

# RENDICONTI *del* SEMINARIO MATEMATICO *della* UNIVERSITÀ DI PADOVA

PAOLO SECCHI

## **A note on the generic solvability of the Navier-Stokes equations**

*Rendiconti del Seminario Matematico della Università di Padova*,  
tome 83 (1990), p. 177-182

[http://www.numdam.org/item?id=RSMUP\\_1990\\_\\_83\\_\\_177\\_0](http://www.numdam.org/item?id=RSMUP_1990__83__177_0)

© Rendiconti del Seminario Matematico della Università di Padova, 1990, tous droits réservés.

L'accès aux archives de la revue « Rendiconti del Seminario Matematico della Università di Padova » (<http://rendiconti.math.unipd.it/>) implique l'accord avec les conditions générales d'utilisation (<http://www.numdam.org/conditions>). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

NUMDAM

*Article numérisé dans le cadre du programme  
Numérisation de documents anciens mathématiques*  
<http://www.numdam.org/>

## A Note on the Generic Solvability of the Navier-Stokes Equations.

PAOLO SECCHI (\*)

### 1. - Introduction.

Let  $\Omega \subset \mathbb{R}^3$  be a bounded domain in  $\mathbb{R}^3$  with boundary  $\partial\Omega$  of class  $C^2$ . Consider the Navier-Stokes equations

$$(1.1) \quad \begin{aligned} u' + (u \cdot \nabla)u - \Delta u + \nabla\pi &= f && \text{in } Q_T \equiv (0, T) \times \Omega, \\ \operatorname{div} u &= 0 && \text{in } Q_T, \\ u &= 0 && \text{on } \Sigma_T \equiv (0, T) \times \partial\Omega, \\ u(0) &= u_0 && \text{in } \Omega, \end{aligned}$$

with some  $T > 0$ . By a strong solution  $(u, \nabla\pi)$  of (1.1) we mean a solution with

$$\begin{aligned} u &\in W_p^{2,1}(Q_T) \equiv L^p(0, T; W_p^2(\Omega)^3) \cap W_p^1(0, T; L^p(\Omega)^3), \\ \nabla\pi &\in L^p(Q_T) \equiv L^p(0, T; L^p(\Omega)^3) \end{aligned}$$

for some  $p$  with  $2 \leq p < \infty$ . Let  $J_p^{2-2/p}(\Omega)$  denote the closure in the norm of  $W_p^{2-2/p}(\Omega)^3$  of the set of smooth finite solenoidal vectors equal

(\*) Indirizzo dell'A.: Dipartimento di Matematica Pura ed Applicata, Università di Padova, via Belzoni 7, 35131 Padova, Italy.

to zero on  $\partial\Omega$ . Consider  $u_0 \in \mathcal{J}_p^{2-2/p}(\Omega)$  and  $f \in L^p(Q_T)$ . Then it is known that these assumptions on the data assure the existence of a local in time unique strong solution of (1.1) (see [4]). The existence of strong solutions for arbitrary  $T > 0$  is an important open problem. Therefore it is interesting to know properties of the set

$$R(u_0) = \{f \in L^p(Q_T) / (1.1) \text{ has a unique strong solution } (u, \nabla\pi) \text{ with data } u_0, f\}$$

for a fixed initial value  $u_0 \in \mathcal{J}_p^{2-2/p}(\Omega)$ . It is not known whether or not  $R(u_0) = L^p(Q_T)$ ; however it is interesting to prove some density properties of this set, since this gives information about how many  $f$  do exist such that (1.1) is strongly solvable. In this concern H. Sohr and W. von Wahl [3] have proved the following interesting result: *the set  $R(u_0) \subset L^p(Q_T)$  is dense in the norm of  $L^s(0, T; L^q(\Omega)^3)$  for all  $s, q \in (1, \infty)$  with  $4 < 2/s + 3/q$  (see also [2] for a weaker previous result).* Their result is proved by a regularization procedure for (1.1) using an approximation of Yosida type and an estimate of the non-linear term  $(u \cdot \nabla)u$  using the exponent  $p = 5/4$  (see [3]). The aim of the present note is to prove the same result with a completely different method. We use an approximation method due to H. Beirão da Veiga [1] plus Sobolev imbedding and Hölder inequality. This approach is particularly simple and so we think it is of interest, even if the result is not new. Denote by  $|\cdot|_p$  the norm in  $L^p(\Omega)^3$  and by  $\|\cdot\|_{s,q,T}$  the norm in  $L^s(0, T; L^q(\Omega)^3)$ . Our result reads as follows

