, 4, 25], glaciology [13], image restoration [17], and in the theory of deformation plasticity under longitudinal shear (anti-plane strain deformation) [3, 14, 23, 24, 26].

In this paper, we study the problem of periodic homogenization associated with the solutions u_ϵ to the problems

$$-\operatorname{div} \left(A \left(\frac{x}{\epsilon}, \nabla u_\epsilon \right) \right) = f \text{ on } \Omega, \quad u_\epsilon \in W_0^{1,p_1}(\Omega), \quad (1.2)$$

where Ω is a bounded open subset of \mathbb{R}^n , $f \in W^{-1,q_2}(\Omega)$, and $1/p_1 + 1/q_2 = 1$. The differential operator on the left-hand side of (1.2) is the $p_\epsilon(x)$ -Laplacian. All solutions are understood in the usual weak sense [30].

It was shown in Chapter 15 of [30] that $\{u_\epsilon\}_{\epsilon>0}$ converges weakly in $W^{1,p_1}(\Omega)$ to the solution u of the homogenized problem

$$-\operatorname{div} (b(\nabla u)) = f \text{ on } \Omega, \quad u \in W_0^{1,p_1}(\Omega), \quad (1.3)$$

where the monotone map $b : \mathbb{R}^n \rightarrow \mathbb{R}^n$ (independent of f and Ω) can be obtained by solving an auxiliary problem for the operator (1.2) on a periodicity cell.

The idea of homogenization is intimately related to the Γ -convergence of a suitable family of energy functionals I_ϵ as $\epsilon \rightarrow 0$ [30]. Here the connection is natural in that the family of boundary value problems (1.2) correspond to the Euler equations of the associated energy functionals I_ϵ and the solutions u_ϵ are their minimizers. The homogenized solution is precisely the minimizer of the Γ -limit of the sequence $\{I_\epsilon\}_{\epsilon>0}$. The connections between Γ limits and homogenization for the power-law materials studied here can be found in Chapter 15 of [30]. The explicit formula for the Γ -limit of the associated energy functionals for layered materials was obtained recently in [22].

The homogenization result found in Chapter 5 of [30] shows that the *average* of the error incurred in approximating $\{\nabla u_\epsilon\}_{\epsilon>0}$ in terms of ∇u , where u is the solution of (1.3) decays to 0. Then again, the presence of large local fields either electric or mechanical often precede the onset of material failure (see, [16]). The goal of our analysis is to develop tools for quantifying the effect of load transfer between length scales inside heterogeneous media. To this end, we present a new corrector result that approximates, inside each phase, ∇u_ϵ up to an error that converges to zero strongly in the norm (see Sect. 2.2.1).

The corrector result is then used to develop new tools that provide lower bounds on the local gradient field intensity inside micro-structured media. The bounds are expressed in terms of the L^q -norms of gradients of the solutions of the local corrector problems. These results provide a lower bound on the amplification of the macroscopic gradient field by the microstructure see, Section 2.2.2. These bounds provide a rigorous way to assess the effect of field concentrations generated by the microgeometry without having to compute the actual solution u_ϵ . In [19], similar lower bounds were established for field concentrations for mixtures of linear electrical conductors in the context of two scale convergence.

In this paper, the corrector results are presented for layered materials (Fig. 2) and for dispersions of inclusions embedded inside a host medium (Fig. 1). For the dispersed microstructures the included material is taken to have the lower power-law exponent than that of the host phase. The reason we use dispersed and layered microstructures is that in both cases we are able to show that the homogenized solution lies in $W_0^{1,p_2}(\Omega)$,

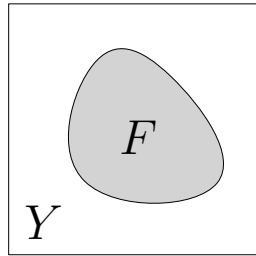


FIGURE 1. Unit cell: dispersed microstructure.

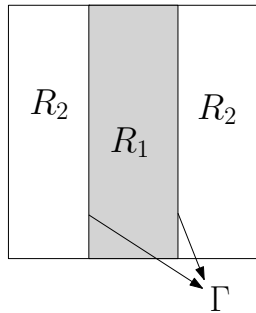


FIGURE 2. Unit cell: layered material.

see Theorem 2.5. Possible extensions of this work include the study of other microstructures for which this higher order integrability condition of the homogenized solution is satisfied. The higher order integrability is used to provide an algorithm for building correctors and construct a sequence of strong approximations to the gradients inside each material, see Theorem 2.7. When the host phase has a lower power-law exponent than the included phase, one can only conclude that the homogenized solution lies in $W_0^{1,p_1}(\Omega)$ and the techniques developed here do not apply.

The presentation of the paper is organized as follows. In Section 2, we state the problem and formulate the main results. Section 3 contains technical lemmas and integral inequalities for the correctors used to prove the main results. Section 4 contains the proof of the main results. The Appendix contains all proofs of lemmas stated in Section 3 and some remarks related to the proof of Theorem 2.7 found in Section 4.

2. STATEMENT OF THE PROBLEM AND MAIN RESULTS

2.1. Notation

In this paper we consider two nonlinear power-law materials periodically distributed inside a domain $\Omega \subset \mathbb{R}^n$. The periodic mixture is described as follows. We introduce the unit period cell $Y = (0, 1)^n$ of the microstructure. Let F be an open subset of Y of material 1, with smooth boundary ∂F , such that $\overline{F} \subset Y$. The function $\chi_1(y) = 1$ inside F and 0 outside and $\chi_2(y) = 1 - \chi_1(y)$. We extend $\chi_1(y)$ and $\chi_2(y)$ by periodicity to \mathbb{R}^n and the ϵ -periodic mixture inside Ω is described by the oscillatory characteristic functions $\chi_1^\epsilon(x) = \chi_1(x/\epsilon)$ and $\chi_2^\epsilon(x) = \chi_2(x/\epsilon)$. Here we will consider the case where F is given by a simply connected inclusion embedded inside a host material (see Fig. 1). A distribution of such inclusions is commonly referred to as a periodic dispersion of inclusions.

We also consider layered materials. For this case the representative unit cell consists of a layer of material 1, denoted by R_1 , sandwiched between layers of material 2, denoted by R_2 . The interior boundary of R_1 is denoted by Γ (see Fig. 2). Here $\chi_1(y) = 1$ for $y \in R_1$ and 0 in R_2 , and $\chi_2(y) = 1 - \chi_1(y)$.

We denote by $\theta_1 = \int_Y \chi_1(y)dy$ and $\theta_2 = 1 - \theta_1$ the volume fractions of phase 1 and phase 2 inside the composite.

On the unit cell Y , the constitutive law for the nonlinear material is given by (1.1) with exponents p_1 and p_2 satisfying $1 < p_1 \leq p_2 \leq 2$ or $1 < p_1 \leq 2 \leq p_2$. Their Hölder conjugates are denoted by $q_2 = p_1/(p_1 - 1)$ and $q_1 = p_2/(p_2 - 1)$ respectively. For $i = 1, 2$, $W_{\text{per}}^{1,p_i}(Y)$ denotes the set of all functions $u \in W^{1,p_i}(Y)$ with mean value zero that have the same trace on the opposite faces of Y . Each function $u \in W_{\text{per}}^{1,p_i}(Y)$ can be extended by periodicity to a function of $W_{\text{loc}}^{1,p_i}(\mathbb{R}^n)$.

The Euclidean norm and the scalar product in \mathbb{R}^n are denoted by $|\cdot|$ and (\cdot, \cdot) , respectively. If $A \subset \mathbb{R}^n$, $|A|$ denotes the Lebesgue measure and $\chi_A(x)$ denotes its characteristic function.

The constitutive law for the ϵ -periodic composite is described by $A_\epsilon(x, \xi) = A(x/\epsilon, \xi)$, for every $\epsilon > 0$, for every $x \in \Omega$, and for every $\xi \in \mathbb{R}^n$.

We have [7] that A fulfills the following conditions:

- (1) For all $\xi \in \mathbb{R}^n$, $A(\cdot, \xi)$ is Y -periodic and Lebesgue measurable.
- (2) $|A(y, 0)| = 0$ for all $y \in \mathbb{R}^n$.
- (3) Continuity: for almost every $y \in \mathbb{R}^n$ and for every $\xi_i \in \mathbb{R}^n$ ($i = 1, 2$) we have

$$|A(y, \xi_1) - A(y, \xi_2)| \leq C |\xi_1 - \xi_2|^{\alpha(y)} (1 + |\xi_1| + |\xi_2|)^{p(y)-1-\alpha(y)} \tag{2.1}$$

where $\alpha(y) = \chi_1(y)\alpha_1(y) + \chi_2(y)\alpha_2(y)$ and

$$\alpha_i(y) = \begin{cases} p_i - 1 & \text{if } 1 \leq p_i \leq 2 \\ 1 & \text{if } p_i \geq 2. \end{cases}$$

- (4) Monotonicity: for almost every $y \in \mathbb{R}^n$ and for every $\xi_i \in \mathbb{R}^n$ ($i = 1, 2$) we have

$$\begin{aligned} (A(y, \xi_1) - A(y, \xi_2), \xi_1 - \xi_2) &\geq C |\xi_1 - \xi_2|^{\beta(y)} (|\xi_1| + |\xi_2|)^{p(y)-\beta(y)} \\ &\geq C |\xi_1 - \xi_2|^{\beta(y)} (1 + |\xi_1| + |\xi_2|)^{p(y)-\beta(y)} \end{aligned} \tag{2.2}$$

where $\beta(y) = \chi_1(y)\beta_1(y) + \chi_2(y)\beta_2(y)$ and

$$\beta_i(y) = \begin{cases} 2 & \text{if } 1 \leq p_i \leq 2 \\ p_i & \text{if } p_i \geq 2. \end{cases}$$

The structure conditions for A given by (2.1) and (2.2) recover the ones stated in [15] where $\alpha_i(y) = 1$ and $\beta_i(y) = p_i$ for $i = 1, 2$. In the context of this paper, the analysis when the exponents p_1 and p_2 are in the regime between 1 and 2 becomes more involved. This particular set of structure conditions (or related variants) are used, for example, in [5, 6, 9, 10].

2.2. Dirichlet boundary value problem

We consider the following Dirichlet boundary value problem

$$\begin{cases} -\text{div}(A_\epsilon(x, \nabla u_\epsilon)) = f \text{ on } \Omega, \\ u_\epsilon \in W_0^{1,p_1}(\Omega). \end{cases} \tag{2.3}$$

where $f \in W^{-1,q_2}(\Omega)$.

The following homogenization result holds.

Theorem 2.1 (homogenization theorem (see Chap. 15 of [30])). *As $\epsilon \rightarrow 0$, the solutions u_ϵ of (2.3) converge weakly to u in $W^{1,p_1}(\Omega)$, where u is the solution of*

$$-\operatorname{div}(b(\nabla u)) = f \text{ on } \Omega, \quad (2.4)$$

$$u \in W_0^{1,p_1}(\Omega); \quad (2.5)$$

and the function $b : \mathbb{R}^n \rightarrow \mathbb{R}^n$ is defined for all $\xi \in \mathbb{R}^n$ by

$$b(\xi) = \int_Y A(y, P(y, \xi)) dy, \quad (2.6)$$

where $p : \mathbb{R}^n \times \mathbb{R}^n \rightarrow \mathbb{R}^n$ is defined by

$$P(y, \xi) = \xi + \nabla v_\xi(y), \quad (2.7)$$

where v_ξ is the solution to the cell problem:

$$\begin{cases} \int_Y (A(y, \xi + \nabla v_\xi), \nabla w) dy = 0, \text{ for every } w \in W_{\text{per}}^{1,p_1}(Y), \\ v_\xi \in W_{\text{per}}^{1,p_1}(Y). \end{cases} \quad (2.8)$$

Remark 2.2. The following *a priori* bound is satisfied

$$\sup_{\epsilon > 0} \left(\int_\Omega \chi_1^\epsilon(x) |\nabla u_\epsilon(x)|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |\nabla u_\epsilon(x)|^{p_2} dx \right) \leq C < \infty, \quad (2.9)$$

where C does not depend on ϵ . The proof of this bound is given in Lemma 3.5.

Remark 2.3. The function b , defined in (2.6), is continuous and monotone (see Lem. 3.7).

