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# THE SOLVABILITY OF NON $L^2$ SOLVABLE OPERATORS

NILS DENCKER

## 1. INTRODUCTION

Lerner proved in [4] that there are first order pseudodifferential operators of principal type satisfying condition  $(\Psi)$ , that are not solvable in  $L^2$  in any neighborhood of the origin. This was quite unexpected, since for first order differential operators of principal type, condition  $(\Psi)$  is equivalent to local  $L^2$  solvability.

In this paper, we shall show that the counterexamples in [4] are locally solvable in  $C^\infty$ , and that we lose at most one derivative in the estimate for the adjoint operators. In some cases we only lose  $\varepsilon$  derivatives in the estimate, for any  $\varepsilon > 0$ .

By local solvability in  $L^2$  we mean that the equation  $Pu = f$  has a local solution  $u \in L^2(\mathbf{R}^n)$  for any  $f \in L^2(\mathbf{R}^n)$  satisfying a finite number of compatibility conditions. We say that  $P$  is locally solvable in  $C^\infty$  if the equation has a solution  $u \in \mathcal{D}'$  for any  $f \in C^\infty$  satisfying a finite number of compatibility conditions. Recall that an operator is of principal type if the Hamilton field  $H_p$  of the principal symbol  $p$  is independent of the Liouville vector field.

Condition  $(\Psi)$  means that the imaginary part of the principal symbol does not change sign from  $-$  to  $+$  along the oriented bicharacteristics of the real part, see Definition 26.4.6 in [2]. This condition is invariant under multiplication of the principal symbol by non-vanishing factors.

It was conjectured by Nirenberg and Treves [5] that condition  $(\Psi)$  was equivalent to local solvability for operators of principal type, and they proved this in several cases. The necessity of  $(\Psi)$  for local solvability in the  $C^\infty$  category was proved by Moyers in two dimensions and by Hörmander in general, see Corollary 26.4.8 in [2]. In the analytic category, the sufficiency of condition  $(\Psi)$  for solvability of microdifferential operators acting on microfunctions was proved by Trépreau [6]. The sufficiency of  $(\Psi)$  for local  $L^2$  solvability for first order pseudodifferential operators in two dimensions, was proved by Lerner [3].

For differential operators, condition  $(\Psi)$  is equivalent to condition  $(P)$ , which rules out any sign changes of the imaginary part of the principal symbol along the bicharacteristics of the real part. The sufficiency of  $(P)$  for local  $L^2$  solvability for first order pseudodifferential operators was proved by Nirenberg and Treves [5] in the case when the principal symbol is real analytic, and by Beals and Fefferman [1] in the general case.

## 2. STATEMENT OF RESULTS

We shall consider the following type of operators, which includes the operators Lerner used in his counter-examples. First, let  $(t, x) \in \mathbf{R} \times \mathbf{R}^n$ ,  $n \geq 2$ , and

$$(2.1) \quad P = D_t + i \sum_{\nu \in \mathbf{Z}_+} Q_\nu(t, x_1, D_x) + R(t, x, D_x)$$

where  $R(t, x, D_x) \in C^\infty(\mathbf{R}, \Psi_{1,0}^0(T^*\mathbf{R}^n))$  and  $\sum_\nu Q_\nu(t, x_1, D_x) \in C^\infty(\mathbf{R}, \Psi_{1,0}^1)$  is on the form

$$(2.2) \quad Q_\nu(t, x_1, D_x) = \alpha_\nu(t)(D_{x_1} + H(t)\nu^k W(\nu^k x_1))\Psi_\nu(D_x), \quad \nu \in \mathbf{Z}_+.$$

Here  $0 \leq \alpha_\nu(t) \in C^\infty(\mathbf{R})$  uniformly, such that  $0 \notin \text{supp } \alpha_\nu$  and  $\alpha_\nu(t)H(t)$  is non-decreasing with  $H(t)$  the Heaviside function,  $0 \leq W(x_1) \in C^\infty(\mathbf{R})$  and  $k > 0$ . We also have  $0 \leq \Psi_\nu(\xi) \in S_{1,0}^0(T^*\mathbf{R}^n)$  uniformly, having non-overlapping interiors of the supports and  $0 < c \leq |\xi|2^{-\nu} \leq C$  in  $\text{supp } \Psi_\nu$ . Since  $0 \notin \text{supp } \alpha_\nu$  we may write  $\alpha_\nu(t)H(t) \equiv \alpha_\nu(t)\beta_\nu(t)$ , where  $\beta_\nu(t) \in C^\infty$  (but not uniformly) such that  $0 \leq \beta_\nu(t) \leq 1$  and  $0 \leq \partial_t \beta_\nu$ . We find that  $\sum_\nu \nu^k W(\nu^k x_1) \Psi_\nu(D_x) \in C^\infty(\mathbf{R}, \Psi_{1,0}^\varepsilon)$ , for any  $\varepsilon > 0$ . Since  $0 \leq \alpha_\nu(t)$  and  $W(\nu^k x_1) \Psi_\nu(\xi) \geq 0$ , it is clear that  $P$  satisfies condition  $(\Psi^*)$ , i. e., the adjoint  $P^*$  satisfies condition  $(\Psi)$ . In what follows, we shall suppress the  $t$  dependence and write  $S^m$  instead of  $C^\infty(\mathbf{R}, S^m)$  for example. We shall use the classical calculus of pseudo-differential operators, but with the general metrics and weights of the Weyl calculus. For notation and calculus results, see chapter 18 in

We define the norms

$$(2.3) \quad \|u\|_{(s,k)}^2 = \int |\hat{u}(\xi)|^2 \langle \xi \rangle^{2s} (\log \langle \xi \rangle + 1)^{2k} d\xi \quad s, k \in \mathbf{R},$$

where  $\langle \xi \rangle^2 = 1 + |\xi|^2$ . Then  $\|u\|_{(s,0)} \cong \|u\|_{(s)}$ , the usual Sobolev norm, and  $\forall s, k \in \mathbf{R}$  we have

$$(2.4) \quad c_{k,\varepsilon} \|u\|_{(s-\varepsilon)} \leq \|u\|_{(s,k)} \leq C_{k,\varepsilon} \|u\|_{(s+\varepsilon)} \quad \forall \varepsilon > 0.$$

We find that  $\|u\|_{(s,k)}$  is equivalent to  $\sum_\nu \langle \xi_\nu \rangle^{2s} (\log \langle \xi_\nu \rangle + 1)^{2k} \|\psi_\nu(D_x)u\|^2$  if  $\{\psi_\nu(\xi)\}_\nu$  is a partition of unity:  $\sum_\nu |\psi_\nu|^2 = 1$  such that  $\langle \xi \rangle \approx \langle \xi_\nu \rangle$  only varies with a fixed factor in  $\text{supp } \psi_\nu$ .

