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## P. PEISKER <br> On the numerical solution of the first biharmonic equation

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# ON THE NUMERICAL SOLUTION OF THE FIRST BIHARMONIC EQUATION (*) 

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Communcated by O. Pironneau


#### Abstract

We constder a mixed fintte element discrettzation of the biharmontc problem Following Glowinski and Pironneau the original indefinite linear system is transformed into a positive definte one for the unknown boundary value $\lambda=\left.\Delta u\right|_{\partial \Omega}$. This system is solved by a conjugate gradıent method We establish a preconditioning and prove that the number of iteration steps required for a given accuracy is independent of the mesh size.


Résumé. - On considere une méthode d'éléments fins muxtes pour le problème de Dirichlet de l'opérateur biharmonique Comme Glowinskl, Pironneau, on transforme le problème origınal, quı est indéfint, en un problème défint postuf pour la trace $\lambda=\left.\Delta u\right|_{\partial \Omega}$. Ce problème est résolu par la méthode du gradıent conjugué On établit une méthode de précondttonnement et on démontre que le nombre d'ttératons pour réduıre l'erreur d'un facteur fixe ne dépend pas du paramètre de discrétisation.

## 1. INTRODUCTION

We consider the numerical solution of the biharmonic equation

$$
\begin{equation*}
\Delta^{2} u=f \quad \text { in } \quad \Omega, \quad u=\frac{\partial u}{\partial n}=0 \quad \text { on } \quad \partial \Omega \tag{1.1}
\end{equation*}
$$

where $\Omega$ is a convex, polygonal domain in the plane. Suppose that the boundary value $\lambda^{*}=\left.\Delta u\right|_{\partial \Omega}$ is known. Then (1.1) is splitted into two separated Poisson equations. An initial guess $\lambda^{(0)}$ for the boundary value may be iteratively improved using the following procedure for $k=0,1,2, \ldots$
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[^0]Given $\lambda^{(k)}$. Then solve

$$
\begin{align*}
\Delta \phi^{(k)} & =f \text { in } \Omega, & \left.\phi^{(k)}\right|_{\partial \Omega} & =\lambda^{(k)} \\
\Delta u^{(k)} & =\phi^{(k)} \text { in } \Omega, & \left.u^{(k)}\right|_{\partial \Omega} & =0  \tag{1.2}\\
\lambda^{(k+1)} & =\lambda^{(k)}+\delta \frac{\partial u^{(k)}}{\partial n}, & & \delta
\end{align*}>0
$$

This method is known as coupled equation approach in the finite differences context (see e.g. [7]). In the framework of finite elements the discrete analogue of (1.2) was first studied by Ciarlet \& Glowinski [6], see also [3] for numerical experiments, and further improved by Glowinski \& Pironneau [8].

Given $\lambda$, denote by $\left(\phi_{\lambda}, u_{\lambda}\right)$ the solution of $(1.2 a, b)$ with right hand side $f=0$. A linear mapping $L$ is defined by

$$
\lambda \mapsto\left(\phi_{\lambda}, u_{\lambda}\right) \mapsto-\frac{\partial u_{\lambda}}{\partial n} .
$$

Glowinski \& Pironneau [8] observed that the operator $L$ is $H^{-1 / 2}(\Gamma)$-elliptic for a smooth domain. The corresponding discrete operator $L_{h}$ reflects this property. Specifically, the matrix $L_{h}$ is positive definite and the spectral condition number $\kappa\left(L_{h}\right)=\lambda_{\max }\left(L_{h}\right) / \lambda_{\min }\left(L_{h}\right)$ grows as $h^{-1}$ [8], where $h$ is a mesh parameter. The discrete system is solved by the method of conjugate gradients [1].

In order to speed up the convergence, Glowinski \& Pironneau have already suggested to use the $H^{-1 / 2}(\Gamma)$-ellipticity for preconditioning. Following this idea we will provide a preconditioning matrix $C_{h}$, such that the resulting condition number becomes independent of the mesh size. The matrix $C_{h}$ is based on the inverse of the square root of a discretization of $-d^{2} / d s^{2}$ with homogeneous boundary conditions on each line segment $\Gamma_{k}$ of $\partial \Omega$.

The proof, which is postponed to the last sections, has the following structure. First, we will generalize the properties of the continuous operator, mentioned above, to the case of a convex, polygonal domain. Here, the dual spaces $H^{-1 / 2}\left(\Gamma_{i}\right)$ of $H_{00}^{1 / 2}\left(\Gamma_{i}\right), \Gamma_{i}$ being a line segment of $\partial \Omega$, are involved. Specifically, we distinguish between $\sum_{i} H^{-1 / 2}\left(\Gamma_{i}\right)$ and $H^{-1 / 2}(\Gamma)$. Next, we will show that the properties of $L$ carry over to the discrete operator $L_{h}$. Finally, we will prove that the inner product $\lambda_{h}^{T} C_{h} \lambda_{h}$ induces a norm which is equivalent to the $\sum_{i} H^{-1 / 2}\left(\Gamma_{i}\right)$-norm on the finite element space.

Numerical experiments which confirm the theoretical results are included.

## 2. PRELIMINARIES

We study a finite element discretization, which is based on the mixed variable formulation of (1.1) :

$$
\begin{array}{ll}
\int_{\Omega} \phi \psi d x-\int_{\Omega} \nabla u \nabla \psi d x=0, & \forall \psi \in H^{1}(\Omega)  \tag{2.1}\\
\int_{\Omega} \nabla \phi \nabla v d x=-\int_{\Omega} f v d x, & \forall v \in H_{0}^{1}(\Omega)
\end{array}
$$

In the numerical solution, the Sobolev spaces $H^{1}(\Omega)$ and $H_{0}^{1}(\Omega)$ are replaced by finite dimensional subspaces $X_{h} \subset H^{1}(\Omega)$ and $X_{0 h}=$ $X_{h} \cap H_{0}^{1}(\Omega)$. Specifically, let $X_{h}$ be the finite element space of continuous, piecewise linear polynomials on the given regular triangulation $\mathfrak{C}_{h}$ of $\Omega$. Let $R_{h}$ denote the $p$-dimensional subspace of $X_{h}$ spanned by those basis functions, which are associated with nodes on the boundary. Then

$$
\begin{equation*}
X_{h}=X_{0 h} \oplus R_{h} . \tag{2.2}
\end{equation*}
$$

Identifying each finite element function via the nodal basis with the associated coefficient vector, the discrete problem which corresponds to (2.1) is written in matrix-vector notation as

$$
\left(\begin{array}{ccc}
M_{11} & M_{12} & B_{0}  \tag{2.3}\\
M_{21} & M_{22} & -T^{T} \\
B_{0} & -T & 0
\end{array}\right)\binom{\varphi_{h}}{u_{h}}=\binom{0}{-f_{h}}
$$

The square-matrix $B_{0}$ represents the discretization of the Poisson equation with Dirichlet boundary condition.

With respect to the decomposition (2.2) we write $\varphi_{h}$ as $\varphi_{h}^{T}=\left(\varphi_{h 0}, \lambda_{h}\right)$. After eliminating the variables $\varphi_{h 0}$ and $u_{h}$ in (2.3) we obtain a positive definite linear system

$$
\begin{equation*}
L_{h} \lambda_{h}=b_{h}, \tag{2.4}
\end{equation*}
$$

where

$$
L_{h}=\left(T^{T} B_{0}^{-1}, I_{p}\right)\left(\begin{array}{ll}
M_{11} & M_{12}  \tag{2.5}\\
M_{21} & M_{22}
\end{array}\right)\binom{B_{0}^{-1} T}{I_{p}}
$$

Since the matrix $L_{h}$ is only given implicitly, we prefer an iterative procedure for the solution of (2.4), especially the method of conjugate gradients. Given $\lambda_{h} \in R_{h}$, the evaluation $L_{h} \lambda_{h}$ requires the solution of two discrete

Poisson equations. For this purpose a multigrid algorithm [9] is well suited (see [4, 11]).