Remark 2.4. Since the solution v_ξ of (2.8) can be extended by periodicity to a function of $W_{\text{loc}}^{1,p_1}(\mathbb{R}^n)$, then (2.8) is equivalent to $-\operatorname{div}(A(y, \xi + \nabla v_\xi(y))) = 0$ over $D'(\mathbb{R}^n)$, i.e.,

$$-\operatorname{div}(A(y, P(y, \xi))) = 0 \text{ in } D'(\mathbb{R}^n) \text{ for every } \xi \in \mathbb{R}^n. \quad (2.10)$$

Moreover, by (2.8), we have

$$\int_Y (A(y, P(y, \xi)), P(y, \xi)) dy = \int_Y (A(y, P(y, \xi)), \xi) dy = (b(\xi), \xi). \quad (2.11)$$

For $\epsilon > 0$, define $P_\epsilon : \mathbb{R}^n \times \mathbb{R}^n \rightarrow \mathbb{R}^n$ by

$$P_\epsilon(x, \xi) = P\left(\frac{x}{\epsilon}, \xi\right) = \xi + \nabla v_\xi\left(\frac{x}{\epsilon}\right), \quad (2.12)$$

where v_ξ is the unique solution of (2.8). The functions P and P_ϵ are easily seen to have the following properties

$$P(\cdot, \xi) \text{ is } Y\text{-periodic and } P_\epsilon(x, \xi) \text{ is } \epsilon\text{-periodic in } x, \quad (2.13)$$

$$\int_Y P(y, \xi) dy = \xi, \quad (2.14)$$

$$P_\epsilon(\cdot, \xi) \rightharpoonup \xi \text{ in } L^{p_1}(\Omega; \mathbb{R}^n) \text{ as } \epsilon \rightarrow 0, \quad (2.15)$$

$$P(y, 0) = 0 \text{ for almost every } y, \tag{2.16}$$

$$A\left(\frac{\cdot}{\epsilon}, P_\epsilon(\cdot, \xi)\right) \rightharpoonup b(\xi) \text{ in } L^{q_2}(\Omega; \mathbb{R}^n), \text{ as } \epsilon \rightarrow 0. \tag{2.17}$$

We now state the higher order integrability properties of the homogenized solution for periodic dispersions of inclusions and layered microgeometries.

Theorem 2.5. *Given a periodic dispersion of inclusions or a layered material then the solution u of (2.4) belongs to $W_0^{1,p_2}(\Omega)$.*

Remark 2.6. The proof of Theorem 2.5 [15] uses a variational approach and considers the homogenized Lagrangian associated with $b(\xi)$ defined in (2.6). The integrability of the homogenized solution u of (2.4) is determined by the growth of the homogenized Lagrangian with respect to its argument, which follows from the regularity of the Lagrangian. For periodic dispersed and layered microstructures no Lavrentiev phenomenon occurs. The proof of the regularity of the homogenized Lagrangian for periodic dispersed microstructure can be found in Chapter 14 of [30] and for layered microstructure in [15]. Both proofs are valid for $p_1, p_2 \geq 1$

2.2.1. *Statement of the corrector theorem*

We now describe the family of correctors that provide a strong approximation of the sequence $\{\chi_i^\epsilon \nabla u_\epsilon\}_{\epsilon > 0}$ in the $L^{p_i}(\Omega, \mathbb{R}^n)$ norm, for $i = 1, 2$. We denote the rescaled period cell with side length $\epsilon > 0$ by Y_ϵ and write $Y_\epsilon^k = \epsilon k + Y_\epsilon$, where $k \in \mathbb{Z}^n$. In what follows it is convenient to define the index set $I_\epsilon = \{k \in \mathbb{Z}^n : Y_\epsilon^k \subset \Omega\}$. For $\varphi \in L^{p_2}(\Omega; \mathbb{R}^n)$, we define the local average operator M_ϵ associated with the partition $Y_\epsilon^k, k \in I_\epsilon$ by

$$M_\epsilon(\varphi)(x) = \sum_{k \in I_\epsilon} \chi_{Y_\epsilon^k}(x) \frac{1}{|Y_\epsilon^k|} \int_{Y_\epsilon^k} \varphi(y) dy; \text{ if } x \in \bigcup_{k \in I_\epsilon} Y_\epsilon^k. \tag{2.18}$$

The family of approximations of the identity map M_ϵ has the following properties (for a proof, see, for example [29]):

1. For $i = 1, 2, \|M_\epsilon(\varphi) - \varphi\|_{L^{p_i}(\Omega; \mathbb{R}^n)} \rightarrow 0$ as $\epsilon \rightarrow 0$.
2. $M_\epsilon(\varphi) \rightarrow \varphi$ a.e. on Ω .
3. From Jensen’s inequality we have $\|M_\epsilon(\varphi)\|_{L^{p_i}(\Omega; \mathbb{R}^n)} \leq \|\varphi\|_{L^{p_2}(\Omega; \mathbb{R}^n)}$, for every $\varphi \in L^{p_2}(\Omega; \mathbb{R}^n)$ and $i = 1, 2$.

The strong approximation to the sequence $\{\chi_i^\epsilon \nabla u_\epsilon\}_{\epsilon > 0}$ is given by the following corrector theorem.

Theorem 2.7 (corrector theorem). *Let $f \in W^{-1,q_2}(\Omega)$, let u_ϵ be the solutions to the problem (2.3), and let u be the solution to problem (2.4). Then, up to a subsequence, for periodic dispersions of inclusions and for layered materials, we have*

$$\int_\Omega |\chi_i^\epsilon(x) P_\epsilon(x, M_\epsilon(\nabla u)(x)) - \chi_i^\epsilon(x) \nabla u_\epsilon(x)|^{p_i} dx \rightarrow 0, \tag{2.19}$$

as $\epsilon \rightarrow 0$, for $i = 1, 2$.

The proof of Theorem 2.7 is given in Section 4.1.

2.2.2. *Lower bounds on the local amplification of the macroscopic field*

We show lower bounds on the L^q norm of the gradient fields inside each material that are given in terms of the correctors presented in Theorem 2.7. We begin by presenting a general lower bound that holds for the composition of the sequence $\{\chi_i^\epsilon \nabla u_\epsilon\}_{\epsilon > 0}$ with any non-negative Carathéodory function. Recall that $\psi : \Omega \times \mathbb{R}^n \rightarrow \mathbb{R}$ is a Carathéodory function if $\psi(x, \cdot)$ is continuous for almost every $x \in \Omega$ and if $\psi(\cdot, \lambda)$ is measurable in x for every $\lambda \in \mathbb{R}^n$. The lower bound on the sequence obtained by the composition of $\psi(x, \cdot)$ with $\chi_i^\epsilon(x) \nabla u_\epsilon(x)$ is given by

Theorem 2.8. For all Carathéodory functions $\psi \geq 0$ and measurable sets $D \subset \Omega$ we have

$$\int_D \int_Y \psi(x, \chi_i(y) P(y, \nabla u(x))) \, dy dx \leq \liminf_{\epsilon \rightarrow 0} \int_D \psi(x, \chi_i^\epsilon(x) \nabla u_\epsilon(x)) \, dx.$$

If the sequence $\{\psi(x, \chi_i^\epsilon(x) \nabla u_\epsilon(x))\}_{\epsilon > 0}$ is weakly convergent in $L^1(\Omega)$, then the inequality becomes an equality.

In particular, for $\psi(x, \lambda) = |\lambda|^q$ with $q > 1$, we have

$$\int_D \int_Y \chi_i(y) |P(y, \nabla u(x))|^q \, dy dx \leq \liminf_{\epsilon \rightarrow 0} \int_D \chi_i^\epsilon(x) |\nabla u_\epsilon(x)|^q \, dx. \quad (2.20)$$

Theorem 2.8 together with (2.20) provide explicit lower bounds on the gradient field inside each material. It relates the local excursions of the gradient inside each phase $\chi_i^\epsilon \nabla u_\epsilon$ to the average gradient ∇u through the multiscale quantity given by the corrector $P(y, \nabla u(x))$. It is clear from (2.20) that the $L^q(Y \times \Omega; \mathbb{R}^n)$ integrability of $P(y, \nabla u(x))$ provides a lower bound on the $L^q(\Omega; \mathbb{R}^n)$ integrability of ∇u_ϵ .

The proof of Theorem 2.8 is given in Section 4.2.

3. TECHNICAL LEMMAS

In this section we state some technical a priori bounds and convergence properties for the sequences P_ϵ defined in (2.12), ∇u_ϵ , and $A_\epsilon(x, P_\epsilon(x, \nabla u_\epsilon))$ that are used in the proof of the main results of this paper. The proofs of these lemmas can be found in the Appendix.

Lemma 3.1. For every $\xi \in \mathbb{R}^n$ we have

$$\int_Y \chi_1(y) |P(y, \xi)|^{p_1} \, dy + \int_Y \chi_2(y) |P(y, \xi)|^{p_2} \, dy \leq C (1 + |\xi|^{p_1} \theta_1 + |\xi|^{p_2} \theta_2) \quad (3.1)$$

and by a change of variables, we obtain

$$\int_{Y_\epsilon} \chi_1^\epsilon(x) |P_\epsilon(x, \xi)|^{p_1} \, dx + \int_{Y_\epsilon} \chi_2^\epsilon(x) |P_\epsilon(x, \xi)|^{p_2} \, dx \leq C (1 + |\xi|^{p_1} \theta_1 + |\xi|^{p_2} \theta_2) |Y_\epsilon|. \quad (3.2)$$

Lemma 3.2. For every $\xi_1, \xi_2 \in \mathbb{R}^n$ we have

– For $1 < p_1 \leq p_2 \leq 2$:

$$\begin{aligned} & \int_{Y_\epsilon} \chi_1^\epsilon(x) |P_\epsilon(x, \xi_1) - P_\epsilon(x, \xi_2)|^{p_1} \, dx + \int_{Y_\epsilon} \chi_2^\epsilon(x) |P_\epsilon(x, \xi_1) - P_\epsilon(x, \xi_2)|^{p_2} \, dx \\ & \leq C \left[\theta_1^{\frac{1}{3-p_1}} |\xi_1 - \xi_2|^{\frac{p_1}{3-p_1}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_1}{3-p_1}} \right. \\ & \quad + \theta_2^{\frac{p_1}{2p_2-p_1p_2+p_1}} |\xi_1 - \xi_2|^{\frac{p_1p_2}{2p_2-p_1p_2+p_1}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{p_2(2-p_1)}{2p_2-p_1p_2+p_1}} \\ & \quad + \theta_1^{\frac{p_2}{2p_1-p_1p_2+p_2}} |\xi_1 - \xi_2|^{\frac{p_1p_2}{2p_1-p_1p_2+p_2}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{p_1(2-p_2)}{2p_1-p_1p_2+p_2}} \\ & \quad \left. + \theta_2^{\frac{1}{3-p_2}} |\xi_1 - \xi_2|^{\frac{p_2}{3-p_2}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_2}{3-p_2}} \right] |Y_\epsilon|; \end{aligned} \quad (3.3)$$

– For $1 < p_1 \leq 2 \leq p_2$:

$$\begin{aligned}
 & \int_{Y_\epsilon} \chi_1^\epsilon(x) |P_\epsilon(x, \xi_1) - P_\epsilon(x, \xi_2)|^{p_1} dx + \int_{Y_\epsilon} \chi_2^\epsilon(x) |P_\epsilon(x, \xi_1) - P_\epsilon(x, \xi_2)|^{p_2} dx \\
 & \leq C \left[\theta_1^{\frac{1}{3-p_1}} |\xi_1 - \xi_2|^{\frac{p_1}{3-p_1}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_1}{3-p_1}} \right. \\
 & \quad + \theta_2^{\frac{p_1}{2p_2-p_1}} |\xi_1 - \xi_2|^{\frac{p_1 p_2}{2p_2-p_1}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2(p_2-p_1)}{2p_2-p_1}} + \theta_1 |\xi_1 - \xi_2|^{p_1} \\
 & \quad \left. + \theta_2^{\frac{1}{p_2-1}} |\xi_1 - \xi_2|^{\frac{p_2}{p_2-1}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{p_2-2}{p_2-1}} \right] |Y_\epsilon|.
 \end{aligned} \tag{3.4}$$

Remark 3.3. Note the two “extra” terms in (3.3) of Lemma 3.2 where there is a “mixing” of the exponents p_1 and p_2 which do not appear in the corresponding property of P_ϵ given by Lemma 5.2 in [15]. Since Lemma 3.2 is used to prove Lemma 3.4, these two terms appear again in (3.6) and therefore in the proof of Theorem 2.7 and the proof of Lemma 4.1 used to prove Theorem 2.8.

