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ON A CONSERVATIVE UPWIND FINITE ELEMENT SCHEME FOR CONVECTIVE DIFFUSION EQUATIONS (*)

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Abstract — The purpose of this paper is to present a new class of upwind finite element schemes for convective diffusion equations and to give error analysis. These schemes based on an integral formula have the following advantages:

(i) They are effective particularly in the case when the convection is dominated,
(ii) Solutions obtained by them satisfy a discrete conservation law,
(iii) Solutions obtained by a scheme with a particular choice satisfy a discrete maximum principle (under suitable conditions).

We show that the finite element solutions converge to the exact one with rate $O(h)$ in $L^2(0, T, H^1(\Omega))$ and $L^\infty(0, T, L^2(\Omega))$.

Resumé — Le but de cet article est de présenter une classe nouvelle de schémas d'éléments finis conservatifs et décentrés pour des équations de diffusion avec convection, et de donner des estimations d'erreur. Les schémas, qui sont basés sur une formule intégrale, ont les avantages suivants:

(i) Ils sont effectifs surtout dans le cas où la convection est dominante,
(ii) Des solutions obtenues par eux satisfont une loi de conservation discrète,
(iii) Des solutions obtenues par un schéma particulier satisfont au principe du maximum discret (sous des conditions convenables).

On montre que les solutions obtenues par éléments finis convergent vers la solution exacte en $O(h)$ dans $L^2(0, T, H^1(\Omega))$ et $L^\infty(0, T, L^2(\Omega))$.

INTRODUCTION

Consider the convective diffusion equation

$$\frac{\partial u}{\partial t} = d \Delta u - \nabla \cdot (bu) + f \quad \text{in} \quad \Omega \times (0, T), \quad (0.1)$$

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where $\Omega$ is a bounded domain in $\mathbb{R}^n$. The solution $u(x, t)$ of (0.1) subject to the free boundary condition

$$d \frac{\partial u}{\partial \nu} - \mathbf{b} \cdot \mathbf{u} = 0 \quad \text{on} \quad \partial \Omega \times (0, T)$$

satisfies the mass-conservation law

$$\int_\Omega u(x, t) \, dx = \int_\Omega u^0(x) \, dx + \int_0^t \int_\Omega f(x, t) \, dx ,$$

(0.2)

where $u^0$ is an initial value. Furthermore, if the initial value $u^0$ and the source term $f$ are non-negative, so is $u$. The purpose of this paper is to present an upwind finite element scheme for (0.1) whose solution satisfies a discrete mass-conservation law. Namely, this scheme is effective even in the case when the convection is dominated

$$0 < d \ll |\mathbf{b}|$$

(0.3)

and it gives a non-negative and discrete mass-conservative solution when the initial value and the source term are non-negative.

The two properties of the exact solution discussed above, the mass-conservation law and the non-negativity, are important from a physical point of view. Since the solution $u$ stands for the density of a substance in the diffusion process, it must be non-negative and the total mass is unchangeable without any source. In the case where the convection is not dominated, it is not difficult to obtain the numerical solution satisfying these two properties. In case (0.3), however, the conventional finite element method is not effective for obtaining non-negative solutions (cf. [11, 12]). To obtain an effective scheme in the case (0.3) it is required to consider a suitable approximation for the drift term $\mathbf{V} \cdot (\mathbf{b} u)$, something like the upwind finite differencing. One of the authors [12] considered an upwind finite element scheme, whose key point was to choose an “upwind element” according to the direction of the flow for approximating the drift term. This scheme is effective in the case (0.3) and has a good feature (see Concluding Remarks), but it is not sufficient for our purpose since it does not satisfy the discrete mass conservation law.

Our scheme is based on an integral formula of the drift term on the “barycentric domain”. After applying the Green formula to the integral on the barycentric domain, we approximate the drift term by considering the upwind nodal point. For other techniques to handle the case (0.3) by the finite element method, we refer to [1, 7, 8, 9, 10, 11, 13]. For finite difference methods sharing the discrete conservation law and the non-negativity of the solution in $n = 1$ we refer to [6].
The present paper is composed of five sections. In § 1 we present a conservative upwind finite element scheme and give two theorems (theorems 1.1 and 1.2). Theorem 1.1 states that the finite element solutions of our scheme share the two properties (a discrete mass conservation law and the non-negativity). This theorem is proved in § 2. Theorem 1.2 is concerned with $L^2$-convergence of the finite element solutions. In § 3 after showing the fundamental properties of our approximation (lemma 3.2), we prove theorem 1.2. In § 4 we consider the case when $\text{div } b = 0$ and $b \cdot v |_{\partial \Omega} = 0$. Furthermore, if $f = 0$, the solution of (0.1) satisfies the maximum principle. We show that our scheme with a special choice also share this principle (theorem 4.1). We also discuss $L^\infty$-convergence of the finite element solutions. In § 5 we give some concluding remarks.

In this paper we use the following function spaces. For $1 \leq p < +\infty$ and a non-negative integer $m$,

$$| u |_{m,p,\Omega} = \sum_{|\beta| = m} \left\{ \int_{\Omega} | D_\beta^p u |^p \, dx \right\}^{1/p},$$

$$\| u \|_{m,p,\Omega} = \sum_{j=0}^m | u |_{j,p,\Omega},$$

$W^m_p(\Omega) = \{ u; u is measurable in \Omega, \| u \|_{m,p,\Omega} < +\infty \}$,

$H^m(\Omega) = W^m_2(\Omega)$.

For $0 < \alpha \leq 1$ and a non-negative integer $m$,

$$| u |_{m+\alpha,\infty,\Omega} = \sup \left\{ \frac{| D_\beta^p u(x) - D_\beta^p u(y) |}{\| x - y \|^{\alpha}}; \| \beta \| = m, x, y \in \Omega \right\},$$

$$| u |_{m,\infty,\Omega} = \sup \{ | D_\beta^p u(x) |; \| \beta \| = m, x \in \Omega \},$$

$$\| u \|_{m,\infty,\Omega} = \sum_{j=0}^m | u |_{j,\infty,\Omega},$$

$$\| u \|_{m+\alpha,\infty,\Omega} = \| u \|_{m,\infty,\Omega} + | u |_{m+\alpha,\infty,\Omega},$$

$C^m(\overline{\Omega}) = \{ u; u is continuously differentiable up to order $m$ in $\overline{\Omega}$ \},$

$C^{m+\alpha}(\overline{\Omega}) = \{ u; u \in C^m(\overline{\Omega}), \| u \|_{m+\alpha,\infty,\Omega} < +\infty \}.$

Let $X$ be a Banach space with norm $\| . \|_X$.

