## Annales de l'I. H. P., Section B

Robert J. Elliott<br>Allanus H. Tsoi<br>Time reversal of non-Markov point processes

Annales de l'I. H. P., section B, tome 26, no 2 (1990), p. 357-373
[http://www.numdam.org/item?id=AIHPB_1990__26_2_357_0](http://www.numdam.org/item?id=AIHPB_1990__26_2_357_0)
© Gauthier-Villars, 1990, tous droits réservés.
L'accès aux archives de la revue «Annales de l'I. H. P., section B » (http://www.elsevier.com/locate/anihpb) implique l'accord avec les conditions générales d'utilisation (http://www.numdam.org/conditions). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

## Numdam

# Time reversal of non-Markov point processes 

Robert J. ELLIOTT ( ${ }^{1}$ ) and Allanus H. TSOI ( ${ }^{( }$)<br>Department of Statistics and Applied Probability, University of Alberta, Edmonton, Alberta, Canada T6G 2G1

Abstract. - Time reversal is considered for a standard Poisson process, a point process with Markov intensity and a point process with a predictable intensity. In the latter case an analog of the Fréchet derivative for functionals of a Poisson process is introduced and used in techniques of integration-by-parts to obtain formulate similar to those of Föllmer in the Wiener space situation.

Key words : Point processes, Poisson process, predictable intensity, non-Markov, integra-tion-by-parts, Fréchet derivative.

Résumé. - Le retournement du temps est considéré pour un processus de Poisson, un processus ponctuel avec intensité markovienne et un processus ponctuel avec intensité prévisible. Pour le dernier cas, nous introduisons une sorte de dérivée Fréchet pour les fonctionnels d'un processus de Poisson et l'utilisons dans les méthodes d'intégration par parties pour obtenir des formules qui sont similaires à celles de Föllmer pour la situation brownienne.

[^0]
## 1. INTRODUCTION

The time reversal of stochastic processes has been investigated for some years. One motivation comes from quantum theory, and this is discussed in the book of Nelson [11]. The time reversal of Markov diffusions is treated in, for example, the papers of Elliott and Anderson [4], and Haussman and Pardoux [8]. However, the first discussion of time reversal for a non-Markov process on Wiener space appears in the paper by Föllmer [7], in which he uses an integration-by-parts formula related to the Malliavin calculus.

In the present paper an analog of the Fréchet derivative is introduced for functionals of a Poisson process. The integration-by-parts formula on Poisson space, see [6], is formulated in terms of this derivative and counterparts of Föllmer's formulae are obtained.

In Section 2 the time reversed form of the standard Poisson process is derived. Section 3 considers a point (counting) process N with Markov intensity $h\left(\mathrm{~N}_{t}\right)$, so that $\mathrm{Q}_{t}=\mathrm{N}_{\mathrm{t}}-\int_{0}^{t} h\left(\mathrm{~N}_{s}\right) d s$ is a martingale, and obtains the reverse time decomposition of Q for $t \in(0,1]$. Finally, in Section 4, the situation when $h$ is predictable is considered using the "Fréchet" derivative and integration-by-parts techniques mentioned above.

## 2. TIME REVERSAL UNDER THE ORIGINAL MEASURE

Consider a standard Poisson process $\mathrm{N}=\left\{\mathrm{N}_{t}: 0 \leqq t \leqq 1\right\}$ on $(\Omega, \mathscr{F}, \mathrm{P})$. We take $\mathrm{N}_{0}=0$. Let $\left\{\mathscr{F}_{t}\right\}$ be the right-continuous, complete filtration generated by N . Let $\mathrm{G}_{t}^{0}=\sigma\left\{\mathrm{N}_{s}: t \leqq s \leqq 1\right\}$ and $\left\{\mathrm{G}_{t}\right\}$ be the left-continuous completion of $\left\{\mathrm{G}_{t}^{0}\right\}$.

The following result is well known; see, for example, Theorem 2.6 in [9]. For completeness we sketch the proof.

Theorem 2.1. - Under $\mathrm{P}, \mathrm{N}$ is a reverse time $\mathrm{G}_{t}$-quasimartingale, and it has the decomposition:

$$
\mathrm{N}_{t}=\mathrm{N}_{1}+\mathrm{M}_{t}-\int_{t}^{1} \frac{\mathrm{~N}_{s}}{s} d s
$$

where M is a reverse time $\mathrm{G}_{t}$-martingale.
Proof. - Since N is Markov, we have, for $\varepsilon>0$,

$$
\begin{align*}
\mathrm{E}\left[\mathrm{~N}_{t-\varepsilon}-\mathrm{N}_{t} \mid \mathrm{G}_{t}\right] & =\mathrm{E}\left[\mathrm{~N}_{t-\varepsilon}-\mathrm{N}_{t} \mid \mathrm{N}_{t}\right] \\
& =-\frac{\varepsilon}{t} \mathrm{~N}_{t} \tag{2.1}
\end{align*}
$$

(see [5] and [10]). Thus

$$
\int_{0}^{t} \mathrm{E}\left|\mathrm{E}\left[\mathrm{~N}_{s-\varepsilon}-\mathrm{N}_{s} \mid \mathrm{G}_{s}\right]\right| d s=O(\varepsilon)
$$

By Stricker's theorem [12], $\mathrm{N}_{t}$ is a reverse time $\mathrm{G}_{t}$-quasimartingale. Considering approximate Laplacians we see it has the decomposition

$$
\begin{equation*}
\mathrm{N}_{t}=\mathrm{N}_{1}+\mathrm{M}_{t}+\int_{t}^{1} \alpha_{s} d s \tag{2.2}
\end{equation*}
$$

where from (2.1) and (2.2), for almost all $t$

$$
\begin{aligned}
\alpha_{t} & =\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \int_{t-\varepsilon}^{t} \mathrm{E}\left[\alpha_{s} \mid \mathrm{G}_{t}\right] d s \\
& =\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}\left[\mathrm{~N}_{t-\varepsilon}-\mathrm{N}_{t} \mid \mathrm{G}_{t}\right] \\
& =-\frac{N_{t}}{t} .
\end{aligned}
$$

## 3. TIME REVERSAL AFTER A CHANGE OF MEASURE: THE MARKOV CASE

Consider a process $h_{t}=h\left(\mathrm{~N}_{t}\right)$ which satisfies: There exist positive constants A, $\mathrm{K}>0$ such that $0<\mathrm{A}<h\left(\mathrm{~N}_{t}\right) \leqq \mathrm{K}$ for all $t$, a.s.