Lemma 3.4. *Let φ be such that*

$$\sup_{\epsilon > 0} \left\{ \int_\Omega \chi_1^\epsilon(x) |\varphi(x)|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |\varphi(x)|^{p_2} dx \right\} < \infty,$$

and let Ψ be a simple function of the form

$$\Psi(x) = \sum_{j=0}^m \eta_j \chi_{\Omega_j}(x), \tag{3.5}$$

with $\eta_j \in \mathbb{R}^n \setminus \{0\}$, $\Omega_j \subset\subset \Omega$, $|\partial\Omega_j| = 0$, $\Omega_j \cap \Omega_k = \emptyset$ for $j \neq k$ and $j, k = 1, \dots, m$; and set $\eta_0 = 0$ and $\Omega_0 = \Omega \setminus \bigcup_{j=1}^m \Omega_j$. Then

– For $1 < p_1 \leq p_2 \leq 2$:

$$\begin{aligned}
 & \limsup_{\epsilon \rightarrow 0} \sum_{i=1}^2 \int_\Omega \chi_i^\epsilon(x) |P_\epsilon(x, M_\epsilon \varphi) - P_\epsilon(x, \Psi)|^{p_i} dx \\
 & \leq C \limsup_{\epsilon \rightarrow 0} \left[\left(\int_\Omega \chi_1^\epsilon(x) |\varphi - \Psi|^{p_1} dx \right)^{\frac{1}{3-p_1}} \left(|\Omega| + \int_\Omega \chi_1^\epsilon(x) |\varphi|^{p_1} dx \right. \right. \\
 & \quad \left. \left. + \int_\Omega \chi_2^\epsilon(x) |\varphi|^{p_2} dx + \int_\Omega \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |\Psi|^{p_2} dx \right)^{\frac{2-p_1}{3-p_1}} \right. \\
 & \quad \left. + \left(\int_\Omega \chi_2^\epsilon(x) |\varphi - \Psi|^{p_2} dx \right)^{\frac{p_1}{p_2-p_1 p_2+p_1}} \left(|\Omega| + \int_\Omega \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |\varphi|^{p_2} dx \right. \right. \\
 & \quad \left. \left. + \int_\Omega \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |\Psi|^{p_2} dx \right)^{\frac{p_2(2-p_1)}{p_2-p_1 p_2+p_1}} \right] \tag{3.6}
 \end{aligned}$$

$$\begin{aligned}
& + \left(\int_{\Omega} \chi_1^\epsilon(x) |\varphi - \Psi|^{p_1} dx \right)^{\frac{p_2}{p_1 - p_1 p_2 + p_2}} \left(|\Omega| + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx \right. \\
& + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx \left. \right)^{\frac{p_1(2-p_2)}{p_1 - p_1 p_2 + p_2}} \\
& + \left(\int_{\Omega} \chi_2^\epsilon(x) |\varphi - \Psi|^{p_2} dx \right)^{\frac{1}{3-p_2}} \left(|\Omega| + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx \right. \\
& \left. + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx \right)^{\frac{2-p_2}{3-p_2}} \Bigg];
\end{aligned}$$

– For $1 < p_1 \leq 2 \leq p_2$:

$$\begin{aligned}
& \limsup_{\epsilon \rightarrow 0} \sum_{i=1}^2 \int_{\Omega} \chi_i^\epsilon(x) |P_\epsilon(x, M_\epsilon \varphi) - P_\epsilon(x, \Psi)|^{p_i} dx \\
& \leq C \limsup_{\epsilon \rightarrow 0} \left[\left(\int_{\Omega} \chi_1^\epsilon(x) |\varphi - \Psi|^{p_1} dx \right)^{\frac{1}{3-p_1}} \left(|\Omega| + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx \right. \right. \\
& \quad \left. \left. + \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx \right)^{\frac{2-p_1}{3-p_1}} \right. \\
& \quad \left. + \left(\int_{\Omega} \chi_2^\epsilon(x) |\varphi - \Psi|^{p_2} dx \right)^{\frac{p_1}{2p_2-p_1}} \left(|\Omega| + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx \right. \right. \\
& \quad \left. \left. + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx \right)^{\frac{2(p_2-p_1)}{2p_2-p_1}} \right. \\
& \quad \left. + \int_{\Omega} \chi_1^\epsilon(x) |\varphi - \Psi|^{p_1} dx + \left(\int_{\Omega} \chi_2^\epsilon(x) |\varphi - \Psi|^{p_2} dx \right)^{\frac{1}{p_2-1}} \left(|\Omega| + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx \right. \right. \\
& \quad \left. \left. + \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx \right)^{\frac{p_2-2}{p_2-1}} \right].
\end{aligned} \tag{3.7}$$

Lemma 3.5. *Let u_ϵ be the solution to (2.3). Then the a priori bound (2.9) holds.*

Lemma 3.6. *If the microstructure is dispersed or layered, we have that*

$$\sup_{\epsilon > 0} \left\{ \int_{\Omega} \chi_i^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x))|^{p_i} dx \right\} \leq C < \infty, \text{ for } i = 1, 2.$$

We use Lemma 3.1 to prove structure conditions of b (2.6) in the following lemma.

Lemma 3.7. *The function b , defined in (2.6), satisfies the following structure properties: for every $\xi_1, \xi_2 \in \mathbb{R}^n$*

1. *Monotonicity:*

$$(b(\xi_2) - b(\xi_1), \xi_2 - \xi_1) \geq 0. \tag{3.8}$$

2. *Continuity: there exists a positive constant C such that*

– For $1 < p_1 \leq p_2 \leq 2$:

$$\begin{aligned}
|b(\xi_1) - b(\xi_2)| \leq C & \left[|\xi_2 - \xi_1|^{\frac{p_1-1}{3-p_1}} (1 + \theta_1 |\xi_1|^{p_1} + \theta_2 |\xi_1|^{p_2} + \theta_1 |\xi_2|^{p_1} + \theta_2 |\xi_2|^{p_2})^{\frac{(2-p_1)(p_1-1)}{p_1(3-p_1)}} \right. \\
& \left. + |\xi_2 - \xi_1|^{\frac{p_2-1}{3-p_2}} (1 + \theta_1 |\xi_1|^{p_1} + \theta_2 |\xi_1|^{p_2} + \theta_1 |\xi_2|^{p_1} + \theta_2 |\xi_2|^{p_2})^{\frac{(2-p_2)(p_2-1)}{p_2(3-p_2)}} \right];
\end{aligned} \tag{3.9}$$

– For $1 < p_1 \leq 2 \leq p_2$:

$$\begin{aligned}
 |b(\xi_1) - b(\xi_2)| \leq C & \left[|\xi_2 - \xi_1|^{\frac{p_1-1}{3-p_1}} (1 + \theta_1 |\xi_1|^{p_1} + \theta_2 |\xi_1|^{p_2} + \theta_1 |\xi_2|^{p_1} + \theta_2 |\xi_2|^{p_2})^{\frac{(2-p_1)(p_1-1)}{p_1(3-p_1)}} \right. \\
 & \left. + |\xi_2 - \xi_1|^{\frac{1}{p_2-1}} (1 + \theta_1 |\xi_1|^{p_1} + \theta_2 |\xi_1|^{p_2} + \theta_1 |\xi_2|^{p_1} + \theta_2 |\xi_2|^{p_2})^{\frac{p_2-2}{p_2-1}} \right]. \tag{3.10}
 \end{aligned}$$

These structure conditions for b are different to the ones obtained in [15] and their proofs require different techniques, for example, the use of (A.4) to obtain (3.9). These structure conditions (3.9) and (3.10) will be used in the proof of Theorem 2.7.

Lemma 3.8. For all $j = 0, \dots, m$, we have that $\int_{\Omega_j} |(A_\epsilon(x, P_\epsilon(x, \eta_j)), \nabla u_\epsilon(x))| dx$ and $\int_{\Omega_j} \left| (A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, \eta_j)) \right| dx$ are uniformly bounded with respect to ϵ .

Lemma 3.9. As $\epsilon \rightarrow 0$, up to a subsequence, $(A_\epsilon(\cdot, P_\epsilon(\cdot, \eta_j)), \nabla u_\epsilon(\cdot))$ converges weakly to a function $g_j \in L^1(\Omega_j; \mathbb{R})$, for all $j = 0, \dots, m$. In a similar way, up to a subsequence, $(A_\epsilon(\cdot, \nabla u_\epsilon(\cdot)), P_\epsilon(\cdot, \eta_j))$ converges weakly to a function $h_j \in L^1(\Omega_j; \mathbb{R})$, for all $j = 0, \dots, m$.

4. PROOF OF MAIN RESULTS

4.1. Proof of the corrector theorem

We are now in the position to give the proof of Theorem 2.7. We present the proof for the case when $1 < p_1 \leq p_2 \leq 2$, for $1 < p_1 \leq 2 \leq p_2$ the proof is very similar and the correspondig formulas can be found in the appendix, in Section A.9

Proof. Let $u_\epsilon \in W_0^{1,p_1}(\Omega)$ be the solutions of (2.3). By (2.2), Lemmas 3.5 and 3.6 we have that

$$\begin{aligned}
 & \int_{\Omega} [\chi_1^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)|^{p_1} + \chi_2^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)|^{p_2}] dx \\
 \leq C & \left[(|\Omega| + \int_{\Omega} \chi_1^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x))|^{p_1} dx + \int_{\Omega} \chi_1^\epsilon(x) |\nabla u_\epsilon(x)|^{p_1} dx)^{\frac{2-p_1}{2}} \right. \\
 & \times \left(\int_{\Omega} \chi_1^\epsilon(x) (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \right)^{\frac{p_1}{2}} \\
 & + \left(\int_{\Omega} \chi_2^\epsilon(x) (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \right)^{\frac{p_2}{2}} \\
 & \left. \times \left(\int_{\Omega} \chi_2^\epsilon(x) (1 + |P_\epsilon(x, M_\epsilon \nabla u(x))|^{p_2} + |\nabla u_\epsilon(x)|^{p_2}) dx \right)^{\frac{2-p_2}{2}} \right] \\
 \leq C & \sum_{i=1}^2 \left(\int_{\Omega} (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \right)^{\frac{p_i}{2}}.
 \end{aligned}$$

To prove Theorem 2.7, we show that

$$\begin{aligned} & \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u(x))) - A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, M_{\epsilon}\nabla u(x)) - \nabla u_{\epsilon}(x)) \, dx \\ &= \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u)), P_{\epsilon}(x, M_{\epsilon}\nabla u)) \, dx - \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u)), \nabla u_{\epsilon}) \, dx \\ & \quad - \int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}), P_{\epsilon}(x, M_{\epsilon}\nabla u)) \, dx + \int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}), \nabla u_{\epsilon}) \, dx \end{aligned}$$

goes to 0, as $\epsilon \rightarrow 0$. This is done in four steps.

In what follows, we use the following notation

$$\xi_{\epsilon}^k = \frac{1}{|Y_{\epsilon}^k|} \int_{Y_{\epsilon}^k} \nabla u \, dx.$$

Step 1. Let us prove that

$$\int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u)), P_{\epsilon}(x, M_{\epsilon}\nabla u)) \, dx \rightarrow \int_{\Omega} (b(\nabla u), \nabla u) \, dx \quad (4.1)$$

as $\epsilon \rightarrow 0$.

Proof. From (2.11) and (2.18), we obtain

$$\begin{aligned} \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u(x))), P_{\epsilon}(x, M_{\epsilon}\nabla u(x))) \, dx &= \sum_{k \in I_{\epsilon}} \int_{Y_{\epsilon}^k} \left(A\left(\frac{x}{\epsilon}, P\left(\frac{x}{\epsilon}, \xi_{\epsilon}^k\right)\right), P\left(\frac{x}{\epsilon}, \xi_{\epsilon}^k\right) \right) \, dx \\ &= \epsilon^n \sum_{k \in I_{\epsilon}} \int_Y (A(y, P(y, \xi_{\epsilon}^k)), P(y, \xi_{\epsilon}^k)) \, dy \\ &= \sum_{k \in I_{\epsilon}} \int_{\Omega} \chi_{Y_{\epsilon}^k}(x) (b(\xi_{\epsilon}^k), \xi_{\epsilon}^k) \, dx \\ &= \int_{\Omega} (b(M_{\epsilon}\nabla u(x)), M_{\epsilon}\nabla u(x)) \, dx. \end{aligned}$$

By (3.9) in Lemma 3.7, Hölder's inequality, Theorem 2.5, and Jensen's inequality, we have

$$\int_{\Omega} |b(M_{\epsilon}\nabla u(x)) - b(\nabla u(x))|^{q_1} \, dx \leq C \left[\left(\int_{\Omega} |M_{\epsilon}\nabla u - \nabla u|^{p_2} \, dx \right)^{\frac{p_1-1}{(p_2-1)(3-p_1)}} + \left(\int_{\Omega} |M_{\epsilon}\nabla u - \nabla u|^{p_2} \, dx \right)^{\frac{1}{3-p_2}} \right].$$

From Property 1 of M_{ϵ} , we obtain that

$$b(M_{\epsilon}\nabla u) \rightarrow b(\nabla u) \text{ in } L^{q_1}(\Omega; \mathbb{R}^n), \text{ as } \epsilon \rightarrow 0. \quad (4.2)$$

Now, (4.1) follows from (4.2) since $M_{\epsilon}\nabla u \rightarrow \nabla u$ in $L^{p_2}(\Omega; \mathbb{R}^n)$, so

$$\begin{aligned} \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u(x))), P_{\epsilon}(x, M_{\epsilon}\nabla u(x))) \, dx &= \int_{\Omega} (b(M_{\epsilon}\nabla u(x)), M_{\epsilon}\nabla u(x)) \, dx \\ &\rightarrow \int_{\Omega} (b(\nabla u(x)), \nabla u(x)) \, dx, \text{ as } \epsilon \rightarrow 0. \quad \square \end{aligned}$$

Step 2. We now show that

$$\int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u(x))), \nabla u_{\epsilon}(x)) \, dx \rightarrow \int_{\Omega} (b(\nabla u(x)), \nabla u(x)) \, dx \quad (4.3)$$

as $\epsilon \rightarrow 0$.