$C^m(0, T; X) = \{ u; u is continuously differentiable up to order $m$ as a function from $[0, T]$ into $X$ \},$

$$\| u \|_{C^m(0, T; X)} = \sum_{j=0}^m \max \{ \| D_i^j u(t) \|_X; t \in [0, T] \},$$

vol. 15, n° 1, 1981
$C^{m+a}(0, T; X) = \{ u ; u \in C^m(0, T; X), D_t^m u(t) \text{ is Hölder continuous with exponent } \alpha \text{ as a function from } [0, T] \text{ into } X \}$,

$$\| u \|_{C^{m+a}(0, T; X)} = \| u \|_{C^m(0, T; X)} + \sup \left\{ \frac{\| D_t^m u(t) - D_t^m u(s) \|_X}{| t - s |^{a}} ; t, s \in [0, T] \right\}.$$  

Let $X_i$, $i = 1, \ldots, m$, be Banach spaces. The norm of Banach space $Z = \bigcap_{i=1}^m X_i$ is given by

$$\| u \|_Z = \sum_{i=1}^m \| u \|_{X_i}.$$  

We use $c$ as a generic positive constant independent of $h$, and we denote by $c(A_1, \ldots, A_m)$ a positive constant dependent on $A_i$, $i = 1, \ldots, m$.

§ 1. RESULTS IN THE GENERAL CASE

Let $\Omega$ be a simply connected bounded domain in $\mathbb{R}^n$ with a $C^3$-class boundary $\Gamma$ or a polyhedral domain in $\mathbb{R}^n$, and $T$ be a fixed positive number. Consider the convective-diffusion problem,

$$\frac{\partial u}{\partial t} = d \Delta u - \nabla.(bu) + f \quad \text{in} \quad Q = \Omega \times (0, T),$$

$$\frac{\partial u}{\partial n} - (b \cdot \nu) u = 0 \quad \text{on} \quad \Sigma = \Gamma \times (0, T),$$

$$u = u^0 \quad \text{in} \quad \Omega \text{ at } t = 0,$$

where $d > 0$ is a given diffusion constant, $\nu$ is the unit outer normal to $\Gamma$, $b = (b_1(x), b_2(x), \ldots, b_n(x)) \in C^{0+1}(\overline{\Omega})$ is a given flow, $f \in C(0, T; L^2(\Omega))$ is a given source function, $u^0 = u^0(x) \in L^2(\Omega)$ is a given initial function and

$$\Delta = \sum_{i=1}^n \frac{\partial^2}{\partial x_i^2} , \quad \nu = (\partial/\partial x_1, \ldots, \partial/\partial x_n).$$

As was mentioned in the introduction, the solution $u$ of (1.1) satisfies the mass conservation law (0.2) and $u$ is non-negative if so are $f$ and $u^0$. Our aim is to present a finite element scheme, effective also in the case (0.3), whose solution is discrete mass conservative and non-negative (if so are $f$ and $u^0$).

We first triangulate $\overline{\Omega}$ and obtain a set of closed $n$-simplices $\{ T_j \}_{j=1}^{N_E}$ and a set of nodal points $\{ P_i \}_{i=1}^{N_P}$ satisfying the following conditions:

(i) the interiors of $T_i$ and $T_j$, $i \neq j$, are disjoint,

(ii) any one of the sides of $T_i$ is a side of another $n$-simplex $T_j$ or a portion of the boundary of the polyhedron $\bigcup_{j=1}^{N_E} T_j$.
(iii) every nodal point $P_i$ is a vertex of an $n$-simplex,

(iv) all the nodal points lying on the boundary of the polyhedron $\bigcup_{j=1}^{N_E} T_j$ exist on $\Gamma$.

Define $h(T_j)$, $\rho(T_j)$, $h$, $\kappa$ and $\Omega_h$:

$h(T_j) =$ the diameter of the smallest ball containing $T_j$,

$\rho(T_j) =$ the diameter of the largest ball contained in $T_j$,

$h = \max \{ h(T_j); j = 1, \ldots, N_E \}$,

$\kappa =$ the minimum perpendicular length of all the simplices,

$\Omega_h =$ the interior of the polyhedron $\bigcup_{j=1}^{N_E} T_j$.

Denoting by $\mathcal{T}_h(=\{ T_j \}_{j=1}^{N_E})$ a triangulation of $\Omega$ satisfying the above conditions ($N_E$ and $N_h$ may, of course, vary depending on a triangulation), we consider a family of triangulation $\{ \mathcal{T}_h \}$, $h \downarrow 0$.

**DEFINITION 1.1:** (i) We say that $\{ \mathcal{T}_h \}$ is $\gamma$-regular if there exists a constant $\gamma (> 1)$ such that

$h(T_k) \leq \gamma \rho(T_k)$ for any $T_k \in \mathcal{T}_h \in \{ \mathcal{T}_h \}$.

(ii) We say that $\{ \mathcal{T}_h \}$ is of acute type if

$\sigma(T_k) \leq 0$ for any $T_k \in \mathcal{T}_h \in \{ \mathcal{T}_h \}$,

where

$\sigma(T_k) = \max \{ \cos (\nabla \lambda_i, \nabla \lambda_j); 0 \leq i < j \leq n \}$,

and $\lambda_i, i = 0, \ldots, n$, are the barycentric coordinates with respect to the vertices of the $n$-simplex $T_k$.

**Remark 1.1:** (i) Obviously it holds that $\kappa < h$.

(ii) In the case when $\Omega$ is a polyhedral domain, we can take

$\Omega_h = \Omega$.

(iii) In $n = 2$, a family of triangulation $\{ \mathcal{T}_h \}$ is $\gamma$-regular and of acute type if and only if every angle $\theta$ of any triangle $T_j \in \mathcal{T}_h$ satisfies that

$\theta_0 \leq \theta \leq \pi/2$,

where $\theta_0$ is a positive angle independent of $h$. 

vol. 15, no 1, 1981
With each nodal point $P_i$, $i = 1, \ldots, N_p$, we associate functions $\phi_{ih} \in H^1(\Omega_h)$ and $\overline{\phi}_{ih} \in L^2(\Omega_h)$ such that:

(i) $\phi_{ih}$ is linear on each triangle and $\phi_{ih}(P_j) = \delta_{ij}$ for $i, j = 1, 2, \ldots, N_p$,

(ii) $\overline{\phi}_{ih}$ is the characteristic function of the barycentric domain $D_i$ associated with $P_i$, i.e.,

$$D_i = \bigcup_k \{ D^k_i \cap T_k \in \mathcal{T}_h \text{ such that } P_i \text{ is a vertex of } T_k \},$$

where

$$D^k_i = \bigcap_{j=1}^n \{ x \in T_k \text{ and } \lambda_{i,j}(x) \leq \lambda_i(x) \}$$

and $\lambda_1, \lambda_{i,1}, \ldots, \lambda_{i,n}$ are the barycentric coordinates with respect to $P_i, P_{i,1}, \ldots, P_{i,n}$, the vertices of $T_k$.

