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ASYMPTOTIC TIME-BEHAVIOR FOR WEIGHTED SCALAR CONSERVATION LAWS (*)

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Abstract. — For weighted scalar conservation laws introduced as model equations for gas flows in axisymmetric coordinates or in a nozzle, the asymptotic time-behavior of a entropy weak solution is obtained according to the behaviors of the weight and flux functions, thanks to the explicit formula previously derived in [10].

Résumé. — Considérant une loi de conservation scalaire avec poids qui modélise l'évolution d'un gaz en géométrie axisymétrique ou dans une tuyère, on utilise la formule explicite obtenue dans [10] pour préciser le comportement asymptotique de la solution faible entropique suivant le comportement de la fonction-flux et de la fonction-poids de l'équation.

1. INTRODUCTION : WEIGHTED SCALAR CONSERVATION LAWS

We are interested in *weighted scalar nonlinear hyperbolic conservation laws* in the half space

$$(1.1) \quad \frac{\partial}{\partial t} (r(x) u(x, t)) + \frac{\partial}{\partial x} (r(x) f(u(x, t))) = 0, \quad x > 0, \quad t > 0$$

with a \mathcal{C}^2 convex flux-function $f: \mathbb{R} \rightarrow \mathbb{R}$ and a \mathcal{C}^2 positive weight-function $r:]0, \infty[\rightarrow]0, \infty[$. Such equations are considered as model equations for gas dynamics in axisymmetric coordinates where the weight-function satisfies

$$(H.1) \quad r \in L^\infty(0, 1) \quad \text{and} \quad \frac{1}{r} \in L^\infty(1, \infty)$$

(for example take $r(x) = x^\alpha$, $\alpha = 0, 1, 2$), or as model equations for gas

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flows in a nozzle where there exists two positive constant r_- and r_+ such that

$$(H.2) \quad 0 < r_- \leq r(\cdot) \leq r_+ < +\infty .$$

For equations of the form (1.1), an uniqueness and existence result was proved in Le Floch-Nedelec [10] thanks to an explicit formula generalizing the one of Lax [7]. Moreover, a suitable formulation of boundary condition (at $x = 0$) was proposed. Refer also to previous works of Whitham [17] and Schonbek [15]. In this paper, we look for the asymptotic time-behavior of the solution of (1.1). We generalize well-known results of time-behavior for equations (1.1) without weight-function (that is $r(\cdot) \equiv 1$). The rate of convergence is specified according to the properties of the weight-function and the flux-function. For classical results on time-behavior of conservation laws, we refer to Lax [7], Dafermos [2], Conway [1], Liu-Pierre [13]... We present our result of time-behavior in the following section 2. Before, we detail some important properties of the mixed problem associated to the weighted scalar conservation law (1.1). The proofs can be found in [8]-[11].

We look for weak solutions (that is in the sense of distributions) of (1.1) satisfying an initial condition

$$(1.2) \quad u(x, 0) = u_0(x) , \quad x > 0$$

with an initial data $u_0 : \mathbb{R}_+ \rightarrow \mathbb{R}$. And for the sake of uniqueness, one adds an entropy condition (Lax [7], Oleinik [14], Kruskov [6])

$$(1.3) \quad u(x - 0, t) \geq u(x + 0, t) , \quad x > 0, t > 0 .$$

We impose a zero boundary condition at $x = 0$, since we are only interested in this paper in the rate of convergence and not in the nonlinear aspect of the boundary condition. Concerning boundary conditions, see [10] and also Le Floch [8], [9], Dubois-Le Floch [3, 4].

