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# Regularity properties of semilinear boundary problems in Besov and Triebel-Lizorkin spaces 

Jon Johnsen

## 1. Summary

For simplicity's sake the following two model problems are considered on a bounded open set $\Omega \subset \mathbb{R}^{n}$, where $n \geq 2$ and $\Gamma:=\partial \Omega$ is $C^{\infty}$-smooth: first there is the Dirichlét problem

$$
\begin{align*}
-\Delta u+u \partial_{x_{1}} u=f & \text { in } \Omega, \\
\gamma_{0} u=\varphi & \text { on } \Gamma . \tag{1.1}
\end{align*}
$$

Here $\gamma_{0} u=\left.u\right|_{\Gamma}$ and $-\Delta u=-\left(\partial_{x_{1}}^{2}+\cdots+\partial_{x_{n}}^{2}\right) u$. Secondly there is the corresponding Neumann problem

$$
\begin{align*}
-\Delta u+u \partial_{x_{1}} u=f & \text { in } \Omega, \\
\gamma_{1} u=\varphi & \text { on } \Gamma, \tag{1.2}
\end{align*}
$$

where $\gamma_{1} u=\gamma_{0}(\vec{n} \cdot \operatorname{grad} u)$ with $\vec{n}$ denoting the unit outward normal vectorfield near $\Gamma$. For the stationary Navier-Stokes equations and other problems, see Theorem 1.3 and Section 6.

The regularity of the solution $u(x)$ is studied here together with the question of carrying over weak solutions to other spaces. To obtain a unified treatment of various well-known scales of function spaces, the Besov spaces $B_{p, q}^{s}$ are considered together with the Triebel-Lizorkin spaces $F_{p, q}^{s}$; hereby $s \in \mathbb{R}$ and $p$ and $q \in] 0, \infty]$ in general, although $p<\infty$ is required throughout for the $F_{p, q}^{s}$ spaces.

Among the various identifications, recall eg that $B_{\infty, \infty}^{s}=C_{*}^{s}$ for $s>0$ (the Hölder-Zygmund spaces); $B_{p, p}^{s}=W_{p}^{s}$ for $s \in \mathbb{R}_{+} \backslash \mathbf{N}, 1<p<\infty$ (Sobolev-Slobodetskiĭ); $F_{p, 2}^{s}=H_{p}^{s}$ for $s \in \mathbb{R}, 1<p<\infty$ (Bessel-potentials) so in particular this encompasses the $W_{p}^{k}$ and $L_{p} ; F_{p, 2}^{0}=h_{p}$ for $0<p<\infty$ (local Hardy space). The scales coincide when $p=q$, so $B_{2,2}^{s}=F_{2,2}^{s}=H^{s}$ is the usual Sobolev space for $s \in \mathbb{R}$.

On $\mathbb{R}^{n}$ the spaces are defined by means of Littlewood-Paley decompositions, $B_{p, q}^{s}(\bar{\Omega})=r_{\Omega} B_{p, q}^{s}\left(\mathbb{R}^{n}\right)$ etc denotes the restriction to $\Omega$; on $\Gamma$ local coordinates are used. A concise review of the definition and the properties of the Besov and Triebel-Lizorkin spaces is given in [8], so details are omitted here; for a proper exposition the reader is referred to the books of $H$. Triebel $[\mathbf{1 4}, \mathbf{1 5}]$ and to Theorems 3.6 and 3.7 in M. Yamazaki's article [16].

[^0]For the Dirichlét problem above there is the following result:
THEOREM 1.1. Let $u(x)$ in $F_{p, q}^{s}(\bar{\Omega})$ be a solution of (1.1) for data $f(x)$ in $F_{r, o}^{t-2}(\bar{\Omega})$ and $\varphi(x)$ in $B_{r, r}^{t-\frac{1}{r}}(\Gamma)$, and suppose that

$$
\begin{align*}
& s>\max \left(\frac{1}{2}, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)  \tag{1.3}\\
& t>\max \left(\frac{1}{2}, \frac{n}{r}-1+\frac{1}{2} \delta_{n 2}\right) \tag{1.4}
\end{align*}
$$

Then $u(x)$ is also an element of $F_{r, o}^{t}(\bar{\Omega})$.
Analogously, if $u \in B_{p, q}^{s}(\bar{\Omega}), f \in B_{r, o}^{t-2}(\bar{\Omega})$ and $\varphi \in B_{r, o}^{t-\frac{1}{r}}(\Gamma)$, then (1.3)-(1.4) imply that $u \in B_{r, o}^{t}(\bar{\Omega})$.

The conditions (1.3)-(1.4) in the theorem are natural, for both $\gamma_{0}$ and $B(v):=v \partial_{1} v$ make sense on $B_{p, q}^{s}(\bar{\Omega})$ and $F_{p, q}^{s}(\bar{\Omega})$ when (1.3) holds. Actually $B(\cdot)$ is even 'better behaved' on these spaces than $-\Delta$ then; this is made precise below by taking a specific $\delta=\delta(s, p)$ such that $\delta>0$.