**THEOREM 2.1.** Let  $P$  be given by (2.1). Then, for any  $s \in \mathbf{R}$  there exists positive  $T_s$  and  $C_s$  such that

$$(2.5) \quad \int \|u\|_{(s)}^2(t) dt \leq C_s T^2 \int \|Pu\|_{(s,2k)}^2(t) dt$$

if  $u \in \mathcal{S}$  has support where  $|t| \leq T \leq T_s$ .

Thus, we obtain for any  $s \in \mathbf{R}$  that

$$(2.6) \quad \int \|u\|_{(s)}^2(t) dt \leq C_{s,\varepsilon} T^2 \int \|Pu\|_{(s+\varepsilon)}^2 ds \quad \forall \varepsilon > 0$$

if  $u \in \mathcal{S}$  has support where  $|t| \leq T \leq T_s$ . This shows that  $P^*$  is locally solvable in  $C^\infty$ , with loss of  $\varepsilon$  derivatives,  $\forall \varepsilon > 0$ .

We shall also consider the following operators, which includes the operators Lerner used in his counter-example with homogeneous symbols. Let

$$(2.7) \quad P = D_t + i \sum_{\nu \in J} Q_\nu(t, x, D_x) + R(t, x, D_x)$$

where  $J$  is a subset of  $\mathbf{Z}_+$  and  $\sum_\nu Q_\nu(t, x, D_x) \in \Psi_{1,0}^1$  is given by

$$(2.8) \quad Q_\nu(t, x, D_x) = \alpha_\nu(t) C(D_x) \chi_\nu(x_2) (D_{x_1} + H(t) \nu^k W(\nu^k x_1) 2^{-\nu} D_{x_2}) \quad \nu \in J.$$

Here we have the same conditions on  $\alpha_\nu$ ,  $W$  and  $R$  as before. Also,  $0 \leq C(\xi)$  is homogeneous, supported where  $|\xi_1| \leq C\xi_2$  and  $0 \leq \chi_\nu(x_2) \in S(1, dx_2^2)$  uniformly with non-overlapping supports. In fact, there exists a function  $\mu(\nu)$  on  $\mathbf{Z}_+$  such that  $\mu(\nu) \leq C_N \nu^N$ , for some  $N > 0$ , and there exists  $\tilde{\chi}_\nu \in S(1, \mu^2(\nu) dx_2^2)$  uniformly, with disjoint supports such that  $0 \leq \tilde{\chi}_\nu(x_2) \leq 1$  and  $\tilde{\chi}_\nu = 1$  on  $\text{supp } \chi_\nu$ . As before, we find that  $P$  satisfies condition  $(\Psi^*)$ .

THEOREM 2.2. Let  $P$  be given in (2.7). Then, for every  $s \in \mathbf{R}$  we find  $T_s > 0$  and  $C_s > 0$  such that

$$(2.9) \quad \int \|u\|_{(s)}^2(t) dt \leq C_s T^2 \int \|Pu\|_{(s+1)}^2(t) dt \quad \forall s$$

if  $u \in \mathcal{S}$  has support where  $|t| \leq T \leq T_s$ .

Thus  $P^*$  is locally solvable in  $C^\infty$ , with loss of one derivative. The theorems are going to be proved in the next sections.

### 3. PROOF OF THEOREM 2.1

Clearly, by conjugating with  $\langle D_x \rangle^s$  we may assume that  $s = 0$ , which only changes  $R(t, x, D_x) \in \Psi_{1,0}^0$  (dependingly on  $s$ ). Next, we shall eliminate  $R(t, x, D_x)$ . We choose  $E_\pm(t, x, D_x) \in \Psi_{1,0}^0$  with principal symbols

$$(3.1) \quad e_\pm(t, x, \xi) = \exp(\pm \int_0^t iR(t, x, \xi) dt),$$

such that  $E_- E_+ \cong E_+ E_- \cong \text{Id}$  modulo  $\Psi^{-\infty}$ . Then by conjugating with  $E_\pm$  we obtain  $R \in \Psi_{1,0}^{-1}$ , but this changes  $Q_\nu$  into

$$(3.2) \quad Q_\nu(t, x, D_x) = \alpha_\nu(t) \left( (D_{x_1} + H(t)\nu^k W(\nu^k x_1)) \Psi_\nu(D_x) + \varrho_\nu(t, x, D_x) \right)$$

where  $\{ \varrho_\nu(t, x, \xi) \}_\nu \in S_{1,0}^0$ . Since we may skip terms in  $\Psi^{-1}$  in  $P$  in the estimate (2.5), we may assume that  $\text{supp } \varrho_\nu \subseteq \text{supp } \Psi_\nu$ .

We shall localize in  $S_{1/2,0}^0$  in order to separate the different  $Q_\nu$  terms. Let  $\{ \phi_j(\xi) \}_j \in S_{1/2,0}^0$  be a partition of unity such that  $\phi_j$  is supported where  $|\xi - \xi_j| \leq c\langle \xi_j \rangle^{1/2}$ , and  $\text{supp } \phi_j$  is connected,  $\forall j$ . Let  $J \subset \mathbf{Z}_+$  be the set of those  $j$  for which  $\text{supp } \phi_j$  intersects  $\cap_\nu \text{supp } \Psi_\nu$ . Since the principal symbol of  $\sum_\nu Q_\nu \in \Psi_{1,0}^1$  vanishes of infinite order somewhere in  $\text{supp } \phi_j$  when  $j \in J$ , and  $\phi_j(\xi) \in S_{1/2,0}^0$ , we find that

$$(3.3) \quad \phi_j(D_x)Pu = \phi_j(D_x)D_t u + R_j(t, x, D_x)u$$

with  $\{ R_j \}_{j \in J} \in \Psi_{1,0}^0$  (with values in  $\ell^2$ ). We have

$$(3.4) \quad \int \|\phi_j(D_x)u\|^2(t) dt \leq CT^2 \int \|D_t \phi_j(D_x)u\|^2(t) dt \\ \leq CT^2 \int \|\phi_j(D_x)Pu\|^2(t) + \|R_j u\|^2(t) dt$$

for  $j \in J$ . Since  $\sum_{j \in J} \|R_j u\|^2 \leq C\|u\|^2$ , we get the result for small enough  $T$ , providing that we also have an estimate for the other terms.