Define the family $\left\{\Lambda_{t}, 0 \leqq t \leqq 1\right\}$ of exponentials:

$$
\Lambda_{t}=\prod_{0 \leqq u \leqq t}\left(1+\left(h\left(\mathrm{~N}_{u-}\right)-1\right) \Delta \mathrm{N}_{u}\right) \exp \left(\int_{0}^{t}\left(1-h\left(\mathrm{~N}_{u-}\right)\right) d u\right)
$$

Then $\Lambda$ is an $\left(\mathscr{F}_{t}\right)$-martingale under $\mathbf{P}$, and is the unique solution of the equation

$$
\Lambda_{t}=1+\int_{0}^{t} \Lambda_{u-}\left(h\left(\mathrm{~N}_{u-}\right)-1\right)\left(d \mathrm{~N}_{u}-d u\right)
$$

Define a new probability measure $\mathrm{P}^{\boldsymbol{h}}$ by

$$
\frac{d \mathrm{P}^{h}}{d \mathrm{P}}=\Lambda_{1}
$$

Then under $\mathrm{P}^{h}$, the process $\mathrm{H}_{t}=\mathrm{N}_{t}-\int_{0}^{t} h\left(\mathrm{~N}_{u_{-}}\right) d u$ is an $\left(\mathscr{F}_{t}\right)$-martingale (see [3]). Let $\beta(t)=\int_{0}^{t} h\left(\mathrm{~N}_{u_{-}}\right) d u$ so that $\beta$ is positive and increasing in $t$ because $h$ is positive. Write

$$
\begin{aligned}
\psi(t) & =\beta^{-1}(t), \\
\mathrm{N}_{t}^{\prime} & =\mathrm{N}_{\psi(t)}, \\
\mathscr{F}_{t}^{\prime} & =\mathscr{F}_{\psi(t)} .
\end{aligned}
$$

Lemma 3.1. - $\left(\mathrm{N}_{t}^{\prime}\right)$ is a Poisson process under $\left(\Omega, \mathscr{F},\left(\mathscr{F}_{t}^{\prime}\right), \mathrm{P}^{h}\right)$.
Proof. - Since $\mathrm{H}_{t}=\mathrm{N}_{t}-\beta(t)$ is an $\left(\mathscr{F}_{t}\right)$-martingale under $\mathrm{P}^{h}$, $\mathrm{H}_{t}^{\prime}=\mathrm{H}_{\psi(t)}=\mathrm{N}_{\psi(t)}-t$ is an $\left(\mathscr{F}_{t}^{\prime}\right)$-martingale under $\mathrm{P}^{h}$. By Itô's rule,

$$
\begin{aligned}
\mathrm{H}^{\prime 2} & =2 \int_{0}^{t} \mathrm{H}_{s-}^{\prime} d \mathrm{H}_{s}^{\prime}+\sum_{s \leqq t}\left(\Delta \mathrm{~N}_{\psi(s)}\right)^{2} \\
& =2 \int_{0}^{t} \mathrm{H}_{s-}^{\prime} d \mathrm{H}_{s}^{\prime}+\mathrm{N}_{\psi(t)} .
\end{aligned}
$$

Hence $\mathrm{H}_{\psi(t)}^{2}-t$ is also an $\left(\mathscr{F}_{t}^{\prime}\right)$-martingale under $\mathrm{P}^{h}$. Therefore, $\left\{\mathrm{N}_{t}^{\prime}\right\}$ is Poisson by Lévy's characterization (Theorem 12.31 in [2]).

Lemma 3.2. - N is Markov under $\mathrm{P}^{h}$.
Proof. - Consider any $\varphi \in \mathrm{C}_{0}^{\infty}(\mathbb{R})$. For $t \geqq s$, by Bayes' formula,

$$
\begin{aligned}
\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right) \mid \mathscr{F}_{s}\right] & =\frac{\mathrm{E}\left[\Lambda_{t} \varphi\left(\mathrm{~N}_{t}\right) \mid \mathscr{F}_{s}\right]}{\mathrm{E}\left[\Lambda_{t} \mid \mathscr{F}_{\mathrm{s}}\right]} \\
& =\mathrm{E}\left[\Lambda_{s}^{t} \varphi\left(\mathrm{~N}_{t}\right) \mid \mathscr{F}_{s}\right] \\
& =\mathrm{E}\left[\Lambda_{s}^{t} \varphi\left(\mathrm{~N}_{t}\right) \mid \mathrm{N}_{s}\right],
\end{aligned}
$$

because N is Markov under P , where

$$
\Lambda_{s}^{t}=\prod_{s<u \leqq t}\left(1+\left(h\left(\mathrm{~N}_{u}\right)-1\right) \Delta \mathrm{N}_{u}\right) \exp \left(\int_{s}^{t}\left(1-h\left(\mathrm{~N}_{u}\right)\right) d u\right)
$$

Hence

$$
\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right) \mid \mathscr{F}_{s}\right]=\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right) \mid \mathrm{N}_{s}\right]
$$

and N is Markov under $\mathrm{P}^{h}$.
Note that

$$
\begin{equation*}
\mathrm{H}_{t}=\mathrm{H}_{1}+\mathrm{N}_{t}-\mathrm{N}_{1}+\int_{t}^{1} h\left(\mathrm{~N}_{s}\right) d s . \tag{3.1}
\end{equation*}
$$

Thus $H_{t}$ is a reverse time $\mathrm{G}_{t}$-quasimartingale under $\mathrm{P}^{h}$ if and only if $\mathrm{N}_{t}$ is. To determine the reverse time decomposition we again investigate the approximate Laplacians, as in [4].