If one denotes $\mathcal{A}_{D}=\binom{-\Delta}{\gamma_{0}}$ and $\mathcal{B}(u)=\binom{B(u)}{0}$, problem (1.1) becomes $\mathcal{A}_{D} u+\mathcal{B}(u)=\binom{f}{\varphi}$. Then, if $\delta(s, p)>0$ and $\delta(t, r)>0$, ie if $\mathcal{B}(\cdot)$ respects the direct regularity properties of $\mathcal{A}_{D}$ at $(s, p, q)$ and $(t, r, o)$, the theorem asserts that $\mathcal{B}(\cdot)$ also respects the inverse regularity properties of $\mathcal{A}_{D}$ at these two parameters. Moreover, this holds for both of the $B_{p, q}^{s}$ and $F_{p, q}^{s}$ scales.

For the Neumann problem $\mathcal{A}_{N} u+\mathcal{B}(u)=\binom{f}{\varphi}$, where $\mathcal{A}_{N}=\binom{-\Delta}{\gamma_{1}}$, there is

Theorem 1.2. Let $u(x)$ in $F_{p, q}^{s}(\bar{\Omega})$ be a solution of (1.2) for data $f \in$ $F_{r, o}^{t-2}(\bar{\Omega})$ and $\varphi \in B_{r, r}^{t-1-\frac{1}{r}}(\Gamma)$, and suppose that

$$
\begin{equation*}
s>\max \left(\frac{1}{p}+1, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right), \quad t>\max \left(\frac{1}{r}+1, \frac{n}{r}-1+\frac{1}{2} \delta_{n 2}\right) \tag{1.5}
\end{equation*}
$$

Then $u(x)$ belongs to $F_{r, o}^{t}(\bar{\Omega})$. The analogous result holds in the $B_{p, q}^{s}(\bar{\Omega})$ spaces.

It turns out that Theorem 1.2 is rather more complicated to prove than Theorem 1.1. The reason for this is that the requirement $s>\frac{1}{2}$ is replaced by $s>\frac{1}{p}+1$ (because of $\gamma_{1}$ ), which is 'bad' in its $p$-dependence. Roughly speaking, this means that if $F_{p, q}^{s}(\bar{\Omega})+F_{r, o}^{t}(\bar{\Omega}) \subset F_{p_{1}, q_{1}}^{s_{1}}(\bar{\Omega})$, then $\left(s_{1}, p_{1}, q_{1}\right)$ need not satisfy (1.5) even if both $(s, p, q)$ and $(t, r, o)$ do so. As outlined in Section 5 below the fine theory of pointwise multiplication provides estimates of $B(\cdot)$, that may be used to overcome the difficulties.

Instead of the model problems above, the methods may be applied to eg the stationary Navier-Stokes equations. For each of the five boundary conditions considered in [7] one finds regularity results for the solutions that correspond to either Theorem 1.1 or Theorem 1.2. See [11, Thm. 5.5.5] or [10] for this.

In addition the existence of weak solutions of the Dirichlét problem may be carried over to the $B_{p, q}^{s}$ and $F_{p, q}^{s}$ spaces in this way. In more details the
problem is:

$$
\begin{align*}
-\Delta u+\sum_{j=1}^{n} u \partial_{j} u+\operatorname{grad} \mathbf{p} & =f & & \text { in } \Omega,  \tag{1.6}\\
\operatorname{div} u & =g & & \text { in } \Omega, \\
\gamma_{0} u & =\varphi & & \text { on } \Gamma .
\end{align*}
$$

Here the solution $(u, \mathbf{p})$ and the data $(f, g, \varphi)$ are sought such that

$$
\begin{gather*}
u \in B_{p, q}^{s}(\bar{\Omega})^{n}, \quad \mathbf{p} \in B_{p, q}^{s-1}(\bar{\Omega}) \\
f \in B_{p, q}^{s-2}(\bar{\Omega})^{n}, \quad g \in B_{p, q}^{s-1}(\bar{\Omega}), \quad \varphi \in B_{p, q}^{s-\frac{1}{p}}(\Gamma)^{n} \tag{1.7}
\end{gather*}
$$

for $s>\max \left(\frac{1}{2}, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)$; observe that the problem in (1.1) may serve as a model problem for (1.6). For the $F_{p, q}^{s}$ spaces the requirement is the same, but again $B_{p, q}^{s-\frac{1}{p}}(\Gamma)$ should be replaced by $B_{p, p}^{s-\frac{1}{p}}(\Gamma)$.