Thus we only have to consider the case when  $\text{supp } \phi_j$  does not intersect  $\cap_\nu \text{supp } \Psi_\nu$ , i. e.  $j \notin J$ . Since  $\text{supp } \phi_j$  is connected, we find that  $\text{supp } \phi_j$  is contained in the interior of  $\text{supp } \Psi_\nu$  for some unique  $\nu = \nu_j$  when  $j \notin J$ . Observe that this gives  $|\xi_j| \approx 2^{\nu_j}$  in  $\text{supp } \phi_j$ . Clearly, since  $\text{supp } Q_\nu \subseteq \text{supp } \Psi_\nu$  we have  $P\phi_j(D_x)u = P_{\nu_j}\phi_j(D_x)u$  where we define

$$(3.5) \quad P_\nu = D_t + iQ_\nu(t, x_1, D_x).$$

Now we use the following

**Lemma 3.1.** Let  $P_\nu$  be given by (3.5). Then we find

$$(3.6) \quad \int \|u\|^2(t)(\nu^{2k}\alpha_\nu(t) + 1) dt \leq CT^2 \nu^{4k} \int \|P_\nu u\|^2(t)(\nu^{2k}\alpha_\nu(t) + 1)^{-1} dt$$

uniformly in  $\nu$ , if  $u \in \mathcal{S}$  has support in  $|t| \leq T$ , for  $T$  small enough.

By substituting  $\phi_j(D_x)u$ , taking  $\nu = \nu_j$  in (3.6), and replacing  $P_{\nu_j}$  by  $P$ , we obtain for  $j \notin J$  that

$$(3.7) \quad \int \|\phi_j(D_x)u\|^2(t) dt \leq CT^2 \nu_j^{4k} \int \|P\phi_j(D_x)u\|^2(t) dt \\ \leq CT^2 \nu_j^{4k} \int \|\phi_j(D_x)Pu\|^2(t) + \|[P, \phi_j(D_x)]u\|^2(t) dt.$$

Now  $\{\nu_j^{2k}[P, \phi_j(D_x)]\}_{j \notin J} \in \Psi_{1/2,0}^{\varepsilon-1/2}$  with values in  $\ell^2$ ,  $\forall \varepsilon > 0$ . In fact, we find that  $\sum_{\nu} \nu^k W(\nu^k x_1) \Psi_{\nu}(D_x) \in C^{\infty}(\mathbf{R}, \Psi_{1,0}^{\varepsilon})$  and  $\{\nu_j^{2k} \phi_j(\xi)\}_{j \notin J} \in S_{1/2,0}^{\varepsilon}$ ,  $\forall \varepsilon > 0$ , since  $\phi_j(\xi)$  is supported where  $|\xi| \approx 2^{\nu_j}$  when  $j \notin J$ . Thus by summing up (3.4) and (3.7) we obtain (2.5) for  $s = 0$  and small enough  $T$ . This completes the proof of Theorem 2.1.

*Proof.* [Proof of Lemma 3.1] We may assume  $\nu$  is fixed in what follows. In the proof, we are going to localize in  $|\xi_1| \geq \nu^{2k}$ . For that purpose we use the metric

$$(3.8) \quad g_{\nu} = \nu^{2k}|dx|^2 + |d\xi|^2/(\nu^{4k} + \xi_1^2) \quad \nu \in \mathbf{Z}_+$$

which is uniformly slowly varying,  $\sigma$  temperate and

$$(3.9) \quad g_{\nu}/g_{\nu}^{\sigma} = h_{\nu}^2 = \nu^{2k}/(\nu^{4k} + \xi_1^2)$$

which makes  $h_{\nu}^{-2} = |\xi_1|^2 \nu^{-2k} + \nu^{2k} \geq 2|\xi_1|$ . We find that  $Q_{\nu} \in \text{Op } S(h_{\nu}^{-2}, g_{\nu})$  but  $\nu^k W(\nu^k x_1) \in S(h_{\nu}^{-1}, g_{\nu})$  uniformly.

Now we localize with  $\chi_0(\xi_1) = \chi(\xi_1 \nu^{-2k}) \in S(1, g_{\nu})$  where  $\chi \in C_0^{\infty}$  is equal to 1 near 0, and with  $\chi_{\pm}(\xi_1) = H(\pm \xi_1)(1 - \chi_0(\xi_1)) \in S(1, g_{\nu})$  which has support where  $\pm \xi_1 > c\nu^{2k}$  so that  $\chi_0 + \chi_+ + \chi_- \equiv 1$ . We also choose non-negative  $\tilde{\chi}_{\pm}(\xi_1)$  and  $\tilde{\chi}_0(\xi_1) \in S(1, g_{\nu})$  such  $\tilde{\chi}_{\pm} \chi_{\pm} = \chi_{\pm}$  and  $\tilde{\chi}_0 \chi_0 = \chi_0$ . This can be done so that  $\tilde{\chi}_{\pm}$  have support where  $\pm \xi_1 > c_0 \nu^{2k}$ ,  $c_0 > 0$ , and  $\tilde{\chi}_0$  has support where  $|\xi_1| \leq C\nu^{2k}$ .