We will be concerned with preconditioning techniques in the application of the cg-algorithm. Given a positive definite matrix $C_{h}$, the condition number of $L_{h}$ with respect to $C_{h}$ is given by $\kappa\left(C_{h}^{-1} L_{h}\right)$. Specifically, let $0<\alpha_{h}<\beta_{h}$ be constants such that

$$
\begin{equation*}
\alpha_{h} \lambda_{h}^{T} C_{h} \lambda_{h} \leqslant \lambda_{h}^{T} L_{h} \lambda_{h} \leqslant \beta_{h} \lambda_{h}^{T} C_{h} \lambda_{h} \tag{2.6}
\end{equation*}
$$

then $\kappa\left(C_{h}^{-1} L_{h}\right) \leqslant \beta_{h} / \alpha_{h}$. We will provide a preconditioning matrix $C_{h}$ such that the constants in (2.6) are independent of the mesh size $h$.

The finite element solution $\varphi_{\lambda_{h}} \in X_{h}$ of

$$
\begin{equation*}
\int_{\Omega} \nabla \varphi_{\lambda_{h}} \nabla v_{h} d x=0 \quad \forall v_{h} \in X_{0 h}, \quad \varphi_{\lambda_{h}}-\lambda_{h} \in X_{0 h} \tag{2.7}
\end{equation*}
$$

is called discrete harmonic. The coefficient vector is given by $\varphi_{\lambda_{h}}=$ $\left(\begin{array}{cc}B_{0}^{-1} & T \\ & I_{p}\end{array}\right) \lambda_{h}$. Hence, from (2.5) it follows that

$$
\begin{equation*}
\lambda_{h}^{T} L_{h} \lambda_{h}=\left\|\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)}^{2} \tag{2.8}
\end{equation*}
$$

Therefore, in order to prove (2.6) we will be concerned with a priori estimates of the $L_{2}$-norm of discrete harmonic functions in terms of their boundary values.

## 3. THE PRECONDITIONING MATRIX

When using the method of conjugate gradients for the solution of the linear system $L_{h} x=b$, the number of iteration steps required for a given accuracy grows as $\sqrt{\kappa\left(L_{h}\right)}=O\left(h^{-1 / 2}\right)$. In order to speed up the convergence, preconditioning techniques have turned out to be useful.

We shall now construct a preconditioning matrix $C_{h}$ such that the condition number $\kappa\left(C_{h}^{-1} L_{h}\right)$ is bounded independently of the mesh size. Since $\Omega$ is assumed to be polygonal, the boundary $\Gamma$ of $\Omega$ consists of a finite number of straight lines $\Gamma_{k}, 1 \leqslant k \leqslant r$. Let $p_{k}$ denote the number of the interior nodes of $\Gamma_{k}$. The number of all boundary nodes is equal to $p=\sum_{k=1}^{r}\left(p_{k}+1\right)$.

The $p_{k} \times p_{k}$-matrices

$$
\begin{equation*}
D_{k}=\operatorname{tridiag}[-1,2,-1] \tag{3.1}
\end{equation*}
$$

correspond to the usual three-point approximation of the differential operator $-d^{2} / d x^{2}$ with homogeneous Dirichlet boundary conditions on $\Gamma_{k}$. The eigenvectors and eigenvalues of these Toeplitz-matrices are explicitly given by

$$
\theta_{l}^{(k)}=\left(\frac{2}{p_{k}+1}\right)^{1 / 2}\left[\sin \left(l \frac{j \pi}{p_{k}+1}\right)\right]_{j=1}^{p_{k}}
$$

and

$$
\lambda_{l}^{(k)}=4 \sin ^{2}\left(\frac{l \pi}{p_{k}+1}\right) .
$$

Therefore, $D_{k}$ admits the factorization

$$
D_{k}=Q_{k} \Lambda_{k} Q_{k}^{T}
$$

where $Q_{k}=\left[\theta_{1}^{(k)}, \theta_{2}^{(k)}, \ldots, \theta_{p_{k}}^{(k)}\right] \quad$ is unitary and $\Lambda_{k}=\operatorname{diag}\left(\lambda_{1}^{(k)}\right.$, $\lambda_{2}^{(k)}, \ldots, \lambda_{p_{k}}^{(k)}$. The powers $D_{k}^{s}, s \in \mathbb{R}$, are defined by

$$
\begin{equation*}
D_{k}^{s}=Q_{k} \Lambda_{k}^{s} Q_{k}^{T} \tag{3.2}
\end{equation*}
$$

Using Fast Fourier Transform (FFT), the evaluation of $D_{k}^{s} x$ requires only $O\left(p_{k} \ln p_{k}\right)$ arithmetic operations [16], provided that $p_{k}=s .2^{t}$ with $s$ being small.

We shall also need the tridiagonal $p_{k} \times p_{k}$-mass matrices

$$
\begin{equation*}
M_{k}=\left(\int_{\Gamma_{k}} \psi_{i} \psi_{j} d x\right)_{i j=1}^{p_{k}} \tag{3.3}
\end{equation*}
$$

on $T_{k}$, where $\psi_{i}(x)$ denotes the piecewise linear nodal basis function, which satisfies $\psi_{i}\left(x_{j}^{k}\right)=\delta_{i j}, 1 \leqslant i \leqslant p_{k}$ for the nodes $x_{j}^{k}$ on $\Gamma_{k}$. Set

$$
\begin{equation*}
C_{k}:=M_{k} D_{k}^{-1 / 2} M_{k} \tag{3.4}
\end{equation*}
$$

For preconditioning, we choose the $p \times p$ matrix $C_{h}$, which has block diagonal form :

$$
C_{h}=\left[\begin{array}{lllllll}
C_{1} & & & & & &  \tag{3.5}\\
& h^{2} & & & & & \\
& & C_{2} & & & & \\
& & & h^{2} & & & \\
& & & & \ddots & & \\
& & & & & C_{r} & \\
& & & & & & h^{2}
\end{array}\right]
$$

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The application of the preconditioning (3.5) requires for $k=1,2, \ldots, r$ two real sine transformations and the solution of two linear systems with the tridiagonal mass matrix $M_{k}$. Since $M_{k}$ is spectrally equivalent to $I_{p_{k}}$, one might expect at first glance, that $M_{k} D_{k}^{-1 / 2} M_{k}$ and $D_{k}^{-1 / 2}$ are spectrally equivalent, too. Indeed, if the meshpoints are distributed equidistantly, then the associated matrices

$$
D_{k}=\operatorname{tridiag}[-1,2,-1], \quad M_{k}=\frac{h_{k}}{6} \text { tridiag }[1,4,1]
$$

$h_{k}=1 /\left(p_{k}+1\right)$, have the same eigenvectors. Thus, the matrices $M_{k}$ and $D_{k}^{-1 / 4}$ commute, i.e.

$$
M_{k} D_{k}^{-1 / 2} M_{k}=D_{k}^{-1 / 4} M_{k}^{2} D_{k}^{-1 / 4}
$$

and we obtain

$$
\frac{1}{9} D_{k}^{-1 / 2} \leqslant \frac{1}{h_{k}^{2}} M_{k} D_{k}^{-1 / 2} M_{k} \leqslant D_{k}^{-1 / 2}
$$

However, if the meshpoints on $\Gamma_{k}$ are not distributed equidistantly, then $D_{k}$ and $M_{k}$ do not commute. In this case the conjecture is not always true, as is illustrated by the following example.

Example: Consider the matrices

$$
D=\left(\begin{array}{ll}
a & 0 \\
0 & 1
\end{array}\right), \quad M=\left(\begin{array}{ll}
2 & 1 \\
1 & 2
\end{array}\right)
$$

with $a \gg 1$. The diagonal dominant matrix $M$ has the eigenvalues $\lambda_{1}=1$ and $\lambda_{2}=3$. Set $x=(1,-2)^{T}$. Then

$$
x^{T} D x=a+4, \quad x^{T} M D M x=9 .
$$

Now we state our main result.
THEOREM 3.1: Let $\Omega$ be convex and let $C_{h}$ be the preconditioning matrix (3.5). Then the condition number $\kappa\left(C_{h}^{-1} L_{h}\right)$ of $L_{h}$ with respect to $C_{h}$ is bounded independently of the mesh size $h$.

We finish this section with the following remark. The condition number $\kappa\left(C_{h}^{-1} L_{h}\right)$ depends on the interior angles $\omega_{k}$ of the polygon $\Omega$ and grows as $\max \left\{\frac{1}{\sin ^{2} \omega_{k}}\right\}$. Angles, which are close to zero or $\pi$ spoil the condition number. Specifically, Theorem 3.1 does not apply if $\Omega=\Omega_{h}$ is the polygonal approximation of a smooth domain. In that case another preconditioning is appropriate : Consider on the whole boundary $\Gamma$ the three-point approxi-
mation $D_{h}$ of $I-d^{2} / d x^{2}$ with periodical boundary conditions [10] and set $C_{h}=M_{h} D_{h}^{-1 / 2} M_{h}$.