Proof. Let $\delta > 0$. From Theorem 2.5 we have $\nabla u \in L^{p_2}(\Omega; \mathbb{R}^n)$ and there exists a simple function Ψ satisfying the assumptions of Lemma 3.4 such that

$$\|\nabla u - \Psi\|_{L^{p_2}(\Omega; \mathbb{R}^n)} \leq \delta. \quad (4.4)$$

Let us write

$$\begin{aligned} \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u(x))), \nabla u_{\epsilon}(x)) \, dx &= \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, \Psi)), \nabla u_{\epsilon}) \, dx \\ &\quad + \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u)) - A_{\epsilon}(x, P_{\epsilon}(x, \Psi)), \nabla u_{\epsilon}) \, dx. \end{aligned}$$

We first show that

$$\int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) \, dx \rightarrow \int_{\Omega} (b(\Psi(x)), \nabla u(x)) \, dx \text{ as } \epsilon \rightarrow 0.$$

We have

$$\int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) \, dx = \sum_{j=0}^m \int_{\Omega_j} (A_{\epsilon}(x, P_{\epsilon}(x, \eta_j)), \nabla u_{\epsilon}(x)) \, dx.$$

Now from (2.17), we have that $A_{\epsilon}(\cdot, P_{\epsilon}(\cdot, \eta_j)) \rightarrow b(\eta_j) \in L^{q_2}(\Omega_j; \mathbb{R}^n)$, and by (2.10), $\int_{\Omega_j} (A_{\epsilon}(x, P_{\epsilon}(x, \eta_j)), \nabla \varphi(x)) \, dx = 0$, for $\varphi \in W_0^{1, p_1}(\Omega_j)$.

Take $\varphi = \delta u_{\epsilon}$, with $\delta \in C_0^{\infty}(\Omega_j)$ to get

$$0 = \int_{\Omega_j} (A_{\epsilon}(x, P_{\epsilon}(x, \eta_j)), (\nabla \delta)u_{\epsilon}) \, dx + \int_{\Omega_j} (A_{\epsilon}(x, P_{\epsilon}(x, \eta_j)), (\nabla u_{\epsilon})\delta) \, dx.$$

Taking the limit as $\epsilon \rightarrow 0$, and using the fact that $u_{\epsilon} \rightharpoonup u$ in $W_0^{1, p_1}(\Omega)$ and (2.17), we have by Lemma 3.9 that

$$\int_{\Omega_j} g_j(x)\delta(x) \, dx = \lim_{\epsilon \rightarrow 0} \int_{\Omega_j} (A_{\epsilon}(x, P_{\epsilon}(x, \eta_j)), (\nabla u_{\epsilon})\delta) \, dx = \int_{\Omega_j} (b(\eta_j), (\nabla u)\delta) \, dx.$$

Therefore, we may conclude that $g_j = (b(\eta_j), \nabla u)$, so

$$\sum_{j=0}^n \int_{\Omega_j} (A_{\epsilon}(x, P_{\epsilon}(x, \eta_j)), \nabla u_{\epsilon}(x)) \, dx \rightarrow \sum_{j=0}^n \int_{\Omega_j} (b(\eta_j), \nabla u(x)) \, dx, \text{ as } \epsilon \rightarrow 0.$$

Thus, we get

$$\int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) \, dx \rightarrow \int_{\Omega} (b(\Psi(x)), \nabla u(x)) \, dx, \text{ as } \epsilon \rightarrow 0.$$

On the other hand, let us estimate

$$\int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon}\nabla u(x))) - A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) \, dx.$$

By (2.1), Hölder's inequality, and (2.9) we obtain

$$\begin{aligned} & \left| \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon} \nabla u(x))) - A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) \, dx \right| \\ & \leq C \left[\left(\int_{\Omega} \chi_1^{\epsilon}(x) |P_{\epsilon}(x, M_{\epsilon} \nabla u) - P_{\epsilon}(x, \Psi)|^{p_1} \, dx \right)^{\frac{p_1-1}{p_1}} \right. \\ & \quad \left. + \left(\int_{\Omega} \chi_2^{\epsilon}(x) |P_{\epsilon}(x, M_{\epsilon} \nabla u) - P_{\epsilon}(x, \Psi)|^{p_2} \, dx \right)^{\frac{p_2-1}{p_2}} \right]. \end{aligned} \quad (4.5)$$

Applying Lemma 3.4 and (4.4) to (4.5), we discover that

$$\begin{aligned} & \limsup_{\epsilon \rightarrow 0} \left| \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon} \nabla u(x))) - A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) \, dx \right| \\ & \leq C \sum_{i=1}^2 \left[\delta^{\frac{p_1}{3-p_1}} + \delta^{\frac{p_1 p_2}{2p_2 - p_2 p_1 + p_1}} + \delta^{\frac{p_1 p_2}{2p_1 - p_2 p_1 + p_2}} + \delta^{\frac{p_2}{3-p_2}} \right]^{\frac{p_i-1}{p_i}}, \end{aligned} \quad (4.6)$$

where C is independent of δ . Since δ is arbitrary we conclude that the limit on the left hand side of (4.6) is equal to 0.

Finally, using the continuity of b (3.9) in Lemma 3.7, Theorem 2.5, and Hölder's inequality, we obtain

$$\left| \int_{\Omega} (b(\nabla u(x)) - b(\Psi(x)), \nabla u(x)) \, dx \right| \leq C \left[\delta^{\frac{p_2(p_1-1)}{(p_2-1)(3-p_1)}} + \delta^{\frac{p_2}{3-p_2}} \right]^{\frac{1}{q_1}},$$

where C does not depend on δ .

Step 2 is proved noticing that δ can be taken arbitrarily small. \square

Step 3. We will show that

$$\int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, M_{\epsilon} \nabla u(x))) \, dx \rightarrow \int_{\Omega} (b(\nabla u(x)), \nabla u(x)) \, dx \quad (4.7)$$

as $\epsilon \rightarrow 0$.

Proof. Let $\delta > 0$. As in the proof of Step 2, assume Ψ is a simple function satisfying assumptions of Lemma 3.4 and such that $\|\nabla u - \Psi\|_{L^{p_2}(\Omega; \mathbb{R}^n)} < \delta$.

Let us write

$$\begin{aligned} \int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, M_{\epsilon} \nabla u(x))) \, dx &= \int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, \Psi(x))) \, dx \\ &\quad + \int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, M_{\epsilon} \nabla u(x)) - P_{\epsilon}(x, \Psi(x))) \, dx. \end{aligned}$$

We first show that

$$\int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, \Psi(x))) \, dx \rightarrow \int_{\Omega} (b(\nabla u(x)), \Psi(x)) \, dx.$$

We start by writing

$$\int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, \Psi(x))) \, dx = \sum_{j=0}^m \int_{\Omega_j} (A_{\epsilon}(x, \nabla u_{\epsilon}(x)), P_{\epsilon}(x, \eta_j)) \, dx.$$

From Lemma 3.9, up to a subsequence, $(A_\epsilon(\cdot, \nabla u_\epsilon), P_\epsilon(\cdot, \eta_j))$ converges weakly to a function $h_j \in L^1(\Omega_j; \mathbb{R})$, as $\epsilon \rightarrow 0$.

By Theorem 2.1, we have $A_\epsilon(\cdot, \nabla u_\epsilon) \rightharpoonup b(\nabla u) \in L^{q_2}(\Omega; \mathbb{R}^n)$ and

$$-\operatorname{div}(A_\epsilon(x, \nabla u_\epsilon)) = f = -\operatorname{div}(b(\nabla u)).$$

From (2.15), p_ϵ satisfies $P_\epsilon(\cdot, \eta_j) \rightharpoonup \eta_j$ in $L^{p_1}(\Omega_j, \mathbb{R}^n)$.

Arguing as in Step 2, we find that $(A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, \eta_j)) \rightharpoonup (b(\nabla u(x)), \eta_j)$ in $D'(\Omega_j)$, as $\epsilon \rightarrow 0$.

Therefore, we may conclude that $h_j = (b(\nabla u), \eta_j)$, and hence,

$$\sum_{j=0}^n \int_{\Omega_j} (A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, \eta_j)) \, dx \rightarrow \sum_{j=0}^n \int_{\Omega_j} (b(\nabla u(x)), \eta_j) \, dx, \text{ as } \epsilon \rightarrow 0.$$

Thus, we get

$$\int_{\Omega} (A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, \Psi(x))) \, dx \rightarrow \int_{\Omega} (b(\nabla u(x)), \Psi(x)) \, dx, \text{ as } \epsilon \rightarrow 0.$$

Moreover, applying Hölder's inequality, (2.1), and (2.9) we have

$$\left| \int_{\Omega} (A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - P_\epsilon(x, \Psi(x))) \, dx \right| \leq C \sum_{i=1}^2 \left(\int_{\Omega} \chi_i^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u) - P_\epsilon(x, \Psi)|^{p_i} \, dx \right)^{\frac{1}{p_i}}.$$

As in the proof of Step 2 we see that

$$\limsup_{\epsilon \rightarrow 0} \left| \int_{\Omega} (A_\epsilon(x, \nabla u_\epsilon), P_\epsilon(x, M_\epsilon \nabla u) - P_\epsilon(x, \Psi)) \, dx \right| \leq C \sum_{i=1}^2 \left[\delta^{\frac{p_1}{3-p_1}} + \delta^{\frac{p_1 p_2}{2p_2 - p_2 p_1 + p_1}} + \delta^{\frac{p_1 p_2}{2p_1 - p_2 p_1 + p_2}} + \delta^{\frac{p_2}{3-p_2}} \right]^{\frac{1}{p_i}},$$

where C does not depend on δ .

Hence, proceeding as in Step 2, we find that

$$\begin{aligned} & \limsup_{\epsilon \rightarrow 0} \left| \int_{\Omega} (A_\epsilon(x, \nabla u_\epsilon), P_\epsilon(x, M_\epsilon \nabla u)) \, dx - \int_{\Omega} (b(\nabla u), \nabla u) \, dx \right| \\ & \leq C \left[\sum_{i=1}^2 \left(\delta^{\frac{p_1}{3-p_1}} + \delta^{\frac{p_1 p_2}{2p_2 - p_2 p_1 + p_1}} + \delta^{\frac{p_1 p_2}{2p_1 - p_2 p_1 + p_2}} + \delta^{\frac{p_2}{3-p_2}} \right)^{\frac{1}{p_i}} + \delta \|b(\nabla u)\|_{L^{q_2}(\Omega, \mathbb{R}^n)} \right], \end{aligned}$$

where C is independent of δ . Now since δ is arbitrarily small, the proof of Step 3 is complete. □

Step 4. Finally, let us prove that

$$\int_{\Omega} (A_\epsilon(x, \nabla u_\epsilon(x)), \nabla u_\epsilon(x)) \, dx \rightarrow \int_{\Omega} (b(\nabla u(x)), \nabla u(x)) \, dx, \text{ as } \epsilon \rightarrow 0. \tag{4.8}$$

Proof. Since

$$\int_{\Omega} (A_\epsilon(x, \nabla u_\epsilon), \nabla u_\epsilon) \, dx = \langle -\operatorname{div}(A_\epsilon(x, \nabla u_\epsilon)), u_\epsilon \rangle = \langle f, u_\epsilon \rangle, \tag{4.9}$$

$$\int_{\Omega} (b(\nabla u), \nabla u) \, dx = \langle -\operatorname{div}(b(\nabla u)), u \rangle = \langle f, u \rangle, \tag{4.10}$$

and $u_\epsilon \rightharpoonup u$ in $W^{1,p_1}(\Omega)$, the result follows immediately. □

Finally, Theorem 2.7 follows from (4.1), (4.3), (4.7) and (4.8). □

4.2. Proof of the lower bound on the amplification of the macroscopic field by the microstructure

The sequence $\{\chi_i^\epsilon(x)\nabla u_\epsilon(x)\}_{\epsilon>0}$ has a Young measure $\nu^i = \{\nu_x^i\}_{x\in\Omega}$ associated to it (see Thm. 6.2 and the discussion following in [21]), for $i = 1, 2$.

As a consequence of Theorem 2.7 proved in the previous section, we have that

$$\left\| \chi_i^\epsilon(x)P\left(\frac{x}{\epsilon}, M_\epsilon(\nabla u)(x)\right) - \chi_i^\epsilon(x)\nabla u_\epsilon(x) \right\|_{\mathbf{L}^{p_i}(\Omega; \mathbb{R}^n)} \rightarrow 0,$$

as $\epsilon \rightarrow 0$, which implies that the sequences

$$\left\{ \chi_i^\epsilon(x)P\left(\frac{x}{\epsilon}, M_\epsilon(\nabla u)(x)\right) \right\}_{\epsilon>0} \quad \text{and} \quad \{ \chi_i^\epsilon(x)\nabla u_\epsilon(x) \}_{\epsilon>0}$$

share the same Young measure (see Lem. 6.3 of [21]), for $i = 1, 2$.

The next lemma identifies the Young measure ν^i . The lemma is proven for $1 < p_1 \leq p_2 \leq 2$; the proof for the case when $1 < p_1 \leq 2 \leq p_2$ follows in a similar way.