Theorem 3.3.

$$
\begin{equation*}
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\mathrm{~N}_{t-\varepsilon}-\mathrm{N}_{t} \mid \mathrm{G}_{t}\right]=-\mathrm{E}^{h}\left[\left.h\left(\mathrm{~N}_{t}-1\right) \frac{\mathrm{N}_{t}}{\int_{0}^{t} h\left(\mathrm{~N}_{u}\right) d u} \right\rvert\, \mathrm{N}_{t}\right] \tag{3.2}
\end{equation*}
$$

Proof. - By Lemma 3.2,

$$
\mathrm{E}^{h}\left[\mathrm{~N}_{t}-\mathrm{N}_{t-\varepsilon} \mid \mathrm{G}_{t}\right]=\mathrm{E}^{h}\left[\mathrm{~N}_{t}-\mathrm{N}_{t-\varepsilon} \mid \mathrm{N}_{t}\right] .
$$

Consider a bounded, differentiable function $\varphi$ on $\mathbb{R}$ and its restriction to Z (the range of N ). Now

$$
\varphi\left(\mathrm{N}_{t}\right)=\varphi\left(\mathrm{N}_{t-\varepsilon}\right)+\int_{t-\varepsilon}^{t}\left(\varphi\left(\mathrm{~N}_{s-}+1\right)-\varphi\left(\mathrm{N}_{s-}\right)\right) d \mathrm{~N}_{s}
$$

So

$$
\begin{aligned}
\varphi\left(\mathrm{N}_{t}\right)\left(\mathrm{N}_{t}-\mathrm{N}_{t-\varepsilon}\right) & =\int_{t-\varepsilon}^{t}\left(\mathrm{~N}_{s-}-\mathrm{N}_{t-\varepsilon}\right)\left(\varphi\left(\mathrm{N}_{s-}+1\right)-\varphi\left(\mathrm{N}_{s-}\right)\right) d \mathrm{~N}_{s} \\
& +\int_{t-\varepsilon}^{t} \varphi\left(\mathrm{~N}_{s-}\right) d \mathrm{~N}_{s}+\sum_{t-\varepsilon<s \leqq t} \Delta \varphi\left(\mathrm{~N}_{s}\right) \Delta \mathrm{N}_{s} \\
= & \int_{t-\varepsilon}^{t}\left(\mathrm{~N}_{s-}-\mathrm{N}_{t-\varepsilon}\right)\left(\varphi\left(\mathrm{N}_{s-}+1\right)-\varphi\left(\mathrm{N}_{s-}\right)\right) d \mathrm{~N}_{s} \\
& \quad+\int_{t-\varepsilon}^{t} \varphi\left(\mathrm{~N}_{s-}+1\right) d \mathrm{~N}_{s}
\end{aligned}
$$

Since

$$
\begin{aligned}
\mathrm{H}_{t} & =\mathrm{N}_{t}-\int_{0}^{t} h\left(\mathrm{~N}_{s}\right) d s \\
& =\mathrm{N}_{t}-\int_{0}^{t} h\left(\mathrm{~N}_{s_{-}}\right) d s
\end{aligned}
$$

is a martingale under $\mathrm{P}^{h}$,

$$
\begin{align*}
& \mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right)\left(\mathrm{N}_{t}-\mathrm{N}_{t-\varepsilon}\right)\right] \\
&=\mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t}\left(\mathrm{~N}_{s-}-\mathrm{N}_{t-\varepsilon}\right)( \right.\left.\left.\left(\mathrm{N}_{s-}+1\right)-\varphi\left(\mathrm{N}_{s-}\right)\right) h\left(\mathrm{~N}_{s-}\right) d s\right] \\
&+\mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t} \varphi\left(\mathrm{~N}_{s-}+1\right) h\left(\mathrm{~N}_{s-}\right) d s\right] . \tag{3.3}
\end{align*}
$$

Now, if $|\varphi| \leqq C$,

$$
\begin{aligned}
& \left|\mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t}\left(\mathrm{~N}_{s-}-\mathrm{N}_{t-\varepsilon}\right)\left(\varphi\left(\mathrm{N}_{s-}+1\right)-\varphi\left(\mathrm{N}_{s-}\right)\right) h\left(\mathrm{~N}_{s-}\right) d s\right]\right| \\
& \leqq 2 \mathrm{KC} \int_{t-\varepsilon}^{t} \mathrm{E}^{h}\left[\left|\mathrm{~N}_{s-}-\mathrm{N}_{t-\varepsilon}\right|\right] d s \\
& \leqq 2 \mathrm{KC} \int_{t-\varepsilon}^{t} \mathrm{E}^{h}\left[\left|\mathrm{~N}_{s-}-\mathrm{N}_{t-\varepsilon}-\int_{t-\varepsilon}^{s-} h\left(\mathrm{~N}_{u-}\right) d u\right|\right] \\
& +\mathrm{E}^{h}\left[\left|\int_{t-\varepsilon}^{s-} h\left(\mathrm{~N}_{u-}\right) d u\right|\right] d s \\
& \leqq 2 \mathrm{KC} \int_{t-\varepsilon}^{t}\left\{\left[\mathrm{E}^{h}\left|\mathrm{~N}_{s-}-\mathrm{N}_{t-\varepsilon}-\int_{t-\varepsilon}^{s-} h\left(\mathrm{~N}_{u-}\right) d u\right|^{2}\right]^{1 / 2}+\mathrm{K} \varepsilon\right\} d s \\
& \leqq 2 \mathrm{KC} \int_{t-\varepsilon}^{t}\left\{\mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t} h\left(\mathrm{~N}_{u-}\right) d u\right]^{1 / 2}+\mathrm{K} \varepsilon\right\} d s \\
& \leqq 2 \mathrm{KC} \int_{t-\varepsilon}^{t}\left((\mathrm{~K} \varepsilon)^{1 / 2}+\mathrm{K} \varepsilon\right) d s \leqq \mathrm{~K}^{\prime} \varepsilon^{3 / 2}+\mathrm{K}^{\prime \prime} \varepsilon^{2} .
\end{aligned}
$$