Let $V_h$ be the linear span of $\phi_{ih}$, $i = 1, 2, \ldots, N_p$, and let $-\Delta$ be a lumping operator from $V_h$ into $L^2(\Omega_h)$ defined by

$$v_h \mapsto \overline{v}_h = \sum_{i=1}^{N_p} v_h(P_i) \overline{\phi}_{ih}.$$ 

We now define three bilinear forms $(., .)_h, a_h(., .)$ and $b_h(., .)$ from $V_h \times V_h$ into $\mathbb{R}^1$. The first two are defined by

$$(u_h, v_h)_h = \int_{\Omega_h} \overline{u}_h(x) \overline{v}_h(x) \, dx,$$

$$a_h(u_h, v_h) = \sum_{i=1}^n \int_{\Omega_h} \frac{\partial u_h}{\partial x_i}(x) \overline{\partial v_h}{\partial x_i}(x) \, dx.$$ 

To define $b_h$ we prepare the following. Let $P_i$ and $P_j$ be adjoining nodal points. Let $\Gamma_{ij}$ be the intersection of the boundaries $\partial D_i$ and $\partial D_j$ and let

$$\gamma_{ij} = \text{mes} \Gamma_{ij} \quad \text{(the measure of } \Gamma_{ij}).$$

Let $\beta_{ij}$ be an approximation of $\int_{\Gamma_{ij}} b(x') \cdot v_{ij}(x') \, dx'$, where $v_{ij}$ is the unit outer normal vector to $\Gamma_{ij}$ considered as the boundary of $D_i$ (Therefore, $v_{ij} = - v_{ij}$). Suppose that $\beta_{ij}$ satisfy

$$\beta_{ij} + \beta_{ji} = 0,$$

$$| \beta_{ij} | \leq \| b \|_{0, \infty, \Omega} \gamma_{ij},$$

$$\left| \int_{\Gamma_{ij}} b(x') \cdot v_{ij}(x') \, dx' - \beta_{ij} \right| \leq c \| b \|_{0+1, \infty, \Omega} h'\gamma(T_k),$$

RAIRO Analyse numérique/Numerical Analysis
where $T_k$ is an $n$-simplex containing the side $P_i P_j$. Then, $b_h$ is defined by

$$b_h(u_h, v_h) = \sum_{i=1}^{N_p} v_h(P_i) \sum_{j \in \Lambda_i} \{ \beta_{ij}^+ u_h(P_i) - \beta_{ij}^- u_h(P_j) \}$$

(1.6)

where $\beta_{ij}^+ = \max (\beta_{ij}, 0)$, $\beta_{ij}^- = \max (-\beta_{ij}, 0)$ and $\Lambda_i$ is a set of nodal points defined by

$$\Lambda_i = \{ P_j; P_j, 1 \leq j \leq N_p, \text{is adjacent to} \ P_i \} .$$

Let $\tau$ be a time mesh and $N_T = [T/\tau]$.

Define an operator $Q_h$ from $L^1(\Omega)$ into $V_h$ by

$$Q_h v = \frac{1}{\text{mes } D_i} \int_{D_i} v(x) \, dx \phi_{in} .$$

(1.7)

Now, our finite element scheme corresponding to (1.1) is as follows:
Find $\{ u_h^k; k = 0, ..., N_T \} \subset V_h$ such that

$$(D_t u_h^k, \phi_h)_h = - d a_h(u_h^k, \phi_h) - b_h(u_h^k, \phi_h)$$

$$+ \int_{\Omega_h} f(x, k\tau) \phi_h(x) \, dx \quad \text{for all } \phi_h \in V_h, k = 0, ..., N_T - 1 , \quad (1.8a)$$

$$u_h^0 = Q_h u^0 , \quad (1.8b)$$

where $D_t$ is the forward difference operator defined by

$$D_t u_h^k = \frac{u_h^{k+1} - u_h^k}{\tau} .$$

**Remark 1.2**: (i) We give a concrete way to determine $\beta_{ij}$. Let $\{ T_k \}$ be a set of $n$-simplices containing the side $P_i P_j$. Let $G_k$ be the centroide of $T_k$. Let $\Gamma_{ij}$ be the intersection of $\Gamma_i$ and $T_k$ and let $n_{ij}$ be the unit outer normal vector to $\Gamma_{ij}$ (see figure 1 in the case $n = 2$). We set

$$\beta_{ij} = \sum_k b(G_k) n_{ij} \text{mes } \Gamma_{ij} .$$

(1.9)

It is not difficult to see that (1.9) satisfies (1.3) $\sim$ (1.5).

(ii) In general, the relation $\Omega_h \subset \Omega$ does not hold. Therefore $f$ and $u^0$ in (1.8) should be extended to $\Omega_h - \Omega$. But the way of extension is not significant since we shall show the convergence of rate $h$ and the width of skin $\Omega_h - \Omega$ is of order $h^2$. For example, even the extention by zero is available.
THEOREM 1.1 : (i) Equation (1.8) has a unique solution $u_h$ which satisfies the discrete mass conservation law,

$$
(u_h^k, 1)_h = \int_{\Omega_h} u^0(x) \, dx + \tau \sum_{j=0}^{k-1} \int_{\Omega_h} f(x, j\tau) \, dx \quad \text{for} \quad k = 1, \ldots, N_T. \quad (1.10)
$$

(ii) Suppose that the triangulation is of acute type and that $\tau$ and $\kappa$ satisfy the condition

$$
\tau \leq \frac{\kappa^2}{(n + 1) d + c_n \kappa \| b \|_{0, \infty, \Omega}}, \quad (1.11)
$$

where $c_n$ is a positive constant defined by (2.11) ($c_2 = 4$, $c_3 = 6$). Then, if $u^0$ and $f$ are non-negative, so is $u_h$.

Remark 1.3 : If $u^0 \in C(\overline{\Omega})$ and $f \in C(0, T; C(\overline{\Omega}))$, we can replace (1.8) by

$$(D_t u_h^k, \phi_h)_h = - d a_h(u_h^k, \phi_h) - b_h(u_h^k, \phi_h) + (I_h f(k\tau), \phi_h)_h$$

for all $\phi_h \in V_h$, $k = 0, 1, \ldots, N_T - 1$, \quad (1.12a)

$$
u_h^0 = I_h u^0, \quad (1.12b)$$

where $I_h$ is an interpolating operator from $C(\overline{\Omega})$ into $V_h$ defined by

$$I_h v = \sum_{j=1}^{N_T} v(P_j) \phi_{ik}.$$

In this case we have in place of (1.10),

$$(u_h^k, 1)_h = (I_h u^0, 1)_h + \tau \sum_{j=0}^{k-1} (I_h f(k\tau), 1)_h \quad \text{for} \quad k = 1, \ldots, N_T. \quad (1.13)$$
THEOREM 1.2: Suppose \( \{ T_h \} \) be a \( \gamma \)-regular family of triangulation of \( \Omega \). Suppose also that \( \tau \) and \( \kappa \) satisfy the condition

\[
\tau \leq \begin{cases} 
\frac{\kappa^2}{(n + 1) d} (1 - \varepsilon) & \text{if } \{ T_h \} \text{ is of acute type}, \\
\frac{2 \kappa^2}{(n + 1)^2 d} (1 - \varepsilon) & \text{otherwise},
\end{cases}
\]

(1.13)

where \( \varepsilon < 1 \) is a positive number independent of \( h \). If the solution \( u \) of (1.1) belongs to \( Z_1 = C^{1+0.5}(0, T; L^2(\Omega)) \cap C^1(0, T; H^1(\Omega)) \cap C(0, T; H^m(\Omega)) \), \( m > n/2 \), we have

\[
\max \left\{ \| e_h^k \|_{0,2,\Omega_h} \middle| k = 0, \ldots, N_T \right\}, \left\{ \tau \sum_{k=0}^{N_T-1} \left\| \frac{e_h^{k+1} + e_h^k}{2} \right\|_{1,2,\Omega_h} \right\}^{1/2} \leq c(\gamma, d, \varepsilon, \Omega, n, m, \| b \|_{0+1,\infty,\Omega}) h \| u \|_{Z_1},
\]

(1.14)

where \( e_h^k = u_h^k - I_h u(\kappa \tau) \) and \( u_h^k \) is a solution of (1.8).

