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#### RESONANCE FUNCTIONS OF TWO-BODY SCHRÖDINGER OPERATORS

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We consider the Schrödinger operator  $-\Delta + V$  in  $L^2(\mathbb{R}^n)$ ,  $n \geq 3$ , where V is a short-range, dilation-analytic potential in an angle  $S_\alpha$ . A resonance  $\lambda_0$  appears as a discrete eigenvalue of the complex-dilated Hamiltonian [2], a pole of the S-matrix [3] and as a pole of the analytically continued resolvent, acting from an exponentially weighted space to its dual [4,5]. In [2] resonance functions are obtained as square-integrable eigenfunctions of the complex-dilated Hamiltonian, corresponding to the eigenvalue  $\lambda_0$ , in [5] they are defined as certain exponentially growing solutions f of the Schrödinger equation  $(-\Delta + V - \lambda_0)f = 0$ . In [6] it is proved that for a dilation-analytic multiplicative potential V with resonance  $\lambda_0$ , the resonance functions of [2] and [5] are simply the restrictions of one analytic,  $L^2(S^{n-1})$ -valued function f on  $S_\alpha$  to rays  $e^{i\phi}IR^+$  with  $2\phi > -Arg\lambda_0$  and to  $IR^+$ , respectively.

Moreover, the precise asymptotic behaviour  $f(z)=\frac{ik_0z}{2}$   $\frac{1-n}{z}(\tau+0(|z|^{-\epsilon}))$  with  $\tau\in L^2(S^{n-1})$ , where  $k_0^2=\lambda_0$ , is established together with asymptotics for f'(z). These imply exponential decay in time of resonance states, defined as suitably cut-off resonance functions, as proved in [8].

In this note we shall give a brief account of results on resonance functions, referring for details to [5] and [6].

### 1. Analytic continuation of resolvent and S-matrix

We introduce the weighted  $L^2$ -spaces  $L^2_{\delta,b} = L^2_{\delta,b}(\mathbb{R}^n)$  for  $\delta,b\in\mathbb{R}$  by

$$L_{\delta,b}^{2} = \{f \mid \|f\|_{\delta,b}^{2} = \int_{\mathbb{R}^{n}} |f(x)|^{2} (1+r^{2})^{\delta} e^{2br} dx < \infty \}$$

where  $x \in \mathbb{R}^n$ , r = |x|. The weighted Sobolev spaces are defined by

$$H_{\delta,b}^2 = \{f \mid \|f\|_{2,\delta,b}^2 = \sum_{|\alpha| \le 2} \|D^{\alpha}f\|_{\delta,b}^2 < \infty \}$$

We set  $L_{\delta,0}^2 = L_{\delta}^2$ ,  $H_{\delta,0}^2 = H^2$  and  $h = L^2(S^{n-1})$ ,  $S^{n-1} = \{x \in {\rm IR}^n \mid |x| = 1\}$ . We assume that the dimension  $n \ge 3$ 

$$\mathfrak{C}^+ = \{k \in \mathfrak{C} \mid \text{Im } k > 0\}, \widetilde{\mathfrak{C}}^+ = \overline{\mathfrak{C}^+} \setminus \{0\}.$$

 $B(H_1,H_2)$  and  $C(H_1,H_2)$  denote the spaces of bounded and compact operators from  $H_1$  into  $H_2$ , respectively.

The free Hamiltonian  $H_0$  in  $L^2$  is defined for  $u \in \mathcal{D}_{H_0} = H^2$  by  $H_0 u = -\Delta u$  with resolvent  $R_0(k) = (H_0 - k^2)^{-1} \in \mathcal{B}(L^2)$  for  $k \in \mathfrak{C}^+$ .