Concerning the existence of solutions when $g=0$ there is:
Theorem 1.3. Let $\Omega \subset \mathbb{R}^{n}$, where $n=2$ or 3 , be a $C^{\infty}$-smooth open bounded set, and let $\Omega$ be connected with finitely many components of $\Gamma$, ie $\Gamma=\Gamma_{1} \cup \cdots \cup \Gamma_{N}$.

Suppose that the data $(f, 0, \varphi)$ belong to the spaces indicated in (1.7) for a parameter $(s, p, q)$ satisfying one of the following conditions:
(1) $s>\max \left(1, \frac{n}{p}+1-\frac{n}{2}\right)$;
(2) $s>1, s=\frac{n}{p}+1-\frac{n}{2}$ and $q \leq 2$;
(3) $s=1$ and $p \geq 2 \geq q$.

Assume in addition that $\int_{\Gamma_{j}} \vec{n} \cdot \varphi=0$ for $j=1, \ldots, N$.
Then there exists a solution ( $u, \mathbf{p}$ ) of (1.6) as in (1.7) above.
For the $F_{p, q}^{s}$ spaces the analogous result holds (for any $\left.\left.q \in\right] 0, \infty\right]$ in (2)).
The special case with $(s, p, q)=(1,2,2)$ is identical to the classical result on weak solutions, cf [13]. As a particular case the theorem gives a solvability theory in the Hölder-Zygmund spaces $C_{*}^{s}(\bar{\Omega})$ for $s>1$.

In addition solutions may be constructed by successive approximations for any $s>\max \left(\frac{1}{2}, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)$ in (1.7) provided only that $\int_{\Omega} g=\int_{\Gamma} \vec{n} \cdot \varphi$ and that the norms of the data are small enough; for the present spaces, this is elaborated in [11]. In comparison Theorem 1.3 asserts that when $g=0$ and $s$ is sufficiently large (plus some stricter conditions on $\Omega$ and $\varphi$ ), then solutions exist for arbitrarily large data.

In view of this even the $C_{*}^{s}$ result should be new.
The purpose of this paper is only to indicate the proofs of the theorems; a detailed exposition is in preparation [10]. The results are based on $[\mathbf{8}, \mathbf{9}]$.

## 2. The pseudo-differential boundary operators

For an efficient treatment of the problems in (1.1), (1.2) and (1.6) one can utilise the calculus of pseudo-differential boundary operators of L. Boutet de Monvel [1] for the linear parts. An extension of this calculus to the $B_{p, q}^{s}$ and $F_{p, q}^{s}$ scales may be found in $[8,11]$ (with the results of J. Franke (partially contained) in [2] and the $H_{p}^{s}$ and $B_{p, p}^{s}$ versions of G. Grubb [4] as forerunners).

Introductions to the calculus may be found in [5, Sect. 2] and [3, Sect. 1.1 ff], or [4, Sect. 4], so here it is recalled that the generic object to study is a Green operator

$$
\mathcal{A}=\left(\begin{array}{cc}
P_{\Omega}+G & K  \tag{2.1}\\
T & S
\end{array}\right): \begin{gathered}
C^{\infty}(\bar{\Omega})^{N} \\
\oplus
\end{gathered} \rightarrow \begin{gathered}
C^{\infty}(\bar{\Omega})^{N^{\prime}} \\
C^{\infty}(\Gamma)^{M}
\end{gathered} \rightarrow \begin{gathered}
\oplus \\
C^{\infty}(\Gamma)^{M^{\prime}}
\end{gathered},
$$

whereby $P_{\Omega}=r_{\Omega} P e_{\Omega}$ denotes the truncation to $\Omega$ of a pseudo-differential operator on $\mathbb{R}^{n} ; T$ is a trace operator, $K$ a Poisson operator and $S$ is a pseudo-differential operator on $\Gamma$; finally $G$ is a singular Green operator.

To assure that $P_{\Omega}\left(C^{\infty}(\bar{\Omega})^{N}\right) \subset C^{\infty}(\bar{\Omega})^{N^{\prime}}$ the so-called transmission condition at $\Gamma$ is imposed on $P$ (cf the elementary exposition in [5, Sect. 1]). More precisely, the results in [8] have been established for the space-uniformly estimated calculus, for which the Hörmander class $S_{1,0}^{d}\left(\mathbb{R}^{n} \times \mathbb{R}^{n}\right)$ is the basic symbol class on $\mathbb{R}^{n}$; this version of the calculus has been introduced systematically in [6]. Hence $P$ is required to satisfy the uniform two-sided transmission condition at $\Gamma$, and for $\mathcal{A}$ of the described kind the main result in [8] is:

Theorem 2.1. Suppose all entries in $\mathcal{A}$ have order $d \in \mathbb{Z}$ and that both $T$ and $P_{\Omega}+G$ are of class $r \in \mathbb{Z}$. Then there is continuity of

$$
\left.\mathcal{A}: \begin{array}{c}
B_{p, q}^{s}(\bar{\Omega})^{N} \\
\oplus
\end{array} \rightarrow \begin{array}{c}
B_{p, q}^{s-d}(\bar{\Omega})^{N^{\prime}}  \tag{2.3}\\
B_{p, q}^{s-\frac{1}{p}}(\Gamma)^{M}
\end{array}\right)
$$

for $s>r+\max \left(\frac{1}{p}-1, \frac{n}{p}-n\right)$. In both cases, boundedness can only hold for $s<r+\max (\ldots)$ if both $\operatorname{class}(T)$ and $\operatorname{class}\left(P_{\Omega}+G\right)$ are $<r$.