Our assumptions concerning the initial data are as follows

$$(H.3) \quad r \cdot u_0 \in L^1(\mathbb{R}_+) , \quad r \cdot f(u_0) \in L^\infty(\mathbb{R}_+) .$$

And, the flux-function satisfies for simplifications

$$(H.4) \quad f'' > 0 , \quad f(0) = f'(0) = 0 .$$

(for example, take $f(u) = \frac{u^2}{2}$). Then, the problem (1.1)-(1.3) admits one and only one solution $u(\cdot, \cdot)$ such that $r \cdot u$ belongs to $L^\infty(\mathbb{R}_t^+ ; L^1(\mathbb{R}_x^+))$ and $rf(u)$ to $L^\infty(\mathbb{R}^+ \times \mathbb{R}^+)$. It satisfies also the following *stability properties* for almost every $t \geq 0$: the L^∞ -stability

$$(1.4) \quad \|r(\cdot) f(u(\cdot, t))\|_{L^\infty(\mathbb{R}_x^+)} \leq \|r(\cdot) f(u_0(\cdot))\|_{L^\infty(\mathbb{R}_x^+)}$$

and the L^1 -semi-group property

$$(1.5) \quad \|(u(\cdot, t) - v(\cdot, t)) \cdot r(\cdot)\|_{L^1(\mathbb{R}_x^+)} \leq \|(u_0(\cdot) - v_0(\cdot)) \cdot r(\cdot)\|_{L^1(\mathbb{R}_x^+)}$$

for two solutions u and v corresponding to two initial data u_0 and v_0 respectively. The property (1.5) generalizes a previous result of Keyfitz [4].

Moreover, take $a = f'$ and let $f_+^{-1} : \mathbb{R}_+ \rightarrow \mathbb{R}_+$ and $f_-^{-1} : \mathbb{R}_- \rightarrow \mathbb{R}_-$ be the two inverse functions of the convex function f . For each $(x, t) \in \mathbb{R}_+ \times \mathbb{R}_+$, we can define — thanks to the hypotheses (H.1) or (H.2) (see [10]) — the function :

$$\mathbb{R}_+ \times \{-1, +1\} \ni (c, \varepsilon) \mapsto y(c, \varepsilon) \in \mathbb{R}^+$$

by the algebraic relation

$$(1.6a) \quad t = \int_{y(c, \varepsilon)}^x \frac{d\xi}{a\left(f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)\right)}$$

and the function $G : \mathbb{R}_+ \times \{-1, 1\} \ni (c, \varepsilon) \rightarrow G(c, \varepsilon)$ as follows :

$$(1.6b) \quad G(c, \varepsilon) = \int_0^{y(c, \varepsilon)} u_0(\xi) r(\xi) d\xi - c \cdot t + \int_{y(c, \varepsilon)}^x f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right) r(\xi) d\xi.$$

Then the value $u(x, t)$ of the solution u of (1.1)-(1.3) is provided by the following *explicit formula*

$$(1.7) \quad u(x, t) = f_{\varepsilon(x, t)}^{-1}\left(\frac{c(x, t)}{r(x)}\right)$$

where $(c(x, t), \varepsilon(x, t))$ minimizes the function G .

2. ASYMPTOTIC TIME-BEHAVIOR

Using the explicit formula (1.6)-(1.7), we are able as in Lax [7] to get an uniform decay in power of t for the solution u of problem (1.1)-(1.3), according to the behavior of the function f . Assuming that

$$(2.1) \quad k_- \cdot |v|^{p-2} \leq f''(v) \leq k_+ \cdot |v|^{p-2}, \quad \forall v \in \mathbb{R}; p \geq 2, 0 < k_- \leq k_+$$

we have :

THEOREM Under hypotheses (H 1) or (H 2) and (H 3)-(H 4), the solution $u(.,.)$ of (1.1)-(1.3), when the time t tends to infinity, decreases with the following rate

$$(2.2) \quad |u(x, t)| \leq \left(\frac{kM}{r(x)} \right)^{\frac{1}{p}} \cdot \frac{1}{t^{1/p}} \quad t > 0, x > 0$$

where the constants k and M depend only on the flux-function and the initial data respectively

$$k = 2p \frac{k_+}{k_-^2}$$

and

$$M = \|r \cdot u_0\|_{L^1(\mathbb{R}_+)} \quad \blacksquare$$

In the case of flows in axisymmetric coordinates

$$r(x) = x^\alpha, \quad \alpha \geq 0$$

the estimation (2.2) becomes

$$(2.3) \quad |u(x, t)| \leq \frac{c}{x^{\alpha/p}} \cdot \frac{1}{t^{1/p}} \quad (c > 0)$$

Note that, when $\alpha > 0$, this inequality (2.3) provides also the behavior of u when $x \rightarrow 0+$ and $x \rightarrow +\infty$