Lemma 4.1. *For all $\phi \in C_0(\mathbb{R}^n)$ and for all $\zeta \in C_0^\infty(\mathbb{R}^n)$, we have*

$$\int_{\Omega} \zeta(x) \int_{\mathbb{R}^n} \phi(\lambda) d\nu_x^i(\lambda) dx = \int_{\Omega} \zeta(x) \int_Y \phi(\chi_i(y)P(y, \nabla u(x))) dy dx. \quad (4.11)$$

Proof. To prove (4.11), we will show that given $\phi \in C_0(\mathbb{R}^n)$ and $\zeta \in C_0^\infty(\mathbb{R}^n)$,

$$\lim_{\epsilon \rightarrow 0} \int_{\Omega} \zeta(x) \phi\left(\chi_i^\epsilon(x)P\left(\frac{x}{\epsilon}, M_\epsilon(\nabla u)(x)\right)\right) dx = \int_{\Omega} \zeta(x) \int_Y \phi(\chi_i(y)P(y, \nabla u(x))) dy dx. \quad (4.12)$$

We consider the difference

$$\begin{aligned} & \left| \int_{\Omega} \zeta(x) \phi\left(\chi_i\left(\frac{x}{\epsilon}\right)P\left(\frac{x}{\epsilon}, M_\epsilon(\nabla u)(x)\right)\right) dx - \int_{\Omega} \zeta(x) \int_Y \phi(\chi_i(y)P(y, \nabla u(x))) dy dx \right| \\ & \leq \left| \sum_{k \in I_\epsilon} \int_{Y_\epsilon^k} \zeta(x) \phi\left(\chi_i\left(\frac{x}{\epsilon}\right)P\left(\frac{x}{\epsilon}, \xi_\epsilon^k\right)\right) dx - \int_{\Omega_\epsilon} \zeta(x) \int_Y \phi(\chi_i(y)P(y, \nabla u(x))) dy dx \right| \\ & \quad + C|\Omega \setminus \Omega_\epsilon|. \end{aligned} \quad (4.13)$$

Note that the term $C|\Omega \setminus \Omega_\epsilon|$ goes to 0, as $\epsilon \rightarrow 0$. Now set x_ϵ^k to be the center of Y_ϵ^k . On the first integral use the change of variables $x = x_\epsilon^k + \epsilon y$, where y belongs to Y , and since $dx = \epsilon^n dy$, we get

$$\begin{aligned} & \left| \sum_{k \in I_\epsilon} \int_{Y_\epsilon^k} \zeta(x) \phi\left(\chi_i\left(\frac{x}{\epsilon}\right)P\left(\frac{x}{\epsilon}, \xi_\epsilon^k\right)\right) dx - \sum_{k \in I_\epsilon} \int_{Y_\epsilon^k} \zeta(x) \int_Y \phi(\chi_i(y)P(y, \nabla u(x))) dy dx \right| \\ & = \left| \sum_{k \in I_\epsilon} \epsilon^n \int_Y \zeta(x_\epsilon^k + \epsilon y) \phi(\chi_i(y)P(y, \xi_\epsilon^k)) dy - \sum_{k \in I_\epsilon} \int_{Y_\epsilon^k} \zeta(x) \int_Y \phi(\chi_i(y)P(y, \nabla u(x))) dy dx \right|. \end{aligned}$$

Applying Taylor's expansion for ζ , we have

$$\begin{aligned} & \leq \left| \sum_{k \in I_\epsilon} \int_{Y_\epsilon^k} \int_Y (\zeta(x) + CO(\epsilon)) [\phi(\chi_i(y)P(y, \xi_\epsilon^k)) - \phi(\chi_i(y)P(y, \nabla u(x)))] dy dx \right| + CO(\epsilon) \\ & \leq \left| \int_{\Omega_\epsilon} |\zeta(x)| \int_Y |\phi(\chi_i(y)P(y, M_\epsilon \nabla u(x))) - \phi(\chi_i(y)P(y, \nabla u(x)))| dy dx \right| + CO(\epsilon). \end{aligned}$$

Because of the uniform Lipschitz continuity of ϕ , we get

$$\leq C \left| \int_{\Omega_\epsilon} |\zeta(x)| \int_Y |P(y, M_\epsilon \nabla u(x)) - P(y, \nabla u(x))| dy dx \right| + CO(\epsilon).$$

By Hölder's inequality twice and Lemma 3.2, we have

$$\begin{aligned} &\leq C \left\{ \left[\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_1}{3-p_1}} \right. \right. \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{2-p_1}{3-p_1}} dx \\ &\quad + \int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_1 p_2}{2p_2 - p_1 p_2 + p_1}} \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{p_2(2-p_1)}{2p_2 - p_1 p_2 + p_1}} dx \\ &\quad + \int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_1 p_2}{2p_1 - p_1 p_2 + p_2}} \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{p_1(2-p_2)}{2p_1 - p_1 p_2 + p_2}} dx \\ &\quad \left. + \int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_2}{3-p_2}} \right. \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{2-p_2}{3-p_2}} dx \left. \right]^{1/p_1} \\ &\quad + \left[\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_1}{3-p_1}} \right. \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{2-p_1}{3-p_1}} dx \\ &\quad + \int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_1 p_2}{2p_2 - p_1 p_2 + p_1}} \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{p_2(2-p_1)}{2p_2 - p_1 p_2 + p_1}} dx \\ &\quad + \int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_1 p_2}{2p_1 - p_1 p_2 + p_2}} \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{p_1(2-p_2)}{2p_1 - p_1 p_2 + p_2}} dx \\ &\quad \left. + \int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{\frac{p_2}{3-p_2}} \right. \\ &\quad \times (1 + |M_\epsilon \nabla u(x)|^{p_1} + |M_\epsilon \nabla u(x)|^{p_2} + |\nabla u(x)|^{p_1} + |\nabla u(x)|^{p_2})^{\frac{2-p_2}{3-p_2}} dx \left. \right]^{1/p_2} \Big\} \\ &\quad + CO(\epsilon). \end{aligned}$$

Applying Hölder's inequality, Jensen's inequality and Theorem 2.5, we get

$$\begin{aligned}
&\leq C \left\{ \left[\left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_1} dx \right)^{\frac{1}{3-p_1}} + \left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_2} dx \right)^{\frac{p_1}{2p_2-p_1p_2+p_1}} \right. \right. \\
&\quad + \left. \left. \left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_1} dx \right)^{\frac{p_2}{2p_1-p_1p_2+p_2}} \left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_2} dx \right)^{\frac{1}{3-p_2}} \right]^{1/p_1} \right. \\
&\quad + \left[\left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_1} dx \right)^{\frac{1}{3-p_1}} + \left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_2} dx \right)^{\frac{p_1}{2p_2-p_1p_2+p_1}} \right. \\
&\quad + \left. \left. \left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_1} dx \right)^{\frac{p_2}{2p_1-p_1p_2+p_2}} + \left(\int_{\Omega_\epsilon} |M_\epsilon \nabla u(x) - \nabla u(x)|^{p_2} dx \right)^{\frac{1}{3-p_2}} \right]^{1/p_2} \right\} \\
&\quad + CO(\epsilon).
\end{aligned}$$

Finally, from the approximation property of M_ϵ in Section 2.2.1, as $\epsilon \rightarrow 0$, we obtain (4.12). \square

Lemma 4.2 (see Thm. 6.11 in [21]). *The sequence $\{\chi_i^\epsilon(x)\nabla u_\epsilon(x)\}_{\epsilon>0}$ and the Young measure $\nu^i = \{\nu_x^i\}_{x \in \Omega}$ associated to it satisfy*

$$\int_D \int_{\mathbb{R}^n} \psi(x, \lambda) d\nu_x^i(\lambda) dx \leq \liminf_{\epsilon \rightarrow 0} \int_D \psi(x, \chi_i^\epsilon(x)\nabla u_\epsilon(x)) dx,$$

for all Carathéodory functions $\psi \geq 0$ and measurable subset $D \subset \Omega$.

Using Lemmas 4.1 and 4.2, we obtain

$$\begin{aligned}
\lim_{\epsilon \rightarrow 0} \int_{\Omega} \zeta(x) \phi \left(\chi_i^\epsilon(x) P \left(\frac{x}{\epsilon}, M_\epsilon(\nabla u)(x) \right) \right) dx &= \int_{\Omega} \zeta(x) \int_Y \phi(\chi_i(y) P(y, \nabla u(x))) dy dx \\
&= \int_{\Omega} \zeta(x) \int_{\mathbb{R}^n} \phi(\lambda) d\nu_x^i(\lambda) dx \\
&\leq \lim_{\epsilon \rightarrow 0} \int_{\Omega} \zeta(x) \phi(\chi_i^\epsilon(x)\nabla u_\epsilon(x)) dx,
\end{aligned}$$

for all $\phi \in C_0(\mathbb{R}^n)$ and for all $\zeta \in C_0^\infty(\mathbb{R}^n)$, which concludes the proof of Theorem 2.8.

APPENDIX A

The proofs presented here are for the case when $1 < p_1 \leq p_2 \leq 2$. The proofs for the lemmas for the case when $1 < p_1 \leq 2 \leq p_2$ follow in a similar way. The letter C will represent a generic positive constant independent of ϵ , and it can take different values from one line to the other.

A.1. Proof of Lemma 3.1

Let $\xi \in \mathbb{R}^n$. By (2.2) we have that

$$(A(y, P(y, \xi)), P(y, \xi)) \geq C(\chi_1(y) |P(y, \xi)|^{p_1} + \chi_2(y) |P(y, \xi)|^{p_2}).$$

Integrating both sides over Y , using (2.1), and Young's inequality, we get

$$\int_Y \chi_1(y) |P(y, \xi)|^{p_1} dy + \int_Y \chi_2(y) |P(y, \xi)|^{p_2} dy \leq C \left[\frac{\delta^{q_2} \int_Y \chi_1(y) |P(y, \xi)|^{p_1} dy}{q_2} + \frac{\delta^{-p_1} \int_Y \chi_1(y) |\xi|^{p_1} dy}{p_1} \right. \\ \left. + \frac{\delta^{q_1} \int_Y \chi_2(y) |P(y, \xi)|^{p_2} dy}{q_1} + \frac{\delta^{-p_2} \int_Y \chi_2(y) |\xi|^{p_2} dy}{p_2} \right].$$

Doing some algebraic manipulations, we obtain

$$(1 - C\delta^{q_2}) \int_Y \chi_1(y) |P(y, \xi)|^{p_1} dy + (1 - C\delta^{q_1}) \int_Y \chi_2(y) |P(y, \xi)|^{p_2} dy \leq C [(\delta^{-p_1} |\xi|^{p_1} \theta_1 + \delta^{-p_2} |\xi|^{p_2} \theta_2)].$$

On choosing an appropriate δ , we finally obtain (3.1).

A.2. Proof of Lemma 3.2

Using Hölder's inequality, we have

$$\int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1} dy + \int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2} dy \\ \leq C \sum_{i=1}^2 \left[\left(\int_Y \frac{\chi_i(y) |P(y, \xi_1) - P(y, \xi_2)|^2}{(1 + |P(y, \xi_1)| + |P(y, \xi_2)|)^{2-p_i}} dy \right)^{\frac{p_i}{2}} \right. \\ \left. \times \left(\int_Y \chi_i(y) (1 + |P(y, \xi_1)| + |P(y, \xi_2)|)^{p_i} dy \right)^{\frac{2-p_i}{2}} \right].$$

By Lemma 3.1, (2.2), and (2.8) we get

$$\leq C \sum_{i=1}^2 \left[\left(\int_Y |A(y, P(y, \xi_1)) - A(y, P(y, \xi_2))| |\xi_1 - \xi_2| dy \right)^{\frac{p_i}{2}} \right. \\ \left. \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_i}{2}} \right].$$

By (2.1) and Hölder's inequality, we have

$$\leq C \left[\left(\int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1} dy \right)^{\frac{p_1-1}{2}} \theta_1^{\frac{1}{2}} |\xi_1 - \xi_2|^{\frac{p_1}{2}} \right. \\ \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_1}{2}} \\ + \left(\int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2} dy \right)^{\frac{p_1(p_2-1)}{2p_2}} \theta_2^{\frac{p_1}{2p_2}} |\xi_1 - \xi_2|^{\frac{p_1}{2}} \\ \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_1}{2}} \\ \left. + \left(\int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1} dy \right)^{\frac{p_2(p_1-1)}{2p_1}} \theta_1^{\frac{p_2}{2p_1}} |\xi_1 - \xi_2|^{\frac{p_2}{2}} \right]$$

$$\begin{aligned} & \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_2}{2}} \\ & + \left(\int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2} dy \right)^{\frac{p_2-1}{2}} \theta_2^{\frac{1}{2}} |\xi_1 - \xi_2|^{\frac{p_2}{2}} \\ & \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_2}{2}} \Big]. \end{aligned}$$

Applying Young's inequality, with $r_1 = \frac{2}{p_1-1}$, $r_2 = \frac{2p_2}{p_1(p_2-1)}$, $r_3 = \frac{2p_1}{p_2(p_1-1)}$, and $r_4 = \frac{2}{p_2-1}$.

