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Partial Differential Equations

Some asymptotic properties for solutions of one-dimensional advection–diffusion equations with Cauchy data in $L^p(\mathbb{R})$

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Abstract

We state and discuss a number of fundamental asymptotic properties of solutions $u(\cdot, t)$ to one-dimensional advection–diffusion equations of the form $u_t + f(u)_x = (a(u)u_x)_x$, $x \in \mathbb{R}$, $t > 0$, assuming initial values $u(\cdot, 0) = u_0 \in L^p(\mathbb{R})$ for some $1 \leq p < \infty$.

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Résumé

Quelques propriétés asymptotiques des solutions des équations d’advection–diffusion unidimensionnelles aux données initiales dans $L^p(\mathbb{R})$. Nous établissons plusieurs propriétés asymptotiques fondamentales des solutions $u(\cdot, t)$ des équations d’advection–diffusion du type $u_t + f(u)_x = (a(u)u_x)_x$, $x \in \mathbb{R}$, $t > 0$, aux données initiales dans l’espace de Lebesgue $L^p(\mathbb{R})$, où $1 \leq p < \infty$. Pour citer cet article : P. Braz e Silva, P.R. Zingano, C. R. Acad. Sci. Paris, Ser. I 342 (2006).

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1. Introduction

We examine some properties of solutions $u(\cdot, t)$ to the Cauchy problem for scalar advection–diffusion equations

$$u_t + f(u)_x = (a(u)u_x)_x, \quad x \in \mathbb{R}, \quad t > 0, \quad (1a)$$

$$u(\cdot, 0) = u_0 \in L^p(\mathbb{R}), \quad 1 \leq p < \infty, \quad (1b)$$

where $a(\cdot)$ and $f(\cdot)$ are given smooth functions. We assume that $a(u) \geq \mu > 0$ for some fixed constant μ and all values of u concerned. Under these conditions, it is known that problem (1) admits a unique, smooth (classical), globally defined solution $u(\cdot, t) \in C^0([0, \infty[, L^p(\mathbb{R}))$, which is bounded for $t > 0$ and satisfies (1b) in L^p sense, that is, $\|u(\cdot, t) - u_0\|_{L^p(\mathbb{R})} \rightarrow 0$ as $t \rightarrow 0$, see e.g. [2–4] and Section 2 below. Several additional properties of $u(\cdot, t)$ are given next.

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2. Decay estimates

We first state an important energy-type inequality for $u(\cdot, t)$.

Theorem 2.1. Assume $a(\cdot), f(\cdot) \in C^1$ with $a(\cdot)$ bounded below by some constant $\mu > 0$. Then, for each $q \geq \max\{p, 2\}$, the solution $u(\cdot, t) \in C^0([0, \infty[, L^p(\mathbb{R}))$ of the Cauchy problem (1) satisfies

$$\begin{aligned} T^{\frac{q}{2p}} \|u(\cdot, T)\|_{L^q(\mathbb{R})}^q + q(q-1)\mu \int_0^T t^{\frac{q}{2p}} \int_{\mathbb{R}} |u(x, t)|^{q-2} |u_x(x, t)|^2 dx dt \\ \leq 2 \left(\frac{q}{2p} \right)^{\frac{1}{2}(\frac{q}{p}+1)} \left(1 - \frac{1}{q} \right)^{-\frac{1}{2}(\frac{q}{p}-1)} \|u_0\|_{L^p(\mathbb{R})}^q \mu^{-\frac{1}{2}(\frac{q}{p}-1)} T^{\frac{1}{2}} \end{aligned} \quad (2)$$

for all $T > 0$.

Theorem 2.1 can be proved adapting the method discussed in [5] to our present needs. Decay estimates for $u(\cdot, t)$ are readily obtained from inequality (2). Indeed, by the maximum principle and choosing $q = 4p$, inequality (2) yields the supnorm estimate

$$\|u(\cdot, t)\|_{L^\infty(\mathbb{R})} \leq C_p \|u_0\|_{L^p(\mathbb{R})} (\mu t)^{-\frac{1}{2p}}, \quad C_p := 4^{\frac{1}{p}} \left(4 - \frac{1}{p} \right)^{-\frac{1}{2p}} \quad (3)$$

for all $t > 0$. We note that the minimal value for C_p is not known; its particular value given here is not optimal. Now, since $\|u(\cdot, t)\|_{L^p(\mathbb{R})} \leq \|u_0\|_{L^p(\mathbb{R})}$ for all $t \geq 0$, it follows from a simple interpolation estimate that

$$\|u(\cdot, t)\|_{L^r(\mathbb{R})} \leq C_p^{1-\frac{p}{r}} \|u_0\|_{L^p(\mathbb{R})} (\mu t)^{-\frac{1}{2}(\frac{1}{p}-\frac{1}{r})}, \quad \forall t > 0 \quad (4)$$

for all $p \leq r \leq \infty$. Therefore, solutions decay in L^r for any $r > p$. Using standard estimates for fundamental solutions of linear parabolic problems (see [1,2]), one obtains decay rates for the derivatives of $u(\cdot, t)$ as well, e.g.

$$\|u_x(\cdot, t)\|_{L^r(\mathbb{R})} \leq C(p, r, \mu, K_p, t_0) \|u_0\|_{L^p(\mathbb{R})} t^{-\frac{1}{2}(\frac{1}{p}-\frac{1}{r})-\frac{1}{2}}, \quad \forall t \geq t_0 \quad (5)$$

for each $t_0 > 0$. Similar bounds hold for the other derivatives. Here, the constant $C(p, r, \mu, K_p, t_0)$ depends on the particular functions $a(\cdot)$ and $f(\cdot)$, on the values of p, r, μ, t_0 , and on K_p , a bound for $\|u_0\|_{L^p(\mathbb{R})}$, i.e., $K_p > 0$ chosen such that

$$\|u_0\|_{L^p(\mathbb{R})} \leq K_p. \quad (6)$$

3. Asymptotic behavior, $p = 1$

For $t \gg 1$, more detailed behavior of $u(\cdot, t)$ can be obtained from Theorem 3.1 below. This theorem follows from the estimates given in Section 2. Here, we assume $a, f \in C^2$, with f Hölder continuous at 0. We also assume $a(u) \geq \mu > 0$ for all u , as before.