§ 2. PROOF OF THEOREM 1.1

In this section, we prove theorem 1.1. We first transform (1.8) into the following form (2.1). Substituting each base function \( \phi_{ih} \) into the test function \( \phi_h \) of (1.8) and dividing both sides by \( m_t = (1, \phi_{ih})_h = \text{mes } D_i \), we obtain

\[
\frac{u_h^{k+1}(P_i) - u_h^k(P_i)}{\tau} = -\sum_{j=1}^{N_p} \frac{m_t}{m_t} \int_{D_i} f(x, \kappa \tau) \phi_{ih} dx + \frac{1}{m_t} \int_{\Omega_h} f(x, k \tau) \phi_{ih} dx,
\]

(2.1a)

\[
u_h^0(P_i) = \frac{1}{m_t} \int_{D_i} u^0(x) dx,
\]

(2.1b)

where

\[
a_{ij} = a_h(\phi_{j,k}, \phi_{i,h}), \quad b_{ij} = b_h(\phi_{j,k}, \phi_{i,h}).
\]

(2.2)

LEMMA 2.1: Let \( \kappa_i \) be the minimum perpendicular length of all the \( n \)-simplices containing a vertex \( P_i \).

(i) We have, for \( i, j = 1, \ldots, N_p \),

\[
\sum_{i=1}^{N_p} a_{ij} = \sum_{i=1}^{N_p} b_{ij} = 0,
\]

(2.3)
\[ b_{ij} \leq 0 \quad \text{if} \quad i \neq j \quad \text{and} \quad 0 \leq \frac{b_{ii}}{m_i} \leq \frac{c_n}{\kappa_i} \| b \|_{0, \infty, \Omega}, \quad (2.4) \]

where \( c_n \) is a positive constant defined by (2.11) \((c_2 = 4, c_3 = 6)\).

(ii) Suppose the triangulation is of acute type. Then,

\[ a_{ij} \leq 0 \quad \text{if} \quad i \neq j \quad \text{and} \quad 0 < \frac{a_{ii}}{m_i} \leq \frac{n + 1}{\kappa_i^2}. \quad (2.5) \]

**Proof**: The assertions concerning \( a_{ij} \) are now well-known (cf. [4], [5, lemma 1]). From (2.2) we have

\[ b_{ii} = \sum_{k \in \Lambda_i} \beta_{ik}^+, \quad b_{ij} = - \sum_{k \in \Lambda_i} \beta_{ik}^- \delta_{jk} \quad (i \neq j). \quad (2.6) \]

Noting (1.3), we obtain

\[ \sum_{i=1}^{N_k} b_{ij} = \sum_{k \in \Lambda_j} \beta_{jk}^+ - \sum_{i \neq j} \sum_{k \in \Lambda_i} \beta_{ik}^- \delta_{jk} = \sum_{k \in \Lambda_j} \beta_{kj}^- - \sum_{i \in \Lambda_j} \beta_{ij}^- = 0. \]

The first part of (2.4) is a direct consequence of (2.6).

From (1.4) and (2.6) we have

\[ \frac{b_{ii}}{m_i} \leq \| b \|_{0, \infty, \Omega} \sum_{j \in \Lambda_i} \gamma_{ij}/m_i. \quad (2.7) \]

Let \( T_k \) be an \( n \)-simplex containing \( P_i \). Let \( P_j \) be another vertex of \( T_k \) and let \( R_{ij} \) be the set of the other \( n - 1 \) vertices. We denote \( T_k \) by \( T_k = \mathcal{S}[P_i, P_j, R_{ij}] \). Let \( M_{ij} \) be the midpoint of side \( P_i P_j \). Define three \((n - 1)\)-simplices \( S_{0j}^k, S_{1j}^k, S_{2j}^k \) by

\[ S_{0j}^k = \mathcal{S}[P_i, R_{ij}], \quad S_{1j}^k = \mathcal{S}[M_{ij}, R_{ij}], \quad S_{2j}^k = \mathcal{S}[P_j, R_{ij}]. \]

Since \( M_{ij} \) is the midpoint, we have

\[ \text{mes } S_{0j}^k + \text{mes } S_{2j}^k \geq 2 \text{mes } S_{1j}^k. \quad (2.8) \]

We set

\[ \pi_n = \text{mes } (\Gamma_{ij} \cap S_{1j}^k)/\text{mes } S_{1j}^k. \quad (2.9) \]

Obviously \( \pi_n \) is a constant which depends only on \( n \).
Noting that
\[ \text{mes } T_k \geq \frac{\kappa_i}{2 n} \{ \text{mes } S^k_{0j} + \text{mes } S^k_{2j} \}, \]
we have by (2.8) and (2.9),
\[ \text{mes } (\Gamma_{ij} \cap S^k_{1j}) = \pi_n \text{mes } S^k_{1j} \leq \frac{\pi_n n}{\kappa_i} \text{mes } T_k. \]

Since \( P_j \) is an arbitrary vertex of \( T_k \) except \( P_b \), we obtain
\[ \sum_j \text{mes } (\Gamma_{ij} \cap S^k_{1j}) \leq \frac{\pi_n n^2}{\kappa_i} \text{mes } T_k. \quad (2.10) \]

Summing up (2.10) with respect to \( k \), we have
\[ \sum_{j \in A_i} \gamma_{ij} \leq \frac{\pi_n n^2}{\kappa_i} \sum_k \text{mes } T_k = \frac{\pi_n n^2(n + 1)}{\kappa_i} m_i. \]

By setting
\[ c_n = \pi_n n^2(n + 1), \quad (2.11) \]
we get (2.4). Since \( \pi_2 = 1/3, \pi_3 = 1/6 \), we obtain \( c_2 = 4, c_3 = 6 \).
\[ \text{q.e.d.} \]

Proof of theorem 1.1: Multiplying \( \tau m_i \) on both sides of (2.1a) and summing up over all the nodal points \( P_b \), we obtain in virtue of (2.3),
\[ \sum_{i=1}^{N_k} m_i u^{k+1}_h(P_i) - \sum_{i=1}^{N_k} m_i u^k_h(P_i) = \tau \int_{\Omega_h} f(x, k\tau) \, dx, \]
which is equivalent to the formula
\[ (u^{k+1}_h, 1)_h - (u^k_h, 1)_h = \tau \int_{\Omega_h} f(x, k\tau) \, dx. \quad (2.12) \]

Summing up (2.12) and noting that
\[ (u^0_h, 1)_h = \int_{\Omega_h} u^0(x) \, dx, \]
we obtain (1.10).
From (2.1a) we have

\[ u^{k+1}_h(P_i) = \left( 1 - \tau \frac{da_{ii} + b_{ii}}{m_i} \right) u^k_h(P_i) - \tau \sum_{j \neq i} \frac{da_{ij} + b_{ij}}{m_i} u^k_h(P_j) \]

\[ + \frac{\tau}{m_i} \int_{\Omega_h} f(x, \kappa t) \phi_h \, dx \quad \text{for} \quad i = 1, \ldots, N_T - 1. \quad (2.13) \]

By (2.4), (2.5) and (1.11) all the coefficients of \( u^k_h \) of the right-hand side are non-negative. Therefore, if \( u^0 \) and \( f \) are non-negative, so is \( u_h \). \qquad \text{q.e.d.}

§ 3. PROOF OF THEOREM 1.2

In this section we prove theorem 1.2. It can be proved along similar lines to [12, theorem 1.2] if some properties of \( b_h \) are shown. Therefore main efforts are devoted to show them (lemma 3.2).