When all symbols are poly-homogeneous and $\mathcal{A}$ is elliptic, the theorem applies also to any parametrix $\widetilde{\mathcal{A}}$, and it was shown in [4, Thm. 5.4] that $\widetilde{\mathcal{A}}$ can be taken of class $r-d$. In general this result is best possible because $\mathcal{A}$ is a parametrix of $\tilde{\mathcal{A}}$. (An exception is when $\mathcal{A}$ itself only contains a negligible part of class $r$.)

With obvious modifications the theorem also holds for multi-order and multi-class operators (of the Douglis-Nirenberg type) or when either $M$ or $M^{\prime}=0$; see [8, Thm. 5.2]. As examples there are then $\mathcal{A}_{D}$ and $\mathcal{A}_{N}$; throughout

$$
\widetilde{\mathcal{A}}_{D}:=\left(\begin{array}{ll}
R_{D} & K_{D} \tag{2.4}
\end{array}\right)=\mathcal{A}_{D}^{-1}
$$

will serve as a special choice of parametrix of $\mathcal{A}_{D}$.
For convenience $\mathbb{D}_{k}$, with $k \in \mathbb{Z}$, will denote the admissible parameters $(s, p, q)$ for which the inequality

$$
\begin{equation*}
s>k+\max \left(\frac{1}{p}-1, \frac{n}{p}-n\right) \tag{2.5}
\end{equation*}
$$

holds. Equivalently this means that $s>k-1+\frac{1}{p}+(n-1)\left(\frac{1}{p}-1\right)_{+}$.

## 3. Product estimates

The bilinear operator $B(v, w)=v \partial_{1} w$, that has been used above with $B(v):=B(v, v)$, is analysed as the composite

$$
\begin{equation*}
(v, w) \mapsto\left(v, \partial_{1} w\right) \mapsto \pi\left(v, \partial_{1} w\right), \tag{3.1}
\end{equation*}
$$

where $\pi(f, g)$ denotes $f(x) \cdot g(x)$. More precisely, $\pi(\cdot, \cdot)$ is the following generalisation, that eg allows $s>\frac{1}{2}$ (instead of $s>1$ ) in (1.3):

Definition 3.1. For $u$ and $v \in \mathcal{S}\left(\mathbb{R}^{n}\right)$ let, with $\psi_{k}(\xi)=\psi\left(2^{-k} \xi\right)$,

$$
\begin{equation*}
\pi(u, v)=\lim _{k \rightarrow \infty} \mathcal{F}^{-1}\left(\psi_{k} \hat{u}\right) \cdot \mathcal{F}^{-1}\left(\psi_{k} \hat{v}\right) \tag{3.2}
\end{equation*}
$$

whenever the limit, calculated in $\mathcal{D}^{\prime}\left(\mathbb{R}^{n}\right)$ : (i) exists for each $\psi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$ equal to 1 near 0 , and (ii) is independent of such $\psi$ 's.

This product has been studied in [9], where it is shown that it fills a part of the gap between two immediate meanings of 'pointwise multiplication': $\pi(f, u)=f u$ for $f \in \mathcal{O}_{M}$ and $u \in \mathcal{S}^{\prime}$, and $\pi\left(f_{0}, f_{1}\right)=f_{0} \cdot f_{1}$ when the $f_{j}$ lie in $L_{p_{j}}^{\text {loc }} \cap \mathcal{S}^{\prime}$ such that $0 \leq \frac{1}{p_{0}}+\frac{1}{p_{1}} \leq 1$ (so that $f_{0} \cdot f_{1} \in L_{1}^{\text {loc }}$ ).

Moreover, for an open set $\Omega \subset \mathbb{R}^{n}$ there is a restriction to $\Omega$ defined as

$$
\begin{equation*}
\pi_{\Omega}(u, v)=\lim _{k \rightarrow \infty} r_{\Omega}\left(\mathcal{F}^{-1}\left(\psi_{k} \hat{u}_{1}\right) \mathcal{F}^{-1}\left(\psi_{k} \hat{v}_{1}\right)\right) \tag{3.3}
\end{equation*}
$$

when the limit exists in $\mathcal{D}^{\prime}(\Omega)$ and satisfies (i) and (ii) for some $u_{1}$ and $v_{1}$ in $\mathcal{S}^{\prime}$ such that $r_{\Omega} u_{1}=u$ and $r_{\Omega} v_{1}=v$. (The existence of such a pair $\left(u_{1}, v_{1}\right)$ implies that the limit exists, equals $\pi_{\Omega}(u, v)$ and fulfils (i) and (ii) for any other pair restricting to $(u, v)$.)