Thus from (3.3),

$$
\begin{align*}
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right)\left(\mathrm{N}_{t}-\mathrm{N}_{t-\varepsilon}\right)\right]=\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t-}+1\right)\right. & \left.h\left(\mathrm{~N}_{t-}\right)\right] \\
& =\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}+1\right) h\left(\mathrm{~N}_{t}\right)\right] \tag{3.4}
\end{align*}
$$

However,

$$
\begin{aligned}
& \mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}+1\right) h\left(\mathrm{~N}_{t}\right)\right]=\mathrm{E}^{h} {\left[\varphi\left(\mathrm{~N}_{\psi(\beta(t))}+1\right) h\left(\mathrm{~N}_{\psi(\beta(t))}\right)\right] } \\
&=\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{\beta(t)}^{\prime}+1\right) h\left(\mathrm{~N}_{\beta(t)}^{\prime}\right)\right] \\
&=\mathrm{E}^{h}\left[\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{\beta(t)}^{\prime}+1\right) h\left(\mathrm{~N}_{\beta(t)}^{\prime}\right) \mid \beta(t)\right]\right]
\end{aligned}
$$

And

$$
\begin{aligned}
& \mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{\beta(t)}^{\prime}+1\right) h\left(\mathrm{~N}_{\beta(t)}^{\prime}\right) \mid \beta(t)\right] \\
&=\sum_{k=0}^{\infty} \varphi(k+1) h(k) \frac{\beta(t)^{k} e^{-\beta(t)}}{k!} \\
&= \sum_{t=0}^{\infty} \varphi(l) h(l-1) \frac{\beta(t)^{l} e^{-\beta(t)}}{l!} \frac{l}{\beta(t)} \\
&= \mathrm{E}^{h}\left[\left.\varphi\left(\mathrm{~N}_{\beta(t)}^{\prime}\right) h\left(\mathrm{~N}_{\beta(t)}^{\prime}-1\right) \frac{\mathrm{N}_{\beta(t)}^{\prime}}{\beta(t)} \right\rvert\, \beta(t)\right] \\
&=\mathrm{E}^{h}\left[\left.\varphi\left(\mathrm{~N}_{t}\right) h\left(\mathrm{~N}_{t}-1\right) \frac{\mathrm{N}_{t}}{\beta(t)} \right\rvert\, \beta(t)\right] .
\end{aligned}
$$

Hence,

$$
\begin{equation*}
\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}+1\right) h\left(\mathrm{~N}_{t}\right)\right]=\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right) h\left(\mathrm{~N}_{t}-1\right) \frac{\mathrm{N}_{t}}{\int_{0}^{t} h\left(\mathrm{~N}_{u}\right) d u}\right] \tag{3.5}
\end{equation*}
$$

Thus from (3.4) and (3.5),

$$
\lim _{\varepsilon \downarrow 0} \mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right) \frac{\left(\mathrm{N}_{t}-\mathrm{N}_{t-\varepsilon}\right)}{\varepsilon}\right]=\mathrm{E}^{h}\left[\varphi\left(\mathrm{~N}_{t}\right) h\left(\mathrm{~N}_{t}-1\right) \frac{\mathrm{N}_{t}}{\int_{0}^{t} h\left(\mathrm{~N}_{u}\right) d u}\right]
$$

or

$$
\lim _{\varepsilon \downarrow 0} \mathrm{E}^{h}\left[\left.\frac{\mathrm{~N}_{t-\varepsilon}-\mathrm{N}_{t}}{\varepsilon} \right\rvert\, \mathrm{G}_{t}\right]=-\mathrm{E}^{h}\left[\left.h\left(\mathrm{~N}_{t}-1\right) \frac{\mathrm{N}_{t}}{\int_{0}^{t} h\left(\mathrm{~N}_{u}\right) d u} \right\rvert\, \mathrm{N}_{t}\right] .
$$

By Theorem 3.3 and an argument similar to that in [4], we see that N , and hence H , is a reverse time $\mathrm{G}_{t}$-quasimartingale under $\mathrm{P}^{h}$, and it has the decomposition

$$
\begin{equation*}
\mathrm{H}_{t}=\mathrm{H}_{1}+\mathrm{M}_{t}+\int_{t}^{1} \alpha_{t} d_{t} . \tag{3.6}
\end{equation*}
$$

Moreover, we have the following expression for $\alpha_{t}$ :
Theorem 3.4. - The integrand $\alpha_{t}$ that appears in (3.6) is given by

$$
\alpha_{t}=h\left(\mathrm{~N}_{t}\right)-\mathrm{E}^{h}\left[\left.h\left(\mathrm{~N}_{t}-1\right) \frac{\mathrm{N}_{t}}{\int_{0}^{t} h\left(\mathrm{~N}_{u}\right) d u} \right\rvert\, \mathrm{N}_{t}\right]
$$

Proof. - From (3.1) and (3.6),

$$
\begin{aligned}
& \mathrm{E}^{h}\left[\mathrm{H}_{t-\varepsilon}-\mathrm{H}_{t} \mid \mathrm{G}_{t}\right]=\mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t} \alpha_{s} d s \mid \mathrm{G}_{t}\right] \\
&=\mathrm{E}^{h}\left[\mathrm{~N}_{t-\varepsilon}-\mathrm{N}_{t} \mid \mathrm{G}_{t}\right]+\mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t} h\left(\mathrm{~N}_{s}\right) d s \mid \mathrm{G}_{t}\right]
\end{aligned}
$$

Thus for almost all $t$

$$
\alpha_{t}=\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t} \alpha_{s} d s \mid \mathrm{G}_{t}\right]=\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\mathrm{~N}_{t-\varepsilon}-\mathrm{N}_{t} \mid \mathrm{G}_{t}\right]+h\left(\mathrm{~N}_{t}\right) .
$$

From Theorem 3.3, $\alpha_{t}$ has the stated form.