## 4. FURTHER PRECONDITIONING. THE BIHARMONIC EQUATION IN A SQUARE

The preconditioning matrix (3.5) has block diagonal form, each block $C_{k}$ corresponding to a line segment $\Gamma_{k}$. Therefore, the condition number of $L_{h}$ with respect to the block diagonal part of $L_{h}$ is again independent of the mesh size. In general, solving a linear system with the block diagonal part of $L_{h}$ is still expensive. However, if $\Omega$ is the unit square, then the block diagonal part is easy to invert, as we will describe in the remainder of this section. Thus, in this special case, another preconditioning matrix is available.

We approximate the biharmonic problem by piecewise linear elements on a subdivision of $\Omega$ into Courant's triangles $K_{h}$ of length $h, h=1 /(n+1)$. Using the quadrature rule

$$
\int_{K_{h}} \varphi(x) d x \approx \frac{\text { meas }\left(K_{h}\right)}{3}\left(\varphi\left(x_{1}\right)+\varphi\left(x_{2}\right)+\varphi\left(x_{3}\right)\right),
$$

when evaluating integrals, the mass matrix $M$ in (2.3) is replaced by the diagonal matrix $\tilde{M}=2 \mathrm{diag}(M)$. The same discretization results when the 13-point finite difference approximation is used. Inserting $\tilde{M}$ into (2.5) and neglecting the equations corresponding to the four corner points, we obtain

$$
\begin{equation*}
L_{h}=\frac{1}{2} I+T B_{0}^{-2} T^{T} \tag{4.1}
\end{equation*}
$$

We decompose the boundary space $R_{h}$ as

$$
R_{h}=R_{h}^{1} \oplus R_{h}^{2}
$$

where $R_{h}^{1}$ is spanned by those basis functions, which are associated with nodes on the lower and upper part of the boundary and $R_{h}^{2}$ is defined analogously. With respect to this decomposition the ( $4 n \times 4 n$ )-matrix $L_{h}$ has $2 \times 2$-block structure. The preconditioning by the block diagonal part diag ( $L_{11}, L_{22}$ ) is investigated. Each block $L_{i i}$ corresponds to the biharmonic problem with $\Delta u$ rather than $u_{n}$ specified on two opposite sides of the square $\Omega$. Bjorstad [2] has observed that this problem is easy to solve, since separation of the variables is possible. Assuming that $\Delta u$ is specified at the left and right part of the boundary, we choose a row-wise ordering of the
nodes. The resulting linear equations of the full problem have the following structure

$$
\left(\begin{array}{ccc}
h^{2} I_{n^{2}} & 0 & B_{0}  \tag{4.2}\\
0 & \frac{1}{2} h^{2} I_{2 n} & -U^{T} \otimes I_{n} \\
B_{0} & -U \otimes I_{n} & 0
\end{array}\right)\left(\begin{array}{l}
\varphi_{0} \\
\varphi_{1} \\
u
\end{array}\right)=\left(\begin{array}{c}
0 \\
0 \\
-h^{2} f_{h}
\end{array}\right)
$$

with $U=\left[e_{1}, e_{n}\right]$. After eliminating the variable $\varphi$, we obtain the positive definite system

$$
A u=h^{4} f_{h}
$$

with

$$
\begin{equation*}
A=B_{0}^{2}+2\left(U U^{T} \otimes I\right) \tag{4.3}
\end{equation*}
$$

Since the inverse of $L_{11}$ can be expressed via the inverse of the matrix $A$ as

$$
\begin{equation*}
L_{11}^{-1}=\frac{2}{h}\left[2\left(U^{T} \otimes I\right) A^{-1}(U \otimes I)-I\right] \tag{4.4}
\end{equation*}
$$

we will study the solution of a linear system with the matrix $A$. The discretization $B_{0}$ of the two-dimensional Laplacian on the unit square can be expressed via the approximation of the one-dimensional Laplacian

$$
D=\operatorname{tridiag}[-1,2,-1]
$$

as

$$
B_{0}=I \otimes D+D \otimes I
$$

with $I=I_{n}$. Using the spectral decomposition of $D$, i.e.

$$
D=Q \Lambda Q
$$

with $\Lambda=\operatorname{diag}\left(\lambda_{i}\right)$ and $Q=\left[\theta_{1}, \theta_{2}, \ldots, \theta_{n}\right]$ as defined in $\S 3$, the matrix $A$ can be written as

$$
\begin{aligned}
A & =I \otimes D^{2}+2(D \otimes D)+D^{2} \otimes I+2\left(U U^{T} \otimes I\right) \\
& =(I \otimes Q)\left[I \otimes \Lambda^{2}+2(D \otimes \Lambda)+\left(D^{2} \otimes I\right)+2\left(U U^{T} \otimes I\right)\right](I \otimes Q)
\end{aligned}
$$

each block corresponding to a row. Using the permutation $P$, which converts the row-wise ordering into a column-wise ordering, Bjørstad [2] obtains

$$
\begin{array}{r}
A=(I \otimes Q) P\left[\Lambda^{2} \otimes I+2(\Lambda \otimes D)+I \otimes D^{2}+2\left(I \otimes U U^{T}\right)\right] \times  \tag{4.5}\\
\times P(I \otimes Q)=:(I \otimes Q) P S P(I \otimes Q)
\end{array}
$$

The matrix $S=\operatorname{diag}\left(S_{i}\right)$ is block diagonal with $(n \times n)$-matrices

$$
S_{i}=\lambda_{i}^{2} I+2 \lambda_{i} D+D^{2}+2 U U^{T}
$$

having bandwidth $d=2$ and therefore being easily invertable.
Inserting (4.5) into (4.4), we obtain

$$
L_{11}^{-1}=\frac{2}{h}\left[2\left(U^{T} \otimes Q\right) P S^{-1} P(U \otimes Q)-I\right]
$$

Therefore, the evaluation of $L_{11}^{-1} r_{1}$ involves four sine transformations of length $n$ and the solution of $a$ linear system with the $n^{2} \otimes n^{2}$-matrix $S$, which is pentadiagonal.

## 5. NUMERICAL RESULTS

We will provide some numerical results for the biharmonic equation on the unit square. We use Courant's triangulation with triangles of length $h, h, \sqrt{2} h$. In order to determine the unknown boundary value $\lambda=\left.\Delta u\right|_{\partial \Omega}$, we solve the system of linear equations

$$
\begin{equation*}
L_{h} \lambda_{h}=b_{h} \tag{5.1}
\end{equation*}
$$

using the method of conjugate gradients. The evaluation of $L_{h}$ requires the solution of two discrete Poisson equations. Since $\Omega$ is the unit square, we have used Buneman's algorithm for this purpose.

The performance of the preconditioning techniques is studied by choosing the right hand side as

$$
b_{h}(x, y)=\left\{\begin{array}{lc}
\sin (\pi x)+\sin \left(\pi x h^{-1} / 2\right) & \text { on } \quad\{(x, 0), 0 \leqslant x \leqslant 1\} \\
0 & \text { on } \quad \partial \Omega \backslash\{(x, 0), 0 \leqslant x \leqslant 1\}
\end{array}\right.
$$

which is a superposition of low and high frequencies. The starting value is $\lambda_{h}^{(0)}=0$. The iteration is terminated, if the relative error of the residuum with respect to the Euclidian norm is less than $\varepsilon$, i.e.

$$
R_{k}=\frac{\left\|L_{h} \lambda_{h}^{(k)}-b_{h}\right\|}{\left\|b_{h}\right\|} \leqslant \varepsilon
$$

Without preconditioning the number of iteration steps required to gain a given accuracy $\varepsilon$ is bounded by $O\left(h^{-1 / 2} \log \varepsilon^{-1}\right)$. This is confirmed by the following table.
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Table 1
Number of cg-iterations for a given accuracy $\varepsilon$

| mesh size $h$ |  |  |  |
| :---: | ---: | :---: | :---: |
| accuracy $\varepsilon$ | $1 / 16$ | $1 / 32$ | $1 / 64$ |
| $10^{-3}$ | 7 | 9 | 12 |
| $10^{-6}$ | 12 | 16 | 21 |

The next table shows the independence of the number of pcg-iterations on the mesh size $h$, when using the preconditioning (3.5).

