# Journées ÉQUATIONS AUX DÉRIVÉES PARTIELLES 

## Dimitri G. Vassiliev

Characteristic properties of distributions associated with the wave group
Journées Équations aux dérivées partielles (1992), p. 1-14
<http://www.numdam.org/item?id=JEDP_1992 $\qquad$ A15_0>
© Journées Équations aux dérivées partielles, 1992, tous droits réservés.
L'accès aux archives de la revue «Journées Équations aux dérivées partielles » (http://www. math.sciences.univ-nantes.fr/edpa/) implique l'accord avec les conditions générales d'utilisation (http://www.numdam.org/conditions). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

# CHARACTERISTIC PROPERTIES OF DISTRIBUTIONS ASSOCIATED WITH THE WAVE GROUP 

D.G.Vassiliev

Let us consider the spectral problem

$$
\begin{gather*}
A v=\lambda^{2 m} v  \tag{A}\\
\left.\left(B^{(j)} v\right)\right|_{\partial M}=0, \quad j=1,2, \ldots, m
\end{gather*}
$$

where $\lambda>0$ is the spectral parameter, $A$ is a positive self-adjoint elliptic linear differential operator of order $2 m(m \in \mathbb{N})$ acting on half-densities on a compact $n$-dimensional ( $n \geq 2$ ) manifold $M$ with boundary $\partial M$ or without boundary $(\partial M=\emptyset)$. The $B^{(j)}$ are "boundary" linear differential operators describing regularly elliptic (see [1]) boundary conditions in the case $\partial M \neq \emptyset$.

Under our assumptions the differential operator $A$ initially defined on

$$
D(A)=\left\{v \in C^{\infty}(M):\left.\left(B^{(j)} v\right)\right|_{\partial M}=0, \quad j=1,2, \ldots, m\right\}
$$

has a self-adjoint closure $\mathcal{A}$ in $L_{2}(M)$ with a domain of definition

$$
D(\mathcal{A})=\left\{v \in H^{2 m}(M):\left.\left(B^{(j)} v\right)\right|_{\partial M}=0, \quad j=1,2, \ldots, m\right\}
$$

where $H^{2 m}(M)$ denotes the Sobolev space of half-densities belonging to $L_{2}(M)$ together with all their partial derivatives of order $\leq 2 m$. It is well known [1] that the operator $\mathcal{A}$ has a positive discrete spectrum $0<\nu_{1} \leq \nu_{2} \leq \ldots$ accumulating to $+\infty$ (we numerate the eigenvalues taking their multiplicities into account). The numbers $\lambda_{k}=\nu_{k}^{1 / 2 m}, k=1,2, \ldots$ may be interpreted as the eigenvalues of the operator $\mathcal{A}^{(1 / 2 m)}$. It is also well known that the respective eigenfunctions (more precisely, half-densities) $v_{k}$ are infinitely smooth on $M$, satisfy (A), (B) and form an orthonormal basis in $L_{2}(M)$.

Note that whereas $\mathcal{A}$ is an operator in the full sense of the word $A$ is not really an operator: it is more correct to call $A$ a "differential expression". Nevertheless, for the sake of simplicity we speak in both cases of "operators", distinguishing $\mathcal{A}$ and $A$ by different script.

This paper is devoted to the study of the operator $\exp \left(-i t \mathcal{A}^{1 /(2 m)}\right)$, $t \in\left(T_{-}, T_{+}\right)$. Information on this operator is essential for obtaining two-term spectral asymptotics for the eigenvalue problem (A), (B), see [2-8].

1. Functional spaces and basic notation. In this section we define precisely the functional spaces of our distributions in order to prepare the ground for the formulation and subsequent proof of Theorem 1.

Local coordinates on $M$ will be denoted by $x=\left(x_{1}, x_{2}, \ldots, x_{n}\right)$ or by $y=$ $\left(y_{1}, y_{2}, \ldots, y_{n}\right)$. In a neighborhood of $\partial M$ we will use only special coordinate systems of the type $x=\left(x^{\prime}, x_{n}\right)$, where $x^{\prime}=\left(x_{1}, x_{2}, \ldots, x_{n-1}\right)$ are coordinates on $\partial M$ (the boundary coordinates) and $x_{n} \geq 0$ is the "normal" coordinate in the sense that $\partial M=\left\{x_{n}=0\right\}$.

Following Schwartz [9, Sect. 3.7] and Hörmander [10, vol. 1, Sect. 2.3] we denote by $\mathcal{E}(M)$ the vector space of infinitely differentiable (up to the boundary!) complexvalued half-densities $v(x)$ on $M$ equipped with the usual $C^{\infty}$-topology defined by the semi-norms

$$
v \rightarrow \sum_{p=1}^{q} \sum_{|\alpha| \leq k} \max _{x \in M^{(p)}}\left|\partial_{x}^{\alpha}\left(\chi_{p} v\right)\right|
$$

where $k$ ranges over all integers $\geq 0$ and $1=\sum_{p=1}^{q} \chi_{p}(x), \chi_{p} \in C^{\infty}(M)$, $\operatorname{supp} \chi_{p} \subset M^{(p)}$, is some partition of unity on $M$ with local coordinates $x$ in coordinate maps $M^{(p)}$. We denote by $\mathcal{E}_{B}(M)$ the subspace of $\mathcal{E}(M)$ consisting of all the half-densities which satisfy the boundary conditions

$$
\begin{equation*}
\left.\left(B^{(j)} A^{r} \bar{v}\right)\right|_{\partial M}=0, \quad j=1,2, \ldots, m, \quad r=0,1,2, \ldots ; \tag{1}
\end{equation*}
$$

the topology on $\mathcal{E}_{B}(M)$ is taken to be the same as on $\mathcal{E}(M)$. By $\mathcal{E}^{\prime}(M), \mathcal{E}_{B}^{\prime}(M)$ we denote the dual spaces of $\mathcal{E}(M), \mathcal{E}_{B}(M)$ respectively (i.e. spaces of linear continuous functionals on $\mathcal{E}(M), \mathcal{E}_{B}(M)$ ) equipped with the dual (strong) topology generated by the initial topology on $\mathcal{E}(M), \mathcal{E}_{B}(M)$, see [9, Sect. 3.2, 3.3 and 3.7]. Obviously, $\mathcal{E}^{\prime}(M) \subset \mathcal{E}_{B}^{\prime}(M)$ because $\mathcal{E}_{B}(M) \subset \mathcal{E}(M)$. The value of the functional (distribution) $u$ on the test half-density $v$ will be denoted by $\langle u, v\rangle_{x}$ with the subscript $x$ emphasizing the variable in which the distribution is acting.