## 4. TIME REVERSAL AFTER A CHANGE OF MEASURE: THE NON-MARKOV CASE

This section involves an integration by parts for Poisson processes which is effected by using a Girsanov transformation to change the intensity and then compensating by a time change. In contrast, the integration by parts considered in [1] is obtained by introducing a perturbation of the size of the jumps. The topic is further investigated in [6].

Suppose $\left\{\mathrm{N}_{t}: 0 \leqq t \leqq 1\right\}$ is a Poisson process with jump times $\mathrm{T}_{1} \wedge 1, \ldots, \mathrm{~T}_{n} \wedge 1, \ldots$ Let $\left\{u_{t}\right\}$ be a real predictable process satisfying $\left\{u_{t}\right\}$ is positive and bounded a.s.

For $\varepsilon>0$, consider the family of exponentials:

$$
\Lambda_{t}^{\varepsilon}=\prod_{0 \leqq s \leqq t}\left(1+\varepsilon u_{s} \Delta \mathrm{~N}_{s}\right) \exp \left(-\int_{0}^{t} \varepsilon u_{s} d s\right)
$$

Then $\left\{\Lambda_{t}^{\varepsilon}\right\}$ is an $\left\{\mathscr{F}_{t}\right\}$-martingale with $\mathrm{E}\left[\Lambda_{t}^{\varepsilon}\right]=1$ (see [6]). Define a probability measure $\mathrm{P}^{\varepsilon}$ on $\mathscr{F}_{1}$ by

$$
\frac{d \mathbf{P}^{\varepsilon}}{d \mathrm{P}}=\Lambda_{1}^{\varepsilon}
$$

Set

$$
\varphi_{\varepsilon}(t)=\int_{0}^{t}\left(1+\varepsilon u_{s}\right) d s
$$

and write

$$
\begin{gathered}
\psi_{\varepsilon}(t)=\varphi_{\varepsilon}^{-1}(t)=\int_{0}^{t} \frac{1}{1+\varepsilon u_{\psi_{\varepsilon}(s)}} d s \\
\mathscr{F}_{t}^{\varepsilon}=\mathscr{F}_{\psi_{\varepsilon}(t)}
\end{gathered}
$$

Then the process $\mathrm{N}_{t}^{\varepsilon}=\mathrm{N}_{\Psi_{\varepsilon}(t)}$ is Poisson on $\left(\Omega, \mathscr{F},\left(\mathscr{F}_{t}^{\varepsilon}\right), \mathrm{P}^{\varepsilon}\right)$ with jump times $\varphi_{\varepsilon}\left(\mathrm{T}_{1}\right) \wedge 1, \ldots, \varphi_{\varepsilon}\left(\mathrm{T}_{n}\right) \wedge 1, \ldots$ (see [6]).

For $\left\{u_{t}\right\}$ as above, set $\mathrm{U}_{t}=\int_{0}^{t} u_{s} d s$. Suppose $g_{s}(w)$ is an $\left\{\mathrm{F}_{t}\right\}$-predictable function on $[0,1]$. Then for $0 \leqq s \leqq \mathrm{~T}_{1} \wedge 1$,

$$
g_{s}(w)=g(s)
$$

and in general, for $\mathrm{T}_{n-1} \wedge 1<s \leqq \mathrm{~T}_{n} \wedge 1$,

$$
g_{s}(w)=g\left(s, \mathrm{~T}_{1} \wedge 1, \ldots, \mathrm{~T}_{n-1} \wedge 1\right)
$$

Note that by setting $g_{s}(0,0, \ldots)=g(s)$ for $0 \leqq s \leqq \mathrm{~T}_{1} \wedge 1$, $\left.g_{s}\left(\left(s-\mathrm{T}_{1}\right) \vee 0, \ldots,\left(s-\mathrm{T}_{n-1}\right) \vee 0\right), 0,0, \ldots\right)$ for $\mathrm{T}_{n-1} \wedge 1<s \leqq \mathrm{~T}_{n} \wedge 1$, etc., such a $g$ can be written in the form

$$
\begin{equation*}
g_{s}(w)=g_{s}\left(\left(s-\mathrm{T}_{1}\right) \vee 0,\left(s-\mathrm{T}_{2}\right) \vee 0, \ldots\right), \quad s \in[0,1] \tag{4.1}
\end{equation*}
$$

Therefore, we shall consider a predictable function $g$ of this form, and further assume that if

$$
g=g_{s}\left(t_{1}, t_{2}, \ldots\right)
$$

then all the partial derivatives $\frac{\partial g_{s}}{\partial t_{i}}$ exist for all $s$, and there is a constant $K>0$ such that

$$
\begin{equation*}
\left|\frac{\partial g_{s}}{\partial t_{i}}\right|<\mathrm{K} \quad \text { for all } i, \text { and for all } s \tag{4.2}
\end{equation*}
$$

We now define the analog of the Fréchet derivative for functionals of the Poisson process.