**THEOREM A.** *Let  $2 < p < \infty$  and  $u_0 \in \mathcal{J}_p^{2-2/p}(\Omega)$ . Then the set  $R(u_0) \subset L^p(Q_T)$  is dense in  $L^s(0, T; L^q(\Omega)^3)$  for all  $s, q \in (1, \infty)$  with  $4 < 2/s + 3/q$ . Therefore, for every  $f \in L^p(Q_T)$  and every  $\varepsilon > 0$  there exists  $g_\varepsilon \in L^p(Q_T)$  with  $\|g_\varepsilon\|_{s,q,T} \rightarrow 0$  as  $\varepsilon \rightarrow 0$  and such that*

$$\begin{aligned} u' + (u \cdot \nabla)u - \Delta u + \nabla\pi &= f + g_\varepsilon && \text{in } Q_T, \\ \operatorname{div} u &= 0 && \text{in } Q_T, \\ u &= 0 && \text{on } \Sigma_T, \\ u(0) &= u_0 && \text{in } \Omega, \end{aligned}$$

has a unique strong solution  $(u, \nabla\pi)$ .

**2. - Proof of Theorem A.**

Following [1] we define the set of vectors

$$A \equiv \{v \in C^\infty(\overline{Q_T})/v(t) \in C_0^\infty(\Omega)^3, \operatorname{div} v(t) = 0 \text{ in } \Omega \text{ for all } t \in [0, T]\},$$

where  $T > 0$  arbitrary, and consider the linearized system

$$(2.1) \quad \begin{aligned} u' + (v \cdot \nabla)u - \Delta u + \nabla \pi &= f && \text{in } Q_T, \\ \operatorname{div} u &= 0 && \text{in } Q_T, \\ u &= 0 && \text{on } \Sigma_T, \\ u(0) &= u_0 && \text{in } \Omega, \end{aligned}$$

for  $v \in A$ . For convenience define

$$(2.2) \quad \begin{aligned} A(u_0, f) &= |u_0|_2 + \|f\|_{1,2,T}, \\ A_1^2(u_0, f) &\equiv |u_0|_2^2 + 2\|f\|_{1,2,T}^2. \end{aligned}$$

From [4] (Theorem 4.2, p. 487) we have the following preliminary result:

**THEOREM 1.** *Let  $v \in A$ ,  $u_0 \in \dot{J}_p^{2-2/p}(\Omega)$ ,  $f \in L^p(Q_T)$ . Then there exists a unique solution  $(u, \nabla \pi)$  of problem (2.1) such that*

$$(2.3) \quad u \in W_p^{2,1}(Q_T), \quad \nabla \pi \in L^p(Q_T).$$

We quote now the result which gives us the approximating solutions we shall use later. It is proved in [1] (see Theorem 1.6, p. 329) as a consequence of a very interesting general approximation theorem.