$$\begin{aligned} & \leq C \left[\frac{\delta^{\frac{2}{p_1-1}} \int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1} dy}{\frac{2}{p_1-1}} \right. \\ & + \frac{\delta^{\frac{-2}{3-p_1}} \theta_1^{\frac{1}{3-p_1}} |\xi_1 - \xi_2|^{\frac{p_1}{3-p_1}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_1}{3-p_1}}}{\frac{2}{3-p_1}} \\ & + \frac{\delta^{\frac{2p_2}{p_1(p_2-1)}} \int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2} dy}{\frac{2p_2}{p_1(p_2-1)}} \\ & + \left(\frac{\delta^{\frac{-2p_2}{2p_2-p_1p_2+p_1}} \theta_2^{\frac{p_1}{2p_2-p_1p_2+p_1}} |\xi_1 - \xi_2|^{\frac{p_1p_2}{2p_2-p_1p_2+p_1}}}{\frac{2p_2}{2p_2-p_1p_2+p_1}} \right) \\ & \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{p_2(2-p_1)}{2p_2-p_1p_2+p_1}} \\ & + \frac{\delta^{\frac{2p_1}{p_2(p_1-1)}} \int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1} dy}{\frac{2p_1}{p_2(p_1-1)}} \\ & + \left(\frac{\delta^{\frac{-2p_1}{2p_1-p_1p_2+p_2}} \theta_1^{\frac{p_2}{2p_1-p_1p_2+p_2}} |\xi_1 - \xi_2|^{\frac{p_1p_2}{2p_1-p_1p_2+p_2}}}{\frac{2p_1}{2p_1-p_1p_2+p_2}} \right) \\ & \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{p_1(2-p_2)}{2p_1-p_1p_2+p_2}} \\ & + \frac{\delta^{\frac{2}{p_2-1}} \int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2} dy}{\frac{2}{p_2-1}} \\ & + \left. \frac{\delta^{\frac{-2}{3-p_2}} \theta_2^{\frac{1}{3-p_2}} |\xi_1 - \xi_2|^{\frac{p_2}{3-p_2}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_2}{3-p_2}}}{\frac{2}{3-p_2}} \right]. \end{aligned}$$

Straightforward algebraic manipulation delivers

$$\begin{aligned} & \left(1 - \frac{C\delta^{\frac{2}{p_1-1}}}{\frac{2}{p_1-1}} - \frac{C\delta^{\frac{2p_1}{p_2(p_1-1)}}}{\frac{2p_1}{p_2(p_1-1)}} \right) \int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1} dy \\ & + \left(1 - \frac{C\delta^{\frac{2}{p_2-1}}}{\frac{2}{p_2-1}} - \frac{C\delta^{\frac{2p_2}{p_1(p_2-1)}}}{\frac{2p_2}{p_1(p_2-1)}} \right) \int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2} dy \end{aligned}$$

$$\begin{aligned} &\leq C \left[\frac{\delta^{\frac{-2}{3-p_1}} \theta_1^{\frac{1}{3-p_1}} |\xi_1 - \xi_2|^{\frac{p_1}{3-p_1}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_1}{3-p_1}}}{\frac{2}{3-p_1}} \right. \\ &\quad + \left(\frac{\delta^{\frac{-2p_2}{2p_2-p_1p_2+p_1}} \theta_2^{\frac{p_1}{2p_2-p_1p_2+p_1}} |\xi_1 - \xi_2|^{\frac{p_1p_2}{2p_2-p_1p_2+p_1}}}{\frac{2p_2}{2p_2-p_1p_2+p_1}} \right) \\ &\quad \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{p_2(2-p_1)}{2p_2-p_1p_2+p_1}} \\ &\quad + \left(\frac{\delta^{\frac{-2p_1}{2p_1-p_1p_2+p_2}} \theta_1^{\frac{p_2}{2p_1-p_1p_2+p_2}} |\xi_1 - \xi_2|^{\frac{p_1p_2}{2p_1-p_1p_2+p_2}}}{\frac{2p_1}{2p_1-p_1p_2+p_2}} \right) \\ &\quad \times (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{p_1(2-p_2)}{2p_1-p_1p_2+p_2}} \\ &\quad \left. + \frac{\delta^{\frac{-2}{3-p_2}} \theta_2^{\frac{1}{3-p_2}} |\xi_1 - \xi_2|^{\frac{p_2}{3-p_2}} (1 + |\xi_1|^{p_1} \theta_1 + |\xi_2|^{p_1} \theta_1 + |\xi_1|^{p_2} \theta_2 + |\xi_2|^{p_2} \theta_2)^{\frac{2-p_2}{3-p_2}}}{\frac{2}{3-p_2}} \right]. \end{aligned}$$

The result follows on choosing an appropriate δ and doing a change of variables.

A.3. Proof of Lemma 3.4

Let Ψ of the form (3.5). For every $\epsilon > 0$, let us denote by $\Omega_\epsilon = \bigcup_{k \in I_\epsilon} \overline{Y_\epsilon^k}$; and for $j = 0, 1, 2, \dots, m$, we set

$$I_\epsilon^j = \{k \in I_\epsilon : Y_\epsilon^k \subseteq \Omega_j\}, \text{ and } J_\epsilon^j = \{k \in I_\epsilon : Y_\epsilon^k \cap \Omega_j \neq \emptyset, Y_\epsilon^k \setminus \Omega_j \neq \emptyset\}.$$

Furthermore, $E_\epsilon^j = \bigcup_{k \in I_\epsilon^j} \overline{Y_\epsilon^k}$, $F_\epsilon^j = \bigcup_{k \in J_\epsilon^j} \overline{Y_\epsilon^k}$, and as $\epsilon \rightarrow 0$, we have $|F_\epsilon^j| \rightarrow 0$.

Set

$$\xi_\epsilon^k = \frac{1}{|Y_\epsilon^k|} \int_{Y_\epsilon^k} \varphi(y) dy.$$

For ϵ sufficiently small Ω_j ($j \neq 0$) is contained in Ω_ϵ .

From (3.5), (2.18), using the fact that $\Omega_j \subset E_\epsilon^j \cup F_\epsilon^j$, Lemma 3.2, and Hölder’s inequality it follows that

$$\begin{aligned} &\int_\Omega \chi_1^\epsilon(x) |P_\epsilon(x, M_\epsilon \varphi) - P_\epsilon(x, \Psi)|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |P_\epsilon(x, M_\epsilon \varphi) - P_\epsilon(x, \Psi)|^{p_2} dx \\ &\leq C \left[\left(\int_\Omega \chi_1^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_1} dx + \int_\Omega \chi_1^\epsilon(x) |\varphi - \Psi|^{p_1} dx \right)^{\frac{1}{3-p_1}} \right. \\ &\quad \times \left(\int_\Omega \chi_1^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_1} dx + \int_\Omega \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_2} dx \right. \\ &\quad \left. \left. + \int_\Omega \chi_2^\epsilon(x) |\varphi|^{p_2} dx + \int_\Omega \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_\Omega \chi_2^\epsilon(x) |\Psi|^{p_2} dx + |\Omega| \right)^{\frac{2-p_1}{3-p_1}} \right] \tag{A.1} \end{aligned}$$

$$\begin{aligned}
& + \left(\int_{\Omega} \chi_1^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_1} dx + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_2} dx \right. \\
& + \left. \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx + |\Omega| \right)^{\frac{p_2(2-p_1)}{p_2-p_1 p_2+p_1}} \\
& \times \left(\int_{\Omega} \chi_2^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_2} dx + \int_{\Omega} \chi_2^\epsilon(x) |\varphi - \Psi|^{p_2} dx \right)^{\frac{p_1}{p_2-p_1 p_2+p_1}} \\
& + \left(\int_{\Omega} \chi_1^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_1} dx + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_2} dx \right. \\
& + \left. \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx + |\Omega| \right)^{\frac{p_1(2-p_2)}{p_1-p_1 p_2+p_2}} \\
& \times \left(\int_{\Omega} \chi_1^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_1} dx + \int_{\Omega} \chi_1^\epsilon(x) |\varphi - \Psi|^{p_1} dx \right)^{\frac{p_2}{p_1-p_1 p_2+p_2}} \\
& + \left(\int_{\Omega} \chi_1^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_1} dx + \int_{\Omega} \chi_1^\epsilon(x) |\varphi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_2} dx \right. \\
& + \left. \int_{\Omega} \chi_2^\epsilon(x) |\varphi|^{p_2} dx + \int_{\Omega} \chi_1^\epsilon(x) |\Psi|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\Psi|^{p_2} dx + |\Omega| \right)^{\frac{2-p_2}{3-p_2}} \\
& \times \left. \left(\int_{\Omega} \chi_2^\epsilon(x) |M_\epsilon \varphi - \varphi|^{p_2} dx + \int_{\Omega} \chi_2^\epsilon(x) |\varphi - \Psi|^{p_2} dx \right)^{\frac{1}{3-p_2}} \right] \\
& + C \sum_{j=0}^m \left[\left(\int_{F_\epsilon^j} \theta_1 \left| \sum_{k \in J_\epsilon^j} \chi_{Y_\epsilon^k} \xi_\epsilon^k - \eta_j \right|^{p_1} dx \right)^{\frac{1}{3-p_1}} \left(|F_\epsilon^j| + \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_1} \theta_1 dx \right. \right. \\
& + \left. \left. \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_2} \theta_2 dx + |\eta_j|^{p_1} \theta_1 |F_\epsilon^j| + |\eta_j|^{p_2} \theta_2 |F_\epsilon^j| \right)^{\frac{2-p_1}{3-p_1}} \right. \\
& + \left. \left(\int_{F_\epsilon^j} \theta_2 \left| \sum_{k \in J_\epsilon^j} \chi_{Y_\epsilon^k} \xi_\epsilon^k - \eta_j \right|^{p_2} dx \right)^{\frac{p_1}{2p_2-p_1 p_2+p_1}} \left(|F_\epsilon^j| + \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_1} \theta_1 dx \right. \right. \\
& + \left. \left. \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_2} \theta_2 dx + |\eta_j|^{p_1} \theta_1 |F_\epsilon^j| + |\eta_j|^{p_2} \theta_2 |F_\epsilon^j| \right)^{\frac{p_2(2-p_1)}{2p_2-p_1 p_2+p_1}} \right. \\
& + \left. \left(\int_{F_\epsilon^j} \theta_1 \left| \sum_{k \in J_\epsilon^j} \chi_{Y_\epsilon^k} \xi_\epsilon^k - \eta_j \right|^{p_1} dx \right)^{\frac{p_2}{2p_1-p_1 p_2+p_2}} \left(|F_\epsilon^j| + \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_1} \theta_1 dx \right. \right. \\
& + \left. \left. \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_2} \theta_2 dx + |\eta_j|^{p_1} \theta_1 |F_\epsilon^j| + |\eta_j|^{p_2} \theta_2 |F_\epsilon^j| \right)^{\frac{p_1(2-p_2)}{2p_1-p_1 p_2+p_2}} \right. \\
& + \left. \left(\int_{F_\epsilon^j} \theta_2 \left| \sum_{k \in J_\epsilon^j} \chi_{Y_\epsilon^k} \xi_\epsilon^k - \eta_j \right|^{p_2} dx \right)^{\frac{1}{3-p_2}} \left(|F_\epsilon^j| + \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_1} \theta_1 dx \right. \right. \\
& + \left. \left. \int_{F_\epsilon^j} |M_\epsilon \varphi|^{p_2} \theta_2 dx + |\eta_j|^{p_1} \theta_1 |F_\epsilon^j| + |\eta_j|^{p_2} \theta_2 |F_\epsilon^j| \right)^{\frac{2-p_2}{3-p_2}} \right].
\end{aligned}$$

Since $|\partial\Omega_j| = 0$ for $j \neq 0$, we have that $|F_\epsilon^j| \rightarrow 0$ as $\epsilon \rightarrow 0$, for every $j = 0, 1, 2, \dots, m$. By Property (1) of M_ϵ mentioned in Section 2.2.1, we have

$$\int_{\Omega} \chi_i^\epsilon(x) |M_\epsilon \varphi(x) - \varphi(x)|^{p_i} dx \rightarrow 0, \text{ as } \epsilon \rightarrow 0, \text{ for } i = 1, 2.$$

Therefore, taking lim sup as $\epsilon \rightarrow 0$ in (A.1), we obtain (3.6).