Table 2
Number of pcg-iterations with $C_{h} \sim D^{-1 / 2}$

| mesh size $h$ |  |  |  |
| :---: | ---: | ---: | :---: |
| accuracy $\varepsilon$ | $1 / 16$ | $1 / 32$ | $1 / 64$ |
| $\varepsilon=10^{-3}$ | 7 | 7 | 7 |
| $\varepsilon=10^{-6}$ | 12 | 13 | 13 |

Finally, we present the results when using the block diagonal part of $L_{h}$ for preconditioning. Note, however, that this preconditioning is only available in case of a rectangular domain $\Omega$.

Table 3
Number of pcg-iterations with $C_{h}=$ blockdiag $\left(L_{h}\right)$

| mesh size $h$ |  |  |  |
| :---: | :---: | :---: | :---: |
| accuracy $\varepsilon$ | $1 / 16$ | $1 / 32$ | $1 / 64$ |
| $\varepsilon=10^{-3}$ | 6 | 7 | 7 |
| $\varepsilon=10^{-6}$ | 9 | 10 | 10 |

## 6. A PRIORI ESTIMATES

The rest of this paper is concerned with the proof that the condition number $\kappa\left(C_{h}^{-1} L_{h}\right)$ is independent of the mesh size, i.e.

$$
\begin{equation*}
\alpha \lambda_{h}^{T} C_{h} \lambda_{h} \leqslant \lambda_{h}^{T} L_{h} \lambda_{h} \leqslant \beta \lambda_{h}^{T} C_{h} \lambda_{h}, \quad \lambda_{h} \in R_{h} \tag{6.1}
\end{equation*}
$$

with $\alpha, \beta$ being independent of $h$. Here, $C_{h}$ is the preconditioning matrix (3.5).

Let $\varphi_{\lambda_{h}} \in X_{h}$ be the discrete harmonic function with boundary value $\lambda_{h}$, i.e. the solution of (2.7). Using (2.8) and the notation

$$
\left\|\lambda_{h}\right\|_{C_{h}}:=\sqrt{\lambda_{h}^{T} C_{h} \lambda_{h}}
$$

inequality (6.1) is rewritten as

$$
\begin{equation*}
\alpha\left\|\lambda_{h}\right\|_{C_{h}} \leqslant\left\|\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} \leqslant \beta\left\|\lambda_{h}\right\|_{C_{h}} . \tag{6.1'}
\end{equation*}
$$

In order to prove (6.1), we consider the continuous case at first. Let

$$
H_{00}^{1 / 2}\left(\Gamma_{i}\right)=\left[H_{0}^{1}\left(\Gamma_{i}\right), L_{2}\left(\Gamma_{i}\right)\right]_{1 / 2}
$$

denote the interpolation space [10], and let

$$
H^{-1 / 2}\left(\Gamma_{i}\right)=\left(H_{00}^{1 / 2}\left(\Gamma_{i}\right)\right)^{\prime}
$$

denote the dual space. Set

$$
\begin{equation*}
\|\lambda\|_{-1 / 2, \Gamma}:=\left(\sum_{t=1}^{r}\|\lambda\|_{H^{-1 / 2}\left(\Gamma_{t}\right)}^{2}\right)^{1 / 2} \tag{6.2}
\end{equation*}
$$

In the proof of the following theorem, we will make use of a trace theorem given in the appendix.

THEOREM 6.1: Let the polygonal domain $\Omega \subset \mathbb{R}^{2}$ be convex and $\lambda \in H^{1 / 2}(\Gamma)$. Then the $L_{2}$-norm of the harmonic function $\phi_{\lambda}$ can be estimated from below and from above by the $\|\cdot\|_{-1 / 2, \Gamma}$-norm of it's boundary value $\lambda$ :

$$
\begin{equation*}
c_{0}\|\lambda\|_{-1 / 2, \Gamma} \leqslant\left\|\phi_{\lambda}\right\|_{L_{2}(\Omega)} \leqslant c_{1}\|\lambda\|_{-1 / 2, \Gamma} \tag{6.3}
\end{equation*}
$$

Proof: Since $\lambda \in H^{1 / 2}(\Gamma)$, we have $\phi_{\lambda} \in H^{1}(\Omega)$. By partial integration we obtain for all $u \in H_{0}^{1}(\Omega) \cap H^{2}(\Omega)$

$$
\begin{align*}
\int_{\Omega}(-\Delta u) \phi_{\lambda} d x & =\int_{\Omega} \nabla u \nabla \phi_{\lambda} d x-\int_{\Gamma} \frac{\partial u}{\partial n} \lambda d s  \tag{6.4}\\
& =-\int_{\Gamma} \frac{\partial u}{\partial n} \lambda d s .
\end{align*}
$$

We will first prove the second inequality of (6.3). Let $u_{\lambda} \in H_{0}^{1}(\Omega)$ denote the solution of the Poisson equation (1.2b) with right hand side $\phi_{\lambda}$. Since $\Omega$ is convex, regularity theory ensures that $u_{\lambda} \in H_{0}^{1} \cap H^{2}(\Omega)$ and that

$$
\begin{equation*}
\left\|u_{\lambda}\right\|_{H^{2}(\Omega)} \leqslant c_{2}\left\|\phi_{\lambda}\right\|_{L_{2}(\Omega)} \tag{6.5}
\end{equation*}
$$

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Furthermore, the trace theorem given in the appendix states that $\frac{\partial u_{\lambda}}{\partial n} \in H_{00}^{1 / 2}\left(\Gamma_{l}\right)$ and that

$$
\begin{equation*}
\left(\sum_{i=1}^{r}\left\|\frac{\partial u_{\lambda}}{\partial n}\right\|_{H_{00}^{1 / 2}\left(\Gamma_{t}\right)}^{2}\right)^{1 / 2} \leqslant c_{3}\left\|u_{\lambda}\right\|_{H^{2}(\Omega)} \tag{6.6}
\end{equation*}
$$

Inserting (6.6) and (6.5) into (6.4) yields

$$
\left\|\phi_{\lambda}\right\|_{L_{2}(\Omega)}^{2}=\int_{\Omega}\left(-\Delta u_{\lambda}\right) \phi_{\lambda} d x \leqslant c_{2} c_{3}\|\lambda\|_{-1 / 2, \Gamma}\left\|\phi_{\lambda}\right\|_{L_{2}(\Omega)}
$$

which proves the second inequality of (6.3).
Next, let $\rho_{i} \in H_{00}^{1 / 2}\left(\Gamma_{l}\right)$ such that $\left\|\rho_{l}\right\|_{H_{10}^{1 / 2}\left(\Gamma_{l}\right)}=1$ and

$$
\begin{equation*}
\|\lambda\|_{-1 / 2, \Gamma} \leqslant c_{4} \int_{\Gamma_{t}} \lambda \rho_{t} d s \tag{6.7}
\end{equation*}
$$

By the trace theorem, there is $u \in H_{0}^{1}(\Omega) \cap H^{2}(\Omega)$ such that $\left.\frac{\partial u}{\partial n}\right|_{\Gamma_{t}}=\rho_{l}$, $\left.\frac{\partial u}{\partial n}\right|_{\Gamma_{l}}=0, l \neq i$ and

$$
\begin{equation*}
\|u\|_{H^{2}(\Omega)} \leqslant c_{5}\left\|\rho_{l}\right\|_{H_{00}^{1 / 2}\left(\Gamma_{1}\right)}=c_{5} . \tag{6.8}
\end{equation*}
$$

From (6.4) we get

$$
\begin{align*}
\int_{\Gamma_{t}} \rho_{l} \lambda d s & =\int_{\Omega}(-\Delta u) \phi_{\lambda} d x  \tag{6.9}\\
& \leqslant\|u\|_{H^{2}(\Omega)}\left\|\phi_{\lambda}\right\|_{L_{2}(\Omega)}
\end{align*}
$$

Inserting (6.9) and (6.8) into (6.7) yields the first inequality of (6.3).
The estimate for the continuous case may be carried over to the discrete case (see also [12], [13] for similar arguments used in different contexts).