Let $T_{-}<T_{+}$be real numbers, finite or $\pm \infty$. In accordance with traditional notation we denote by $\mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)$the vector space of linear continuous functionals on $C_{0}^{\infty}\left(T_{-}, T_{+}\right)$. The value of the distribution $f \in \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)$on the test function $g \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$will be denoted by $\langle f, g\rangle_{t}$. For the sake of simplicity we, following Hörmander [10, vol. 1, Sect. 2.1], equip $\mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)$with the weak* topology defined by the semi-norms

$$
\mathcal{D}^{\prime}\left(T_{-}, T_{+}\right) \ni f \rightarrow\left|\langle f, g\rangle_{t}\right|
$$

where $g$ is any fixed function from $C_{0}^{\infty}\left(T_{-}, T_{+}\right)$.
By $C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}^{\prime}\left(M_{x}\right)\right), C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}_{B}^{\prime}\left(M_{x}\right)\right)$ we shall denote the class of distributions from $\mathcal{E}^{\prime}\left(M_{x}\right), \mathcal{E}_{B}^{\prime}\left(M_{x}\right)$ respectively infinitely differentiably depending on $(t, y) \in\left(T_{-}, T_{+}\right) \times M_{y}$ as on a parameter (the subscripts $x$ and $y$ are used to distinguish the two copies of the manifold $M$ ). By $C^{\infty}\left(M_{x} \times M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)\right)$we shall denote the class of distributions from $\mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)$infinitely differentiably depending on $(x, y) \in M_{x} \times M_{y}$ as on a parameter. Here infinite differentiability is understood in the strong (Fréchet) sense with account of the respective topologies in $\mathcal{E}^{\prime}\left(M_{x}\right), \mathcal{E}_{B}^{\prime}\left(M_{x}\right), \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)$.

We shall use the notation " $f=O\left(|\lambda|^{-\infty}\right)$ as $\lambda \rightarrow-\infty$ " to describe the fact that the function $f(\lambda) \in C^{\infty}(\mathbb{R})$ vanishes faster than any given negative power of $|\lambda|$ when $\lambda \rightarrow-\infty$. More generally, we shall use this notation for "functions" $f(\lambda, x, y)$ depending on additional parameters $x \in M_{x}, y \in M_{y}$ to describe the fact that $f$ as well as any given derivative of $f$ with respect to $x, y$ vanishes faster than any given negative power of $|\lambda|$ as $\lambda \rightarrow-\infty$ uniformly over $M_{x} \times M_{y}$.

By $\mathcal{F}_{t \rightarrow \lambda}^{-1}(\cdot) \stackrel{\text { def }}{=}(2 \pi)^{-1} \int_{-\infty}^{+\infty} \exp (i t \lambda)(\cdot) d t$ we shall denote the inverse Fourier transform.
2. Main result. The main result of this paper is the following abstract theorem which plays a fundamental role in the asymptotic analysis of higher order ( $m>1$ ) spectral problems. It allows us to avoid the consideration of an ill-posed Cauchy problem (in the variable $t$ ) for the equation $D_{t}^{2 m} u=A_{x} u \quad\left(D_{t} \stackrel{\text { def }}{=}-i \partial / \partial t\right)$; see [6-8] for applications of this theorem.

Theorem 1. Let $T_{-}<0<T_{+}$be real numbers, finite or $\pm \infty$, and let
(2) $u(t, x, y) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}_{B}^{\prime}\left(M_{x}\right)\right) \cap C^{\infty}\left(M_{x} \times M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)\right)$
be a distribution which behaves as a function in the variable $t$ and as a half-density in the variables $x, y$.

If

$$
\begin{equation*}
u(t, x, y)-\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{x} \times M_{y}\right) \tag{3}
\end{equation*}
$$

then

$$
\begin{equation*}
D_{t}^{2 m} u-A_{x} u \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{x} \times M_{y}\right) \tag{4}
\end{equation*}
$$

$$
\begin{gather*}
\left.\left(B_{x}^{(j)} u\right)\right|_{\partial M_{x}} \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times \partial M_{x} \times M_{y}\right), \quad j=1,2, \ldots, m  \tag{5}\\
\mathcal{F}_{t \rightarrow \lambda}^{-1}[g u]=O\left(|\lambda|^{-\infty}\right) \quad \text { as } \quad \lambda \rightarrow-\infty
\end{gather*}
$$

for any $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$.
Inversly, if (4), (5) hold, (6) holds for some $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right), g \not \equiv 0$, and, in addition,

$$
\begin{equation*}
\left.\left(B_{x}^{(j)} A_{x}^{r} u\right)\right|_{\partial M_{x}}=0 \quad \text { at } \quad t=0, \quad j=1,2, \ldots, m, \quad r=0,1,2, \ldots \tag{7}
\end{equation*}
$$

then (3) is fulfilled.
Before proceeeding to the proof of Theorem 1 let us explain in what sense formulas (2)-(7) should be understood (recall that we are dealing with distributions!).

Firstly, (2) means that for any $v(x) \in \mathcal{E}_{B}\left(M_{x}\right), g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$we have the equality

$$
\int_{T_{-}}^{T_{+}}\langle u, v\rangle_{x} g d t=\int_{M_{x}}\langle u, g\rangle_{t} v d x
$$

over all $y \in M_{y}$.
By $\exp \left(-i t \mathcal{A}^{1 /(2 m)}\right), t \in\left(T_{-}, T_{+}\right)$, we denote the bounded operator $L_{2}(M) \rightarrow$ $L_{2}(M)$ (called the unitary exponent) defined by the series

$$
\begin{equation*}
\exp \left(-i t \mathcal{A}^{1 /(2 m)}\right)=\sum_{k=1}^{+\infty} \exp \left(-i t \lambda_{k}\right) v_{k}(x) \int_{M_{y}}(\cdot) \overline{v_{k}(y)} d y \tag{8}
\end{equation*}
$$

where $\lambda_{k}$ are the eigenvalues and $v_{k}$ are the orthonormalized eigenfunctions of the problem (A), (B). In accordance with (8), by $\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y)$ we denote the formal expression
(9) $\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y)=\sum_{k=1}^{+\infty} \exp \left(-i t \lambda_{k}\right) v_{k}(x)\left\langle u(0, x, y), \overline{v_{k}(x)}\right\rangle_{x}$.