Proof of theorem Multiplying (2.1) by $v \in \mathbb{R}$ and integrating lead to

$$(2.4) \quad k_- \cdot \frac{|v|^p}{p} \leq a(v)v - f(v) \leq k_+ \cdot \frac{|v|^p}{p}$$

because of

$$\frac{d}{dv} (a(v)v - f(v)) = a'(v)v = f'' * (v)v$$

Moreover, by two successive integrations of (2.1), we have also

$$(2.5) \quad k_- \cdot \frac{|v|^p}{p(p-1)} \leq f(v) \leq k_+ \cdot \frac{|v|^p}{p(p-1)}$$

Then from (2.4)-(2.5) it results that

$$\frac{k_-}{k_+} (p-1) f(v) \leq a(v)v - f(v) \leq \frac{k_+}{k_-} (p-1) f(v)$$

or, for $v > 0$:

$$(2.6a) \quad \frac{k_-}{k_+} (p - 1) \frac{f(v)}{a(v)} \leq v - \frac{f(v)}{a(v)} \leq \frac{k_+}{k_-} (p - 1) \frac{f(v)}{a(v)}$$

and for $v < 0$:

$$(2.6b) \quad \frac{k_+}{k_-} (p - 1) \frac{f(v)}{a(v)} \leq v - \frac{f(v)}{a(v)} \leq \frac{k_-}{k_+} (p - 1) \frac{f(v)}{a(v)}.$$

Henceforth, the following minoration of the function G defined by (1.6) is an immediate consequence of (2.6) (used with the value $v = f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)$)

$$\begin{aligned} G(c, \varepsilon) &= \int_0^{y(c, \varepsilon)} u_0 r \, d\xi + \int_{y(c, \varepsilon)}^x \left\{ f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right) - \frac{\frac{c}{r(\xi)}}{a f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)} \right\} r(\xi) \, d\xi \\ &\geq -M + \int_{y(c, \varepsilon)}^x \frac{k_-}{k_+} (p - 1) \frac{\frac{c}{r(\xi)}}{a f_\varepsilon^{-1}\left(\frac{c}{r(\xi)}\right)} \cdot r(\xi) \, d\xi \end{aligned}$$

where $M = \|r \cdot u_0\|_{L^1(\mathbb{R}^+)}$. Thus, the inequality

$$(2.7) \quad G(c, \varepsilon) \geq -M + \frac{k_-}{k_+} (p - 1) c \cdot t$$

holds for all $c = 0, \varepsilon = \pm 1, t > 0$.

Furthermore, for each (x, t) in $\mathbb{R}_+ \times \mathbb{R}_+$, the minimum value of the function G — which is by notation attained at $(c(x, t), \varepsilon(x, t))$ — may be majorized by the value of G at $c = 0$:

$$(2.8) \quad G(c(x, t), \varepsilon(x, t)) \leq \int_0^x r u_0 \, d\xi \leq M.$$

Now comparing (2.7) and (2.8), it results the estimation

$$(p - 1) t \frac{k_-}{k_+} c(x, t) \leq 2M$$

or

$$0 \leq c(x, t) \leq \frac{k_+}{k_-} \frac{2M}{(p - 1) t}.$$

Thus in virtue of (1.7), we have proved that $rf(u)$ decreases uniformly in $x \in \mathbb{R}_+$ when $t \rightarrow \infty$:

$$\sup_{x \in \mathbb{R}_+} |r(x) f(u(x, t))| \leq \frac{k_+}{k_-} \frac{2M}{(p-1)t}.$$

It remains to use again (2.5) :

$$k_- \cdot \frac{|u(x, t)|^p}{p(p-1)} \leq f(u(x, t)) \leq \frac{2Mk_+}{k_- \cdot r(x)} \cdot \frac{1}{t}$$

which gives (2.2). ■

When the solution u is bounded — say the weight-function r satisfies the properties (H.2) — it suffices to assume that the inequalities (2.1) hold in the neighborhood of $v = 0$ with one $p_0 \geq 2$ and also that f'' is uniformly bounded :

$$\text{Cte} \leq f'' \leq \text{Cte}'.$$

Namely, thanks to the theorem applied with $p = 2$, we know that the solution u tends to zero uniformly in $x \in \mathbb{R}_+$ when the time goes to infinity

$$\sup_{x \in \mathbb{R}_+} |u(x, t)| \leq \left(\frac{kM}{r_-} \right)^{\frac{1}{2}} \cdot \frac{1}{\sqrt{t}} \quad t > 0.$$