THEOREM 6.2: Let $\Omega$ be convex. Given $\lambda_{h} \in R_{h}$, let $\varphi_{\lambda_{h}}$ be the discrete harmonic function with boundary value $\lambda_{h}$, i.e. the solution of (2.7). Then there are positive constants $\alpha, \beta$, such that

$$
\begin{align*}
\alpha\left(\left\|\lambda_{h}\right\|_{-1 / 2, \Gamma}+h^{1 / 2}\left\|\lambda_{h}\right\|_{L_{2}(\Gamma)}\right) & \leqslant\left\|\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)}  \tag{6.10}\\
& \leqslant \beta\left(\left\|\lambda_{h}\right\|_{-1 / 2, \Gamma}+h^{1 / 2}\left\|\lambda_{h}\right\|_{L_{2}(\Gamma)}\right)
\end{align*}
$$

Proof: Let $\phi_{\lambda_{h}} \in H^{1}(\Omega)$ denote the harmonic function with boundary value $\lambda_{h}$. By the approximation properties of the finite element space $X_{h}$, we see that ([8], p. 184)

$$
\begin{equation*}
\left\|\phi_{\lambda_{h}}-\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} \leqslant c_{1} h^{1 / 2}\left\|\lambda_{h}\right\|_{L_{2}(\Gamma)} \tag{6.11}
\end{equation*}
$$

Therefore, Theorem 6.1 implies

$$
\begin{aligned}
\left\|\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} & \leqslant\left\|\phi_{\lambda_{h}}\right\|_{L_{2}(\Omega)}+\left\|\varphi_{\lambda_{h}}-\phi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} \\
& \leqslant \beta\left(\left\|\lambda_{h}\right\|_{-1 / 2, \Gamma}+h^{1 / 2}\left\|\lambda_{h}\right\|_{L_{2}(\Gamma)}\right),
\end{aligned}
$$

which proves the second inequality of (6.10).
Next we will prove the first inequality of (6.10). Since the triangulation is regular, a simple scaling argument yields

$$
\begin{equation*}
\left\|\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} \geqslant c_{2} h^{1 / 2}\left\|\lambda_{h}\right\|_{L_{2}(\mathrm{\Gamma})} \tag{6.12}
\end{equation*}
$$

Put $\eta:=h^{1 / 2}\left\|\lambda_{h}\right\|_{L_{2}(\Gamma)} /\left\|\lambda_{h}\right\|_{-1 / 2, \Gamma}$. From (6.3) and (6.11) we conclude that

$$
\begin{align*}
\left\|\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} & \geqslant\left\|\phi_{\lambda_{h}}\right\|_{L_{2}(\Omega)}-\left\|\varphi_{\lambda_{h}}-\phi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} \\
& \geqslant c_{3}\left\|\lambda_{h}\right\|_{-1 / 2, \Gamma}-c_{1} h^{1 / 2}\left\|\lambda_{h}\right\|_{L_{2}(\Gamma)}  \tag{6.13}\\
& =\left(c_{3}-c_{1} \eta\right)\left\|\lambda_{h}\right\|_{-1 / 2, \Gamma} .
\end{align*}
$$

Using (6.12) and (6.13) we obtain

$$
\begin{equation*}
\left\|\varphi_{\lambda_{h}}\right\|_{L_{2}(\Omega)} \geqslant \max \left\{c_{2} \eta, c_{3}-c_{1} \eta\right\}\left\|\lambda_{h}\right\|_{-1 / 2, \Gamma} \tag{6.14}
\end{equation*}
$$

Since

$$
\min _{\eta>0} \max \left\{c_{2} \eta, c_{3}-c_{1} \eta\right\}=\frac{c_{3} c_{2}}{c_{1}+c_{2}}:=\alpha>0
$$

we get the result as stated.

## 7. DISCRETE NORMS AND MATRIX-REPRESENTATIONS

In order to apply Theorem 6.2 for the proof of (6.1), we must verify that the norm $\left(\|\cdot\|_{-1 / 2, \Gamma}^{2}+h\|\cdot\|_{L_{2}(\Gamma)}^{2}\right)^{1 / 2}$ is represented by the matrix $C_{h}$ defined by (3.5). More precisely, we will prove that the norms $\|\cdot\|_{C_{h}}$ and $\left(\|\cdot\|_{-1 / 2, \Gamma}^{2}+h\|\cdot\|_{L_{2}(\Gamma)}^{2}\right)^{1 / 2}$ are equivalent on $R_{h}$ with constants being indevol. $22, n^{\circ} 4,1988$
pendent of the mesh size. Let $R_{h}^{0} \subset R_{h}$ denote the subspace consisting of those piecewise linear functions which vanish at the corners of $\partial \Omega$. Then $v_{h} \in R_{h}^{0}$ if and only if $v_{h} \in R_{h}$ and $\left.v_{h}\right|_{\Gamma_{k}} \in H_{0}^{1}\left(\Gamma_{k}\right), 1 \leqslant k \leqslant r$.

### 7.1. Matrix-Representation of the $\boldsymbol{H}_{00}^{1 / 2}\left(\Gamma_{k}\right)$-norm on $\left.R_{h}^{0}\right|_{\Gamma_{k}}$

We denote the nodes on the line segment $\Gamma_{k}={\bar{P}{ }_{k} P_{k+1}}$ by

$$
P_{k}=x_{0}^{k}<x_{1}^{k}<\cdots<x_{p_{k}}^{k}<x_{p_{k}+1}^{k}=P_{k+1} .
$$

Let $\left.v_{h} \in R_{h}^{0}\right|_{\Gamma_{k}}$. Then

$$
\left|v_{h}\right|_{H^{1}\left(\Gamma_{k}\right)}^{2}=\sum_{l=0}^{p_{k}} \frac{1}{\left|x_{l+1}^{k}-x_{l}^{k}\right|}\left(v_{h}\left(x_{l+1}^{k}\right)-v_{h}\left(x_{l}^{k}\right)\right)^{2} .
$$

Since the triangulation is regular, i.e. $\sigma h \leqslant\left|x_{l+1}^{k}-x_{l}^{k}\right| \leqslant h$, the $H^{1}$-norm on $\left.R_{h}^{0}\right|_{\Gamma_{k}}$ is equivalent to

$$
\left[\frac{1}{h} \sum_{l=0}^{p_{k}}\left(v_{h}\left(x_{l+1}^{k}\right)-v_{h}\left(x_{l}^{k}\right)\right)^{2}\right]^{1 / 2},
$$

and the associated bilinear form is represented by the tridiagonal $p_{k} \times p_{k}$-matrix

$$
\begin{equation*}
\frac{1}{h} D_{k}=\frac{1}{h} \operatorname{tridiag}[-1,2,-1] \tag{7.1}
\end{equation*}
$$

By interpolation we obtain the following
Proposition 7.1 : The norms

$$
\left\|v_{h}\right\|_{\left.H_{00}^{1 / 2} \Gamma_{k}\right)} \text { and }\left\|v_{h}\right\|_{1 / 2, \Gamma_{k}}:=\sqrt{v_{h}^{T} D_{k}^{1 / 2} v_{h}}
$$

are equivalent on $\left.R_{h}^{0}\right|_{\Gamma_{k}}$ with constants being independent of the mesh size $h$.
Proof: Let $s \geqslant 0$. Set

$$
\begin{equation*}
\left\|\left\|v_{h}\right\|\right\|_{s, \Gamma_{k}}:=\left[h v_{h}^{T}\left(\frac{1}{h^{2}} D_{k}\right)^{s} v_{h}\right]^{1 / 2} \tag{7.2}
\end{equation*}
$$

Then the imbeddings

$$
i:\left(\left.R_{h}^{0}\right|_{\Gamma_{k}},\| \| \cdot\| \|_{0, \Gamma_{k}}\right) \rightarrow L_{2}\left(\Gamma_{k}\right)
$$

and

$$
i:\left(\left.R_{h}^{0}\right|_{\Gamma_{k}}, \|\left|\left|\left|\left|| |_{1, \Gamma_{k}}\right) \rightarrow H_{0}^{1}\left(\Gamma_{k}\right)\right.\right.\right.\right.
$$

are continuous with constants $c_{0}$ and $c_{1}$, resp. Therefore, the interpolation theorem [10] yields the continuity of $i$ between the interpolated spaces. Specifically, in case $s=1 / 2$ we obtain

$$
\begin{equation*}
\left\|v_{h}\right\|_{H_{00}^{1 / 2}\left(\Gamma_{k}\right)} \leqslant\left(c_{0} c_{1}\right)^{1 / 2}\| \| v_{h} \| \frac{1}{2}, \Gamma_{k},\left.\quad v_{h} \in R_{h}^{0}\right|_{\Gamma_{k}} \tag{7.3}
\end{equation*}
$$