Note that if $u(0, x, y)=\delta(x-y)$ then (9) is the Schwartz kernel of the unitary exponent. The formal expression (9) can be understood in the sense of $C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}_{B}^{\prime}\left(M_{x}\right)\right)$ - distributions as well as $C^{\infty}\left(M_{x} \times M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)\right)-$ distributions. In the first case the value of the distribution $\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y)$ on the test half-density $v(x) \in \mathcal{E}_{B}\left(M_{x}\right)$ is defined as

$$
\begin{align*}
\left\langle\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right)\right. & u(0, x, y), v\rangle_{x}  \tag{10}\\
& =\sum_{k=1}^{+\infty} \exp \left(-i t \lambda_{k}\right) \int_{M_{x}} v_{k}(x) v(x) d x\left\langle u(0, x, y), v_{k}(x)\right\rangle_{x}
\end{align*}
$$

In the second case the value of the distribution $\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y)$ on the test function $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$is defined as

$$
\begin{align*}
&\left\langle\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y), g\right\rangle_{t}  \tag{11}\\
&=\sum_{k=1}^{+\infty} v_{k}(x) \int_{T_{-}}^{T_{+}} \exp \left(-i t \lambda_{k}\right) g(t) d t\left\langle u(0, x, y), v_{k}(x)\right\rangle_{x}
\end{align*}
$$

Elementary integration by parts

$$
\begin{gather*}
\int_{M_{x}} v_{k}(x) v(x) d x=\lambda_{k}^{-s} \int_{M_{x}} v_{k}(x)\left(\overline{A_{x}^{s}} v(x)\right) d x  \tag{12}\\
\int_{T_{-}}^{T_{+}} \exp \left(-i t \lambda_{k}\right) g(t) d t=\left(i \lambda_{k}\right)^{-s} \int_{T_{-}}^{T_{+}} \exp \left(-i t \lambda_{k}\right)\left(\partial_{t}^{s} g(t)\right) d t \tag{13}
\end{gather*}
$$

(obviously we used (1) in deriving (12)) with arbitrary $s \in \mathbb{N}$ proves that the quantities $\int_{M_{x}} v_{k}(x) v(x) d x, \int_{T_{-}}^{T_{+}} \exp \left(-i t \lambda_{k}\right) g(t) d t$ vanish faster than any given
XV-4
negative power of $\lambda_{k}$ as $k \rightarrow+\infty$, and, consequently (since we a priori know from classical works on one-term spectral asymptotics the rough estimate $\lambda_{k} \geq c k^{2 m / n}$, $c>0$ ), faster than any given negative power of $k$. On the other hand, in view of standard embedding theorems and of the fact that a distribution from $\mathcal{E}_{B}^{\prime}\left(M_{x}\right)$ always has finite order (due to the definition of the topology on $\mathcal{E}_{B}\left(M_{x}\right)$ ), the quantities $v_{k}(x),\left\langle u(0, x, y), v_{k}(x)\right\rangle_{x}$ grow not faster than some fixed positive power of $k$. This argument shows that the series (10), (11) are absolutely convergent. Moreover, this convergence is uniform over $\left(T_{-}, T_{+}\right) \times M_{y}, M_{x} \times M_{y}$, and does not suffer as a result of differentiation with respect to $(t, y),(x, y)$; see a more detailed discussion after Lemma 8.

Thus, the left-hand side of (3) can be understood in the sense of $C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}_{B}^{\prime}\left(M_{x}\right)\right)$ - distributions as well as $C^{\infty}\left(M_{x} \times M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)\right)$distributions. The $C^{\infty}$-inclusion here means that there exists a $w(t, x, y) \in$ $C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{x} \times M_{y}\right)$ such that

$$
\begin{aligned}
\left\langle u(t, x, y)-\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y), v\right\rangle_{x} & =\int_{M_{x}} w(t, x, y) v(x) d x \\
\left\langle u(t, x, y)-\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y), g\right\rangle_{t} & =\int_{T_{-}}^{T_{+}} w(t, x, y) g(t) d t
\end{aligned}
$$

for $v(x) \in \mathcal{E}_{B}\left(M_{x}\right), g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$.
The expressions $\left.\left(B_{x}^{(j)} u\right)\right|_{\partial M_{x}}$, and, consequently, the inclusions (5), are understood in the sense of $C^{\infty}\left(\partial M_{x} \times M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)\right)$-distributions:

$$
\begin{equation*}
\left.\left\langle\left.\left(B_{x}^{(j)} u\right)\right|_{\partial M_{x}}, g\right\rangle_{t} \stackrel{\text { def }}{=}\left(B_{x}^{(j)}\langle u, g\rangle_{t}\right)\right|_{\partial M_{x}} \tag{14}
\end{equation*}
$$

for $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$. When (6) holds we shall denote the infinitely differentiable "functions" $\left.\left(B_{x}^{(j)} u\right)\right|_{\partial M_{x}}$ by $b_{j}\left(t, x^{\prime}, y\right), j=1,2, \ldots, m$.

The expression $D_{t}^{2 m} u-A_{x} u$, and, consequently, the inclusion (4), can be understood in the sense of $C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}_{B}^{\prime}\left(M_{x}\right)\right)$-distributions (if we already know that (5) holds) or $C^{\infty}\left(M_{x} \times M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)\right)$-distributions. In the first case

$$
\left\langle D_{t}^{2 m} u-A_{x} u, v\right\rangle_{x} \stackrel{\text { def }}{=} D_{t}^{2 m}\langle u, v\rangle_{x}-\left\langle u, \overline{A_{x}} v\right\rangle_{x}-\left.\sum_{j=1}^{m} \int_{\partial M_{x}} b_{j}\left(C_{x}^{(j)} v\right)\right|_{\partial M_{x}} d x^{\prime}
$$

for $v(x) \in \mathcal{E}_{B}\left(M_{x}\right)$; here the $C^{(j)}$ are "boundary" differential operators defined uniquely (for a given set of $B^{(j)}$ ) by Green's formula

$$
\int_{M}(A u(x)) v(x) d x-\int_{M} u(x)(\bar{A} v(x)) d x=\left.\left.\sum_{j=1}^{m} \int_{\partial M}\left(B^{(j)} u\right)\right|_{\partial M}\left(C^{(j)} v\right)\right|_{\partial M} d x^{\prime}
$$

$\forall u \in \mathcal{E}(M), \forall v \in \mathcal{E}_{B}(M)$. In the second case

$$
\left\langle D_{t}^{2 m} u-A_{x} u, g\right\rangle_{t} \stackrel{\text { def }}{=}\left\langle u, D_{t}^{2 m} g \mathcal{A},\right\rangle_{t}-A_{x}\langle u, g\rangle_{t}
$$

for $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$. In both cases the subscript $x$ is used to emphasize the variable in which the respective differential operators are acting.