Perhaps more importantly, the continuity properties of $\pi(\cdot, \cdot)$ may be obtained by para-multiplication. For spaces over $\Omega$ the definition in (3.3) allows one to carry boundedness over from $\pi(\cdot, \cdot)$ to $\pi_{\Omega}(\cdot, \cdot), \mathrm{cf}[\mathbf{9}, \mathrm{Thm}$. 7.2].

For simplicity's sake only the needed $F_{p, q}^{s}$ results will be recalled. For the Besov spaces it is necessary with a stricter control over the sum-exponents $q$, but in the end this does not affect the results in Theorems 1.1-1.3; hence these technicalities are omitted here.

Theorem 3.2. The product in (3.3) is defined on $F_{p_{0}, q_{0}}^{s_{0}}(\bar{\Omega}) \times F_{p_{1}, q_{1}}^{s_{1}}(\bar{\Omega})$ when

$$
\begin{equation*}
s_{0}+s_{1}>\max \left(0, \frac{n}{p_{0}}+\frac{n}{p_{1}}-n\right), \tag{3.4}
\end{equation*}
$$

and then there is boundedness

$$
\begin{equation*}
\pi_{\Omega}(\cdot, \cdot): F_{p_{0}, q_{0}}^{s_{0}}(\bar{\Omega}) \oplus F_{p_{1}, q_{1}}^{s_{1}}(\bar{\Omega}) \rightarrow F_{p_{2}, q_{2}}^{s_{2}}(\bar{\Omega}) \tag{3.5}
\end{equation*}
$$

if all of the following conditions are fulfilled:

$$
\begin{align*}
& s_{2}<\min \left(s_{0}, s_{1}\right) ;  \tag{3.6}\\
& s_{2}-\frac{n}{p_{2}} \leq \min \left(s_{0}-\frac{n}{p_{0}}, s_{1}-\frac{n}{p_{1}}, s_{0}+s_{1}-\frac{n}{p_{0}}-\frac{n}{p_{1}}\right) \text {; }  \tag{3.7}\\
& s_{2}-\frac{n}{p_{2}}=s_{0}-\frac{n}{p_{0}} \text { and } s_{1}=\frac{n}{p_{1}} \text { hold only if } p_{1} \leq 1 \text {; }  \tag{3.8}\\
& s_{2}-\frac{n}{p_{2}}=s_{1}-\frac{n}{p_{1}} \text { and } s_{0}=\frac{n}{p_{0}} \quad \text { hold only if } p_{0} \leq 1 . \tag{3.9}
\end{align*}
$$

Here it suffices with $s_{2} \leq \min \left(s_{0}, s_{1}\right)$ in (3.6) provided $q_{2} \geq q_{j}$ if $s_{2}=s_{j}$.

For this result the reader is referred to the theorems in [9, Sect.s 6 and 7]. Since $\pi_{\Omega}(\cdot, \cdot)$ is commutative, it may be assumed that $s_{0} \geq s_{1}$, and then the value, $p_{1}^{*}$, of $p_{2}$ for which there can be equality in both (3.6) and (3.7) is given by the formula

$$
\begin{equation*}
\frac{n}{p_{1}^{*}}=\frac{n}{p_{1}}+\left(s_{0}-\frac{n}{p_{0}}\right)_{-}+\left(s_{1}-\frac{n}{p_{1}}-\left(s_{0}-\frac{n}{p_{0}}\right)_{+}\right)_{+} . \tag{3.10}
\end{equation*}
$$

Remark 3.3. In Theorem 3.2 the receiving spaces $F_{p_{2}, q_{2}}^{s_{2}}$ are determined implicitly by (3.6)-(3.9). But, since $\Omega$ is bounded, $F_{p_{1}^{*}, q}^{s_{1}}(\bar{\Omega}) \hookrightarrow F_{p_{2}, q_{2}}^{s_{2}}(\bar{\Omega})$ holds in any case, if $q=q_{1}$ for $s_{0}>s_{1}$ and if $q=\max \left(q_{0}, q_{1}\right)$ for $s_{0}=s_{1}$. Thus the receiving space with $\left(s_{2}, p_{2}, q_{2}\right)=\left(s_{1}, p_{1}^{*}, q\right)$ may be considered as optimal.

## 4. The Dirichlét model problem

This section concerns the proof of Theorem 1.1. Preference will be given to the Triebel-Lizorkin spaces for simplicity, however, everything holds mutatis mutandem for the Besov spaces as well.