Write

$$
g_{s}^{\varepsilon}=g_{s}\left(\left(s-\varphi_{\varepsilon}\left(\mathrm{T}_{1}\right)\right) \vee 0, \ldots,\left(s-\varphi_{\varepsilon}\left(\mathrm{T}_{n}\right)\right) \vee 0, \ldots\right) .
$$

Then

$$
\begin{align*}
&\left.\frac{\partial g_{s}^{\varepsilon}}{\partial \varepsilon}\right|_{\varepsilon=0}=-\sum_{i=1}^{\infty} \frac{\partial}{\partial t_{i}} g_{s}\left(\left(s-\mathrm{T}_{1}\right) \vee 0, \ldots,\left(s-\mathrm{T}_{n}\right) \vee 0, \ldots\right) \\
& \times \int_{0}^{\mathrm{T}_{i}} u_{r} d r \mathrm{I}_{\mathrm{T}_{i}<s} \tag{4.3}
\end{align*}
$$

Define

$$
\mu(d t)=-\sum_{i=1}^{\infty} \frac{\partial g_{s}}{\partial t_{i}} \mathrm{I}_{\mathrm{T}_{i}<s} \delta_{\mathrm{T}_{i}}(d t)
$$

where $\delta_{\mathrm{T}_{i}}$ is the point mass at $\mathrm{T}_{i}$. Then

$$
\begin{aligned}
\left.\frac{\partial g_{s}^{\varepsilon}}{\partial \varepsilon}\right|_{\varepsilon=0} & =\int_{0}^{s} \int_{0}^{t} u_{r} d r \mu(d t) \\
& =\int_{0}^{s} \int_{0}^{s} \mathrm{I}_{0 \leqq r \leqq t \leqq s} u_{r} d r \mu(d t) \\
& =\int_{0}^{s} \mu([r, s]) u_{r} d r \\
& =-\int_{0}^{s} \sum_{i=1}^{\infty} \mathrm{I}_{r \leqq \mathrm{~T}_{i}<s} \frac{\partial g_{s}}{\partial t_{i}} u_{r} d r \\
& =\int_{0}^{s} \mathrm{D} g_{s}(.,[r, s]) u_{r} d r
\end{aligned}
$$

where

$$
\mathrm{D} g_{s}(.,[r, s])=-\sum_{i=1}^{\infty} \mathrm{I}_{r \leqq \mathrm{~T}_{i}<s} \frac{\partial g_{s}}{\partial t_{i}}
$$

Write

$$
\mathrm{D} g_{s}(., \mathrm{U})=\int_{0}^{s} \mathrm{D} g_{s}(.,[r, s]) u_{r} d r
$$

Note that

$$
\begin{equation*}
\mathrm{D} g_{\mathrm{T}_{i}}(., \mathrm{U})=-\sum_{j=1}^{i-1} \frac{\partial g_{\mathrm{T}_{i}}}{\partial t_{j}} \int_{0}^{\mathrm{T}_{j}} u_{\mathrm{r}} d r . \tag{4.4}
\end{equation*}
$$

Definition 4.1. - A process $\left\{g_{s}\right\}$ of the form (4.1) is said to be differentiable if it satisfies (4.2) and (4.3) for all $u$ satisfying (i) and (ii) above, and for all $s$. We call $\mathrm{D} g_{s}(., \mathrm{U})$ the derivative of $g_{s}$ in the direction U. It is of interest to note that this concept of differentiability of a function of a Poisson process is an analog of the Fréchet derivative of a function of a continuous process. See Föllmer [7], where similar formulae arise using the Fréchet derivative.

Now suppose $\left\{h_{s}\right\}$ is a bounded, $\left\{\mathrm{F}_{t}\right\}$-predictable process of the form given by (4.1), which satisfies:
(a) $h$ is differentiable in the sense of Definition 4.1.
(b) $\frac{\partial h_{s}}{\partial s}$ exists, and there exists a constant $\mathrm{A}>0$ such that $\left|\frac{\partial h_{s}}{\partial s}\right|<\mathrm{A}$ for all $s$, a.s.
(c) There are constants $\mathrm{B}>0, \mathrm{C}>0$ such that $0<\mathrm{B}<h_{s}<\mathrm{C}$ for all $s$, a.s.

It is easy to check that $h_{s}=h_{s}\left(\left(s-\mathrm{T}_{1}\right) \vee 0,\left(s-\mathrm{T}_{2}\right) \vee 0, \ldots\right)$ is predictable. Consider the family of exponentials:

$$
\begin{align*}
\mathrm{G}_{t} & =\prod_{0 \leqq s \leqq t}\left(1+\left(h_{s}-1\right) \Delta \mathrm{N}_{s}\right) \exp \left(\int_{0}^{t}\left(1-h_{s}\right) d s\right) \\
& =\left(\prod_{0 \leqq \mathrm{~T}_{i} \leqq t} h_{\mathrm{T}_{i}}\right) \exp \left(\int_{0}^{t}\left(1-h_{s}\right) d s\right) . \tag{4.5}
\end{align*}
$$

Then $\left\{\mathrm{G}_{t}\right\}$ is a martingale with $\mathrm{E}\left[\mathrm{G}_{t}\right]=1$. Since for each fixed $\omega$, if $\mathrm{T}_{n-1}(\omega)<t \leqq \mathrm{~T}_{n}(\omega), \mathrm{G}_{t}$ is a function of $\left(t, \mathrm{~T}_{1}(\omega), \ldots, \mathrm{T}_{n-1}(\omega)\right)$, we see as above that $\mathrm{G}_{t}$ can be considered to be of the form

$$
\mathrm{G}_{t}=\mathrm{G}_{t}\left(\left(t-\mathrm{T}_{1}\right) \vee 0, \ldots,\left(t-\mathrm{T}_{n}\right) \vee 0, \ldots\right) .
$$