For later reference we begin by stating the following lemma.

**Lemma 3.1**: Let \( \{ \mathcal{T}_h \} \) be a \( \gamma \)-regular family of triangulation of \( \Omega \) in \( \mathbb{R}^n \). Let \( T_j \in \mathcal{T}_h \) be any n-simplex and let \( h_j = h(T_j) \).

(i) For every \( v_h \in V_h \) and \( 1 \leq p < +\infty \) we have

\[ \max \{ \| v_h(x) \| ; x \in T_j \} \leq ch_j^{-n/p} \| v_h \|_{0,p,T_j}, \quad (3.1) \]

\[ \max \{ \| v_h(x) - v_h(y) \| ; x, y \in T_j \} \leq ch_j^{1-n/p} \| v_h \|_{1,p,T_j}, \quad (3.2) \]

where \( c = c(\gamma, p, n) \).

(ii) Let \( P \) be any point in \( T_j \) and \( U \) be the intersection of \( T_j \) and any hyperplane through \( P \). Then we have

\[ \int_U \| u(P) - u(x') \| \, dx' \leq c(\gamma, p, n) h_j^{n-p/n} \| u \|_{1,p,T_j}, \quad (3.3) \]

for \( u \in W^1_p(T_j), \quad p > n. \)

**Remark 3.1**: \( u(P) \) in (3.3) is meaningful by Sobolev's imbedding theorem,

\[ \| u \|_{0,\infty,G} \leq c(G) \| u \|_{1,p,G} \quad (p > n), \quad (3.4) \]

where \( G \) is a Lipschitz domain in \( \mathbb{R}^n \).
Proof of lemma 3.1: We only prove (3.3) since (3.1) and (3.2) are well-known results. Let \( \Delta \) be a reference \( n \)-simplex with vertices \( A_0 (0, \ldots, 0), A_1 (1, 0, \ldots, 0), \ldots, A_n (0, \ldots, 1) \) in \( \mathbb{R}_n^c \) and let \( F_j : \mathbb{R}_n^c \to \mathbb{R}^n \) be a linear transformation converting \( \Delta \) onto \( T_j \). We denote

\[
\begin{align*}
    u_0 (\xi) &= u (F_j (\xi)), & P_0 &= F^{-1}_j (P) \quad \text{and} \quad U_0 &= F^{-1}_j (U).
\end{align*}
\]

Let \( \alpha \) be any real number and we set \( u^\alpha = u - \alpha \).

Using (3.4) with \( G = \Delta \), we have

\[
\int_U | u (P) - u (x') | \, dx' = \int_U | u^\alpha (P) - u^\alpha (x') | \, dx'
\]

\[
\leq c h_j^{-1} \int_{U_0} | u_0^\alpha (P_0) - u_0^\alpha (\xi') | \, d\xi'
\]

\[
\leq c h_j^{-1} \| u_0^\alpha \|_{1,p,\Delta}.
\]

Since \( \alpha \) is an arbitrary number, we have

\[
\int_U | u (P) - u (x') | \, dx' \leq c h_j^{-1} \inf \{ \| u_0^\alpha \|_{1,p,\Delta} : \alpha \in \mathbb{R} \}
\]

\[
\leq c h_j^{-1} \| u_0 \|_{1,p,\Delta}
\]

\[
\leq c h_j^{-m+1} \| u_0 \|_{1,p,T_j},
\]

here we have used the fact that the norm of the quotient space \( W^1_p (\Delta) / \mathbb{R} \),

\[
\inf \{ \| v - \alpha \|_{1,p,\Delta} : \alpha \in \mathbb{R} \},
\]

is equivalent to \( |v|_{1,p,\Delta} \) (see [2; theorem 1]).

Bilinear form \( b_h \) satisfies the following fundamental inequalities. Although the inequality (3.7) is required in § 4, we state it here since it leads to (3.6).

Lemma 3.2: Let \( \mathcal{T}_h \) be a \( \gamma \)-regular family of triangulation of \( \Omega \) in \( \mathbb{R}^n \).

(i) For every \( u_h \) and \( \phi_h \in V_h \), we have

\[
| b_h (u_h, \phi_h) | \leq c (\gamma, n) \| b \|_{0,\infty,\Omega} \| u_h \|_{0,2,\Omega_h} \| \phi_h \|_{1,2,\Omega_h}. \tag{3.5}
\]

(ii) For every \( u \in H^m (\Omega), m > n/2, \) and \( \phi_h \in V_h \), we have

\[
| b_h (I_h u, \phi_h) + \int_{\Omega_h} (b u) \cdot (\nabla \phi_h) \, dx |
\]

\[
\leq c (\gamma, \Omega, m, n) \| b \|_{0+1,\infty,\Omega} \| u \|_{m,2,\Omega} \| \phi_h \|_{1,2,\Omega_h}. \tag{3.6}
\]

vol. 15, no 1, 1981
(iii) For every \( u \in W^1_p(\Omega), p > n, \) and \( \phi_h \in V_h, \) we have
\[
\left| b_h(I_h u, \phi_h) + \int_{\Omega_h} (bu) \cdot (\nabla \phi_h) \, dx \right| 
\leq c(\gamma, \Omega, p, n) h \| b \|_{0+1,\infty,\Omega} \| u \|_{1,p,\Omega} \| \phi_h \|_{1,q,\Omega_h},
\] (3.7)
where \( \frac{1}{p} + \frac{1}{q} = 1. \)

Remark 3.2: In (3.6) and (3.7) \( u \) and \( b \) are supposed to be extended smoothly from \( \Omega \) to \( \Omega_h. \)

Proof of lemma 3.2: Let \( \Gamma_h \) be the set consisting of all the internal boundaries \( \Gamma_{ij} \),
\[
\Gamma_h = \{ \Gamma_{ij}; P_j, P_j \in \overline{\Omega_h} \}.
\]
By using \( \Gamma_h \), \( b_h \) can be written as follows,
\[
b_h(u_h, \phi_h) = \sum_{\Gamma_{ij} \in \Gamma_h} \{ \phi_h(P_i) - \phi_h(P_j) \} \{ \beta_{ij}^+ u_h(P_i) - \beta_{ij}^- u_h(P_j) \}.
\] (3.8)
Hence, by (3.2) with \( p = 2, (1.4) \) and (3.1) with \( p = 2, \) we have
\[
\| b_h(u_h, \phi_h) \| \leq c \left\| b \right\|_{0,\infty,\Omega} \sum_{\Gamma_{ij} \in \Gamma_h} \| \phi_h \|_{1,2,T_k} \| u_h \|_{0,2,T_k}
\leq c \left\| b \right\|_{0,\infty,\Omega} \| \phi_h \|_{1,2,\Omega_h} \| u_h \|_{0,2,\Omega_h},
\]
where \( T_k \) is an \( n \)-simplex containing vertices \( P_i \) and \( P_j \). Thus, (3.5) is obtained.
If (3.7) is verified, we can obtain (3.6) immediately by Sobolev's lemma since it holds that
\[
\| u \|_{1,p,\Omega} \leq c \| u \|_{m,2,\Omega} \quad \text{and} \quad \| \phi_h \|_{1,q,\Omega_h} \leq c \| \phi_h \|_{1,2,\Omega_h},
\]
by taking \( p = \frac{2n}{n+2-2m} (> n \geq 2), \) \( q = \frac{p}{p-1} (< 2). \) (In the case \( n + 2 < 2m, \) it suffices to take \( p = n + 1. \))