Firstly, for the linear parts of (1.1), there is boundedness of

$$
\mathcal{A}_{D}:=\binom{-\Delta}{\gamma_{0}}: F_{p, q}^{s}(\bar{\Omega}) \rightarrow \begin{gather*}
F_{p, q}^{s-2}(\bar{\Omega})  \tag{4.1}\\
F_{p, p}^{s-\frac{1}{p}}(\Gamma)
\end{gather*}
$$

for each parameter $(s, p, q)$ with $s>1+\max \left(\frac{1}{p}-1, \frac{n}{p}-n\right)$, ie in $\mathbb{D}_{1}$.
For $\mathcal{A}_{D}$, the calculus asserts that the parametrix $\widetilde{\mathcal{A}}_{D}$ is bounded in the opposite direction in (4.1) for each parameter $(s, p, q) \in \mathbb{D}_{1}$.

Secondly, when the non-linear term $u \partial_{x_{1}} u$ is taken into consideration too, it is found from (1.1) that

$$
\begin{equation*}
u=R_{D} f+K_{D} \varphi-R_{D}\left(u \partial_{x_{1}} u\right) . \tag{4.2}
\end{equation*}
$$

This turns out to be meaningful when $s>\max \left(\frac{1}{p}, \frac{1}{2}, \frac{n}{p}-\frac{1}{2}\right)$ for $n=2$ and for $s>\max \left(\frac{1}{p}, \frac{1}{2}, \frac{n}{p}-1\right)$ when $n \geq 3$, so in general the condition is

$$
\begin{equation*}
s>\max \left(\frac{1}{2}, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right) . \tag{4.3}
\end{equation*}
$$

To obtain this one can derive from Theorem 3.2 that $u \mapsto\left(u, \partial_{x_{1}} u\right) \mapsto u \partial_{x_{1}} u$ is bounded, for some $\delta(s, p)$,

$$
\begin{equation*}
F_{p, q}^{s}(\bar{\Omega}) \rightarrow F_{p, q}^{s}(\bar{\Omega}) \oplus F_{p, q}^{s-1}(\bar{\Omega}) \rightarrow F_{p, q}^{s-2+\delta}(\bar{\Omega}), \tag{4.4}
\end{equation*}
$$

in general when $s>\max \left(\frac{1}{2}, \frac{n}{p}-\frac{n-1}{2}\right)$, cf (3.4). Here the deficit $\delta(s, p)$ measures how much the order of $B(\cdot)$ deviates from the order of $-\Delta$.

It is essential that the non-linear term is more regular than $\mathcal{A}_{D} u$ when (4.3) holds. In fact $\delta(s, p)$ equals $1+\min \left(0, s-\frac{n}{p}\right)$-an increasing function of $s-\frac{n}{p}$ - except that $\delta\left(\frac{n}{p}, p\right)=1-\varepsilon$ for some arbitrary $\varepsilon>0$. Hence $\delta(s, p)>0$ as long as $s>\max \left(\frac{1}{2}, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)$.

Thirdly, after these preparations, an iteration yields that $u \in F_{r, o}^{t}(\bar{\Omega})$ : observe that in (4.2) one has, by (4.4) and Theorem 2.1 applied to $\widetilde{\mathcal{A}}_{D}$, for the summands on the right hand side that

$$
\begin{equation*}
R_{D} f+K_{D} \varphi \in F_{r, o}^{t}(\bar{\Omega}), \quad R_{D}\left(u \partial_{x_{1}} u\right) \in F_{p, q}^{s+\delta(s, p)}(\bar{\Omega}) . \tag{4.5}
\end{equation*}
$$

By determination of $F_{p_{1}, q_{1}}^{s_{1}} \supset F_{r, o}^{t}+F_{p, q}^{s+\delta(s, p)}$, it follows that $u \in F_{p_{1}, q_{1}}^{s_{1}}$. Application of (4.4) then gives $R_{D}\left(u \partial_{x_{1}} u\right) \in F_{p_{1}, q_{1}}^{s_{1}+\delta\left(s_{1}, p_{1}\right)}$ etc.

In the case $r=p$ one may take $p_{1}=p$, and, because $\delta(s, p)>0$, the process ends with the conclusion that $u \in F_{p, q}^{s}$ in approximately $|t-s| / \delta(s, p)$ steps (as is well known).

For $r \neq p$ the conclusion follows by consideration of four different cases, namely those with the combinations of $s+\delta(s, p) \gtreqless t$ and $t-\frac{n}{r} \gtreqless s-\frac{n}{p}$. The procedure is far easier to sketch with a diagram than with words, so the reader is referred to Figure 1.


Figure 1. The $p$-dependent iteration (for $n=3$ ).