The interaction Q is assumed to be a symmetric, short-range,  $S_{\alpha} \text{-dilation-analytic operator in } L^2 \text{. Thus, } Q \in \mathbb{C}(H^2_{-\delta_0}, L^2_{\delta_0}) \text{ for some } \delta_0 > \frac{1}{2} \text{, and if } \{U(\rho)\} \text{ is the dilation group on } L^2 \text{ defined by }$ 

$$(U(\rho)f)(x) = \rho^{\frac{11}{2}}f(\rho x)$$

then the function  $Q(\rho) = U(\rho)QU(\rho^{-1})$  on  $\mathbb{R}^+$  has an analytic,  $C(H_{-\delta_0}^2, L_{\delta_0}^2)$ -valued analytic extension to the angle

$$S_{\alpha} = \{ \rho e^{i\phi} \mid \rho > 0, |\phi| < \alpha \}$$

Moreover,  $Q(z) \in C(H^2_{-\delta_0,b}, L^2_{\delta_0,b})$  for all  $b \in \mathbb{R}$ . (This follows from  $Q \in C(H^2_{-\delta_0}, L^2_{\delta_0})$  if Q is local).

The Hamiltonian  $H=H_0+Q$  is self-adjoint on  $\mathcal{D}_H=H^2$  , and associated with H is a self-adjoint, analytic family of type A , H(z) , given by

$$H(z) = z^{-2}H_0 + Q(z)$$
,

and  $H(\rho e^{i\phi}) = U(\rho)H(e^{i\phi})U(\rho^{-1})$ , so  $\sigma(H(z)) = \sigma(H(e^{i\phi}))$  for  $\rho > 0$ ,  $z = \rho e^{i\phi}$ .

We define the operators  $H_z$  and their resolvents  $R_z$ (k) by

$$H_z = H_0 + z^2 Q(z) = z^2 H(z)$$
,  $R_z(k) = (H_z - k^2)^{-1}$ .

We note that  $R_z(zk) = z^{-2}(H(z) - k^2)^{-1}$ .

We have  $\sigma_e(H(z)) = e^{-2i\phi}R^+$  and  $\sigma_d(H(z)) \setminus R \subset \{\lambda \mid -2\phi < Arg \lambda < 0\}$ .

We define  $R(\phi)$  by  $R(\phi)=\{k\mid 0> {\rm Arg}\, k>-\phi$  ,  $k^2\in\sigma_{\tilde{\bf d}}(H(z))\}$  , R=U  $R(\phi)$  . The points  $\lambda=k^2$  ,  $k\in R$  , are called resonances.  $0<\phi<\alpha$ 

For our analysis we need the following result, proved in [5]:

 $\underline{\text{Lemma 1.1}}. \quad \text{For } \delta > 0 \quad \text{let } S_{\alpha}^{\delta} = \{k \in S_{\alpha} \mid \text{Im}(e^{\mathbf{i} (\alpha - \delta)}k) < \epsilon\} \ .$  There exists  $S_{\alpha}$ -dilation-analytic interactions  $V_{\epsilon}$  and  $W_{\epsilon}$  with  $Q = V_{\epsilon} + W_{\epsilon}$ , such that  $H_{0} + V_{\epsilon}$  has no resonances outside  $(S_{\alpha}^{\delta})^{2}$  and  $W_{\epsilon}$  decays faster than any exponential. This holds with  $W_{\epsilon} = g_{\epsilon} Q g_{\epsilon}$ , where  $g_{\epsilon}(\mathbf{r}) = \exp(-\epsilon \mathbf{r}^{\beta})$ ,  $\beta = \frac{\pi}{2\pi}$ , for  $\epsilon$  small.

Using Lemma 1.1 one can prove all results for fixed  $\delta>0$  with  $S_{\alpha}$  replaced by  $S_{\alpha}-S_{\alpha}^{\delta}$ , using the splitting  $Q=V_{\epsilon}+W_{\epsilon}$ , and then let  $\delta \downarrow 0$ . To simplify the presentation, we assume from the outset (although this can strictly speaking not be obtained) that  $H_1=H_0+V$  has no resonances and fix  $\epsilon$ , setting  $g=g_{\epsilon}$ , W=Qg, V=Q-gW. We denote by  $H_{1z}$ ,  $R_{1z}(k)$  etc. the operators obtained by replacing Q by V.