The expressions $\left.\left(B_{x}^{(j)} A_{x}^{r} u\right)\right|_{\partial M_{x}}$ appearing in the left-hand side of (7) are understood in the sense of $C^{\infty}\left(\partial M_{x} \times M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)\right)$-distributions similarly to (14):

$$
\left.\left\langle\left.\left(B_{x}^{(j)} A_{x}^{r} u\right)\right|_{\partial M_{x}}, g\right\rangle_{t} \stackrel{\text { def }}{=}\left(B_{x}^{(j)} A_{x}^{r}\langle u, g\rangle_{t}\right)\right|_{\partial M_{x}}
$$

for $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$. It is easy to see that if (4), (5) hold then

$$
\begin{align*}
&\left.\left(B_{x}^{(j)} A_{x}^{r} u\right)\right|_{\partial M_{x}} \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times \partial M_{x} \times M_{y}\right)  \tag{15}\\
& j=1,2, \ldots, m, \quad r=0,1,2, \ldots
\end{align*}
$$

so the equalities (7) can be understood in the classical sense. We shall denote the infinitely differentiable "functions" $\left.\left(B_{x}^{(j)} A_{x}^{r} u\right)\right|_{\partial M_{x}}$ by $b_{j r}\left(t, x^{\prime}, y\right)$.

Finally, the expression $\mathcal{F}_{t \rightarrow \lambda}^{-1}[g u]$ is the $C^{\infty}\left(\mathbb{R} \times M_{x} \times M_{y}\right.$ )- "function" (more precisely, function in the variable $\lambda$ and half-density in the variables $x, y$ ) defined as $\mathcal{F}_{t \rightarrow \lambda}^{-1}[g u]=(2 \pi)^{-1}\langle u(t, x, y), \exp (i t \lambda) g(t)\rangle_{t}$.
3. Proof of Theorem 1. If (2), (3) hold then (4)-(6) are obviously fulfilled. So we have to prove only the inverse statement of Theorem 1. It is convenient to split this proof into several parts which we shall consider as separate lemmas.

Lemma 2. Let $b_{j r}\left(t, x^{\prime}\right) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times \partial M\right), \quad j=1,2, \ldots, m$, $r=0,1,2, \ldots$, be a set of "functions" which satisfy $b_{j r}\left(0, x^{\prime}\right)=0$. Then there exists $w(t, x) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M\right)$ such that

$$
\begin{gather*}
\left.\left(B^{(j)} A^{r} w\right)\right|_{\partial M}=b_{j r}, \quad j=1,2, \ldots, m, \quad r=0,1,2, \ldots,  \tag{16}\\
\left.w\right|_{t=0}=0
\end{gather*}
$$

Proof of Lemma 2. In order to simplify notation let us renumerate our boundary "functions", boundary operators and their orders with one index $k=j+m r$ :

$$
b_{j+m r} \stackrel{\text { def }}{=} b_{j r}, \quad B^{(j+m r)} \stackrel{\text { def }}{=} B^{(j)} A^{r}, \quad m_{j+m r} \stackrel{\text { def }}{=} m_{j}+2 m r
$$

$j=1,2, \ldots, m, r=0,1,2, \ldots$ Here $0 \leq m_{1}<m_{2}<\ldots<m_{m} \leq 2 m-1$ are the orders of $B^{(1)}, B^{(2)}, \ldots, B^{(m)}$.

Let us consider some local coordinate system $x=\left(x^{\prime}, x_{n}\right)$ and some compact $K \subset \partial M$ which lies inside the chosen coordinate map. Without loss of generality
we shall assume that $\operatorname{supp}_{x^{\prime}} b_{k} \subseteq K$ for all $t \in\left(T_{-}, T_{+}\right)$and $k=1,2, \ldots$; the general case is reduced to this one by a partition of unity.

Each operator $B^{(k)}, \quad k=1,2, \ldots$, can be represented in the form $B^{(k)}=$ $c_{k}\left(x^{\prime}\right) \partial_{x_{n}}^{m_{k}}-\tilde{B}^{(k)}, c_{k}\left(x^{\prime}\right) \neq 0$, where $\tilde{B}^{(k)}$ is a "boundary" differential operator of order $m_{k}$ without the leading conormal derivative:

$$
\tilde{B}^{(k)}=\sum_{p=0}^{m_{k}-1} \tilde{B}^{(k p)} \partial_{x_{n}}^{p}
$$

Here the $B^{(k p)}$ are differential operators in $x^{\prime}$ of order $\leq m_{k}-p$. Without loss of generality we shall assume that $c_{k}\left(x^{\prime}\right)=1$; the general case is reduced to this one by an obvious renormalization of the operators $B^{(k)}$ and the "functions" $b_{k}\left(t, x^{\prime}\right)$.

Let us construct $w$ as a formal Taylor expansion in $x_{n}$

$$
\begin{equation*}
w(t, x) \sim \sum_{k=1}^{+\infty} w_{k}\left(t, x^{\prime}\right) \frac{x_{n}^{m_{k}}}{m_{k}!} \tag{18}
\end{equation*}
$$

Substituting (18) into (16) we obtain an infinite system of differential equations:

$$
\begin{gather*}
w_{1}=b_{1}  \tag{19}\\
w_{k}=\sum_{l=1}^{k-1}\left(\tilde{B}^{\left(k m_{l}\right)} w_{l}\right)+b_{k}, \quad k=2,3, \ldots
\end{gather*}
$$

Due to its triangular structure this system is solved explicitly: (19) gives $w_{1}$ and (20) gives a recurrent procedure for the determination of $w_{k}, k=2,3, \ldots$ Note that we have $\left.w_{k}\right|_{t=0}=0$ and $\operatorname{supp}_{x^{\prime}} w_{k} \subseteq K$ for all $t \in\left(T_{-}, T_{+}\right)$because the $b_{k}$ possess these properties.

It is known that given an arbitrary formal Taylor expansion (18) one can construct an infinitely smooth "function" $w$ with such Taylor coefficients and with support lying in our coordinate map. Moreover, as in our case the Taylor coefficients vanish at $t=0, w$ can be chosen to vanish identically over $M$ at $t=0$. (These simple statements are proved analogously to [11, Proposition 3.5].)