Next, we consider the $L_{2}$-projection $p_{0}$ onto $\left.R_{h}^{0}\right|_{\Gamma_{k}}$. Obviously $p_{0}: L_{2}\left(\Gamma_{k}\right) \rightarrow$ $\left(\left.R_{h}^{0}\right|_{\Gamma_{k}},\| \| \cdot\| \|_{0, \Gamma_{k}}\right.$ ) is continuous with constant $c_{2}$. Furthermore, we claim that $p_{0}: H_{0}^{1}\left(\Gamma_{k}\right) \rightarrow\left(\left.R_{h}^{0}\right|_{\Gamma_{k}},\| \| .\| \|_{1, \Gamma_{k}}\right)$ is continuous. Indeed, let $v \in H_{0}^{1}\left(\Gamma_{k}\right)$. Since $v$ is a continuous function, the interpolant $\left.J_{h} v \in R_{h}^{1}\right|_{\Gamma_{k}}, J_{h} v\left(x_{l}^{k}\right)=$ $v\left(x_{l}^{k}\right)$, is well defined and

$$
\left\|J_{h} v\right\|_{H^{1}\left(\Gamma_{k}\right)} \leqslant c\|v\|_{H^{1}\left(\Gamma_{k}\right)}
$$

Therefore, using approximation properties and inverse inequalities, we have

$$
\begin{aligned}
c_{3}\left\|p_{0} v\right\| \|_{1, \Gamma_{k}} & \leqslant\left\|p_{0} v\right\|_{H^{1}\left(\Gamma_{k}\right)} \\
& \leqslant\|v\|_{H^{1}}+\left\|J_{h} v-v\right\|_{H^{1}}+\left\|p_{0} v-J_{h} v\right\|_{H^{1}} \\
& \leqslant c\|v\|_{H^{1}}+c h^{-1}\left(\left\|p_{0} v-v\right\|_{L_{2}}+\left\|v-J_{h} v\right\|_{L_{2}}\right) \\
& \leqslant c\|v\|_{H^{1}} .
\end{aligned}
$$

Thus, the projection mappings are continuous and from the interpolation theorem we obtain for $s=1 / 2$

$$
\begin{equation*}
\left\|p_{0} v\right\|_{1 / 2, \Gamma_{k}} \leqslant c\|v\|_{H_{00}^{12}\left(\Gamma_{k}\right)} \tag{7.4}
\end{equation*}
$$

### 7.2. Matrix-Representation of the $H^{-1 / 2}\left(\Gamma_{k}\right)$-norm on $\left.R_{h}^{0}\right|_{\Gamma_{k}}$

By definition,

$$
\begin{equation*}
\|u\|_{H^{-1 / 2}\left(\Gamma_{k}\right)}=\sup _{v \in H_{00}^{1 / 2}\left(\Gamma_{k}\right)} \frac{\int_{\Gamma_{k}} u v d s}{\|v\|_{H_{00}^{12}\left(\Gamma_{k}\right)}} . \tag{7.5}
\end{equation*}
$$

If $\left.u \in R_{h}^{0}\right|_{\Gamma_{k}}$, then we will see that it is sufficient to take the supremum over the subspace $\left.R_{h}^{0}\right|_{\Gamma_{k}}$ :
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PROPOSITION 7.2 : The norms

$$
\left\|u_{h}\right\|_{H^{-1 / 2}\left(\Gamma_{k}\right)} \text { and }\left\|\left\|u_{h}\right\|\right\|_{-1 / 2, \Gamma_{k}}:=\sup _{\left.v_{h} \in R_{h}^{0}\right|_{\Gamma_{k}}} \frac{\int_{\Gamma_{k}} u_{h} v_{h} d s}{\left\|v_{h}\right\|_{1 / 2, \Gamma_{k}}}
$$

are equivalent on $\left.R_{h}^{0}\right|_{\Gamma_{k}}$.
Proof: By (7.4) the $L_{2}$-projection

$$
p_{0}: H_{00}^{1 / 2}\left(\Gamma_{k}\right) \rightarrow\left(\left.R_{h}^{0}\right|_{\Gamma_{k}},\| \|\| \|_{1 / 2, \Gamma_{k}}\right)
$$

is continuous. Let $\left.u_{h} \in R_{h}^{0}\right|_{\Gamma_{k}}$ and $v \in H_{00}^{1 / 2}\left(\Gamma_{k}\right)$. Then

$$
\int_{\mathrm{I}_{k}} u_{h} v d s=\int_{\Gamma_{k}} u_{h} p_{0} v d s \leqslant c \frac{\left|\int_{\Gamma_{k}} u_{h} p_{0} v d s\right|}{\left\|p_{0} v\right\|_{1 / 2, \Gamma_{k}}}\|v\|_{\left.H_{0_{0}\left(\Gamma_{k}\right)}\right)} .
$$

This implies

$$
\left\|u_{h}\right\|_{H^{-1 / 2}\left(\Gamma_{k}\right)} \leqslant c \sup _{\left.v_{h} \in R_{h}^{0}\right|_{\Gamma_{k}}} \int_{\Gamma_{k}} u_{h} v_{h} d s v_{h} \|_{1 / 2, \Gamma_{k}} .
$$

The reversed incquality follows from the inclusion $\left.R_{h}^{0}\right|_{\Gamma_{k}} \subset H_{00}^{1 / 2}\left(\Gamma_{k}\right)$.
Using the $p_{k} \times p_{k}$-mass matrix $M_{k}$ and the Euclidian norm $\|\cdot\|$ we can write

$$
\begin{align*}
\left\|\left\|u_{h}\right\|\right\|_{-1 / 2, \Gamma_{k}} & =\sup _{\left.v_{h} \in R_{h}^{0}\right|_{\Gamma_{k}}} \frac{\int_{r_{k}} u_{h} v_{h} d s}{\left\|v_{h}\right\|_{1 / 2, \Gamma_{k}}} \\
& =\sup _{v_{h} \in \mathbb{R}^{p_{k}}} \frac{u_{h}^{T} M_{k} v_{h}}{\left\|D_{k}^{1 / 4} v_{h}\right\|}  \tag{7.6}\\
& =\sup _{z_{h} \in \mathbb{R}^{p_{k}}} \frac{u_{h}^{T} M_{k} D_{k}^{-1 / 4} z_{h}}{\left\|z_{h}\right\|} \\
& =\left\|D_{k}^{-1 / 4} M_{k} u_{h}\right\| .
\end{align*}
$$

Thus, the $H^{-1 / 2}\left(\Gamma_{k}\right)$-norm on $\left.R_{h}^{0}\right|_{\Gamma_{k}}$ is induced by the $p_{k} \times p_{k}$-matrix

$$
\begin{equation*}
C_{k}=M_{k} D_{k}^{-1 / 2} M_{k} . \tag{7.7}
\end{equation*}
$$

### 7.3. Matrix-Representation of the $H^{-1 / 2}\left(\Gamma_{k}\right)$-norm on $\left.R_{h}\right|_{\Gamma_{k}}$

Finally, we consider the general case that $\left.v_{h} \in R_{h}\right|_{\Gamma_{k}}$ does not necessarily vanish at the endpoints of $\Gamma_{k}$. We claim that the norm $\left(\left\|v_{h}\right\|_{H^{-1 / 2}\left(\Gamma_{k}\right)}^{2}+\right.$ $\left.h\left\|v_{h}\right\|_{L_{2}\left(\Gamma_{k}\right)}^{2}\right)^{1 / 2}$ on $\left.R_{h}\right|_{\Gamma_{k}}$ is induced by the $\left(p_{k}+2\right) \times\left(p_{k}+2\right)$-matrix

$$
\left(\begin{array}{ccc}
h^{2} & &  \tag{7.8}\\
& C_{k} & \\
& & h^{2}
\end{array}\right)
$$

Here, we identify each funtion $\left.v_{h} \in R_{h}\right|_{\Gamma_{k}}$ with the vector of nodal values $\left(v_{0},\left(v_{1}, v_{2}, \ldots, v_{p_{k}}\right), v_{p_{k}+1}\right)^{T} \in \mathbb{R}^{p_{k}+2}$.