First we estimate the  $\chi_{\pm}(D_{x_1})u$  terms by Lemma 5.1 with the operator

$$(3.10) \quad P_{\pm} = D_t + Q_{\nu} \tilde{\chi}_{\pm}(D_{x_1}),$$

where

$$(3.11) \quad \pm \text{Re } Q_{\nu} \tilde{\chi}_{\pm}(D_{x_1}) \geq \mp C \quad \text{on } u \in \mathcal{S},$$

by the Fefferman–Phong inequality, where  $\text{Re } F = (F + F^*)/2$ . In fact, the symbol of

$$(3.12) \quad \pm \alpha_{\nu}(t) \text{Re} \left( D_{x_1} + H(t) \nu^k W(\nu^k x_1) \right) \Psi_{\nu}(D_x) \tilde{\chi}_{\pm}(D_{x_1})$$

is bounded from below, modulo terms in  $S(1, g_{\nu})$ . Thus Lemma 5.1 gives (after changing  $t$  to  $-t$  for  $P_-$ )

$$(3.13) \quad \int \|u\|^2(t) dt \leq CT^2 \int \|P_{\pm} u\|^2(t) dt$$

if  $u \in \mathcal{S}$  is supported where  $|t| \leq T$  and  $T$  is small enough. Now, by substituting  $\chi_{\pm}(D_{x_1})u$  into (3.13) and using that  $P_{\pm} \chi_{\pm}(D_{x_1}) = P_{\nu} \chi_{\pm}(D_{x_1})$  and that  $[P_{\nu}, \chi_{\pm}(D_{x_1})] \in \text{Op } S(1, g_{\nu})$  is uniformly  $L^2$  bounded, we find

$$(3.14) \quad \int \|\chi_{\pm}(D_{x_1})u\|^2(t) dt \leq C_0 T^2 \int \|P_{\nu} u\|^2(t) + \|u\|^2(t) dt$$

if  $u \in \mathcal{S}$  is supported where  $|t| \leq T$  and  $T$  is small enough.

Next, we shall estimate  $\|\chi_0(D_{x_1})u\|^2$ . Let

$$(3.15) \quad B_{\nu} = D_{x_1} \Psi_{\nu}(D_x) \tilde{\chi}_0(D_{x_1}) + \beta_{\nu}(t) \left( \nu^k W(\nu^k x_1) \Psi_{\nu}(D_x) \tilde{\chi}_0(D_{x_1}) + \varrho \right) \in \text{Op } S(h_{\nu}^{-1}, g_{\nu}),$$

where  $\varrho > 0$ . Here  $\beta_\nu \in C^\infty$  such that  $0 \leq \beta_\nu(t) \leq 1$ ,  $0 \leq \partial_t \beta_\nu$  and  $\alpha_\nu(t)H(t) \equiv \alpha_\nu(t)\beta_\nu(t)$ . Since  $\nu^k W(\nu^k x_1) \Psi_\nu(D_x) \tilde{\chi}_0(D_{x_1}) \in \text{Op } S(h_\nu^{-1}, g_\nu)$  has positive principal symbol, we find

$$(3.16) \quad \partial_t B_\nu = \partial_t \beta_\nu(t) \left( \nu^k W(\nu^k x_1) \Psi_\nu(D_x) \tilde{\chi}_0(D_{x_1}) + \varrho \right) \geq 0$$

for large enough  $\varrho$ . We also find  $B_\nu \in \text{Op } S(\nu^{2k}, g_\nu)$  uniformly, thus  $\|B_\nu\| \leq C\nu^{2k}$ . Applying Lemma 5.2 on  $\chi_0(D_{x_1})u$ , with  $P_0 = D_t + \alpha_\nu(t)(B_\nu + r_\nu)$ ,  $r_\nu = \varrho_\nu(t, x, D_x) \tilde{\chi}_0(D_{x_1}) - \beta_\nu(t)\varrho$  and  $M = C\nu^{2k}$ , we find

$$(3.17) \quad \int \|\chi_0(D_{x_1})u\|^2(t)(\nu^{2k}\alpha_\nu(t) + 1) dt \leq C_1 \nu^{4k} T^2 \int \|P_0 \chi_0(D_{x_1})u\|^2(t)(\nu^{2k}\alpha_\nu(t) + 1)^{-1} dt$$

if  $u \in \mathcal{S}$  is supported where  $|t| \leq T$  and  $T$  is small enough. As before, we find  $P_0 \chi_0(D_{x_1}) = P_\nu \chi_0(D_{x_1})$  and we have  $[P_\nu, \chi_0(D_{x_1})] = \alpha_\nu(t)f_\nu$ , where  $f_\nu \in \text{Op } S(1, g_\nu)$  is uniformly  $L^2$  bounded. Since

$$(3.18) \quad \nu^{4k} \alpha_\nu^2(t) / (\nu^{2k} \alpha_\nu(t) + 1) \leq \nu^{2k} \alpha_\nu(t) + 1,$$

we obtain

$$(3.19) \quad \int \|\chi_0(D_x)u\|^2(t)(\nu^{2k}\alpha_\nu(t) + 1) dt \\ \leq C_1 T^2 \left( \int \nu^{4k} \|P_\nu u\|^2(t)(\nu^{2k}\alpha_\nu(t) + 1)^{-1} dt + \int \|u\|^2(t)(\nu^{2k}\alpha_\nu(t) + 1) dt \right)$$

if  $u$  is supported where  $|t| \leq T$  and  $T$  is small enough. Combining (3.14) and (3.19), we obtain (3.6) for small enough  $T$ . ■

#### 4. PROOF OF THEOREM 2.2

First, we conjugate with  $\langle D_x \rangle^{s+1/2}$  to reduce to the case  $s = -1/2$  (this only changes  $R(t, x, D_x)$  dependently on  $s$ ). We choose  $E_\pm(t, x, D_x) \in \Psi_{1,0}^0$  with principal symbols

$$(4.1) \quad e_\pm(t, x, \xi) = \exp(\pm \int_0^t iR(t, x, \xi) dt),$$

such that  $E_- E_+ \cong E_+ E_- \cong \text{Id}$  modulo  $\Psi^{-\infty}$ . As before, the calculus gives  $R \in \Psi_{1,0}^{-1}$  for the new operator, but changes  $Q_\nu$  into

$$(4.2) \quad Q_\nu(t, x, D_x) = \alpha_\nu(t) \left( C(D_x) \chi_\nu(x_2) (D_{x_1} + H(t) \nu^k W(\nu^k x_1) 2^{-\nu} D_{x_2}) + \varrho_\nu(t, x, D_x) \right)$$

where  $\varrho_\nu(t, x, \xi) \in S_{1,0}^0$  uniformly, with  $\text{supp } \varrho_\nu \subseteq \text{supp } \chi_\nu$ . Thus, we may assume  $R \equiv 0$  since the term  $CT\|Ru\|_{(1/2)}$  can be estimated by the left hand side of (2.9) for  $s = -1/2$  and small enough  $T$ .