The figure displays the location of $F_{p, q}^{s+\delta(s, p)}$ and four examples of $F_{r, o}^{t}$ corresponding to the subdivision mentioned above. However, the sumexponents $q$ and $o$ are not represented. The sector where $\delta(s, p)>0$, ie $s>\max \left(\frac{1}{2}, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)$, is indicated in dashed line; note that the ' $\times$ ' representing $F_{p, q}^{s+\delta}$ and $F_{r, o}^{t}$ all lie inside the sector.

The arrows in full line indicates embeddings $F_{p, q}^{s+\delta(s, p)} \hookrightarrow F_{p_{1}, q_{1}}^{s_{1}}$ etc: there are Sobolev embeddings down the lines of slope 1 , and to the right along horizontal lines because $\Omega$ has finite measure. The dotted line indicates the improved knowledge of the non-linear term, ie the spaces $F_{p_{1}, q_{1}}^{s_{1}+\delta\left(s_{1}, p_{1}\right)}$ etc.

Note that one of the four cases is trivial since $F_{p, q}^{s+\delta} \hookrightarrow F_{r, o}^{t}$, while another in finitely many steps (indicated by ' 0 ') reduces to this or to one of the cases with either the "sawtooth" or the "staircase" manoeuvres.

Altogether this leads to the proof of Theorem 1.1.
Remark 4.1. The procedure followed above has been used by S. I. Pohožaev, at least in the case with $s=t$ and $r>p, \operatorname{cf}[12]$.

## 5. The Neumann problem

For the Neumann problem in (1.2) the arguments in Section 4 turn out to require more detailed estimates of the non-linear term. The reasons for this will be described in the following.

For the problem in (1.2), one should take $(s, p, q) \in \mathbb{D}_{2}$, for then

$$
\begin{equation*}
\gamma_{1}: F_{p, q}^{s}(\bar{\Omega}) \rightarrow F_{p, p}^{s-1-\frac{1}{p}}(\Gamma), \tag{5.1}
\end{equation*}
$$

is bounded. It is important here that $\gamma_{1}$ can not be continuous from $F_{p, q}^{s}(\bar{\Omega})$ unless $s \geq 2+\max \left(\frac{1}{p}-1, \frac{n}{p}-n\right)$ holds, so the restriction for $(s, p, q)$ can not be essentially improved. For the Green operator $\mathcal{A}_{N}=\left(-\Delta \gamma_{1}\right)$ this means that the class is 2 .

Since $\mathcal{A}_{N}$ is elliptic, the Boutet de Monvel calculus asserts that there exists a parametrix $\tilde{\mathcal{A}}_{N}=\left(\begin{array}{ll}R_{N} & K_{N}\end{array}\right)$ of class 0 - but not lower - that is bounded

$$
\begin{equation*}
\widetilde{\mathcal{A}}_{N}: \stackrel{F_{p, q}^{s-2}(\bar{\Omega})}{\substack{s, 1-\frac{1}{p} \\ F_{p, p}^{s-1}(\Gamma)}} \rightarrow F_{p, q}^{s}(\bar{\Omega}) \tag{5.2}
\end{equation*}
$$

when $(s, p, q) \in \mathbb{D}_{2}$. Hence the class of $R_{N}$ is 0 , so $R_{N}$ can not be extended to an operator that is continuous from $F_{p, q}^{s-2}(\bar{\Omega})$ when $s-2<\max \left(\frac{1}{p}-\right.$ $\left.1, \frac{n}{p}-n\right)$. Again the restriction on $(s, p, q)$ can not be essentially improved.

Contrary to the Dirichlét case above, $\widetilde{\mathcal{A}}_{N}$ is not an inverse, but

$$
\begin{equation*}
\tilde{\mathcal{A}}_{N} \mathcal{A}_{N}=1-\mathcal{R} \tag{5.3}
\end{equation*}
$$

for an operator $\mathcal{R}$ of order $-\infty$ and class $\leq 2$. In fact, this regularising operator may be taken as $\mathcal{R} u=|\Omega|^{-1} \int_{\Omega} u$ by a specific choice of $\widetilde{\mathcal{A}}_{N}$.

Hence (4.2) is replaced by $u=R_{N} f+K_{N} \varphi+\mathcal{R} u-R_{N}\left(u \partial_{x_{1}} u\right)$, where the first three terms belong to $F_{r, o}^{t}(\bar{\Omega})$.

The conditions that make $u \partial_{x_{1}} u$ defined remain the same, of course, so the assumption for $(s, p, q)$ in (4.3) is here replaced by

$$
\begin{equation*}
s>\max \left(\frac{1}{p}+1, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right) \tag{5.4}
\end{equation*}
$$

and similarly for $(t, r, o)$. Cf (1.5).
Now the cases with $t<s+\delta(s, p)$ and $t-\frac{n}{r}>s+\delta-\frac{n}{p}$ are rather more complicated than the corresponding cases for the Dirichlét problem. The reason for this is that the space $F_{p_{1}, q_{1}}^{s_{1}} \supset F_{p, q}^{s+\delta}+F_{r, o}^{t} \ni u$ seemingly may be much too large for an application of the non-linear operator $v \mapsto R_{N}\left(v \partial_{x_{1}} v\right)$ to it.