Using the $L_{2}$-projection $p_{0}$ onto $\left.R_{h}^{0}\right|_{\Gamma_{k}}$, we consider the decomposition

$$
\begin{equation*}
v=\left(v-p_{0} v\right)+p_{0} v \tag{7.9}
\end{equation*}
$$

By the usual approximation properties of the finite element space $\left.R_{h}^{0}\right|_{\Gamma_{k}}$, we obtain :

Lemma 7.3: Let $\left.v \in R_{h}\right|_{\Gamma_{k}}$. Then

$$
\begin{equation*}
\left\|v-p_{0} v\right\|_{H^{-12}\left(\Gamma_{k}\right)} \leqslant h^{1 / 2}\|v\|_{L_{2}\left(\Gamma_{k}\right)} \tag{7.10}
\end{equation*}
$$

Proof: Let $\phi \in H_{00}^{1 / 2}\left(\Gamma_{k}\right)$. Then

$$
\int_{\mathrm{r}_{k}}\left(v-p_{0} v\right) \phi d s=\int_{\mathrm{r}_{k}} v\left(\phi-p_{0} \phi\right) d s
$$

and

$$
\left\|\phi-p_{0} \phi\right\|_{L_{2}\left(\Gamma_{k}\right)} \leqslant c h^{1 / 2}\|\phi\|_{H_{00}^{1 / 2}\left(\Gamma_{k}\right)}
$$

This proves (7.10).
Let $v^{(1)}=\left(v_{1}, v_{2}, \ldots, v_{p_{k}}\right)$ denote the vector of nodal values associated with the interior nodes of $\Gamma_{k}, v^{(2)}:=\left(v_{0}, 0, \ldots, 0, v_{p_{k}+1}\right)^{T}$ and

$$
M_{12} v^{(2)}:=\left(h_{0} v_{0}, 0, \ldots, 0, h_{p_{k}} v_{p_{k}+1}\right) \in \mathbb{R}^{p_{k}}, \quad h_{l}=\left|x_{l+1}^{k}-x_{l}^{k}\right|
$$

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Then the coefficient vector of $\left.p_{0} v \in R_{h}^{0}\right|_{\Gamma_{k}}$ is given by

$$
p_{0} v=M_{k}^{-1}\left(M_{k} v^{(1)}+M_{12} v^{(2)}\right) .
$$

Thus, using (7.6)

$$
\begin{equation*}
\left\|p_{0} v\right\|_{-1 / 2, \Gamma_{k}=\left\|D_{k}^{-1 / 4} M_{k} v^{(1)}+D_{k}^{-1 / 4} M_{12} v^{(2)}\right\| . . . . ~}^{\text {. }} \tag{7.11}
\end{equation*}
$$

The second member of the right hand side of (7.11) can be estimated using

LEMMA 7.4: Let $v^{(2)}=\left(v_{0}, 0, \ldots, 0, v_{p_{k}+1}\right)$. Then

$$
\begin{equation*}
\left\|D_{k}^{-1 / 4} M_{12} v^{(2)}\right\| \leqslant h\left(v_{0}^{2}+v_{p_{k}+1}^{2}\right)^{1 / 2} . \tag{7.12}
\end{equation*}
$$

Proof: Hölder's inequality implies

$$
\begin{equation*}
\left\|D_{k}^{-1 / 4} M_{12} v^{(2)}\right\| \leqslant\left\|M_{12} v^{(2)}\right\|^{1 / 2}\left\|D_{k}^{-1 / 2} M_{12} v^{(2)}\right\|^{1 / 2} \tag{7.13}
\end{equation*}
$$

The solution $u$ of the linear system

$$
D_{k} u=M_{12} v^{(2)}
$$

is given by

$$
u_{j}=h_{0} v_{0}+\frac{h_{p_{k}} v_{p_{k}+1}-h_{0} v_{0}}{p_{k}+1} j, \quad 1 \leqslant j \leqslant p_{k}
$$

Therefore,

$$
\begin{align*}
\left\|D_{k}^{-1 / 2} M_{12} v^{(2)}\right\|^{2} & =u^{T} M_{12} v^{(2)} \\
& =h_{0}^{2} v_{0}^{2}+h_{p_{k}}^{2} v_{p_{k}+1}^{2}-\frac{1}{p_{k}+1}\left(h_{0} v_{0}-h_{p_{k}} v_{p_{k}+1}\right)^{2}  \tag{7.14}\\
& \leqslant c^{2}\left(v_{0}^{2}+v_{p_{k}+1}^{2}\right) .
\end{align*}
$$

Inserting (7.14) into (7.13) completes the proof.
Collecting the previous results we obtain
Proposition $7.5:$ Let $\left.v_{h} \in M_{h}\right|_{\Gamma_{k}}$. Then the norms

$$
\left(\left\|v_{h}\right\|_{H^{-1 / 2}\left(\Gamma_{k}\right)}^{2}+h\left\|v_{h}\right\|_{L_{2}\left(\mathrm{~F}_{k}\right)}^{2}\right)^{1 / 2}, \quad\left(v_{h}^{T}\left(\begin{array}{lll}
h^{2} & & \\
& C_{k} & \\
& & h^{2}
\end{array}\right) v_{h}\right)^{1 / 2}
$$

are equivalent.

Proof: Combining (7.9)-(7.12) yields

$$
\left\|v_{h}\right\|_{H^{-1 / 2}\left(r_{k}\right)}\left\{\begin{array}{l}
\leqslant\left\|D_{k}^{-1 / 4} M_{k} v^{(1)}\right\|+c_{2} h\left(\sum_{j=0}^{p_{k}+1} v_{j}^{2}\right)^{1 / 2} \\
\geqslant\left\|D_{k}^{-1 / 4} M_{k} v^{(1)}\right\|-c_{2} h\left(\sum_{j=0}^{p_{k}+1} v_{j}^{2}\right)^{1 / 2} .
\end{array}\right.
$$

Using the inverse inequality

$$
h\left(\sum_{j=1}^{p_{k}} v_{j}^{2}\right)^{1 / 2} \leqslant c\left\|D_{k}^{-1 / 4} M_{k} v^{(1)}\right\|
$$

completes the proof.

## APPENDIX : A TRACE THEOREM

We will establish the trace theorem which has been used in the proof of the a priori estimates given in section 6. Let $(a, b) \subset \mathbb{R}$. The Sobolev space $H_{00}^{1 / 2}(a, b)$, which is defined to be the interpolation space $\left[L_{2}(a, b), H_{0}^{1}(a, b)\right]_{12}$, has an explicit representation [10]:
(A.1) $H_{00}^{1 / 2}(a, b)=\left\{u \in H^{1 / 2}(a, b) ; \frac{u}{\sqrt{(b-x)(x-a)}} \in L_{2}(a, b)\right\}$,
the interpolation norm being equivalent to

$$
\begin{equation*}
\|u\|_{H_{0}^{112}(a, b)}:=\left(\|u\|_{H^{12}(a, b)}^{2}+\int_{a}^{b} \frac{|u(x)|^{2}}{(b-x)(x-a)} d x\right)^{1 / 2} \tag{A.2}
\end{equation*}
$$

In (A.1), (A.2) we have used the Sobolev space $H^{1 / 2}(a, b)$, which can be defined using the norm
(A.3) $\|u\|_{H^{s}(a, b)}:=\left(\|u\|_{L_{2}(a, b)}^{2}+\int_{a}^{b} \int_{a}^{b} \frac{|u(x)-u(y)|^{2}}{|x-y|^{1+2 s}} d x d y\right)^{1 / 2}$
with $s=1 / 2$. If $H_{0}^{s}(a, b), 0 \leqslant s \leqslant 1$, is defined to be the completion of $C_{0}^{\infty}(a, b)$ with respect to the norm (A.3), then $H_{0}^{1 / 2}(a, b)=H^{1 / 2}(a, b)$ and $H_{00}^{1 / 2}(a, b)$ is strictly contained in $H_{0}^{1 / 2}(a, b)$ [10]. This explains the additional zero in the indexing of $H_{00}^{1 / 2}(a, b)$.
The functions in $H_{00}^{1 / 2}(a, b)$ can be extended by 0 to functions in $H^{1 / 2}(\mathbb{R})$. The following is contained in [10].