Next, we localize in  $x_2$  to separate the different  $Q_\nu$  terms. By assumption there exists  $\tilde{\chi}_\nu(x_2) \in S(1, \mu^2(\nu) dx_2^2)$  uniformly when  $\nu \in J$ , with disjoint supports, such that  $0 \leq \tilde{\chi}_\nu(x_2) \leq 1$  and  $\tilde{\chi}_\nu \chi_\nu = \chi_\nu$ . We also localize in  $\xi$ : let  $\{\psi_j(\xi)\}_j$  and  $\{\phi_j(\xi)\}_j \in S_{1,0}^0$  (with values in  $\ell^2$ ) such that  $\sum_j \psi_j(\xi)^2 = 1$ ,  $\phi_j(\xi)$  and  $\psi_j(\xi)$  are non-negative,  $\phi_j \psi_j = \psi_j$  and  $\psi_j, \phi_j$  are supported where  $0 < c \leq |\xi| 2^{-\nu} \leq C$ . We may also assume that for some fixed  $N > 0$  we have  $\sum_{|j-k| \leq N} \psi_k^2(\xi) \equiv 1$  on  $\text{supp } \psi_j, \forall j$ .

Since  $\tilde{\chi}_\nu \in S(1, \mu^2(\nu) dx_2^2)$  we find that  $\{\psi_j(\xi)\tilde{\chi}_\nu(x_2)\}_{\nu, j}$  is not in a good symbol class. Therefore, we put

$$(4.3) \quad \tilde{\chi}_{0j}(x_2) = 1 - \sum_{\substack{0 < \nu \leq j^2 \\ \nu \in J}} \tilde{\chi}_\nu(x_2).$$

Since  $\psi_j$  is supported where  $|\xi| \approx 2^j$  and  $\mu(\nu) \leq C_N \nu^N$  for some  $N > 0$ , it is easy to see that  $\{\tilde{\chi}_\nu(x_2)\psi_j(\xi)\}_{J \ni \nu \leq j^2}$  and  $\{\tilde{\chi}_{0j}(x_2)\psi_j(\xi)\}_j \in \Psi_{1, \varepsilon}^0, \forall \varepsilon > 0$ . Let

$$(4.4) \quad \alpha_{\nu j}(t) = \sqrt{\alpha_\nu(t) + 2^{-j}} \quad \forall j \in J, \quad \forall \nu,$$

in what follows. Now, we are going to use the following

**Lemma 4.1.** We find that

$$(4.5) \quad \int \sum_{\substack{J \ni \nu \leq j^2 \\ j}} \|\alpha_{\nu j}(t)\tilde{\chi}_\nu(x_2)\psi_j(D_x)u\|^2(t) + \sum_j \|\tilde{\chi}_{0j}(x_2)\psi_j(D_x)u\|^2(t) dt \\ \leq CT \int \sum_{\substack{J \ni \nu \leq j^2 \\ j}} \|\alpha_{\nu j}^{-1}(t)\tilde{\chi}_\nu(x_2)\psi_j(D_x)Pu\|^2(t) \\ + \sum_j \|\tilde{\chi}_{0j}(x_2)\psi_j(D_x)Pu\|^2(t) + \|u\|_{(-1/2)}^2(t) dt.$$

if  $u \in \mathcal{S}$  has support in  $|t| \leq T$  for  $T$  small enough.

Since  $2^{-j/2} \leq \alpha_{\nu j}, |\xi| \approx 2^j$  in  $\text{supp } \psi_j$ , the supports of  $\tilde{\chi}_\nu$  are disjoint and  $\sum_{J \ni \nu \leq j^2} \tilde{\chi}_\nu + \tilde{\chi}_{0j} \equiv 1, \forall j$ , it is easy to see that the left hand side of (4.5) is greater than  $c \int \|u\|_{(-1/2)}^2(t) dt$  for some  $c > 0$ , and the right hand side is less than  $CT \int \|Pu\|_{(1/2)}^2(t) + \|u\|_{(-1/2)}^2(t) dt$ . Thus (4.5) implies (2.9) for the case  $s = -1/2$  for small  $T$ , and completes the proof of Theorem 2.2.

*Proof.* [Proof of Lemma 4.1] Since  $\psi_j(1 - \phi_j) \equiv 0 \forall j$ , the calculus gives that we may replace  $P$  by  $P_j = D_t + i \sum_{\nu \in J} Q_\nu \phi_j(D_x)$  for the terms containing the factor  $\psi_j(D_x)$  in (4.5).

For the terms  $\|\tilde{\chi}_{0j}(x_2)\psi_j(D_x)u\|^2$  we use the fact that  $\nu^k W(\nu^k x_1) 2^{-\nu} D_{x_2} \phi_j(D_x) \in \Psi^{-\infty}$  uniformly when  $(\log |\xi|)^2 \approx j^2 < \nu$ . Thus we use Nirenberg–Treves estimate in [2, Theorem 26.8.1] with  $B = D_{x_1} \phi_j(D_x)$  bounded, and  $0 \leq A \in \Psi_{1,0}^0$  such that

$$(4.6) \quad A \cong \sum_{J \ni \nu > j^2} \alpha_\nu(t) C(D_x) \chi_\nu(x_2) \quad \text{mod } \Psi_{1,0}^{-1}.$$

By perturbing this estimate with  $L^2$  bounded operators, and substituting the term  $\tilde{\chi}_{0j}(x_2)\psi_j(D_x)u$ , we find for small enough  $T$  that

$$(4.7) \quad \int \|\tilde{\chi}_{0j}(x_2)\psi_j(D_x)u\|^2(t) dt \leq CT^2 \int \|\tilde{P}_j \tilde{\chi}_{0j}(x_2)\psi_j(D_x)u\|^2(t) dt \quad \forall j$$

when  $|t| \leq T$  in  $\text{supp } u$ . Here

$$(4.8) \quad \tilde{P}_j = D_t + i \sum_{J \ni \nu > j^2} \alpha_\nu(t) (C(D_x) \chi_\nu(x_2) D_{x_1} + \varrho_\nu(t, x, D_x)) \phi_j(D_x) \\ \cong D_t + i \sum_{J \ni \nu > j^2} Q_\nu \phi_j(D_x) \quad \text{modulo } \Psi^{-\infty}.$$

Thus  $\tilde{P}_j$  satisfies condition (P), i. e., the imaginary part of the principal symbol has no sign changes for fixed  $(x, \xi)$ .