Indeed, whilst $t=s_{1}$ the integral-exponent $p_{1}$ is smaller than $r$, and in many cases $s_{1}<1+\frac{1}{p_{1}}$ holds, although $t>1+\frac{1}{r}$. An example is sketched in Figure 2 below, where the sector determined by (5.4) is indicated by dashes. When $s_{1}$ is close to $\frac{n}{p_{1}}-1$, then the deficit $\delta\left(s_{1}, p_{1}\right)$ is close to 0 , and so eventually

$$
\begin{equation*}
s_{1}-2+\delta\left(s_{1}, p_{1}\right)<\frac{1}{p_{1}}-1 . \tag{5.5}
\end{equation*}
$$

In such cases, since class $\left(R_{N}\right)=0$, the solution operator $R_{N}$ simply does not make sense on $F_{p_{1}, q_{1}}^{s_{1}-2+\delta\left(s_{1}, p_{1}\right)}(\bar{\Omega})$, as recalled after (5.2). By comparison with the Dirichlét problem, the iteration is seemingly unable to begin.


Figure 2. The integral-exponent $p^{*}$.
At this place the fine theory of pointwise multiplication offers a remedy. In fact, one can do better than regarding the non-linear operator $v \mapsto v \partial_{x_{1}} v$ as one of order $2-\delta(s, p)$, as in (4.4) above.

The problem only arises when $\delta(s, p)$ is close to 0 , hence only for $s<\frac{n}{p}$, and then $v \partial_{1} v$ may be seen to factor through a space with smoothness index $\mathbf{s - 1}$. More exactly, with $\frac{n}{p^{*}}=\frac{n}{p}+\left(\frac{n}{p}-s\right)_{+}$, Theorem 3.2 gives that

$$
\begin{equation*}
B(\cdot): F_{p, q}^{s}(\bar{\Omega}) \rightarrow F_{p^{*}, q}^{s-1}(\bar{\Omega}) \tag{5.6}
\end{equation*}
$$

is a bounded non-linear operator for $s>\max \left(\frac{1}{2}, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)$. With $(s, p, q)$ in the subsector given by $s>\max \left(1, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)$, in which the $\left(s_{1}, p_{1}, q_{1}\right)$ above lies, it may be checked that the receiving space $F_{p^{*}, q}^{s-1}$ in (5.6) lies above the critical broken line $s=\max \left(\frac{1}{p}-1, \frac{n}{p}-n\right)$, so that $\left(s-1, p^{*}, q\right) \in \mathbb{D}_{0}$. See [10] or [11] for this.

According to (5.2) this assures that $R_{N}$ may be applied to $F_{p^{*}, q}^{s-1}(\bar{\Omega})$, so, because $s-1-\frac{n}{p^{*}}=s+\delta(s, p)-\frac{n}{p}$, there is (after all) boundedness of

$$
\begin{equation*}
F_{p, q}^{s}(\bar{\Omega}) \xrightarrow{B(\cdot)} F_{p^{*}, q}^{s-1}(\bar{\Omega}) \xrightarrow{R_{N}} F_{p^{*}, q}^{s+1}(\bar{\Omega}) \hookrightarrow F_{p, q}^{s+\delta(s, p)}(\bar{\Omega}) \tag{5.7}
\end{equation*}
$$

for all $s>\max \left(1, \frac{n}{p}-1+\frac{1}{2} \delta_{n 2}\right)$. Evidently this last sector is stable under the forming of the intermediate parameters $\left(s_{1}, p_{1}, q_{1}\right),\left(s_{2}, p_{2}, q_{2}\right), \ldots$.

In principle also the cases with $t>s+\delta$ and $t-\frac{n}{r}<s+\delta-\frac{n}{p}$ need a special argument, but also here (5.7) may be applied. Altogether the iteration used for the Dirichlét problem applies also to the Neumann problem.

## 6. Final remarks

(1) To prove Theorem 1.3, notice that each of the conditions (1)-(3) there implies that the data belong to the spaces considered in Theorem 2.1 of $[13$, App. 1]. Hence there is a weak solution to which the regularity results apply. For details, see [11, Thm. 5.5.5] or [10].
(2) The iteration methods apply also to the von Karman equations for a plate in $\Omega \subset \mathbb{R}^{2}$, or to problems with a suitable semi-linear perturbation of an injectively elliptic Green operator $\left(\begin{array}{cc}P_{\Omega}+G & K \\ T\end{array}\right)$ in the calculus.

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