**THEOREM 2 ([1]).** *Let  $u_0 \in H_0^1(\Omega)$  and  $f \in L^2(0, T; L^2(\Omega)^3)$  be given and let  $1 < q \leq 5/4$ . Then, in correspondence to every  $\varepsilon > 0$ , there exist*

$$\begin{aligned} u_\varepsilon \in A, \quad u_\varepsilon \in L^\infty(0, T; L^2(\Omega)^3) \cap L^2(0, T; H_0^1(\Omega)) \cap W_q^{2,1}(Q_T), \\ \pi_\varepsilon \in L^q(0, T; W_q^1(\Omega)^3) \end{aligned}$$

verifying the system

$$(2.4)_\varepsilon \quad \begin{aligned} u'_\varepsilon + (v_\varepsilon \cdot \nabla) u_\varepsilon - \Delta u_\varepsilon + \nabla \pi_\varepsilon &= f && \text{in } Q_T, \\ \operatorname{div} u_\varepsilon &= 0 && \text{in } Q_T, \\ u_\varepsilon &= 0 && \text{on } \Sigma_T, \\ u_\varepsilon(0) &= u_0 && \text{in } \Omega, \end{aligned}$$

and for which

$$(2.5) \quad \|u_\varepsilon - v_\varepsilon\|_{2,2,T} < \varepsilon.$$

Moreover the following estimates hold

$$(2.6) \quad \begin{aligned} \|u_\varepsilon\|_{\infty,2,T} &\leq A(u_0, f), \\ \|\nabla u_\varepsilon\|_{2,2,T} &\leq A_1(u_0, f) \end{aligned}$$

Estimates (2.6)<sub>1</sub> and (2.6)<sub>2</sub> hold also for  $v_\varepsilon$  and  $\nabla v_\varepsilon$  respectively.

Let now  $u_0, f$  and  $s, q$  as in Theorem A. A combination of Theorems 1 and 2 gives us that the approximating solution given by Theorem 2 satisfies also (2.3). We are now in position to prove our result. We write (2.4)<sub>ε</sub> in the form

$$(2.7) \quad \begin{aligned} u'_\varepsilon + (u_\varepsilon \cdot \nabla) u_\varepsilon - \Delta u_\varepsilon + \nabla \pi_\varepsilon &= f + (u_\varepsilon \cdot \nabla) u_\varepsilon - (v_\varepsilon \cdot \nabla) u_\varepsilon && \text{in } Q_T, \\ \operatorname{div} u_\varepsilon &= 0 && \text{in } Q_T, \\ u_\varepsilon &= 0 && \text{on } \Sigma_T, \\ u_\varepsilon(0) &= u_0 && \text{in } \Omega. \end{aligned}$$

Hence  $(u_\varepsilon, \nabla \pi_\varepsilon)$  is a strong solution of (1.1) with external force  $f + g_\varepsilon$ ,  $g_\varepsilon \equiv (u_\varepsilon \cdot \nabla) u_\varepsilon - (v_\varepsilon \cdot \nabla) u_\varepsilon$ . Observe that, because  $4 < 2/s + 3/q$ , we have  $s < 2, q < 3/2$ ; it follows that  $L^p(Q_T)$ ,  $2 \leq p < \infty$ , is densely contained in  $L^s(0, T; L^q(\Omega)^3)$ . Hence the theorem is proved if we show that  $g_\varepsilon \in L^p(Q_T)$  and  $\|g_\varepsilon\|_{s,q,T} \rightarrow 0$  as  $\varepsilon \rightarrow 0$ . Since  $v_\varepsilon \in \mathcal{A}$  and  $(u_\varepsilon, \nabla \pi_\varepsilon)$  satisfies (2.3), using some Sobolev imbeddings it easily follows that  $g_\varepsilon \in L^p(Q_T)$ . On the other hand, using the Hölder inequality gives

$$(2.8) \quad \|g_\varepsilon\|_{s,q,T} = \|((u_\varepsilon - v_\varepsilon) \cdot \nabla) u_\varepsilon\|_{s,q,T} \leq \|u_\varepsilon - v_\varepsilon\|_{s_1, q_1, T} \|\nabla u_\varepsilon\|_{2,2,T}$$

where  $(1/s_1) + (1/2) = 1/s$ ,  $(1/q_1) + (1/2) = 1/q$ . Since  $1 < s < 2$ ,  $1 < q < 3/2$  with  $2/s + 3/q > 4$  we obtain  $2 < s_1$ ,  $2 < q_1 < 6$  with  $2/s_1 + 3/q_1 > 3/2$ . Let  $s_2, q_2$  be the solution of