Theorem 3.1. Let $u_0 \in L^1(\mathbb{R})$, and let $v(\cdot, t) \in C^0([0, \infty[, L^1(\mathbb{R}))$ be (any) solution of the Burgers equation

$$v_t + f'(0)v_x + f''(0)vv_x = a(0)v_{xx}, \quad x \in \mathbb{R}, \quad t > 0 \quad (7)$$

having the same mass as $u(\cdot, t)$, i.e., $\int_{\mathbb{R}} v(x, 0) dx = \int_{\mathbb{R}} u_0(x) dx$. Then, one has

$$\lim_{t \rightarrow \infty} t^{\frac{1}{2}(1-\frac{1}{r})} \|u(\cdot, t) - v(\cdot, t)\|_{L^r(\mathbb{R})} = 0 \quad (8)$$

for each $1 \leq r \leq \infty$, uniformly in r .

Now, solutions of (7) can be studied in detail through explicit representation formulas obtained with the so-called Hopf–Cole transformation, see [6]. In many cases, these properties can be recast for (1a) using Theorem 3.1 above, as illustrated by the following two results.

Theorem 3.2. Let $u(\cdot, t), \hat{u}(\cdot, t) \in C^0([0, \infty[, L^1(\mathbb{R}))$ be solutions of Eq. (1a) corresponding to initial states $u_0, \hat{u}_0 \in L^1(\mathbb{R})$, respectively, with the same mass. Then, one has

$$\lim_{t \rightarrow \infty} t^{\frac{1}{2}(1-\frac{1}{r})} \|u(\cdot, t) - \hat{u}(\cdot, t)\|_{L^r(\mathbb{R})} = 0, \quad (9)$$

for all $1 \leq r \leq \infty$, uniformly in r .

Theorem 3.3. Let $u(\cdot, t) \in C^0([0, \infty[, L^1(\mathbb{R}))$ be the solution of problem (1) for some given initial state $u_0 \in L^1(\mathbb{R})$ with mass m . Then, for each $1 \leq r \leq \infty$,

$$\lim_{t \rightarrow \infty} t^{\frac{1}{2}(1-\frac{1}{r})} \|u(\cdot, t)\|_{L^r(\mathbb{R})} = \gamma_r(m),$$

where, if $f''(0) \neq 0$, the quantity $\gamma_r(m)$ is given by

$$\gamma_r(m) = \frac{|m|}{\sqrt{4\pi a(0)}} (4a(0))^{\frac{1}{2r}} \frac{2a(0)}{mf''(0)} \left(1 - e^{-\frac{mf''(0)}{2a(0)}}\right) \|\mathcal{F}\|_{L^r(\mathbb{R})},$$

for $\mathcal{F} \in L^1(\mathbb{R}) \cap L^\infty(\mathbb{R})$ defined by

$$\mathcal{F}(x) = 2e^{-x^2} / \left(1 + e^{-\frac{mf''(0)}{2a(0)}} - \left(1 - e^{-\frac{mf''(0)}{2a(0)}}\right) \text{erf}(x)\right).$$

Here, $\text{erf}(x)$ is the error function, and $\gamma_r(0) = 0$. If $f''(0) = 0$, $\gamma_r(m)$ is given by

$$\gamma_r(m) = \frac{|m|}{\sqrt{4\pi a(0)}} \left(\frac{4\pi a(0)}{r}\right)^{\frac{1}{2r}}.$$

The case $r = 1$ is worth explicit mention:

$$\lim_{t \rightarrow \infty} \|u(\cdot, t)\|_{L^1(\mathbb{R})} = |m| \quad \text{and} \quad \lim_{t \rightarrow \infty} \|u(\cdot, t) - \hat{u}(\cdot, t)\|_{L^1(\mathbb{R})} = 0$$

for any two solutions $u(\cdot, t), \hat{u}(\cdot, t)$ of Eq. (1a) transporting the same mass m .

4. Asymptotic behavior, $p > 1$

Lastly, we consider solutions $u(\cdot, t)$ of problem (1) when $p > 1$. In this case, taking cut-off approximations $v_R = u_0 \cdot \chi_{[-R, R]} \in L^1(\mathbb{R}) \cap L^p(\mathbb{R})$ of the given initial data $u_0 \in L^p(\mathbb{R})$, and using results for the case $p = 1$, one obtains the following.

Theorem 4.1. Let $u(\cdot, t) \in C^0([0, \infty[, L^p(\mathbb{R}))$ be the solution of problem (1) corresponding to an initial state $u_0 \in L^p(\mathbb{R})$, $p > 1$. Then, one has

$$\lim_{t \rightarrow \infty} t^{\frac{1}{2}(\frac{1}{p}-\frac{1}{r})} \|u(\cdot, t)\|_{L^r(\mathbb{R})} = 0 \quad (10)$$

for each $p \leq r \leq \infty$, uniformly in r .

As for the heat equation, the convergence to zero in (10) can be arbitrarily slow (for suitable $u_0 \in L^p(\mathbb{R})$ verifying (6), $K_p > 0$ fixed). Therefore, no rates better than (4) can be given in general.

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