Since  $\alpha_\nu \leq C\alpha_{\nu j}$  and  $\text{supp } \varrho_\nu \subseteq \text{supp } \chi_\nu$ , the calculus gives that

$$(4.9) \quad \left\{ [\tilde{P}_j, \tilde{\chi}_{0j}(x_2)\psi_j(D_x)] \right\}_j \cong \left\{ \sum_{\nu > j^2} \alpha_{\nu j}(t) f_{\nu j}(x, D_x) \right\}_j \quad \text{mod } \Psi_{1,\varepsilon}^{-1/2}$$

where  $\{f_{\nu j}\}_{\nu j} \in \Psi_{1,0}^0$  with values in  $\ell^2$ , and  $\text{supp } f_{\nu j} \subseteq \text{supp } \chi_\nu \psi_j$ . In order to estimate these terms we need the following

**Lemma 4.2.** If  $\{f_{\nu j}(x, D_x)\}_{\nu j} \in \Psi_{1,0}^0$  with values in  $\ell^2$ , and  $\text{supp } f_{\nu j} \subseteq \text{supp } \chi_\nu \psi_j$ ,  $\forall \nu j$ , then

$$(4.10) \quad \sum_{\substack{\nu \in J \\ j}} \|\alpha_{\nu j}(t) f_{\nu j}(x, D_x) u\|^2 \leq C \left( \sum_{\nu \leq j^2} \|\alpha_{\nu j}(t) \tilde{\chi}_j(x_2) \psi_j(D_x) u\|^2 \right. \\ \left. + \sum_j \|\tilde{\chi}_{0j}(x_2) \psi_j(D_x) u\| + \|u\|_{(-1/2)}^2 \right)$$

for  $u \in \mathcal{S}$ .

Since  $\tilde{\chi}_{0j} \equiv 0$  on  $\text{supp } \chi_\nu$  when  $J \ni \nu \leq j^2$ , we find that  $\left\{ \tilde{\chi}_{0j}(x_2) \psi_j(D_x) (\tilde{P}_j - P_j) \right\}_j \in \Psi^{-\infty}$ , where as before  $P_j = D_t + i \sum_{\nu \in J} Q_\nu \phi_j(D_x) \in \Psi_{1,0}^1$ . Thus we find

$$(4.11) \quad \int \sum_j \|\tilde{\chi}_{0j}(x_2) \psi_j(D_x) \tilde{P}_j u\|^2(t) dt \\ \leq CT \int \sum_j \|\tilde{\chi}_{0j}(x_2) \psi_j(D_x) P_j u\|^2(t) + \|u\|_{(-1/2)}^2(t) dt.$$

This gives the estimate (4.5) for the terms  $\|\tilde{\chi}_{0j}(x_2) \psi_j(D_x) u\|^2$  for small  $T$ , providing we can estimate the other terms.

As before, we are going to use Lemma 5.2 with  $a(t) = \alpha_\nu(t)$  and

$$(4.12) \quad B_t = \text{Re } C(D_x) \chi_\nu(x_2) \left( D_{x_1} \phi_j(D_x) + \beta_\nu(t) \left( \nu^k W(\nu^k x_1) 2^{-\nu} D_{x_2} \phi_j(D_x) + \varrho \right) \right),$$

where  $\varrho > 0$ . Here  $\beta_\nu \in C^\infty$  such that  $0 \leq \beta_\nu(t) \leq 1$ ,  $0 \leq \partial_t \beta_\nu$  and  $\alpha_\nu(t) H(t) \equiv \alpha_\nu(t) \beta_\nu(t)$ . We have  $\|B_t\| \leq C2^j$ ,  $\partial_t B_t \geq 0$  for large  $\varrho$  and  $R_t \in \Psi^0$ . By substituting  $\tilde{\chi}_\nu(x_2) \psi_j(D_x) u$  in this Lemma, we find for small  $T$  that

$$(4.13) \quad \int \|\tilde{\chi}_\nu(x_2) \psi_j(D_x) u\|^2(t) (2^j \alpha_\nu(t) + 1) dt \\ \leq CT^2 2^{2j} \int \|(D_t + iQ_\nu \phi_j(D_x)) \tilde{\chi}_\nu(x_2) \psi_j(D_x) u\|^2(t) (2^j \alpha_\nu(t) + 1)^{-1} dt$$

when  $J \ni \nu \leq j^2$ , providing  $|t| \leq T$  in  $\text{supp } u$ . This is equivalent to

$$(4.14) \quad \int \|\alpha_{\nu j}(t) \tilde{\chi}_\nu(x_2) \psi_j(D_x) u\|^2(t) dt \\ \leq CT^2 \int \|\alpha_{\nu j}^{-1}(t) (D_t + iQ_\nu \phi_j(D_x)) \tilde{\chi}_\nu(x_2) \psi_j(D_x) u\|^2(t) dt.$$