A.4. Proof of Lemma 3.5

Evaluating u_ϵ in the weak formulation for (2.3), applying Hölder's inequality, and since $f \in W^{-1, q_2}(\Omega)$, we obtain

$$\begin{aligned} s \int_{\Omega} (A_\epsilon(x, \nabla u_\epsilon), \nabla u_\epsilon) dx &= \sigma_1 \int_{\Omega} \chi_1^\epsilon(x) |\nabla u_\epsilon|^{p_1} dx + \sigma_2 \int_{\Omega} \chi_2^\epsilon(x) |\nabla u_\epsilon|^{p_2} dx \\ &= \langle f, u_\epsilon \rangle \leq C \left[\left(\int_{\Omega} \chi_1^\epsilon(x) |\nabla u_\epsilon|^{p_1} dx \right)^{\frac{1}{p_1}} + \left(\int_{\Omega} \chi_2^\epsilon(x) |\nabla u_\epsilon|^{p_2} dx \right)^{\frac{1}{p_2}} \right]. \end{aligned} \quad (\text{A.2})$$

Applying Young's inequality to the last term in (A.2), we obtain

$$\begin{aligned} &\sigma_1 \int_{\Omega} \chi_1^\epsilon(x) |\nabla u_\epsilon|^{p_1} dx + \sigma_2 \int_{\Omega} \chi_2^\epsilon(x) |\nabla u_\epsilon|^{p_2} dx \\ &\leq C \left[\frac{\delta^{p_1}}{p_1} \int_{\Omega} \chi_1^\epsilon(x) |\nabla u_\epsilon|^{p_1} dx + \frac{\delta^{-q_2}}{q_2} + \frac{\delta^{p_2}}{p_2} \int_{\Omega} \chi_2^\epsilon(x) |\nabla u_\epsilon|^{p_2} dx + \frac{\delta^{-q_1}}{q_1} \right]. \end{aligned} \quad (\text{A.3})$$

By rearranging the terms in (A.3), one gets

$$\left(\sigma_1 - C \frac{\delta^{p_1}}{p_1} \right) \int_{\Omega} \chi_1^\epsilon(x) |\nabla u_\epsilon|^{p_1} dx + \left(\sigma_2 - C \frac{\delta^{p_2}}{p_2} \right) \int_{\Omega} \chi_2^\epsilon(x) |\nabla u_\epsilon|^{p_2} dx \leq \frac{\delta^{-q_2}}{q_2} + \frac{\delta^{-q_1}}{q_1}.$$

Therefore, by choosing δ small enough so that $\min \left\{ \sigma_1 - C \frac{\delta^{p_1}}{p_1}, \sigma_2 - C \frac{\delta^{p_2}}{p_2} \right\}$ is positive, one obtains

$$\int_{\Omega} \chi_1^\epsilon(x) |\nabla u_\epsilon(x)|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |\nabla u_\epsilon(x)|^{p_2} dx \leq C.$$

A.5. Proof of Lemma 3.6

Using (2.18), we have

$$\begin{aligned} &\int_{\Omega} \chi_1^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x))|^{p_1} dx + \int_{\Omega} \chi_2^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x))|^{p_2} dx \\ &= \sum_{k \in \mathbf{I}_\epsilon} \left[\int_{Y_\epsilon^k} \chi_1^\epsilon(x) |P_\epsilon(x, \xi_\epsilon^k)|^{p_1} dx + \int_{Y_\epsilon^k} \chi_2^\epsilon(x) |P_\epsilon(x, \xi_\epsilon^k)|^{p_2} dx \right] \\ &\leq C \sum_{k \in \mathbf{I}_\epsilon} \left(|Y_\epsilon^k| + |\xi_\epsilon^k|^{p_1} \theta_1 |Y_\epsilon^k| + |\xi_\epsilon^k|^{p_2} \theta_2 |Y_\epsilon^k| \right) \\ &\leq C \left(|\Omega| + \|\nabla u\|_{\mathbf{L}^{p_1}(\Omega)}^{p_1} + \|\nabla u\|_{\mathbf{L}^{p_2}(\Omega)}^{p_2} \right) < \infty, \end{aligned}$$

where the last three inequalities follow from Lemma 3.1, Jensen's inequality, and Theorem 2.5.

A.6. Proof of Lemma 3.7

We prove properties (3.8) and (3.9) of the homogenized operator b . Property (3.10), which occurs in the case when $1 < p_1 \leq 2 \leq p_2$, follows in a similar way.

A.6.1. Proof of (3.8)

Using (2.8) and (2.2), we have

$$\begin{aligned} (b(\xi_2) - b(\xi_1), \xi_2 - \xi_1) &= \int_Y (A(y, P(y, \xi_2)) - A(y, P(y, \xi_1)), P(y, \xi_2) - P(y, \xi_1)) \, dy \\ &\geq C \int_Y |P(y, \xi_1) - P(y, \xi_2)|^{\beta(y)} (|P(y, \xi_1)| + |P(y, \xi_2)|)^{p(y) - \beta(y)} \, dy \geq 0. \end{aligned}$$

A.6.2. Proof of (3.9)

Note that

$$\begin{aligned} &\int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^2 (|P(y, \xi_1)| + |P(y, \xi_2)|)^{p_1 - 2} \, dy \\ &+ \int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^2 (|P(y, \xi_1)| + |P(y, \xi_2)|)^{p_2 - 2} \, dy \tag{A.4} \\ &\leq C \int_Y (A(y, P(y, \xi_2)) - A(y, P(y, \xi_1)), P(y, \xi_2) - P(y, \xi_1)) \, dy \\ &= C \int_Y (A(y, P(y, \xi_2)) - A(y, P(y, \xi_1)), \xi_2 - \xi_1) \, dy \\ &= (b(\xi_2) - b(\xi_1), \xi_2 - \xi_1) \leq |b(\xi_2) - b(\xi_1)| |\xi_2 - \xi_1| \end{aligned}$$

by (2.2), (2.8), and (2.6). Also, using (2.1), we obtain

$$\begin{aligned} |b(\xi_1) - b(\xi_2)| &\leq \int_Y |A(y, P(y, \xi_1)) - A(y, P(y, \xi_2))| \, dy \\ &\leq C \left[\int_Y \chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1 - 1} \, dy + \int_Y \chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2 - 1} \, dy \right] \\ &= C \left[\int_Y \frac{\chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_1 - 1} (|P(x, \xi_1)| + |P(x, \xi_2)|)^{\frac{(2-p_1)(p_1-1)}{2}}}{(|P(x, \xi_1)| + |P(x, \xi_2)|)^{\frac{(2-p_1)(p_1-1)}{2}}} \, dy \right. \\ &\quad \left. + \int_Y \frac{\chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^{p_2 - 1} (|P(x, \xi_1)| + |P(x, \xi_2)|)^{\frac{(2-p_2)(p_2-1)}{2}}}{(|P(x, \xi_1)| + |P(x, \xi_2)|)^{\frac{(2-p_2)(p_2-1)}{2}}} \, dy \right]. \end{aligned}$$

Using Hölder's inequality with $r_i = \frac{2}{p_i - 1} > 1$ and $\frac{1}{s_i} = 1 - \frac{1}{r_i} = \frac{3-p_i}{2}$, for $i = 1, 2$

$$\begin{aligned} &\leq C \left[\left(\int_Y \frac{\chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^2}{(|P(x, \xi_1)| + |P(x, \xi_2)|)^{2-p_1}} \, dy \right)^{\frac{p_1-1}{2}} \left(\int_Y \chi_1(y) (|P(x, \xi_1)| + |P(x, \xi_2)|)^{\frac{(2-p_1)(p_1-1)}{3-p_1}} \, dy \right)^{\frac{3-p_1}{2}} \right. \\ &\quad \left. + \left(\int_Y \frac{\chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^2}{(|P(x, \xi_1)| + |P(x, \xi_2)|)^{2-p_2}} \, dy \right)^{\frac{p_2-1}{2}} \left(\int_Y \chi_2(y) (|P(x, \xi_1)| + |P(x, \xi_2)|)^{\frac{(2-p_2)(p_2-1)}{3-p_2}} \, dy \right)^{\frac{3-p_2}{2}} \right]. \end{aligned}$$

Using Hölder’s inequality with $r_i = \frac{p_i(3-p_i)}{(p_i-1)(2-p_i)} > 1$, for $i = 1, 2$

$$\begin{aligned} &\leq C \left[\left(\int_Y \frac{\chi_1(y) |P(y, \xi_1) - P(y, \xi_2)|^2}{(|P(x, \xi_1)| + |P(x, \xi_2)|)^{2-p_1}} dy \right)^{\frac{p_1-1}{2}} \left(\int_Y \chi_1(y) (|P(x, \xi_1)| + |P(x, \xi_2)|)^{p_1} dy \right)^{\frac{(2-p_1)(p_1-1)}{2p_1}} \right. \\ &\quad \left. + \left(\int_Y \frac{\chi_2(y) |P(y, \xi_1) - P(y, \xi_2)|^2}{(|P(x, \xi_1)| + |P(x, \xi_2)|)^{2-p_2}} dy \right)^{\frac{p_2-1}{2}} \left(\int_Y \chi_2(y) (|P(x, \xi_1)| + |P(x, \xi_2)|)^{p_2} dy \right)^{\frac{(2-p_2)(p_2-1)}{2p_2}} \right]. \end{aligned}$$

By (A.4) and Lemma (3.1)

$$\leq C \sum_{i=1}^2 \left[|b(\xi_2) - b(\xi_1)|^{\frac{p_i-1}{2}} |\xi_2 - \xi_1|^{\frac{p_i-1}{2}} (1 + \theta_1 |\xi_1|^{p_1} + \theta_2 |\xi_1|^{p_2} + \theta_1 |\xi_2|^{p_1} + \theta_2 |\xi_2|^{p_2})^{\frac{(2-p_i)(p_i-1)}{2p_i}} \right].$$

By Young’s Inequality with $r_i = \frac{2}{p_i-1} > 1$ with $i = 1, 2$

$$\leq C \sum_{i=1}^2 \left[\frac{\delta^{\frac{2}{p_i-1}} |b(\xi_2) - b(\xi_1)|}{\frac{2}{p_i-1}} + \frac{\delta^{\frac{-2}{3-p_i}} |\xi_2 - \xi_1|^{\frac{p_i-1}{3-p_i}} (1 + \theta_1 |\xi_1|^{p_1} + \theta_2 |\xi_1|^{p_2} + \theta_1 |\xi_2|^{p_1} + \theta_2 |\xi_2|^{p_2})^{\frac{(2-p_i)(p_i-1)}{p_i(3-p_i)}}}{\frac{2}{3-p_i}} \right].$$

Rearranging the terms, and taking δ small enough we obtain (3.9)

A.7. Proof of Lemma 3.8

Using Hölder’s inequality, (2.1), and (2.9), we obtain

$$\begin{aligned} &\int_{\Omega_j} |(A_\epsilon(x, P_\epsilon(x, \eta_j)), \nabla u_\epsilon(x))| dx \leq \int_{\Omega_j} |A_\epsilon(x, P_\epsilon(x, \eta_j))| |\nabla u_\epsilon(x)| dx \\ &\leq C \left[\left(\int_{\Omega_j} \chi_1^\epsilon(x) |P_\epsilon(x, \eta_j)|^{p_1} dx \right)^{\frac{1}{q_2}} + \left(\int_{\Omega_j} \chi_2^\epsilon(x) |P_\epsilon(x, \eta_j)|^{p_2} dx \right)^{\frac{1}{q_1}} \right] \\ &\leq C, \text{ where } C \text{ does not depend on } \epsilon. \end{aligned}$$

The proof of the uniform boundedness of $\int_{\Omega_j} |(A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, \eta_j))| dx$ follows in the same manner.

A.8. Proof of Lemma 3.9

We prove the first statement of the lemma, the second statement follows in a similar way. The lemma follows from the Dunford–Pettis theorem (see [8]). To apply this theorem, the following conditions are necessary:

- $\int_{\Omega_j} |(A_\epsilon(x, P_\epsilon(x, \eta_j)), \nabla u_\epsilon(x))| dx$ is uniformly bounded with respect to ϵ , which was proved in Lemma 3.8.
- $(A_\epsilon(\cdot, P_\epsilon(\cdot, \eta_j)), \nabla u_\epsilon(\cdot))$ is equiintegrable for all $j = 0, \dots, m$

Proof. By Theorem 1.5 in [8], we have that $\chi_1^\epsilon(\cdot) |A_\epsilon(\cdot, P_\epsilon(\cdot, \eta_j))|^{q_2}$ and $\chi_2^\epsilon(\cdot) |A_\epsilon(\cdot, P_\epsilon(\cdot, \eta_j))|^{q_1}$ are equiintegrable.

By (2.9), for any $E \subset \Omega$, we have

$$\max_{i=1,2} \left\{ \sup_{\epsilon > 0} \left\{ \left(\int_E \chi_i^\epsilon(x) |\nabla u_\epsilon(x)|^{p_i} dx \right)^{\frac{1}{p_i}} \right\} \right\} \leq C.$$

Let $\lambda > 0$ arbitrary and choose $\lambda_1 > 0$ and $\lambda_2 > 0$ such that $\lambda_1^{\frac{1}{q_2}} + \lambda_2^{\frac{1}{q_1}} < \lambda/C$.

For λ_1 and λ_2 , there exist $\omega(\lambda_1) > 0$ and $\omega(\lambda_2) > 0$ such that for every $E \subset \Omega$ with $|E| < \min\{\omega(\lambda_1), \omega(\lambda_2)\}$, we have

$$\int_E \chi_1^\epsilon(x) |A_\epsilon(x, P_\epsilon(x, \eta_j))|^{q_2} dx < \lambda_1 \quad \text{and} \quad \int_E \chi_2^\epsilon(x) |A_\epsilon(x, P_\epsilon(x, \eta_j))|^{q_1} dx < \lambda_2.$$

Take $\omega = \omega(\lambda) = \min\{\omega(\lambda_1), \omega(\lambda_2)\}$. Then, for all $E \subset \Omega$ with $|E| < \omega$, we have

$$\begin{aligned} \int_E |(A_\epsilon(x, P_\epsilon(x, \eta_j)), \nabla u_\epsilon)| dx &\leq \int_E |A_\epsilon(x, P_\epsilon(x, \eta_j))| |\nabla u_\epsilon| dx \\ &\leq \left(\int_E \chi_1^\epsilon(x) |A_\epsilon(x, P_\epsilon(x, \eta_j))|^{q_2} dx \right)^{\frac{1}{q_2}} \left(\int_E \chi_1^\epsilon(x) |\nabla u_\epsilon|^{p_1} dx \right)^{\frac{1}{p_1}} \\ &\quad + \left(\int_E \chi_2^\epsilon(x) |A_\epsilon(x, P_\epsilon(x, \eta_j))|^{q_1} dx \right)^{\frac{1}{q_1}} \left(\int_E \chi_2^\epsilon(x) |\nabla u_\epsilon|^{p_2} dx \right)^{\frac{1}{p_2}} \\ &\leq C(\lambda_1^{1/q_2} + \lambda_2^{1/q_1}) < \lambda, \end{aligned}$$

for every $\lambda > 0$, and so $(A_\epsilon(\cdot, P_\epsilon(\cdot, \eta_j)), \nabla u_\epsilon)$ is equiintegrable. \square

A.9. Case $1 < p_1 \leq 2 \leq p_2$: Proof of Theorem 2.7 in Section 4

The proof of Theorem 2.7 for the case $1 < p_1 \leq 2 \leq p_2$ is very similar to the one presented in Section 4. Here we will indicate only the main differences in the different parts of the proof.