$$(2.9) \quad 2/s_2 + 3/q_2 = 3/2$$

$$(2.10) \quad (1 - 2/q_1)1/s_2 - (1 - 2/s_1)1/q_2 = 1/s_1 - 1/q_1.$$

The estimates on  $s_1, q_1$  yield  $2 < s_1 < s_2$ ,  $2 < q_1 < q_2 < 6$ . Using the Hölder inequality gives

$$(2.11) \quad \|u_\varepsilon - v_\varepsilon\|_{s_1, q_1, T} \leq \|u_\varepsilon - v_\varepsilon\|_{2, 2, T}^a \|u_\varepsilon - v_\varepsilon\|_{s_2, q_2, T}^b$$

where  $a, b$  must satisfy

$$(2.12) \quad \begin{aligned} a + b &= 1, \\ a/2 + b/q_2 &= 1/q_1, \\ a/2 + b/s_2 &= 1/s_1. \end{aligned}$$

System (2.12) has a solution if and only if the determinant of the corresponding complete matrix

$$\begin{pmatrix} 1 & 1 & 1 \\ 1/2 & 1/q_2 & 1/q_1 \\ 1/2 & 1/s_2 & 1/s_1 \end{pmatrix}$$

is zero. This condition is satisfied since (2.10) holds. Hence we find a solution of (2.12)

$$a = (1/q_1 - 1/q_2)/(1/2 - 1/q_2), \quad b = (1/2 - 1/q_1)/(1/2 - 1/q_2)$$

with  $a, b > 0$ . Since we have

$$1/q_2 = (1 - 2/s_2)/2 + (2/s_2)/6$$

with  $0 < (2/s_2) < 1$  it follows

$$\|u_\varepsilon - v_\varepsilon\|_{q_2} \leq \|u_\varepsilon - v_\varepsilon\|_2^{1 - (2/s_2)} \|u_\varepsilon - v_\varepsilon\|_6^{2/s_2} \leq C_1 \|u_\varepsilon - v_\varepsilon\|_2^{1 - (2/s_2)} |\nabla(u_\varepsilon - v_\varepsilon)|_2^{2/s_2},$$

where  $C_1$  is a positive constant. Integrating in time at the  $s_2$ -th power and the Young's inequality give

$$(2.13) \quad \|u_\varepsilon - v_\varepsilon\|_{s_2, a, T} \leq C_2 [\|u_\varepsilon - v_\varepsilon\|_{\infty, 2, T} + \|\nabla(u_\varepsilon - v_\varepsilon)\|_{2, 2, T}],$$

where  $C_2$  is a positive constant; the right-hand side of (2.13) is bounded because of (2.6). Hence from (2.5), (2.6) (also for  $v_\varepsilon$ ), (2.8), (2.11), (2.13) we obtain

$$\|g_\varepsilon\|_{s, a, T} \leq C_3 \varepsilon^a,$$

where  $C_3$  is a positive constant independent of  $\varepsilon$ . The theorem is proved.

#### REFERENCES

- [1] H. BEIRÃO DA VEIGA, *On the construction of suitable weak solutions to the Navier-Stokes equations via a general approximation theorem*, J. Math. Pures Appl., **64** (1985), pp. 321-334.
- [2] A. V. FURSIKOV, *On some problems of control and results concerning the unique solvability of a mixed boundary value problem for the three-dimensional Navier-Stokes and Euler systems*, Dokl. Akad. Nauk SSSR, **252** (1980), pp. 1066-1070.
- [3] H. SOHR - W. VON WAHL, *Generic solvability of the equations of Navier-Stokes*, Hiroshima Math. J., **17** (1987), pp. 613-625.
- [4] V. A. SOLONNIKOV, *Estimates for the solutions of nonstationary Navier-Stokes equations*, J. Soviet Math., **8** (1977), pp. 467-529.

Manoscritto pervenuto in redazione il 29 Giugno 1989.