Lemma 2 is proved.
Lemma 3. Let $t_{-}<0<t_{+}$be finite real numbers and let $a(t, x) \in$ $C^{\infty}\left(\left[t_{-}, t_{+}\right] \times M\right)$ be a "function" which satisfies

$$
\begin{equation*}
\left.\left(B^{(j)} A^{r} a\right)\right|_{\partial M}=0, \quad j=1,2, \ldots, m, \quad r=0,1,2, \ldots \tag{21}
\end{equation*}
$$

Then there exists $w(t, x) \in C^{\infty}\left(\left[t_{-}, t_{+}\right] \times M\right)$ such that equalities

$$
\begin{gather*}
D_{t}^{2 m} w=A w+a,  \tag{22}\\
\text { XV- } 7
\end{gather*}
$$

$$
\begin{equation*}
\left.\left(B^{(j)} A^{r} w\right)\right|_{\partial M}=0, \quad j=1,2, \ldots, m, \quad r=0,1,2, \ldots, \tag{23}
\end{equation*}
$$

and (17) hold.
Proof of Lemma 3. Set

$$
\begin{equation*}
w(t, x)=\sum_{k=1}^{+\infty} u_{k}(x)\left(w_{k}(t)-w_{k}(0) \cos \lambda_{k} t\right) \tag{24}
\end{equation*}
$$

$$
\begin{equation*}
w_{k}(t)=\sum_{l=1}^{m} w_{k l}^{+}(t)+\sum_{l=m+1}^{2 m} w_{k l}^{-}(t) \tag{25}
\end{equation*}
$$

$$
\begin{equation*}
w_{k l}^{ \pm}(t)=-\frac{i \exp \left(-i t \lambda_{k l}\right)}{2 m \lambda_{k l}^{2 m-1}} \int_{t_{ \pm}}^{t} \exp \left(i \tau \lambda_{k l}\right) a_{k}(\tau) d \tau \tag{26}
\end{equation*}
$$

$$
\begin{equation*}
a_{k}(t)=\int_{M} a(t, x) \overline{v_{k}(x)} d x \tag{27}
\end{equation*}
$$

$$
\begin{equation*}
\lambda_{k l}=\lambda_{k} \exp \left(i \pi(l-1) m^{-1}\right), \quad l=1,2, \ldots, 2 m \tag{28}
\end{equation*}
$$

Recall that by $\lambda_{k}$ and $v_{k}, k=1,2, \ldots$, we denote the eigenvalues and the orthonormalized eigenfunctions of the problem (A), (B). Due to the boundary conditions (21) the quantity $a_{k}(t)$ defined by (27) vanishes faster than any given negative power of $k$ as $k \rightarrow+\infty$ uniformly over $\left[t_{-}, t_{+}\right]$; the same is true for any given derivative of $a_{k}(t)$ with respect to $t$. An elementary analysis of formulas (25), (26), (28) shows that this rapid decay property is inherited by the terms of the series (24): this series converges absolutely, uniformly over $\left[t_{-}, t_{+}\right] \times M$, as well as the series of any given derivatives with respect to $t, x$. Thus (24) defines an infinitely smooth "function" which can be differentiated under the $\sum_{k=1}^{+\infty}$ sign.

Straightforward substitution shows that the constructed $w$ satisfies (22), (23) and (17).

Remark 4. Lemmas 2, 3 and their proofs remain true if the "functions" $b_{j r}$ and $a$ depend smoothly on the additional parameter $y \in M_{y}$. In this case the resulting $w$ will also smoothly depend on $\boldsymbol{y}$.

Lemma 5. It is sufficient to prove the inverse statement of Theorem 1 with

$$
\begin{gather*}
D_{t}^{2 m} u=A_{x} u  \tag{29}\\
\left.\left(B_{x}^{(j)} u\right)\right|_{\partial M_{x}}=0, \quad j=1,2, \ldots, m \tag{30}
\end{gather*}
$$

instead of (4), (5), (7).
Proof of Lemma 5. It follows immediately from (13), Lemma 2 and Remark 4 that we can turn (5), (7) into (30) by adding to $u$ an infinitely smooth with respect to all the variables "function" which vanishes at $t=0$. Obviously such an operation does not spoil (2), (4), (6) and it also has no infiuence on the formula (3) which we are proving.

Take now arbitrary finite real numbers $t_{-}, t_{+}$such that

$$
\begin{equation*}
T_{-}<t_{-}<0<t_{+}<T_{+} . \tag{31}
\end{equation*}
$$

Formulas (4), (30) imply that the "function" $a(t, x, y) \stackrel{\text { def }}{=} D_{t}^{2 m} u-A_{x} u$ satisfies the conditions of Lemma 3. It follows from Lemma 3 and Remark 4 that we can turn (4) into (29) by adding to $u$ an infinitely smooth with respect to all the variables "function" which satisfies (23) in the variable $x$ and which vanishes at $t=0$. Such an operation does not spoil (2), (6), (30) and it does not influence the formula (3) which we are proving, only the time interval becomes smaller $\left(t_{-}, t_{+}\right)$instead of $\left(T_{-}, T_{+}\right)$). Assuming that the inverse statement of Theorem 1 with (4), (5), (7) replaced by $(29),(30)$ is true, we have

$$
\begin{equation*}
u(t, x, y)-\exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y) \in C^{\infty}\left(\left(t_{-}, t_{+}\right) \times M_{x} \times M_{y}\right) \tag{32}
\end{equation*}
$$

As $t_{-}, t_{+}$are arbitrary numbers satisfying (31) formula (32) implies (3).
Lemma 5 is proved.
Thus we have reduced the proof of Theorem 1 to the proof of the following statement.

Let $u(t, x, y)$ be a distribution of the class (2) which behaves as a function in the variable $t$ and as a half-density in the variables $x, y$ and which satisfies (29), (30) and (6) for some $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right), g \not \equiv 0$. Then (3) is fulfilled.