We have by (2.2), Lemmas 3.5 and 3.6 that

$$\begin{aligned} &\int_\Omega [\chi_1^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)|^{p_1} + \chi_2^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)|^{p_2}] dx \\ &\leq C \left[(|\Omega| + \int_\Omega \chi_1^\epsilon(x) |P_\epsilon(x, M_\epsilon \nabla u(x))|^{p_1} dx + \int_\Omega \chi_1^\epsilon(x) |\nabla u_\epsilon(x)|^{p_1} dx)^{\frac{2-p_1}{2}} \right. \\ &\quad \times \left(\int_\Omega \chi_1^\epsilon(x) (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \right)^{\frac{p_1}{2}} \\ &\quad \left. + \int_\Omega \chi_2^\epsilon(x) (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \right] \\ &\leq C \left[\left(\int_\Omega (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \right)^{\frac{p_1}{2}} \right. \\ &\quad \left. + \int_\Omega (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \right]. \end{aligned}$$

Therefore to prove Theorem 2.7 in this case, we also need to show that

$$\begin{aligned} &\int_\Omega (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u(x))) - A_\epsilon(x, \nabla u_\epsilon(x)), P_\epsilon(x, M_\epsilon \nabla u(x)) - \nabla u_\epsilon(x)) dx \\ &= \int_\Omega (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u)), P_\epsilon(x, M_\epsilon \nabla u)) dx - \int_\Omega (A_\epsilon(x, P_\epsilon(x, M_\epsilon \nabla u)), \nabla u_\epsilon) dx \\ &\quad - \int_\Omega (A_\epsilon(x, \nabla u_\epsilon), P_\epsilon(x, M_\epsilon \nabla u)) dx + \int_\Omega (A_\epsilon(x, \nabla u_\epsilon), \nabla u_\epsilon) dx \end{aligned}$$

goes to 0 as $\epsilon \rightarrow 0$ and this is done with the same four steps as in Section 4.

In Step 1, by (3.10), Hölder's inequality, Theorem 2.5, and Jensen's inequality, we obtain

$$\int_{\Omega} |b(M_{\epsilon} \nabla u(x)) - b(\nabla u(x))|^{q_1} dx \leq C \left[\left(\int_{\Omega} |M_{\epsilon} \nabla u - \nabla u|^{p_2} dx \right)^{\frac{p_1-1}{(p_2-1)(3-p_1)}} + \left(\int_{\Omega} |M_{\epsilon} \nabla u - \nabla u|^{p_2} dx \right)^{\frac{1}{(p_2-1)^2}} \right].$$

In Step 2, by (2.1), Hölder's inequality, and (2.9) we obtain (instead of (4.5))

$$\begin{aligned} & \left| \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon} \nabla u(x))) - A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) dx \right| \\ & \leq C \left[\left(\int_{\Omega} \chi_1^{\epsilon}(x) |P_{\epsilon}(x, M_{\epsilon} \nabla u) - P_{\epsilon}(x, \Psi)|^{p_1} dx \right)^{\frac{p_1-1}{p_1}} \right. \\ & \quad \left. + \left(\int_{\Omega} \chi_2^{\epsilon}(x) |P_{\epsilon}(x, M_{\epsilon} \nabla u) - P_{\epsilon}(x, \Psi)|^{p_2} dx \right)^{\frac{1}{p_2}} \right]. \end{aligned} \quad (\text{A.5})$$

Applying Lemma 3.4 and (4.4) to (A.5), we discover that

$$\begin{aligned} & \limsup_{\epsilon \rightarrow 0} \left| \int_{\Omega} (A_{\epsilon}(x, P_{\epsilon}(x, M_{\epsilon} \nabla u(x))) - A_{\epsilon}(x, P_{\epsilon}(x, \Psi(x))), \nabla u_{\epsilon}(x)) dx \right| \\ & \leq C \left[\left(\delta^{\frac{p_1}{3-p_1}} + \delta^{\frac{p_1 p_2}{2p_2-p_1}} + \delta^{p_1} + \delta^{\frac{p_2}{p_2-1}} \right)^{\frac{p_1-1}{p_1}} \right. \\ & \quad \left. + \left(\delta^{\frac{p_1}{3-p_1}} + \delta^{\frac{p_1 p_2}{2p_2-p_1}} + \delta^{p_1} + \delta^{\frac{p_2}{p_2-1}} \right)^{\frac{1}{p_2}} \right], \end{aligned} \quad (\text{A.6})$$

where C is independent of δ . Since δ is arbitrary we conclude that the limit on the left hand side of (A.6) is equal to 0.

Finally, using the continuity of b (3.10) in Lemma 3.7, Theorem 2.5, and Hölder's inequality, we obtain

$$\left| \int_{\Omega} (b(\nabla u(x)) - b(\Psi(x)), \nabla u(x)) dx \right| \leq C \left[\delta^{\frac{p_2(p_1-1)}{(p_2-1)(3-p_1)}} + \delta^{\frac{p_2}{(p_2-1)^2}} \right]^{\frac{1}{q_1}},$$

where C does not depend on δ .

In Step 3, we have

$$\begin{aligned} & \limsup_{\epsilon \rightarrow 0} \left| \int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}), P_{\epsilon}(x, M_{\epsilon} \nabla u) - P_{\epsilon}(x, \Psi)) dx \right| \\ & \leq C \sum_{i=1}^2 \left[\delta^{\frac{p_1}{3-p_1}} + \delta^{\frac{p_1 p_2}{2p_2-p_1}} + \delta^{p_1} + \delta^{\frac{p_2}{p_2-1}} \right]^{\frac{1}{p_i}}, \end{aligned}$$

where C does not depend on δ .

Hence, proceeding as in Step 2, we find that

$$\begin{aligned} & \limsup_{\epsilon \rightarrow 0} \left| \int_{\Omega} (A_{\epsilon}(x, \nabla u_{\epsilon}), P_{\epsilon}(x, M_{\epsilon} \nabla u)) dx - \int_{\Omega} (b(\nabla u), \nabla u) dx \right| \\ & \leq C \left[\sum_{i=1}^2 \left(\delta^{\frac{p_1}{3-p_1}} + \delta^{\frac{p_1 p_2}{2p_2-p_1}} + \delta^{p_1} + \delta^{\frac{p_2}{p_2-1}} \right)^{\frac{1}{p_i}} + \delta \|b(\nabla u)\|_{L^{q_2}(\Omega, \mathbb{R}^n)} \right], \end{aligned}$$

where C is independent of δ . Now since δ is arbitrarily small, the proof of Step 3 is complete.

REFERENCES

- [1] B. Amaziane, S.N. Antontsev, L. Pankratov and A. Piatnitski, Γ -convergence and homogenization of functionals in Sobolev spaces with variable exponents. *J. Math. Anal. Appl.* **342** (2008) 1192–1202.
- [2] S.N. Antontsev and J.F. Rodrigues, On stationary thermo-rheological viscous flows. *Ann. Univ. Ferrara Sez. VII Sci. Mat.* **52** (2006) 19–36.
- [3] C. Atkinson and C.R. Champion, Some boundary-value problems for the equation $\nabla \cdot (|\nabla\varphi|^N \nabla\varphi) = 0$. *Quart. J. Mech. Appl. Math.* **37** (1984) 401–419.
- [4] L.C. Berselli, L. Diening and M. Růžička, Existence of strong solutions for incompressible fluids with shear dependent viscosities. *J. Math. Fluid Mech.* **12** (2010) 101–132.
- [5] A. Braides, V. Chiadò Piat and A. Defranceschi, Homogenization of almost periodic monotone operators. *Ann. Inst. Henri Poincaré, Anal. Non Linéaire* **9** (1992) 399–432.
- [6] J. Byström, Correctors for some nonlinear monotone operators. *J. Nonlinear Math. Phys.* **8** (2001) 8–30.
- [7] J. Byström, Sharp constants for some inequalities connected to the p -Laplace operator. *JIPAM. J. Inequal. Pure Appl. Math.* **6** (2005). Article 56 (electronic) 8.
- [8] B. Dacorogna, Direct methods in the calculus of variations, in *Appl. Math. Sci.* Springer-Verlag, Berlin **78** (1989).
- [9] G. Dal Maso and A. Defranceschi, Correctors for the homogenization of monotone operators. *Differ. Integral Equ.* **3** (1990) 1151–1166.
- [10] N. Fusco and G. Moscarillo, On the homogenization of quasilinear divergence structure operators. *Ann. Mat. Pura Appl.* **146** (1987) 1–13.
- [11] A. Garroni and R.V. Kohn, Some three-dimensional problems related to dielectric breakdown and polycrystal plasticity. *Proc. R. Soc. London Ser. A Math. Phys. Eng. Sci.* **459** (2003) 2613–2625.
- [12] A. Garroni, V. Nesi and M. Ponsiglione, Dielectric breakdown: optimal bounds. *Proc. R. Soc. London Ser. A Math. Phys. Eng. Sci.* **457** (2001) 2317–2335.
- [13] R. Glowinski and J. Rappaz, Approximation of a nonlinear elliptic problem arising in a non-Newtonian fluid flow model in glaciology. *ESAIM: M2AN* **37** (2003) 175–186.
- [14] M. Idiart, The macroscopic behavior of power-law and ideally plastic materials with elliptical distribution of porosity. *Mech. Res. Commun.* **35** (2008) 583–588.
- [15] S. Jimenez and R.P. Lipton, Correctors and field fluctuations for the $p_\epsilon(x)$ -Laplacian with rough exponents. *J. Math. Anal. Appl.* **372** (2010) 448–469.
- [16] A. Kelly and N.H. Macmillan, *Strong Solids. Monographs on the Physics and Chemistry of Materials.* Clarendon Press, Oxford (1986).
- [17] S. Levine, J. Stanich and Y. Chen, *Image restoration via nonstandard diffusion.* Technical report (2004).
- [18] O. Levy and R.V. Kohn, Duality relations for non-Ohmic composites, with applications to behavior near percolation. *J. Stat. Phys.* **90** (1998) 159–189.
- [19] R. Lipton, Homogenization and field concentrations in heterogeneous media. *SIAM J. Math. Anal.* **38** (2006) 1048–1059.
- [20] F. Murat and L. Tartar, H -convergence. In *Topics in the mathematical modelling of composite materials, Progr. Nonlinear Diff. Equ. Appl.* Birkhäuser Boston, Boston, MA **31** (1997) 21–43.
- [21] P. Pedregal, Parametrized measures and variational principles, in *Progress in Nonlinear Differential Equations and their Applications.* Birkhäuser Verlag, Basel (1997).
- [22] P. Pedregal and H. Serrano, Homogenization of periodic composite power-law materials through young measures, in *Multi scale problems and asymptotic analysis. GAKUTO Int. Ser. Math. Sci. Appl.* Gakkōtoshō, Tokyo **24** (2006) 305–310.
- [23] P. Ponte Castañeda and P. Suquet, Nonlinear composites. *Adv. Appl. Mech.* **34** (1997) 171–302.
- [24] P. Ponte Castañeda and J.R. Willis, Variational second-order estimates for nonlinear composites. *Proc. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci.* **455** (1999) 1799–1811.
- [25] M. Růžička, Electrorheological fluids: modeling and mathematical theory. *Lect. Notes Math.* **1748** (2000).
- [26] P. Suquet, Overall potentials and extremal surfaces of power law or ideally plastic composites. *J. Mech. Phys. Solids* **41** (1993) 981–1002.
- [27] D.R.S. Talbot and J.R. Willis, Upper and lower bounds for the overall properties of a nonlinear elastic composite dielectric I. random microgeometry. *Proc. R. Soc. Lond. A* **447** (1994) 365–384.
- [28] D.R.S. Talbot and J.R. Willis, Upper and lower bounds for the overall properties of a nonlinear elastic composite dielectric II. periodic microgeometry. *Proc. R. Soc. Lond. A* **447** (1994) 385–396.
- [29] A.C. Zaenen, *An introduction to the theory of integration.* Publishing Company, North-Holland, Amsterdam (1958).
- [30] V.V. Zhikov, S.M. Kozlov and O.A. Oleinik, *Homogenization of differential operators and integral functionals.* Springer-Verlag, Berlin (1994). Translated from the Russian by G.A. Yosifian.