Basic to our approach is an extended limiting absorption principle proved in [7] and generalized in [5] to non-symmetric, short-range potentials like  $Q_z$ . The idea is to consider  $-\Delta$  and  $-\Delta+Q_z$  as operators  $H_0^{-b}$  and  $H_z^{-b}$  acting in the space  $L_{0,-b}^2$ ,  $b\geq 0$ . The spectrum of  $H_0^{-b}$  and the essential spectrum of  $H_z^{-b}$  coincide with the parabolic region  $P_b=\{k^2\mid |\mathrm{Im}k|\leq b\}$ , and it is then proved that the resolvents  $(H_0^{-b}-(a+ib+i\epsilon)^2)^{-1}$  and  $(H_z^{-b}-(a+ib+i\epsilon)^2)^{-1}$  have boundary values as  $\epsilon \downarrow 0$  in  $B(L_{\delta,-b}^2$ ,  $H_{-\delta-b}^2$ ) for  $\frac{1}{2}<\delta \leq \delta_0$ , except at the so-called singular points.

The singular sets  $\sum_z^c$  ,  $\sum_z^r$  and  $\sum_z$  are defined for z =  $\rho e^{i\phi}$  ,  $\phi$  > 0 , by

$$\sum_{z}^{c} = \{k \in \mathbb{C}^{+} \mid k^{2} = z^{2}\lambda, \lambda \in \sigma_{d}(H(z))\},$$

$$\sum_{z}^{r} = z R \cap IR^{+}$$
 ,  $\sum_{z} = \sum_{z}^{c} \cup \sum_{z}^{r}$  ,

and for  $\phi < 0$  by  $\sum_{z}^{c} = -\frac{\sum_{z}^{c}}{z}$  and similar for  $\sum_{z}^{r}$  and  $\sum_{z}$ . For  $\phi = 0$ ,  $\sum_{z}^{c} = \sum_{z}^{c} \cup \sum_{z}^{r} = \{k \in \widetilde{\mathbb{C}}^{+} \mid k^{2} \in \sigma_{p}(H)\}$ .

The extended limiting absorption principle for  ${\rm H}_{\rm Z}$  can then be formulated as follows:

$$R_z(k) = e^{ikr} R_z(k+i0)e^{-ikr}$$

where

$$R_{z}(k+i0) = \lim_{\epsilon \downarrow 0} R_{z}(k+i\epsilon)$$

in the operator-norm topology of  $\text{B}(\text{L}^2_{\ \delta},\text{H}^2_{-\delta})$  , locally uniformly in k .

For  $f \in L^2_{\delta,-b}$ ,  $u = e^{-ikr}R_z^-(k)e^{ikr}f$  is the unique solution in  $L^2_{\delta,-b}$  of the equation  $(H_z^{-b}-k^2)u=f$ , such that  $\mathcal{D}u \in L^2_{\delta-1,-b}$ , where b=Imk and

$$vu = (v_1 u, \dots, v_n u)$$
,  $v_j = \frac{\partial}{\partial x_j} + \frac{n-1}{2r^2} x_j - ik \frac{x_j}{r}$ 

(the radiation condition).

<u>Proof.</u> We refer to [5] for the proof of the Theorem. It utilizes the result of [7] for  ${\rm H}_0$ , analytic Fredholm theory and control of the singular points using analyticity in  ${\rm k}$  and  ${\rm z}$ .

The trace operators  $T_0(k)$  ,  $T_z(k) \in B(L^2_{\ \delta},h)$  are defined for  $z \in S_{\alpha}$  , by

$$(T_0(\pm k)f)(k,\cdot) = (F_{\pm}f)(k,\cdot), k \in \mathbb{R}^+$$

where

$$(F_{\pm}f)(k,\omega) = (2\pi)^{-\frac{n}{2}} \int_{\mathbb{IR}^{n}} e^{\mp ik\omega \cdot x} f(x) dx ,$$

$$T_{z}(k) = T_{0}(k) (1 - Q_{z}R_{z}(k+i0)) , k \in \mathbb{IR} \setminus \sum_{z}^{r} .$$

We set

$$T_0(k) = T_0(k)e^{ikr}$$
 ,  $T_z^+(k) = T_z(k)e^{ikr}$  .