For the proof of (3.7) we transform the second term of the left-hand side,
\[
\int_{\Omega_h} (bu) \cdot (\nabla \phi_h) \, dx = -\int_{\Omega_h} \nabla \cdot (bu) \phi_h \, dx + \int_{\partial \Omega_h} (b \cdot v) u \phi_h \, dx'
\]
\[
= \left\{ -\int_{\Omega_h} \nabla \cdot (bu) \bar{\phi}_h \, dx + \int_{\partial \Omega_h} (b \cdot v) u \phi_h \, dx' \right\} - \int_{\Omega_h} \nabla \cdot (bu) (\phi_h - \bar{\phi}_h) \, dx
\]
\[
= I_1 + I_2.
\] (3.9)
By the usual way (cf. [3, theorem 5]), we can estimate $I_2$ as follows,

$$|I_2| \leq c(\gamma, p, n) h \| b \|_{0+1,\infty, \Omega} \| u \|_{1,p,\partial \Omega} \| \phi_h \|_{1,q,\Omega_h}. \quad (3.10)$$

Applying the Green formula to the first term of $I_1$, we obtain

$$I_1 = -\sum_k \phi_h(P_k) \int_{\partial \Omega_h} (b \cdot v) u \, dx' + \int_{\partial \Omega_h} (b \cdot v) u \phi_h \, dx'$$

$$= \sum_{\Gamma_{ij} \in \Gamma_h} \{ \phi_h(P_i) - \phi_h(P_j) \} \int_{\Gamma_{ij}} (b \cdot v) u \, dx' + \int_{\partial \Omega_h} (b \cdot v) u \phi_h \, dx'$$

$$= I_{11} + I_{12}, \quad (3.11)$$

where

$$\phi_h(x') = \phi_h(x') - \phi_h(P_j) \quad \text{if} \quad x' \in D_j \cap \partial \Omega_h.$$

Since it holds that

$$\int_{\partial T_k \cap \partial \Omega_h} \phi_h(x') \, dx' = 0,$$

we have

$$I_{12} = \sum_k \int_{\partial \Omega_h \cap \partial T_k} \{ (b \cdot v) u(x') - (b \cdot v) u(M_k) \} \phi_h(x') \, dx'$$

$$\leq ch \| b \|_{0+1,\infty, \Omega} \sum_k \| u \|_{1,p,\Gamma_k} \| \phi_h \|_{1,q,\Omega_h}$$

$$\leq ch \| b \|_{0+1,\infty, \Omega} \| u \|_{1,p,\Omega} \| \phi_h \|_{1,q,\Omega_h},$$

where $M_k$ is the centroide of the face $\partial \Omega_h \cap \partial T_k$ and we have used (3.2) with $p = q$ and (3.3).

Expressing $b_h(I_h u, \phi_h)$ like (3.8), we have

$$b_h(I_h u, \phi_h) + I_{11} = \sum_{\Gamma_{ij} \in \Gamma_h} \{ \phi_h(P_i) - \phi_h(P_j) \} \left( \beta_{ij} - \int_{\Gamma_{ij}} b \cdot v \, dx' \right)$$

$$\times \{ \sigma^+ u(P_i) + \sigma^- u(P_j) \} + \sum_{\Gamma_{ij} \in \Gamma_h} \{ \phi_h(P_i) - \phi_h(P_j) \} \times$$

$$\int_{\Gamma_{ij}} b \cdot v \{ \sigma^+(u(P_i) - u(x')) + \sigma^-(u(P_j) - u(x')) \} \, dx'$$

$$= I_{111} + I_{112}, \quad (3.12)$$

where

$$\sigma^+ = \text{sgn} \beta_{ij}^+, \quad \sigma^- = 1 - \sigma^+.$$
By (3.2) with \( p = q \) and (3.3), we have

\[
| I_{112} | \leq c \| b \|_{0, \infty} \| u \|_{1,p,\Omega} \sum_k h_k^{1+\alpha} \| \phi_h \|_{1,q,T_k} \\
\leq c h \| b \|_{0, \infty} \| u \|_{1,p,\Omega} \left( \sum_k h_k^{1/p} \right) \left( \sum_k \| \phi_h \|_{1,q,T_k} \right)^{1/q}
\]

(3.13)

where \( T_k \) is an \( n \)-simplex containing side \( P \).

By (3.2) with \( p = q \) and (3.3), we have

\[
| I_{111} | \leq c h \| b \|_{0, \infty} \| u \|_{1,p,\Omega} \sum_k \| \phi_h \|_{1,q,T_k} \\
\leq c h \| b \|_{0, \infty} \| u \|_{1,p,\Omega} \| \phi_h \|_{1,q,\Omega_h}.
\]

(3.14)

Combining (3.9) \(-\) (3.14), we obtain (3.7).

q.e.d.

Scheme (1.8) satisfies the following a priori estimate.

**Lemma 3.3:** Let \( \{ T_h \} \) be a \( \gamma \)-regular family of triangulation of \( \Omega \). Suppose that \( \{ w_h^k; k = 0, \ldots, N_T \} \subset \mathbb{V}_h \) satisfies that

\[
(D_{\tau} w_h^k, \phi)_h = - d a_h(w_h^k, \phi) - b_h(w_h^k, \phi) + \delta(h) \theta_k(\phi) \| \phi_h \|_{1,2,\Omega_h}
\]

for \( \phi_h \in \mathbb{V}_h, k = 0, \ldots, N_T - 1 \), (3.15)

where \( \delta(h) \) is a non-negative function of \( h \) and \( \theta_k, k = 0, \ldots, N_T - 1 \), are functionals on \( \mathbb{V}_h \) such that \( | \theta_k | \leq 1 \).

Then, under the condition (1.13), we have

\[
\max \left\{ \| w_h^k \|_{0,2,\Omega_h}; k = 0, \ldots, N_T \right\}, \left\{ \tau \sum_{k=0}^{N_T-1} \frac{w_h^{k+1} + w_h^k}{2} \right\}^{1/2} \leq c \left\{ \| w_h^0 \|_{0,2,\Omega_h} + \delta(h) \right\},
\]

(3.16)

where \( c = c(\gamma, d, \varepsilon, n, T, \| b \|_{0,\infty}) \).