Set

$$
\begin{equation*}
u_{k}(t, y)=\left\langle u(t, x, y), \overline{v_{k}(x)}\right\rangle_{x} \tag{33}
\end{equation*}
$$

In view of (2) we have

$$
\begin{equation*}
u_{k}(t, y) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y}\right) \tag{34}
\end{equation*}
$$

Moreover, formulas (29), (30) and (33) imply $D_{t}^{2 m} u_{k}=\lambda_{k} u_{k}$, and consequently

$$
\begin{gather*}
u_{k}(t, y)=\sum_{l=1}^{2 m} u_{k l}(t, y)  \tag{35}\\
u_{k l}(t, y)=u_{k l}(0, y) \exp \left(-i t \lambda_{k l}\right)=\sum_{p=1}^{2 m} b_{l p}\left(-\lambda_{k}\right)^{1-p} D_{t}^{p-1} u_{k}(t, y)
\end{gather*}
$$

Here $\left\|a_{l p}\right\|$ is the symmetric $2 m$ by $2 m$ matrix with elements

$$
a_{l p}=\exp \left(i \pi(l-1)(p-1) m^{-1}\right), \quad l, p=1,2, \ldots, 2 m
$$

and $\left\|b_{l p}\right\|=\left\|a_{l p}\right\|^{-1}$. The $\lambda_{k l}$ are defined by (28). Note that formula (36) allows to express the functions $u_{k l}(t, y)$ in terms of derivatives of the function $u_{k}(t, y)$; with account of (34) this gives $u_{k l}(t, y) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y}\right)$.

Definition 6. Consider a sequence of "functions" (more precisely, functions in the variable $t$ and half-densities in the variable $y$ )

$$
w_{k}(t, y) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y}\right), \quad k=1,2, \ldots
$$

and let $1=\sum_{p=1}^{q} \chi_{p}(y), \chi_{p} \in C^{\infty}\left(M_{y}\right), \operatorname{supp} \chi_{p} \subset M_{y}^{(p)}$, be some partition of unity on $M_{y}$ with local coordinates $y$ in coordinate maps $M_{y}^{(p)}$. We will say that this sequence increases slowly if for any real numbers $t_{-}, t_{+}$satisfying (31), any multiindex $\alpha \geq 0$ and any integer $r \geq 0$ there exists a natural $s$ such that $k^{-s} \partial_{t}^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) w_{k}(t, y)\right) \rightarrow 0$ as $k \rightarrow+\infty$ uniformly over $t \in\left[t_{-}, t_{+}\right], y \in M_{y}^{(p)}$, $p=1,2, \ldots, q$. We will say that this sequence decreases rapidly if for any real numbers $t_{-}, t_{+}$satisfying (31), any multiindex $\alpha \geq 0$, any integer $r \geq 0$ and any natural $s \quad k^{s} \partial_{t}^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) w_{k}(t, y)\right) \rightarrow 0$ as $k \rightarrow+\infty$ uniformly over $t \in\left[t_{-}, t_{+}\right]$, $y \in M_{y}^{(p)}, p=1,2, \ldots, q$.

Lemma 7. The sequence $u_{k}(t, y), k=1,2, \ldots$, increases slowly.
Proof of Lemma 7. Suppose that the statement of Lemma 7 is false. Then there exist real numbers $t_{-}, t_{+}$satisfying (31), a multiindex $\alpha \geq 0$, an integer $r \geq 0$, a coordinate map $M_{y}^{(p)}$ with local coordinates $y$, a cut-off function $\chi_{p}(y) \in C^{\infty}\left(M_{y}\right)$ with $\operatorname{supp} \chi_{p} \subset M_{y}^{(p)}$, and sequences $k_{s} \in \mathbb{N}, t_{s} \in\left[t_{-}, t_{+}\right], y_{s} \in \operatorname{supp} \chi_{p}$, $s=1,2, \ldots$, such that $k_{s} \rightarrow+\infty$ and

$$
\begin{equation*}
\left.\left(k_{s}\right)^{-s}\left(\partial_{t}^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) u_{k_{s}}(t, y)\right)\right)\right|_{t=t_{s}, y=y_{s}} \nrightarrow 0 \tag{37}
\end{equation*}
$$

as $s \rightarrow+\infty$. Without loss of generality we shall assume that $t_{s} \rightarrow \tilde{t} \in\left[t_{-}, t_{+}\right]$, $y_{s} \rightarrow \tilde{y} \in \operatorname{supp} \chi_{p}$ as $s \rightarrow+\infty$; this can always be achieved by extracting subsequences in view of the compactness of $\left[t_{-}, t_{+}\right] \times \operatorname{supp} \chi_{p}$. Let us introduce the set of half-densities $\mathcal{B}=\left\{\left(k_{s}\right)^{-s} \overline{v_{k_{s}}(x)}: s=1,2, \ldots\right\}$. Obviously $\mathcal{B}$ is a bounded set (see [9, Sect. 3.2 and 3.7]) in $\mathcal{E}_{B}(M)$. Denote by $\tilde{u}(t, x, y)$ the distribution defined by the formula $\langle\tilde{u}, v\rangle_{x}=\partial_{t}^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y)\langle u(t, x, y), v(x)\rangle_{x}\right), v(x) \in \mathcal{E}_{B}\left(M_{x}\right)$. Due to (2) we have $\tilde{u}(t, x, y) \in C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}_{B}^{\prime}\left(M_{x}\right)\right)$. So, in particular, $\tilde{u}(t, x, y)$ is a continuous function of the parameters $t, y$ at the point $t=\tilde{t}$, $y=\tilde{y}$ in the sense of the $\mathcal{E}_{B}^{\prime}\left(M_{x}\right)$ )-topology which means (see [9, Sect. 3.3 and 3.7]) that $\langle\tilde{u}(t, x, y), v(x)\rangle_{x} \rightarrow\langle\tilde{u}(\tilde{t}, x, \tilde{y}), v(x)\rangle_{x}$ as $t \rightarrow \tilde{t}, y \rightarrow \tilde{y}$ uniformly over all $v(x) \in \mathcal{B}$, where $\mathcal{B}$ is an arbitrary bounded set in $\mathcal{E}_{B}(M)$. But this contradicts (37).

Lemma 7 is proved.
Lemma 8. For any $l=1,2, \ldots, 2 m$ the sequence $u_{k l}(t, y), k=1,2, \ldots$, increases slowly.

Proof of Lemma 8. Lemma 8 follows immediately from Lemma 7 and formula (36).