The following result is proved in [5].

We recall the following formulas from the stationary scattering theory:

$$R(k+i0) = R(-k+i0) + \pi i k^{n-2} T^*(k) T(k) , k \in \mathbb{R}^+ \setminus \sum_{r}$$
 (1.2)

$$T(k) = S(k)RT(-k)$$
 (1.3)

where  $(R\tau)(\omega) = \tau(\omega)$  for  $\tau \in h$ .

Inserting (1.3) in (1.2), we obtain

$$R(k+i0) = R(-k+i0) + \pi i k^{n-2} T^*(k) S(k) RT(-k)$$
 (1.4)

The S-matrix S(k) of (H $_0$ ,H) is given for  $k \in \mathbb{R}^+ \setminus \sum_r$  by

$$S(k) = 1 - \pi i k^{n-2} T_0(k) (Q - QR(k + i0)Q) T_0^*(k)$$
 (1.5)

and the S-matrix  $S_1(k)$  of  $(H_0,H_1)$  by (1.5) with Q and R replaced by V and  $R_1$ .

The following result is proved in [3].

Theorem 1.4. The S-matrix S(k) has a meromorphic extension  $\widetilde{S}(k)$  from  $\mathrm{IR}^+$  to  $\mathrm{S}_\alpha$  with poles at R. The S-matrix  $\mathrm{S}_1(k)$  has an analytic extension  $\widetilde{\mathrm{S}}_1(k)$  from  $\mathrm{IR}^+$  to  $\mathrm{S}_\alpha$ . Moreover, for k>0,  $0<\varphi<\alpha$ ,  $\widetilde{\mathrm{S}}_1(k\mathrm{e}^{-i\varphi})=\mathrm{S}_{1,\mathrm{e}^{i\varphi}}(k)$ , where  $\mathrm{S}_{1,\mathrm{e}^{i\varphi}}(k)$  is the S-matrix of  $(\mathrm{H}_0,\mathrm{H}_1,\mathrm{e}^{i\varphi})$  at the point k.

From (1.4) and Theorem 1.2 we obtain

Theorem 1.5. For any b >0 the  $\mathcal{B}(L_{0,b}^2, H_{0,-b}^2)$ -valued function  $\mathcal{R}(k)$  has a meromorphic continuation  $\widetilde{\mathcal{R}}(k)$  from  $\mathfrak{C}^+$  across  $\mathbb{R}^+$  to  $S_{\alpha,b} = \{k \in S_{\alpha} \mid -b < \text{Im} k < 0\}$ , given by

$$\widetilde{R}(k) = R(-k) + \pi i k^{n-2} T^*(\overline{k}) \widetilde{S}(k) T(-k)$$
(1.6)

The  $\mathcal{B}(L_0^2, h^2, h_0^2, -b)$ -valued function  $R_1(k)$  has an analytic continuation  $\widetilde{R}_1(k)$  from  $\mathfrak{C}^+ \setminus \Sigma_{1c}$  across  $\mathbb{R}^+$  to  $S_{\alpha,b}$ , given by (1.6) with R,T and S replaced by  $R_1,T_1$  and  $S_1$ .

The functions  $\widetilde{R}(k)$  and  $\widetilde{R}_1(k)$  are connected by the analytically continued symmetrized resolvent equation

$$\widetilde{R}(k) = \widetilde{R}_1(k) - \widetilde{R}_1(k)g(1 + W\widetilde{R}_1(k)g)^{-1} W\widetilde{R}_1(k)$$
(1.7)

The following result is proved in [5]:

Theorem 1.6.  $\widetilde{R}(k)$  and  $\widetilde{S}(k)$  have the same poles and of the same order in  $S_{\alpha\,,b}$  .

#### 2. Resonance functions

Let  $k_0^2$  be a resonance, and fix  $b > -\mathrm{Im}\,k_0$ . Then  $k_0$  is a pole of  $\widetilde{R}(k) \in \mathcal{B}(L_{0,b}^2, H_{0,-b}^2)$ , defined in Theorem 1.5. Let C be a circle separating  $k_0$  from other poles and let

$$P = -\frac{1}{2\pi i} \int_C \widetilde{R}_2(k) dk^2$$

be the residue of  $\tilde{R}_2(k)$  at  $k_0$ ,  $P \in \mathcal{B}(L_{0,b}^2, H_{0,-b}^2)$  is of finite rank.