Theorem 4.2. $-\left(\mathrm{G}_{t}\right)$ defined in (4.5) is differentiable in the sense of Definition 4.1.
Moreover,

$$
\begin{align*}
& \mathrm{DG}_{1}(., \mathrm{U}) \mathrm{G}_{1}^{-1}=\int_{0}^{1} \gamma_{s} u_{s} \mathrm{G}_{1}^{-1} d s \\
& =\int_{0}^{1} \int_{r}^{1}\left[\frac{\partial h_{s}}{\partial s}+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<s\right\}} \frac{\partial h_{s}}{\partial t_{j}}+\mathrm{D} h_{s}(.,[r, s])\right] \frac{1}{h_{s}} d \mathrm{~N}_{s} u_{r} d r \\
&  \tag{4.6}\\
& \quad-\int_{0}^{1} \int_{r}^{1} \mathrm{D} h_{s}(.,[r, s]) d s u_{r} d r, \quad \text { a.s. }
\end{align*}
$$

where

$$
\gamma_{s}=-\sum_{i=1}^{\infty} \mathrm{I}_{s \leqq \mathrm{~T}_{i} \leqq 1} \frac{\partial}{\partial t_{i}} \mathrm{G}_{1}\left(\left(1-\mathrm{T}_{1}\right) \vee 0, \ldots,\left(1-\mathrm{T}_{n}\right) \vee 0, \ldots\right) .
$$

Proof. - The first identity follows from the definition and properties of the derivative. To determine $\mathrm{DG}_{t}(., \mathrm{U})$ we calculate the derivative of $\mathrm{G}_{t}^{\varepsilon}$ at $\varepsilon=0$. Write

$$
h_{s}^{\varepsilon}=h_{s}\left(\left(s-\varphi_{\varepsilon}\left(\mathrm{T}_{1}\right)\right) \vee 0, \ldots,\left(\mathrm{~s}-\varphi_{\varepsilon}\left(\mathrm{T}_{n}\right)\right) \vee 0, \ldots\right),
$$

$$
\begin{aligned}
\mathrm{G}_{t}^{\varepsilon} & =\prod_{0 \leqq s \leqq t}\left(1+\left(h_{s}^{\varepsilon}-1\right) \Delta \mathrm{N}_{\Psi_{\varepsilon}(s)}\right) \exp \left(\int_{0}^{t}\left(1-h_{s}^{\varepsilon}\right) d s\right) \\
& =\left(\prod_{0 \leqq \varphi_{\varepsilon}\left(\mathrm{T}_{i}\right) \leqq t} h_{\varphi_{\varepsilon}\left(\mathrm{T}_{i}\right)}^{\varepsilon}\right) \exp \left(\int_{0}^{t}\left(1-h_{s}^{\varepsilon}\right) d s\right) \\
& =\left(\prod_{0 \leqq \mathrm{~T}_{i} \leqq \psi_{\varepsilon}(t)} h_{\varphi_{\varepsilon}\left(\mathrm{T}_{i}\right)}^{\varepsilon}\right) \exp \left(\int_{0}^{t}\left(1-h_{s}^{\varepsilon}\right) d s\right) .
\end{aligned}
$$

Then

$$
\begin{equation*}
\log \mathrm{G}_{t}^{\varepsilon}=\sum_{i=1}^{\infty} \mathrm{I}_{\mathrm{T}_{i} \leq \psi_{\varepsilon}(\mathrm{t})} \log h_{\mathrm{q}_{\mathrm{\varepsilon}}\left(\mathrm{~T}_{i}\right)}^{\varepsilon}+\int_{0}^{t}\left(1-h_{s}^{\varepsilon}\right) d s \tag{4.7}
\end{equation*}
$$

Differentiate (4.7) with respect to $\varepsilon$, and then set $\varepsilon=0$, to see

$$
\begin{aligned}
\mathrm{DG}_{t}(., \mathrm{U}) \frac{1}{\mathrm{G}_{t}}= & \sum_{i=1}^{\infty}\left\{\mathrm { I } _ { \mathrm { T } _ { i } \leqq t } \left[\frac{\partial h_{\mathrm{T}_{i}}}{\partial t} \int_{0}^{\mathrm{T}_{\mathrm{i}}} u_{r} d r\right.\right. \\
& \left.\left.+\sum_{j=1}^{i-1} \frac{\partial h_{\mathrm{T}_{i}}}{\partial t_{j}}\left(\int_{0}^{\mathrm{T}_{i}} u_{r} d r-\int_{0}^{\mathrm{T}_{j}} u_{r} d r\right)\right] \frac{1}{h_{\mathrm{T}_{i}}}\right\} \\
& \quad-\int_{0}^{t} \mathrm{D} h_{s}(., \mathrm{U}) d s, \quad \text { a.s. }
\end{aligned}
$$

From (4.4) this is

$$
\begin{gather*}
=\sum_{i=1}^{\infty}\left\{\mathrm { I } _ { \mathrm { T } _ { i } \leq t } \left[\frac{\partial h_{\mathrm{T}_{i}}}{\partial t} \int_{0}^{\mathrm{T}_{i}} u_{r} d r\right.\right. \\
\left.\left.+\sum_{j=1}^{i-1} \frac{\partial h_{\mathrm{T}_{i}}}{\partial t_{j}} \int_{0}^{\mathrm{T}_{i}} u_{r} d r+\mathrm{D} h_{\mathrm{T}_{i}}(., \mathrm{U})\right] \frac{1}{h_{\mathrm{T}_{i}}}\right\} \\
-\int_{0}^{t} \mathrm{D} h_{s}(., \mathrm{U}) d s=\int_{0}^{t}\left[\frac{\partial h_{s}}{\partial s} \int_{0}^{s} u_{r} d r\right. \\
\left.+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<s\right\}} \frac{\partial h_{s}}{\partial t_{j}} \int_{0}^{s} u_{r} d r+\mathrm{D} h_{s}(., \mathrm{U})\right] \frac{1}{h_{s}} d \mathrm{~N}_{s} \\
-\int_{0}^{t} \mathrm{D} h_{s}(., \mathrm{U}) d s . \tag{4.8}
\end{gather*}
$$