So for t sufficiently large, say $t > T_0$, $u(x, t)$ is for each x in \mathbb{R}_+ in the neighborhood where (2.1) holds. And the same proof as previously gives

$$|u(x, t)| \leq \left(\frac{kM}{r_-} \right)^{\frac{1}{p_0}} \cdot \frac{1}{t^{1/p_0}}, \quad \forall x \in \mathbb{R}_+, \forall t > T_0.$$

REFERENCES

- [1] E. D. CONWAY, *The formation and decay of shocks for a conservation law in several dimensions*, Arch. Rat. M.A. 64 (1977) pp. 47-57.
- [2] C. M. DAFERMOS, *Characteristics in hyperbolic conservation law*, in « Nonlinear analysis and mechanics : Heriot-Watt Symposium vol. 1 », Knops Editor (1983).
- [3] F. DUBOIS, Ph. LE FLOCH, *Boundary conditions for nonlinear hyperbolic systems of conservation laws*, Internal Report (1987), École Polytechnique ; J. of Diff. Eq., Vol. 71, No 1, jan. 1988, pp. 93-122.

- [4] F. DUBOIS, Ph. LE FLOCH, *Condition à la limite pour un système de lois de conservation*, Note Compt. Rend. Acad. Sc. Paris, t. 304, Série I, n° 3, pp. 75-78 (1987).
- [5] B. KEYFITZ, *Solutions with shocks, an example of L^1 -contractive semi-group*, Comm. Pure Appl. Math., 24 (1971) pp. 125-132.
- [6] S. N. KRUSKOV, *First order quasi-linear equations in several independent variables*, Math. USSR Sb., 10 (1970) n° 2, pp. 217-243.
- [7] P. D. LAX, *Conservation laws and the mathematical theory of shock waves*, CBMS Ser. Appl. Math., vol. 11, SIAM, Philadelphia (1973).
- [8] Ph. LE FLOCH, *Explicit formula for scalar conservation laws with boundary conditions*, to appear in Math. Meth. in Appl. Sc. (1988) vol. 10.
- [9] Ph. LE FLOCH, *Generalized Riemann problem and boundary conditions for systems of conservation laws*, Thesis (1987) École Polytechnique (France).
- [10] Ph. LE FLOCH, J. C. NEDELEC, *Explicit formula for weighted scalar conservation laws*, Internal Report n° 144 (janv. 1986) of École Polytechnique ; accepted for publication to Transactions of A.M.S.
- [11] Ph. LE FLOCH, J. C. NEDELEC, *Lois de conservation scalaires avec poids*, Note Compt. Rend. Acad. Sc. Paris, t. 301, Série I, n° 17, pp. 1301-1304 (1985).
- [12] Ph. LE FLOCH, P. A. RAVIART, *Un développement asymptotique pour le problème de Riemann généralisé*, Compt. Rend. Acad. Sc. Paris, t. 304, Série I, n° 4, pp. 119-122 (1987) and Ann. Inst. Henri Poincaré, Analyse non linéaire.
- [13] T. P. LIU, M. PIERRE, *Source-solutions and asymptotic behavior in conservation laws*, J. of Diff. Eq. 51, 419-441 (1984).
- [14] O. A. OLEINIK, *Discontinuous solutions of nonlinear differential equations*, A.M.S. Transl., Ser. 2, 26, pp. 95-172 (1963).
- [15] M. E. SCHONBEK, *Existence of solutions to singular conservation laws*, Siam J. Math. Anal., vol. 15, n° 6 (nov. 1984).
- [16] J. A. SMOLLER, *Reaction-Diffusion Equations and Shock Waves*, Springer, Verlag 258 (1983).
- [17] G. B. WHITHAM, *Linear and Non linear Waves*, Wiley Interscience, New York (1974).