Now, it follows from the asymptotic expansion that

$$(4.15) \quad \left\{ [Q_\nu \phi_j(D_x), \tilde{\chi}_\nu(x_2) \psi_j(D_x)] \right\}_{J \in \nu \leq j^2} \cong \left\{ \alpha_\nu(t) \tilde{f}_{\nu j}(t, x, D_x) \right\}_{J \in \nu \leq j^2}$$



modulo  $\Psi_{1,\varepsilon}^{-1/2}$ , where  $\{ \tilde{f}_{\nu j}(t, x, D_x) \}_{\nu j} \in \Psi_{1,0}^0$  with values in  $\ell^2$ ,  $\text{supp } \tilde{f}_{\nu j} \subseteq \text{supp } \chi_\nu \psi_j$ ,  $\forall t$ . Thus, we may estimate the commutator terms by Lemma 4.2:

$$(4.16) \quad \sum_{\substack{J \ni \nu \leq j^2 \\ j}} \|\alpha_{\nu j}(t) \tilde{f}_{\nu j}(t, x, D_x) u\|^2 \\ \leq C \left( \sum_{\nu \leq j^2} \|\alpha_{\nu j} \tilde{\chi}_j \psi_j u\|^2 + \sum_j \|\tilde{\chi}_{0j} \psi_j u\|^2 + \|u\|_{(-1/2)}^2 \right) \quad \forall t.$$

Since the supports of  $\tilde{\chi}_\nu$  are disjoint, and  $\sum_{J \ni \mu \neq \nu} Q_\mu \phi_j(D_x) \in \Psi_{1,0}^1$  uniformly, we obtain that

$$(4.17) \quad \left\{ \tilde{\chi}_\nu(x_2) \psi_j(D_x) \sum_{J \ni \mu \neq \nu} Q_\mu \phi_j(D_x) \right\}_{J \ni \nu \leq j^2} \in \Psi^{-\infty}$$

with values in  $\ell^2$ . Thus we may replace  $D_t + iQ_\nu \phi_j(D_x)$  by  $P_j$  in the estimate, which proves (4.5). ■

*Proof.* [Proof of Lemma 4.2] Since  $\sum_{|j-k| \leq N} \psi_k^2(\xi) \equiv 1$  on  $\text{supp } f_{\nu j}$  and  $\{ f_{\nu j} \}_{\nu j} \in S_{1,0}^0$ , we may use the calculus to write

$$(4.18) \quad \sum_{\nu, j} \|\alpha_{\nu j}(t) f_{\nu j}(x, D_x) u\|^2 \leq \sum_{\substack{\nu, j \\ |k-j| \leq N}} \|\alpha_{\nu j}(t) e_{\nu j k}(x, D_x) \psi_k(D_x) u\|^2 + C \|u\|_{(-1)}^2,$$

where  $\{ e_{\nu j k} \}_{\nu j k} \in \Psi_{1,0}^0$  with values in  $\ell^2$ , and  $\text{supp } e_{\nu j k} \subseteq \text{supp } f_{\nu j} \psi_k$ . Since  $\tilde{\chi}_{0k} + \sum_{\mu \leq k^2} \tilde{\chi}_\mu \equiv 1$ , we find

$$(4.19) \quad \sum_{\substack{\nu, j \\ |k-j| \leq N}} \|\alpha_{\nu j}(t) e_{\nu j k}(x, D_x) \psi_k(D_x) u\|^2 \leq 2 \sum_{\substack{\nu, j \\ |k-j| \leq N}} \|\alpha_{\nu j}(t) e_{\nu j k}(x, D_x) \tilde{\chi}_{0k}(x_2) \psi_k(D_x) u\|^2 \\ + 2 \sum_{\substack{\nu, j \\ |k-j| \leq N}} \left\| \alpha_{\nu j}(t) e_{\nu j k}(x, D_x) \sum_{\mu \leq k^2} \tilde{\chi}_\mu(x_2) \psi_k(D_x) u \right\|^2.$$

By summing up in  $j$  and  $\nu$  we find

$$(4.20) \quad \sum_{\substack{\nu, j \\ |k-j| \leq N}} \|\alpha_{\nu j}(t) e_{\nu j k}(x, D_x) \tilde{\chi}_{0k}(x_2) \psi_k(D_x) u\|^2 \\ \leq C_N \left( \sum_k \|\tilde{\chi}_{0k}(x_2) \psi_k(D_x) u\|^2 + \|u\|_{(-1/2)}^2 \right),$$

since  $\alpha_{\nu j} \leq c$  and  $\{ e_{\nu j k} \}_{\nu j} \in \Psi_{1,0}^0$  with values in  $\ell^2$ , uniformly in  $k$ . Now  $\alpha_{\nu j} \leq C \alpha_{\nu k}$  when  $|j - k| \leq N$  which similarly gives by the calculus

$$(4.21) \quad \sum_{\substack{\nu, j \\ |k-j| \leq N}} \left\| \alpha_{\nu j}(t) e_{\nu j k}(x, D_x) \sum_{\mu \leq k^2} \tilde{\chi}_\mu(x_2) \psi_k(D_x) u \right\|^2 \\ \leq C \sum_{\mu \leq k^2} \|\alpha_{\mu k}(t) \tilde{\chi}_\mu(x_2) \psi_k(D_x) u\|^2 + C \|u\|_{(-1/2)}^2$$

since  $\text{supp } e_{\mu j k} \subseteq \text{supp } \chi_\mu \forall j, k$ . ■

## 5. SOME ESTIMATE LEMMAS

We assume that

$$(5.1) \quad P = D_t + iQ_t + R_t$$

where  $Q_t$  is a closed, densely defined operator on  $L^2(\mathbf{R}^n)$  such that  $\mathcal{S} \subset D(Q_t) \cap D(Q_t^*)$   $\forall t$ ,  $t \mapsto \langle Q_t u, u \rangle$  is continuous for  $u \in \mathcal{S}$ , and

$$(5.2) \quad \operatorname{Re} Q_t \geq -C_1 \quad \text{on } \mathcal{S} \quad \forall t,$$

where  $2 \operatorname{Re} Q_t = Q_t + Q_t^*$ . We also assume that  $\|R_t\| \leq C_0$  on  $L^2(\mathbf{R}^n)$ . Let  $\|u\|$  be the  $L^2$  norm of  $u \in L^2(\mathbf{R}^n)$  and  $\langle u, v \rangle$  the corresponding sesquilinear form.

**Lemma 5.1.** There exists  $T_0 > 0$  and  $C > 0$  such that

$$(5.3) \quad \int \|u\|^2(t) \leq CT^2 \int \|Pu\|^2(t) dt$$

if  $u \in \mathcal{S}$  has support where  $|t| \leq T \leq T_0$ . Here  $T_0$  and  $C$  only depend on  $C_0$  and  $C_1$ .