Proposition A. $1:$ Let $u \in L_{2}(a, b)$ and let $\tilde{u}$ denote the extension of $u$ by zero. Then $\tilde{u} \in H^{1 / 2}(\mathbb{R})$ if and only if $u \in H_{00}^{1 / 2}(a, b)$. In addition, there are positive constants $\sigma_{1}, \sigma_{2}$ such that

$$
\begin{equation*}
\sigma_{1}\|\tilde{u}\|_{H^{1 / 2}(\mathbb{R})} \leqslant\|u\|_{H_{00}^{1 / 2}(a, b)} \leqslant \sigma_{2}\|\tilde{u}\|_{H^{12}(\mathbb{R})} \tag{A.4}
\end{equation*}
$$

for every $u \in H_{00}^{1 / 2}(a, b)$.
If $\Omega$ is smooth, then the trace operator

$$
\begin{aligned}
H^{2}(\Omega) & \rightarrow H^{3 / 2}(\Gamma) \times H^{1 / 2}(\Gamma) \\
u & \mapsto\left(u, \frac{\partial u}{\partial n}\right)
\end{aligned}
$$

is known to be continuous and surjective with continuous right inverse [10]. Generalizations to domains with corners are given by Yakovlev [16]. We only need a special case which is more easily established.

THEOREM A. 2 : Let $\Omega$ be a convex polygon.
a) Let $u \in H_{0}^{1} \cap H^{2}(\Omega)$. Then $\left.\frac{\partial u}{\partial n}\right|_{\Gamma_{i}} \in H_{00}^{1 / 2}\left(\Gamma_{i}\right)$ and

$$
\sum_{i=1}^{r}\left\|\frac{\partial u}{\partial n}\right\|_{H_{00}^{12\left(\Gamma_{i}\right)}} \leqslant \alpha\|u\|_{H^{2}(\Omega)}
$$

b) Let $\rho_{i} \in H_{00}^{1 / 2}\left(\Gamma_{i}\right)$. Then there is a function $u \in H_{0}^{1} \cap H^{2}(\Omega)$ satisfying $\left.\frac{\partial u}{\partial n}\right|_{\Gamma_{i}}=\rho_{i}$ and $\beta\|u\|_{H^{2}(\Omega)} \leqslant \sum_{i=1}^{r}\left\|\rho_{i}\right\|_{H_{00}^{1 / 2}\left(\Gamma_{i}\right)}$.

The constants $\alpha$ and $\beta$ depend on the interior angles $\omega_{i}$ of $\Omega$ and $\alpha / \beta$ grows as $\max _{1 \leqslant i \leqslant r}\left\{\frac{1}{\sin \left(\omega_{i}\right)^{2}}\right\}$.

Proof: Let us consider the special case $\hat{\Omega}=(0,1) \times(0,1)$ at first.
a) If $u \in H_{0}^{1} \cap H^{2}(\hat{\Omega})$, then $\left.\nabla u\right|_{\Gamma} \in H^{1 / 2}(\hat{\Gamma})^{2}$ and

$$
\|\nabla u\|_{H^{1 / 2}(\hat{\Gamma})} \leqslant c\|u\|_{H^{2}(\hat{\Omega})} .
$$

Furthermore, since the derivative $\left.\frac{\partial u}{\partial x}\right|_{\hat{\Gamma}}\left(\right.$ and $\left.\frac{\partial u}{\partial y}\right|_{\hat{\Gamma}}$, resp.) is equal to zero on two opposite sides of the boundary, the result as stated follows from Proposition A.1.
b) Let $\rho \in H_{00}^{1 / 2}(0,1)$. Then $\tilde{\rho} \in H^{1 / 2}(\mathbb{R})$, where $\tilde{\rho}$ is the extension of $\rho$ by zero. Using the results for the half space (see e.g. [10]), there is $u \in H^{2}\left(\mathbb{R}_{+}^{2}\right)$ satisfying

$$
\begin{equation*}
u(x, 0)=0, \quad-\frac{\partial u}{\partial y}(x, 0)=\tilde{\rho}(x) \tag{A.5}
\end{equation*}
$$

and

$$
\|u\|_{H^{2}\left(\mathbb{R}_{+}^{2}\right)} \leqslant c\|\rho\|_{H_{00}^{1 / 2}(0,1)}
$$

We may assume that $u$ vanishes outside $\left(0, \frac{1}{2}\right) \times\left(-\frac{1}{3}, \frac{4}{3}\right)$. For $x>0$ we define by reflection

$$
\left(R_{1} u\right)(x, y):=u(x, y)-\sum_{k=1}^{2} \alpha_{k} u(-k x, y)
$$

where the numbers $\alpha_{1}$ and $\alpha_{2}$ are determined by the two conditions

$$
\begin{equation*}
\frac{\partial^{j}}{\partial x^{j}}\left(R_{1} u\right)(0, y)=0, \quad j=0,1 . \tag{A.6}
\end{equation*}
$$

Note that $R_{1} u$ retains the boundary conditions (A.5) for $x>0$. Using a second reflection at $\left\{(1, y), y \in \mathbb{R}_{+}\right\}$, we obtain a function as stated.

Now we treat the general case. By a partition of unity, we only have to consider the situation in a neighbourhood $U \cap \Omega$ of a convex corner with angle $\omega$. With the help of an affine mapping $F$ we are brought back to a neighbourhood $U \cap \hat{\Omega}$ of a corner with right angle. The affine mapping $F(x)=B x+b, B=\left(\begin{array}{cc}1 & \cos \omega \\ 0 & \sin \omega\end{array}\right)^{-1}$ leads to a correspondence

$$
x \in \Omega \rightarrow \hat{x} \in \hat{\Omega}, \quad u \rightarrow \hat{u} \circ F .
$$

Normals are generally not preserved through affine mappings. However, it is easily seen that

$$
\frac{\partial u}{\partial n}=\frac{\partial \hat{u}}{\partial \hat{n}} \frac{1}{\sin (\omega)} .
$$

Finally, concerning the transformations of the norms we refer to ([5], p. 117).

## REFERENCES

[1] O. Axelsson and V. A. Barker, Finite Element Solution of Boundary Value Problems, Theory and Computation. Academic Press 1984.
[2] P. E. Bjørstad, Fast numerical solution of the biharmonic Dirichlet problem on rectangles, Siam J. Numer. Anal. 20, 59-71 (1983).
[3] J. F. Bourgat, Numerical study of a dual iterative method for solving a finite element approximation of the biharmonic equation, Comput. Methods Appl. Mech. 9, 203-218 (1976).
[4] D. Braess and P. Peisker, On the numerical solution of the biharmonic equation and the role of squaring matrices for preconditioning, IMA Journal of Numerical Analysis 6, 393-404 (1986).
[5] P. G. Ciarlet, The Finite Element Method for Elliptic Problems, North Holland 1978.
[6] P. G. Ciarlet and R. Glowinski, Dual iterative techniques for solving a finite element approximation of the biharmonic equation, Comput. Methods Appl. Mech. 5, 277-295 (1975).
[7] L. W. Ehrlich, Solving the biharmonic equation as a coupled difference equation, Siam J. Numer. Anal. 8, 278-287 (1971).
[8] R. Glowinski and O. Pironneau, Numerical methods for the first biharmonic equation and for the two dimensional Stokes problem, Siam Rev., 167-212 (1979).
[9] W. Hackbusch, Multi-Grid Methods and Applications, Springer Berlin-Heidelberg-New York, Heidelberg 1985.
[10] J. L. Lions and E. Magenes, Non-Homogeneous Boundary Value Problems and Applications I, Springer Berlin-Heidelberg-New York 1972.
[11] P. Peisker, Zwei numerische Verfahren zur Lösung der biharmonischen Gleichung unter besonderer Berücksichtigung der Mehrgitteridee, Dissertation, Bochum 1985.
[12] J. Pitkäranta, Boundary subspaces for the finite element method with Lagrange multipliers, Numer. Math. 33, 273-289 (1979).
[13] R. VERFÜrth, Error estimates for a mixed finite element approximation of the Stokes equations, R.A.I.R.O. Numerical Analysis 18, 175-182 (1984).
[14] O. B. Widlund, Iterative methods for elliptic problems partitioned into substructures and the biharmonic Dirichlet problem, in : Proceedings of the sixth international conference on computing methods in science and engineering held at Versailles, Francc, December, 12-16, 1983.
[15] H. Werner and R. Schaback, Praktische Mathematik II, Springer Berlin-Heidelberg-New York 1979.
[16] G. N. Yakovlev, Boundary properties of functions of class $W_{p}^{l}$ on regions with angular points, Doklady Academy of Sciences of U.S.S.R. 140, 73-76 (1961).


[^0]:    $\mathrm{M}^{2}$ AN Modelisation mathématıque et Analyse numérıque 0399-0516/88/04/655/22/\$ 4.20
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