The space  ${\it F}$  of resonance functions associated with  ${\it k}_0$  is defined by

$$F = \{f \in R_p \mid (-\Delta + Q - k_0^2) f = 0\}$$
.

The following result is proved in [5]:

Theorem 2.1. F is the isomorphic image of  $N(\widetilde{S}^{-1}(k_0))$  and of  $N(1 + W\widetilde{R}_1(k_0)g)$  via the following maps:

$$N(\widetilde{\mathbf{S}}^{-1}(\mathbf{k}_0)) \ni \tau \to \mathbf{T}^*(\overline{\mathbf{k}}_0)\tau = \mathbf{f} \in F$$

$$N(1 + W\widetilde{\mathbf{R}}_1(\mathbf{k}_0)g) \ni \phi \to \widetilde{\mathbf{R}}_1(\mathbf{k}_0)g\phi = \mathbf{f} \in F$$

Remark. From the representation  $f = T^*(\overline{k}_0)\tau$  we conclude by Theorem 1.3 and the uniqueness part of Theorem 1.2 that  $f \in H^2_{-\delta}$ ,  $-b_0 \subset L^2_{\delta-1}$ ,  $-b_0$  for every  $\delta > \frac{1}{2}$  and  $b_0 = -\mathrm{Im}\,k_0$ . A further analysis yields precise asymptotic estimates. We first establish the analyticity properties, using the second isomorphism.

Applying (1.4) to the operator  $H_{1z}$  at a point  $zk_0$  with Arg  $zk_0$  = 0 and noting that by Theorem 1.4,  $S_{1z}(zk_0) = \widetilde{S}_1(k_0)$  we obtain

$$R_{1z}(zk_0 + i0) = R_{1z}(-zk_0 + i0) + \pi i(zk_0)^{n-2}$$

$$T^*_{1z}(\overline{zk_0}) \widetilde{S}_1(k_0) RT_{1z}(-zk_0)$$
(1.7)

By Theorems 1.2 and 1.3 we obtain from (1.7)

Theorem 2.2. The  $B(L_{\delta}^2, H_{-\delta}^2)$ -valued function  $e^{-izk_0r}$   $R_{1z}(zk_0+i0)e^{-izk_0r}$  has an analytic extension from  $\{z\in zk_0\mid IR^+\}$  to  $\{z\in S_{\alpha}\mid Arg\,zk_0<0\}$ , given by

Recalling that  $W_z = Q_z g(rz)$ , where  $g(rz) = \exp\{-\epsilon (rz)^\beta\}$  with  $\beta > 1$ , we obtain from Theorem 2.2

Theorem 2.3. The  $C(L^2)$ -valued function  $W_z R_{1z}(zk_0)g(rz)$  has an analytic continuation from  $\{z \in S_\alpha \mid Argzk_0 > 0\}$  to  $\{z \in S_\alpha \mid Argzk_0 \leq 0\}$ , given by  $W_z \widetilde{R}_{1z}(zk_0)g(rz)$ .

By standard dilation-analytic arguments  $\sigma(W_z\tilde{R}_1(zk_0)g(rz))$  is constant. Let C be a circle separating -1 from the rest of  $\sigma(W_z\tilde{R}_1(zk_0)g(rz))$  and set

$$P(z) = -\frac{1}{2\pi i} \int_{C} (-\lambda + W_z \widetilde{R}_{1z}(zk_0) g(rz))^{-1} d\lambda .$$

Then P(z) is a dilation-analytic  $\mathcal{B}(L^2)$ -valued function of z , and P(z) is a projection on the finite-dimensional algebraic null space of  $1+W_z\widetilde{R}_{1z}(zk_0)g(rz)$ . Let  $\phi\in N(1+W\widetilde{R}_1(k_0)g(rz))$  and pick an  $S_\alpha$ -dilation-analytic vector  $\eta$  in  $L^2$  such that

 $\varphi=P(1)\eta$  . Then  $\varphi(z)=P(z)\eta(z)\in N(1+W_z^{\widetilde{R}}_{1z}(zk_0)g(rz)$  , and  $\varphi(z)$  is dilation-analytic.