Lemma 8 allows us to represent our distribution $u(t, x, y)$ in the form of a series of smooth "functions"

$$
\begin{equation*}
u(t, x, y)=\sum_{k=1}^{+\infty} \sum_{l=1}^{2 m} u_{k l}(t, y) v_{k}(x) \tag{38}
\end{equation*}
$$

with the values of this distribution on the test "functions" $v(x) \in \mathcal{E}_{B}(M)$, $g(t) \in C_{0}^{\infty}\left(T_{-}, T_{+}\right)$given by the natural formulas

$$
\begin{align*}
\langle u(t, x, y), v\rangle_{x} & =\sum_{k=1}^{+\infty} \sum_{l=1}^{2 m} u_{k l}(t, y) \int_{M_{x}} v_{k}(x) v(x) d x  \tag{39}\\
\langle u(t, x, y), g\rangle_{t} & =\sum_{k=1}^{+\infty} \sum_{l=1}^{2 m} v_{k}(x) \int_{T_{-}}^{T_{+}} u_{k l}(t, y) g(t) d t
\end{align*}
$$

It follows from Lemma 8 and formulas (12),

$$
\begin{equation*}
\int_{T_{-}}^{T_{+}} u_{k l}(t, y) g(t) d t=\left(i \lambda_{k l}\right)^{-s} \int_{T_{-}}^{T_{+}} u_{k l}(t, y)\left(\partial_{t}^{s} g(t)\right) d t \tag{41}
\end{equation*}
$$

(cf. (13)) with arbitrary $s \in \mathbb{N}$ that the series (39), (40) converge absolutely, uniformly over $\left(T_{-}, T_{+}\right) \times M_{y}, M_{x} \times M_{y}$, as well as the series of any given derivatives with respect to $(t, y),(x, y)$. This justifies the representation (38) in the sense of $C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{y} ; \mathcal{E}_{B}^{\prime}\left(M_{x}\right)\right)$-distributions as well as $C^{\infty}\left(M_{x} \times\right.$ $M_{y} ; \mathcal{D}^{\prime}\left(T_{-}, T_{+}\right)$)-distributions.

Lemma 9. For any $l \neq 1, m+1$ the sequence $u_{k l}(t, y), k=1,2, \ldots$, decreases rapidly.

Proof of Lemma 9. Assume $m \geq 2$ (otherwise none of the $l=1,2, \ldots, 2 m$ satisfy the condition $l \neq 1, m+1)$. Let us consider first the case $2 \leq l \leq m$. Having fixed $l$ and arbitrary $s, r, \alpha, \chi_{p}$ (in the notation of Definition 6 ) set

$$
\mu_{k}\left(t_{-}, t_{+}\right) \stackrel{\text { def }}{=} \max _{t \in\left[t_{-}, t_{+}\right], y \in M_{y}^{(p)}} k^{s}\left|\partial_{t}^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) u_{k l}(t, y)\right)\right|
$$

where $t_{-}, t_{+}$are arbitrary real numbers satisfying (31). Set $\varepsilon=\left(T_{+}-t_{+}\right) / 2$. Due to the exponential behavior of the "function" $u_{k l}(t, y)$ in the variable $t$ (see (36)) we have

$$
\begin{equation*}
\mu_{k}\left(t_{-}, t_{+}\right)=\exp \left(-\varepsilon \lambda_{k} \sin (\pi / m)\right) \mu_{k}\left(t_{-}, t_{+}+\varepsilon\right) \tag{42}
\end{equation*}
$$

According to Lemma 8 and Definition 6 there exists a natural $\tilde{s}$ such that

$$
\begin{equation*}
k^{-\tilde{s}} \mu_{k}\left(t_{-}, t_{+}+\varepsilon\right) \rightarrow 0 \quad \text { as } \quad k \rightarrow+\infty \tag{43}
\end{equation*}
$$

Combining (42) and (43) we get

$$
\begin{equation*}
\mu_{k}\left(t_{-}, t_{+}\right) \leq c k^{s+\tilde{s}} \exp \left(-\varepsilon \lambda_{k} \sin (\pi / m)\right) \tag{44}
\end{equation*}
$$

where $c=\max _{k \in \mathbb{N}}\left(k^{-\tilde{s}} \mu_{k}\left(t_{-}, t_{+}+\varepsilon\right)\right)$. As the right-hand side of (44) contains an exponential term which is obviously stronger than the term $k^{s+\tilde{s}}$, it follows from (44) that $\mu_{k}\left(t_{-}, t_{+}\right)$vanishes as $k \rightarrow+\infty$. This means rapid decrease in the sense of Definition 6.

The case $m+2 \leq l \leq 2 m$ is handled similarly by estimating $\mu_{k}\left(t_{-}, t_{+}\right)$through $\mu_{k}\left(t_{-}-\varepsilon, t_{+}\right), \varepsilon=\left(t_{-}-T_{-}\right) / 2$.

Lemma 9 is proved.

Lemma 10. The sequence $u_{k(m+1)}(t, y), k=1,2, \ldots$, decreases rapidly.
Proof of Lemma 10. As in the proof of the previous lemma let us fix arbitrary $s, r, \alpha, \chi_{p}$. It follows from (6) that

$$
\begin{equation*}
k^{s}\left(i \lambda_{k}\right)^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) \int_{M_{x}} \mathcal{F}_{t \rightarrow \lambda}^{-1}[g(t) u(t, x, y)] \overline{v_{k}(x)} d x\right) \rightarrow 0 \text { as } \lambda \rightarrow-\infty \tag{45}
\end{equation*}
$$

uniformly over $y \in M_{y}^{(p)}, k \in \mathbb{N}$; here uniformity over $k$ is established by integration by parts in the variable $x$ (similarly to (12)) with account of the boundary conditions (30). Substitution of (38), (36) for $u(t, x, y)$ turns (45) into

$$
\begin{equation*}
\left.k^{s}\left(i \lambda_{k}\right)^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) \sum_{l=1}^{2 m} \mathcal{F}_{t \rightarrow\left(\lambda-\lambda_{k l}\right)}^{-1}[g(t)] u_{k l}(0, y)\right)\right) \rightarrow 0 \text { as } \lambda \rightarrow-\infty \tag{46}
\end{equation*}
$$

Let us choose and fix some $\mu \in \mathbb{R}$ such that $\mathcal{F}_{t \rightarrow \mu}^{-1}[g(t)] \neq 0$ (such a $\mu$ exists because $g \not \equiv 0$ ), and let us relate $\lambda$ and $k$ by the condition $\lambda=\mu-\lambda_{k}$. Then (46) takes the form

$$
\begin{equation*}
\sum_{l=1}^{2 m} g_{k l}(y) \rightarrow 0 \quad \text { as } \quad k \rightarrow+\infty \tag{47}
\end{equation*}
$$

where

$$
\begin{equation*}
g_{k l}(y)=k^{s}\left(i \lambda_{k}\right)^{r} \mathcal{F}_{t \rightarrow\left(\mu-\lambda_{k}-\lambda_{k l}\right)}^{-1}[g(t)]\left(\partial_{y}^{\alpha}\left(\chi_{p}(y) u_{k l}(0, y)\right)\right) \tag{48}
\end{equation*}
$$