**Proof:** Substituting \( \phi_h = \frac{w_h^{k+1} + w_h^k}{2} \) into (3.15), we apply (3.5). After a brief calculation we have

\[
\left\{ \| w_h^k \|_0^2 - \frac{\tau d}{2} | w_h^k |_1^2 \right\} - \left\{ \| w_h^0 \|_0^2 - \frac{\tau d}{2} | w_h^0 |_1^2 \right\}
\]

\( + \left( \frac{d}{4} - \varepsilon_1 \right) \tau \sum_{j=0}^{k-1} | w_h^{j+1} + w_h^j |_1^2 \leq c_1(\varepsilon_1) \left\{ \tau \sum_{j=0}^{k} | w_h^j |_0^2 + \delta(h)^2 \right\},
\]

R.A.I.R.O. Analyse numérique/Numerical Analysis
where $\varepsilon_1 > 0$ is a constant less than $d/4$. Since it holds that for $w_h \in V_h$,

$$|w_h|_1 \leq \begin{cases} \frac{\sqrt{2(n+1)}}{\kappa} \|w_h\|_0 & \text{if } \{\mathcal{T}_h\} \text{ is of acute type}, \\ \frac{n+1}{\kappa} \|w_h\|_0 & \text{otherwise}, \end{cases}$$

(cf. [5, lemma 2]), condition (1.13) and Gronwall's inequality lead to (3.16). 

q.e.d.

Scheme (1.8) approximates the weak form of (1.1) in the following way.

**Lemma 3.4**: Let $\{\mathcal{T}_h\}$ be a $\gamma$-regular family of triangulation of $\Omega$. Suppose that $u$ is a solution of (1.1) belonging to

$$Z = C^{1+\sigma}(0, T ; L^2(\Omega)) \cap C^1(0, T ; H^1(\Omega)) \cap C(0, T ; H^m(\Omega)),$$

$0 < \sigma \leq 1, m > n/2$.

Then we have

$$(D_t I_h u(k\tau), \phi_h)_h = - da_h(I_h u(k\tau), \phi_h) - b_h(I_h u(k\tau), \phi_h)$$

$$+ \int_{\Omega_h} f(x, k\tau) \phi_h(x) \, dx + c_0 \theta_{k}(h + \tau^\sigma) \|\phi_h\|_{1,2,\Omega_h}, \quad (3.17)$$

where $\theta_{k}$ is a number satisfying $|\theta_{k}| \leq 1$ and

$$c_0 = c_0(\gamma, n, m, \|b\|_{0+1,0+1,\Omega}) \|u\|_Z.$$

The proof of lemma 3.4 is like that of [12, (3.19)], except that we use (3.6) for estimating the term $b_h$. So we omit the proof.

**Proof of theorem 1.2**: Since $u$ belongs to $Z_1$, (3.17) is satisfied with $\sigma = 0.5$. Subtracting (3.17) from (1.8), we observe that $\{e_h^k; k = 0, \ldots, N_T\}$ satisfies the assumption of lemma 3.3 with

$$\delta(h) = c(\gamma, d, \varepsilon, \Omega, n, m, \|b\|_{0+1,\infty,\Omega}) \|u\|_{Z_1}$$

by the fact

$$\tau^{1/2} \leq c(d) \kappa \leq c(d) h.$$

Noting that

$$\|e_h^0\|_{0,2,\Omega_h} = \|Q_h u^0 - I_h u^0\|_{0,2,\Omega_h} \leq c h \|u^0\|_{1,2,\Omega},$$

we obtain (1.14). 

q.e.d.
§ 4. THE CASE WHEN DIV $\mathbf{b} = 0$

We now consider the case when (1.1) satisfies the additional conditions

$$\text{div } \mathbf{b} = 0 \quad \text{in } \Omega \quad \text{and} \quad \mathbf{b} \cdot \mathbf{v} = 0 \quad \text{on } \Gamma.$$  \hfill (4.1)

Furthermore, if there is no source

$$f = 0,$$  \hfill (4.2)

the problem (1.1) is reduced to

$$\begin{aligned}
    \frac{\partial u}{\partial t} &= d \Delta u - (\mathbf{b} \cdot \nabla) u \quad \text{in } Q, \\
    \frac{\partial u}{\partial \mathbf{n}} &= 0 \quad \text{on } \Sigma, \\
    u &= u^0 \quad \text{in } \Omega \text{ at } t = 0.
\end{aligned}$$  \hfill (4.3)

The solution $u$ of (4.3) satisfies the maximum principle

$$\min \{ u^0(\Omega); x \in \overline{\Omega} \} \leq u(x, t) \leq \max \{ u^0(x); x \in \overline{\Omega} \} \quad \text{for } (x, t) \in \overline{Q}. \quad (4.4)$$

Let us show that a suitable choice of $\beta_{ij}$ in the scheme (1.8) enables us to obtain approximate solutions possessing not only discrete mass conservation law (1.10) but also discrete maximum principle (4.7).

For each boundary element $T_k \in \mathcal{F}_h$, i.e., a face of $T_k$ is a portion of $\partial \Omega_h$, we correspond an curved element $\tilde{T}_k$ with a corresponding portion of $\Gamma$. If $T_k$ is not a boundary element, we set $\tilde{T}_k = T_k$.

Thus, we obtain $\{ \tilde{T}_k \}$ such that the interiors of $\tilde{T}_i$ and $\tilde{T}_j$, $i \neq j$, are disjoint and that

$$\bigcup_{k=1}^{N_E} \tilde{T}_k = \overline{\Omega}. \hfill (4.5)$$

We define $\tilde{D}_i$ by replacing $T_k$ by $\tilde{T}_k$ in (1.2). Likewise we define $\tilde{D}_i$. Then, the interiors of $\tilde{D}_i$ and $\tilde{D}_j$, $i \neq j$, are disjoint and it holds that

$$\bigcup_{i=1}^{N_p} \tilde{D}_i = \overline{\Omega}. \quad (4.5)$$

Defining $\tilde{\Gamma}_{ij}$ by

$$\tilde{\Gamma}_{ij} = \partial \tilde{D}_i \cap \partial \tilde{D}_j,$$
we set

\[ \beta_{ij} = \int_{\tilde{\Gamma}_{ij}} \mathbf{b}(x') \cdot \mathbf{v}_j(x') \, dx'. \]  

(4.6)

\begin{figure}[h]
\centering
\includegraphics[width=0.5\textwidth]{figure2.png}
\caption{\( \tilde{\Gamma}_{ij} \) in the case \( n = 2 \).}
\end{figure}

**Theorem 4.1:** Assume the conditions of theorem 1.1 (ii). Then, the solution \( u_h \) of (1.8) with (4.6) satisfies the discrete mass conservation law (1.10) and the maximum principle

\[
\min \{ u^0(x); x \in \overline{\Omega} \} \leq u_h^k(x) \leq \max \{ u^0(x); x \in \overline{\Omega} \}
\]

for \( x \in \overline{\Omega}_h, k = 0, \ldots, N_T \).  

(4.7)

In this case we can show uniform convergence of \( u_h \) to \( u \).