*Proof.* We only need to prove the estimate (5.1) for  $R_t \equiv 0$ , since we may perturb it with  $L^2$  bounded terms for small  $T$ . We find

$$(5.4) \quad \langle Q_t u, u \rangle \geq -C_1 \|u\|^2 \quad \forall t$$

when  $u \in \mathcal{S}$ . Since  $iP = \partial_t - Q_t$ , this gives

$$(5.5) \quad \begin{aligned} \|u\|^2(t) &= - \int_t^T 2 \operatorname{Re} \langle \partial_t u, u \rangle(t) dt \\ &= - \int_t^T 2 \operatorname{Re} \langle iPu, u \rangle(t) - \int_t^T 2 \operatorname{Re} \langle Q_t u, u \rangle(t) dt \\ &\leq - \int_t^T 2 \operatorname{Re} \langle iPu, u \rangle(t) dt + 2C_1 \int_t^T \|u\|^2(t) dt \end{aligned}$$

when  $u \in \mathcal{S}$ , and  $u \equiv 0$  when  $t \geq T$ .

By integrating in  $t$  we find

$$(5.6) \quad \int_{-T}^T \|u\|^2(t) dt \leq 4T \int_{-T}^T \operatorname{Im} \langle Pu, u \rangle(t) dt + 4C_1 T \int_{-T}^T \|u\|^2(t) dt$$

By using the Cauchy–Schwarz inequality we obtain

$$(5.7) \quad 2 \langle Pu, u \rangle \leq \lambda \|u\|^2/T + \|Pu\|^2 T/\lambda \quad \forall \lambda > 0.$$

This gives

$$(5.8) \quad (1 - 4CT - 2\lambda) \int \|u\|^2 \leq 2T^2/\lambda \int \|Pu\|^2 dt,$$

which gives (5.3) when  $T_0 \leq 1/16C$  and  $\lambda \leq 1/4$ . ■

The next case we shall consider is

$$(5.9) \quad P = D_t + ia(t)(B_t + R_t)$$

where  $0 \leq a(t) \leq C_0$ ,  $B_t$  and  $\partial_t B_t$  are self-adjoint and bounded,  $\partial_t B_t \geq 0$  and  $\|R_t\| \leq C_1$  on  $L^2(\mathbf{R}^n)$ . We also assume that there exists a constant  $M > 0$  such that

$$(5.10) \quad \|B_t\| \leq M \quad \forall t$$

$$(5.11) \quad \|[B_s, B_t]\| \leq M \quad \forall s, t.$$

**Lemma 5.2.** There exists  $T_0 > 0$  and  $C > 0$  such that

$$(5.12) \quad \int \|u\|^2(t)(a(t) + M^{-1}) dt \leq CT^2 \int \|Pu\|^2(t)(a(t) + M^{-1})^{-1} dt$$

if  $u \in \mathcal{S}$  has support where  $|t| \leq T \leq T_0$ . Here  $C_0$  and  $T_0$  are independent of  $M$ , and only depend on  $C_0$  and  $C_1$ .

*Proof.* First we consider the case  $a(t) \geq M^{-1} > 0$ . Then (5.12) is equivalent to the estimate:

$$(5.13) \quad \int \|u\|^2(t)a(t) dt \leq CT^2 \int \|Pu\|^2(t) dt/a(t)$$

if  $u \in \mathcal{S}$  has support where  $|t| \leq T$  is small enough. Introducing  $s = \int_0^t a(t) dt$  as a new time variable and  $P_0 = D_s + iB_t$ , we find that it suffices to prove

$$(5.14) \quad \int \|u\|^2(s) ds \leq CT^2 \int \|P_0u\|^2(s) ds$$

if  $u \in \mathcal{S}$  has support where  $|t| \leq T$ , which implies  $|s| \leq CT$ . In fact, we may then perturb the estimate with the  $L^2$  bounded term  $iR_tu$  for small  $T$ .

Now  $[P_0^*, P_0] = 2\partial_s B_t \geq 0$ , which implies

$$(5.15) \quad \|P_0u\|^2 - \|P_0^*u\|^2 = \langle [P_0^*, P_0]u, u \rangle \geq 0.$$

Since  $\|D_su\|^2 \leq 2(\|P_0u\|^2 + \|P_0^*u\|^2)$ , we find

$$(5.16) \quad \int \|u\|^2(s) ds \leq C_0T^2 \int \|D_su\|^2(s) ds \leq 4CT^2 \int \|P_0u\|^2(s) ds$$

if  $u \in \mathcal{S}$  has support where  $|s| \leq CT$ . This proves (5.13) in the case  $a(t) \geq M^{-1}$ .

Next we consider the case  $a(t) \geq 0$ . In order to reduce to the case  $a \geq M^{-1}$  we conjugate with  $E_t$  solving

$$(5.17) \quad \begin{cases} \partial_t E_t = -E_t B_t / M \\ E_0 = \text{Id}. \end{cases}$$

This gives bounds on  $\|E_t\|$  and  $\|E_t^{-1}\|$  when  $t$  is bounded (independently of  $M$ ), and the conjugation transforms  $P$  into

$$(5.18) \quad \tilde{P} = D_t + i(a(t) + M^{-1})B_t + a(t)\tilde{R}_t = D_t + i(a(t) + M^{-1})(B_t + S_t)$$

where  $\tilde{R}_t = iE_t^{-1}[B_t + R_t, E_t] + iR_t$  and  $S_t = a(t)\tilde{R}_t/(a(t) + M^{-1})$  are uniformly bounded on  $L^2(\mathbf{R}^n)$  for bounded  $t$ . In fact, if  $F_r = [B_t, E_r]$ ,  $\forall r$ , then

$$(5.19) \quad \partial_r F_r = E_r[B_r, B_t]/M - F_r B_r/M$$

and  $F_0 \equiv 0$ , thus  $F_t = [B_t, E_t]$  is bounded on  $L^2(\mathbf{R}^n)$  for bounded  $t$  (independently of  $M$ ). By using (5.13) with  $\tilde{P}$  and  $a(t) + M^{-1}$ , we obtain (5.12). ■