Formula (48) can also be rewritten as

$$
g_{k l}(y)=\frac{\exp \left(-i \pi r(l-1) m^{-1}\right)}{2 \pi} \int_{T_{-}}^{T_{+}} \exp \left(i t\left(\mu-\lambda_{k}\right)\right) g(t)\left(\partial_{t}^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) u_{k l}(t, y)\right)\right) d t
$$

in the case $l \neq 1, m+1$ the integral in this formula is easily estimated with the help of Lemma 9 which gives

$$
\begin{equation*}
g_{k l}(y) \rightarrow 0 \quad \text { as } \quad k \rightarrow+\infty, \quad l \neq 1, m+1 \tag{49}
\end{equation*}
$$

Let us examine now formula (48) in the case $l=1$. As $\lambda_{k 1}=\lambda_{k}$ the sequence $\mathcal{F}_{t \rightarrow\left(\mu-\lambda_{k}-\lambda_{k 1}\right)}^{-1}[g(t)]$ decreases rapidly (the usual property of the Fourier transform of a $C_{0}^{\infty}$-function), and the presence of slowly increasing factors $k^{s}$, $\left(i \lambda_{k}\right)^{r}, \partial_{y}^{\alpha}\left(\chi_{p}(y) u_{k l}(0, y)\right)$ ) (see Lemma 8) in the right-hand side of (48) cannot spoil this rapid decrease. Thus

$$
\begin{equation*}
g_{k_{1}}(y) \rightarrow 0 \quad \text { as } \quad k \rightarrow+\infty \tag{50}
\end{equation*}
$$

Formulas (47)-(50) and $\lambda_{k(m+1)}=-\lambda_{k}$ imply that
$g_{k(m+1)}(y)=k^{s}\left(i \lambda_{k}\right)^{r} \mathcal{F}_{t \rightarrow \mu}^{-1}[g(t)]\left(\partial_{y}^{\alpha}\left(\chi_{p}(y) u_{k(m+1)}(0, y)\right)\right) \rightarrow 0 \quad$ as $\quad k \rightarrow+\infty$.
Dividing the latter formula by the non-zero constant $\mathcal{F}_{t \rightarrow \mu}^{-1}[g(t)]$ and multiplying it by the function $\exp \left(i t \lambda_{k}\right)$ (which obviously has unit modulus) we arrive at

$$
\begin{equation*}
k^{s}\left(\partial_{t}^{r} \partial_{y}^{\alpha}\left(\chi_{p}(y) u_{k(m+1)}(t, y)\right)\right) \rightarrow 0 \quad \text { as } \quad k \rightarrow+\infty \tag{51}
\end{equation*}
$$

Formula (51) is uniform over $y \in M_{y}^{(p)}$ because the initial formula (45) had this property and all our arguments leading from (45) to (51) preserved uniformity in $y$. Moreover, as the time dependence in $u_{k(m+1)}(t, y)$ is trivial (purely imaginary exponent) formula (51) is uniform over all $t \in \mathbb{R}$. So (51) means rapid decrease of the functional sequence $u_{k(m+1)}(t, y), k=1,2, \ldots$, in the sense of Definition 6.

Lemma 10 is proved.
Now it only remains to rearrange (38) as

$$
\begin{align*}
u(t, x, y)= & \exp \left(-i t \mathcal{A}_{x}^{1 /(2 m)}\right) u(0, x, y)  \tag{52}\\
& +\sum_{k=1}^{+\infty} \sum_{l=2}^{2 m}\left(u_{k l}(t, y)-u_{k l}(0, y) \exp \left(-i t \lambda_{k}\right)\right) v_{k}(x)
\end{align*}
$$

Due to Lemmas 9,10 the infinite sum in the right-hand side of (52) defines a $C^{\infty}\left(\left(T_{-}, T_{+}\right) \times M_{x} \times M_{y}\right)$-"function". So the inclusion (3), and with it our Theorem 1 , is proved.

## References

1. J. L. Lions and E. Magenes, Problèmes aux limites non homogènes et applications, vol. 1, Dunod, Paris, 1968; English transl., Springer-Verlag, Berlin and New York, 1972.
2. J. J. Duistermaat and V. W. Guillemin, The spectrum of positive elliptic operators and periodic bicharacteristics, Invent. Math. 29 (1975), 39-79.
3. V. Ya. Ivriur, On the second term of the spectral asymptotics for the Laplace-Beltrami operator on manifolds with boundary, Funktsional. Anal. i Prilozhen. 14 (1980), no. 2, 25-34; English transl. in Functional Anal. Appl. 14 (1980), 98-106.
4. Victor Ivrǐ, Precise spectral asymptotics for elliptic operators acting in fiberings over manifolds with boundary, Lecture Notes in Math., vol. 1100, Springer-Verlag, Berlin and New York, 1984.
5. R. Melrose, Weyl's conjecture for manifolds with concave boundary, Proc. Sympos. Pure Math., vol. 36, Amer. Math. Soc., Providence, R.I., 1980, pp. 257-273.
6. D. G. Vasil'ev, Binomial asymptotics of the spectrum of a boundary value problem, Funktsional. Anal. i Prilozhen. 17 (1983), no. 4, 79-81; English transl. in Functional Anal. Appl. 17 (1984), 309-311.
7. D. G. Vasil'ev, Two-term asymptotics of the spectrum of a boundary value problem under an interior reflection of general form, Funktsional. Anal. i Prilozhen. 18 (1984), no. 1, 1-13; English transl. in Functional Anal. Appl. 18 (1984), 267-277.
8. D. G. Vasil'ev, Asymptotics of the spectrum of a boundary value problem, Trudy Moskov. Mat. Obshch. 49 (1986), 167-237; English transl. in Trans. Moscow Math. Soc. (1987), 173-245. 9. L.Schwartz, Théorie des distributions, Hermann, Paris, 1966.
9. Lars Hörmander, The analysis of linear partial differential operators, vol. 1, SpringerVerlag, Berlin and New York, 1983.
10. M. A. Shubin, Pseudodifferential operators and spectral theory, "Nauka", Moscow, 1980; English transl., Springer-Verlag, Berlin and New York, 1986.
D.Vassiliev

Mathematics Division
School of Mathematical and Physical Sciences
University of Sussex
Falmer, Brighton BN1 9QH, England
E-mail address: D.Vassiliev@ central.sussex.ac.uk