**Theorem 4.2:** Let \( \{ \mathbf{F}_h \} \) be a \( \gamma \)-regular family. Assume the conditions of theorem 1.1 (ii). If the solution \( u \) of (4.3) belongs to

\[
Z_2 = C^{1+0.5}(0, T; C(\overline{\Omega})) \cap C^1(0, T; C^{1+1}(\overline{\Omega})) \cap C(0, T; C^{2+1}(\overline{\Omega})) ,
\]

we have

\[
\max \{ \| \mathbf{e}_h^k \|_{0, \infty, \Omega_h}; k = 0, \ldots, N_T \} \leq c h \| u \|_{Z_2} , \]

(4.8)

where

\[
\mathbf{e}_h^k = u_h^k - I_h u(\kappa T) \quad \text{and} \quad c = c(\gamma, d, \Omega, n; \| \mathbf{b} \|_{0+1, \infty, \Omega}) .
\]

For the proofs of these theorems we prepare the following lemma.
Lemma 4.1: We have

\[ \sum_{j=1}^{N_x} a_{ij} = 0, \quad \sum_{j=1}^{N_x} b_{ij} = 0 \quad \text{for} \quad i = 1, \ldots, N_p. \quad (4.9) \]

Proof: The first part is trivial. By (2.6), (4.1) and (4.6), we have

\[
\sum_{j=1}^{N_x} b_{ij} = \sum_{k \in \Lambda_i} \beta_{ik}^+ - \sum_{j \neq i} \sum_{k \in \Lambda_i} \beta_{ik}^- \delta_{jk} \\
= \sum_{k \in \Lambda_i} \beta_{ik} \\
= \int_{\partial B_i} b(x') \cdot v(x') \, dx' \\
= \int_{\partial B_i} \text{div} \, b(x) \, dx \\
= 0. \quad \text{q.e.d.}
\]

Proof of theorem 4.1: Since the coefficients of \( u_h^k \) of (2.13) are non-negative by (1.11), we obtain (4.7) by (4.9). q.e.d.

For the proof of uniform convergence we first consider an elliptic problem:

\[
- d \Delta v + (b \cdot \nabla) v + \mu v = f \quad \text{in} \quad \Omega, \quad (4.10a) \\
\frac{\partial v}{\partial n} = g \quad \text{on} \quad \Gamma, \quad (4.10b)
\]

where \( \mu \) is a real constant. We approximate (4.10) by the scheme

\[
d a_h(v_h, \phi_h) + b_h(v_h, \phi_h) + \mu(v_h, \phi_h)_h = \int_{\Omega_h} f \phi_h \, dx + \int_{\partial \Omega_h} g \phi_h \, dx' \\
\text{for any} \quad \phi_h \in V_h. \quad (4.11)
\]

Proposition 4.1: Let \( \{ \mathcal{T}_h \} \) be a \( \gamma \)-regular family. Suppose \( \{ \mathcal{T}_h \} \) is of acute type. Then, there exists a constant \( \mu_0(\gamma, n, \| b \|_{0, \infty, \Omega}) \) such that for every \( \mu \geq \mu_0 \) we have

\[
\| v_h - I_h v \|_{0, \infty, \Omega_h} \leq c h \| v \|_{2, p, \Omega}, \quad (4.12)
\]

where

\[
c = c(\gamma, \Omega, n, p, \mu, \| b \|_{0+1, \infty, \Omega}),
\]

\( v_h \) is the solution of (4.11) and \( v \) is the solution of (4.10) belonging to \( W^2_p(\Omega) \), \( p > n \).
CONVECTIVE DIFFUSION EQUATIONS

The proof of proposition 4.1 is like that of [13, theorem 3.1] whose keypoints were those corresponding to the estimate (3.7) and lemma 4.1. Therefore we omit it.

Uniform convergence of finite element solutions of parabolic problems is proved if the following two conditions are satisfied (cf. [13, 14]):

(i) Uniform convergence of finite element solutions of the corresponding elliptic problem.

(ii) Non-negativity of the scheme for the parabolic problem.

Proposition 4.1 ensures (i) and theorem 1.1 (ii) does (ii). Thus theorem 4.2 is obtained.

Remark 4.1: To obtain a scheme whose solution satisfies the discrete mass conservation law (1.10) and the (discrete) maximum principle (4.7), we have to take $\beta_{ij}$ determined by (4.6). If we define $b_h$ by

$$b_h(u_h, v_h) = \sum_{i=1}^{N_x} v_h(P_i) \sum_{j \in \Lambda_i} \beta_{ij}(u_h(P_i) - u_h(P_j)), \quad (4.13)$$

where $\beta_{ij}$ satisfy only (1.4) and (1.5), the solutions obtained by this scheme satisfy the maximum principle (4.7) and converge uniformly to the exact solution (theorem 4.2).

This result can be shown by using the fact that $b_{ij}$ derived from (4.13) satisfy (4.9) since it holds that

$$b_{ii} = \sum_{k \in \Lambda_i} \beta_{ik}, \quad b_{ij} = \begin{cases} -\beta_{ij} & \text{if } j \in \Lambda_i, \\ 0 & \text{otherwise}. \end{cases}$$

Especially if we take $\beta_{ij}$ by (4.6), (4.13) is equivalent to (1.6).

§ 5. CONCLUDING REMARKS

We have assumed hitherto that the flow $b$ does not depend on time $t$. Now we consider the case $b = b(x, t)$. Let $\{ b_k \}$, $k = 0, \ldots, N_T$, be a set of bilinear forms on $V_h \times V_h$ defined by

$$b_k(u_h, v_h) = \sum_{i=1}^{N_x} v_h(P_i) \sum_{j \in \Lambda_i} \{ (\beta_{ij}^k)^+ u_h(P_i) - (\beta_{ij}^k)^- u_h(P_j) \},$$
where $\beta_{ij}^k k = 0, \ldots, N_T$, satisfy that

$$\beta_{ij}^0 + \beta_{ji}^k = 0,$$

$$|\beta_{ij}^k| \leq \|b\|_{0,\infty,Q} \gamma_{ij},$$

$$\left| \int_{\Gamma_{ij}} b(x', k\tau).v_{ij}(x') \, dx' - \beta_{ij}^k \right| \leq c \|b\|_{0+1,\infty,Q} h^n(T_i),$$

where $T_i$ is an $n$-simplex containing nodal points $P_i$ and $P_j$.

If we replace $b_h$ in (1.8a) by $b_h^N$, theorems 1.1 and 1.2 are valid also in this case.

A feature of the bilinear form $b_h$ is that the difference scheme derived from this form is not locally consistent in $L^{\infty}$-sense, i.e.,

$$\sum_{j=1}^{N_T} \frac{b_{ij}}{m_i} u(P_j) = \frac{1}{m_i} b_h(I_h u, \phi_h) \rightarrow \nabla \cdot (bu)(P_i) \quad \text{as} \quad h \downarrow 0,$$

where $u$ is a smooth function. But $b_h(I_h u, \phi_h)$ approximates

$$-\int_{\Omega_h} (bu).\nabla \phi_h \, dx$$

in the sense (3.7), which is sufficient for obtaining convergence of the finite element solutions. On the other hand an upwind finite element scheme proposed in [12] is locally consistent in $L^{\infty}$-sense. So it is easy to apply it to first order hyperbolic problems (see [15]).

The rates of convergence given in theorems 1.2 and 4.2 are best possible in this type of approximations since the upwind bilinear form corresponds to a one-sided difference approximation.

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