We now obtain, using the second isomorphism of Theorem 2.1,

Theorem 2.4. Let  $f\in F$ . Then there exists an  $S_\alpha$ -dilation-analytic,  $H^2_{-\delta}$ -valued function  $\chi(z)$ , such that  $f=e^{ik_0r}\chi(1)$  and for Arg z  $k_0>0$ 

$$f(z) = e^{ik_0 zr} \chi(z) \in N(H(z) - k_0^2)$$
.

Moreover,  $\chi(z) \notin L^2_{\delta-1}$  for all  $z \in S_{\alpha}$ ,  $\delta > \frac{1}{2}$ .

Proof. Define f(z) by

$$f(z) = \begin{cases} izk_0^r + -izk_0^r \\ e^{-izk_0^r} (zk_0) & e^{-izk_0^r} \\ izk_0^r + (zk_0)^r -izk_0^r \\ e^{-izk_0^r} (zk_0)^r + e^{-izk_0^r} (zk_0)^r e^{$$

where  $R_{1z}^+(zk_0)$  is defined similarly to  $R_{1z}^-(zk_0)$ , replacing -b by b and  $e^{\pm iar}$  with  $e^{\mp iar}$  in Theorem 1.2. Clearly, f(z) is continuous for  $zk_0\in \mathbb{R}^+$ . By Theorem 1.2 and 2.2,  $\chi(z)=e^{-izk_0r}f(z)$  is an analytic  $H_{-\delta}^2$ -valued function in  $S_\alpha$ .

It follows from the uniqueness part of Theorem 1.2 that  $\chi(z) \not\in L^2_{\delta-1} \quad \text{for } \operatorname{Im} zk_0 < 0 \text{ . The fact that } \chi(z) \not\in L^2_{\delta-1} \quad \text{for } \operatorname{Im} zk_0 \geqq 0$  then follows by the next Lemma, proved in [6]:

Lemma 2.5. Let  $\chi(z)$  be an  $S_{\alpha}$ -dilation-analytic vector, and define  $h(\phi)$  for  $\phi \in (-\alpha, \alpha)$  by

$$h(\varphi) = \inf\{s \mid \chi(e^{i\varphi}) \in L_{-s}^2\}$$
.

Then either  $h(\varphi) \equiv -\infty$  or  $h(\varphi) > -\infty$  and h is convex in  $(-\alpha, \alpha)$  .

Using this Lemma together with a recent result of Agmon [1], giving the precise asymptotic behaviour of f(z) for  $\text{Arg}\,zk_0 > 0$ , we finally obtain the desired asymptotic estimates of f and f'. We refer to [6] for the proof.

Theorem 2.6. Assume that Q is an  $S_{\alpha}$ -dilation-analytic multiplicative potential such that  $|Q(z)(x)| \le C|x|^{-1-\epsilon}$  for  $z \in S_{\alpha}$  and  $|x| \ge R$ . Let  $f \in F$ . Then f is an analytic, h-valued function  $f(z,\cdot)$  on  $S_{\alpha}$  of the form

$$f(z,\cdot) = e^{ik_0 z} \frac{1-n}{2} g(z,\cdot)$$

where

$$g(z,\cdot) = \tau + 0(|z|^{-\epsilon})$$

$$g'(z,\cdot) = 0(|z|^{-1-\epsilon})$$

uniformly in any smaller angle  $S_\alpha^*$  for some  $\epsilon>0$  . Moreover,  $\tau\in N(\widetilde{S}^{-1}(k_0)) \text{ and } f=CT^*(\overline{k}_0)\tau\ ,$ 

$$C = k_0^{\frac{n-1}{2}} (-i)^{\frac{1-n}{2}} (2\pi)^{\frac{1}{2}}$$
.

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