(Formally, the differentiation of the indicator functions $\mathrm{I}_{\mathrm{T}_{i} \leqq \psi_{\varepsilon}(t)}$ introduces Dirac measures $\delta\left(t-\mathrm{T}_{i}\right)$ However, $\mathrm{P}\left(\mathrm{T}_{i}=t\right)=0$ and we later will take
expectations, so these can be ignored.) From (4.8),

$$
\begin{aligned}
& \mathrm{DG}_{1}(., \mathrm{U}) \mathrm{G}_{1}^{-1}=\int_{0}^{1}\left\{\frac{\partial h_{s}}{\partial s} \int_{0}^{s} u_{r} d r+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<s\right\}} \frac{\partial h_{s}}{\partial t_{j}} \int_{0}^{s} u_{r} d r\right. \\
& \left.+\int_{0}^{s} \mathrm{D} h_{s}(.,[r, s]) u_{r} d r\right\} \frac{1}{h_{s}} d \mathrm{~N}_{s}-\int_{0}^{1} \int_{0}^{s} \mathrm{D} h_{s}(.,[r, s]) u_{r} d r d s \\
& =\int_{0}^{1} \int_{0}^{1} \mathrm{I}_{0 \leqq r \leqq s \leqq 1}\left\{\frac{\partial h_{s}}{\partial s} u_{r} \frac{1}{h_{s}}+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<s\right\}} \frac{\partial h_{s}}{\partial t_{j}} u_{r} \frac{1}{h_{s}}\right. \\
& \left.+\mathrm{D} h_{s}(.,[r, s]) u_{r} \frac{1}{h_{s}}\right\} d r d N_{s}-\int_{0}^{1} \int_{0}^{1} \mathrm{I}_{0 \leqq r \leqq s \leqq 1} \mathrm{D} h_{s}(.,[r, s]) u_{r} d r d s \\
& =\int_{0}^{1} \int_{r}^{1}\left[\frac{\partial h_{s}}{\partial s}+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<s\right\}} \frac{\partial h_{s}}{\partial t_{j}}+\mathrm{D} h_{s}(.,[r, s])\right] \frac{1}{h_{s}} d \mathrm{~N}_{s} u_{r} d r \\
& -\int_{0}^{1} \int_{r}^{1} \mathrm{D} h_{s}(.,[r, s]) d s u_{r} d r
\end{aligned}
$$

which is (4.6).
Consider the family of exponentials defined by (4.5) and define a new probability measure $\mathrm{P}^{h}$ on $\mathscr{F}_{1}$ by:

$$
\frac{d \mathrm{P}^{h}}{d \mathrm{P}}=\mathrm{G}_{1}
$$

Then (see [3]) the process

$$
\begin{align*}
\mathrm{Z}_{t} & =\mathrm{N}_{t}-\int_{0}^{t} h_{s} d s  \tag{4.9}\\
& =\mathrm{Q}_{t}-\int_{0}^{t}\left(h_{s}-1\right) d s
\end{align*}
$$

where $\mathrm{Q}_{t}=\mathrm{N}_{t}-t$, is an $\left(\mathscr{F}_{t}\right)$-martingale under $\mathrm{P}^{h}$. We want to show that $\mathrm{Z}_{t}$ is a reverse time $\mathrm{G}_{t}$-quasimartingale under $\mathrm{P}^{h}$, having the decomposition

$$
\begin{equation*}
\mathrm{Z}_{t}=\mathrm{Z}_{1}+\mathrm{M}_{t}+\int_{t}^{1} \alpha_{s} d s \tag{4.10}
\end{equation*}
$$

From (4.9), we can write

$$
\mathrm{Z}_{t}=\mathrm{Z}_{1}+\mathrm{Q}_{t}-\mathrm{Q}_{1}+\int_{t}^{1}\left(h_{s}-1\right) d s
$$

Now for almost all $t$

$$
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\int_{t-\varepsilon}^{t}\left(h_{s}-1\right) d s \mid \mathrm{G}_{t}\right]=\mathrm{E}^{h}\left[h_{t}-1 \mid \mathrm{G}_{t}\right]
$$

Hence, to show that $Z_{t}$ has the decomposition given by (4.10), it again suffices to consider approximate Laplacien as in [4] and show that

$$
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\mathrm{Q}_{t-\varepsilon}-\mathrm{Q}_{t} \mid \mathrm{G}_{t}\right]
$$

exists.
Theorem 4.3. - For almost all $t \in[0,1]$

$$
\begin{equation*}
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\mathrm{Q}_{t}-\mathrm{Q}_{t-\varepsilon} \mid \mathrm{G}_{t}\right]=\frac{1}{t} \mathrm{E}^{h}\left[\mathrm{Q}_{t}+a_{t} \mid \mathrm{G}_{t}\right]-\mathrm{E}^{h}\left[b_{t} \mid \mathrm{G}_{t}\right] \tag{4.11}
\end{equation*}
$$

where

$$
\begin{aligned}
& a_{t}=\int_{0}^{t} \int_{s}^{1}\left[\frac{\partial h_{r}}{\partial r}+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<r\right\}} \frac{\partial h_{r}}{\partial t_{j}}\right. \\
&\left.+\mathrm{D} h_{r}(.,[s, r])\right] \frac{1}{h_{r}} d \mathrm{~N}_{r} d s-\int_{0}^{t} \int_{s}^{1} \mathrm{D} h_{r}(.,[s, r]) d r d s
\end{aligned}
$$

and

$$
b_{t}=\int_{t}^{1}\left[\frac{\partial h_{r}}{\partial r}+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathbf{T}_{j}<\boldsymbol{r}\right\}} \frac{\partial h_{r}}{\partial t_{j}}+\mathrm{D} h_{r}(.,[t, r])\right] \frac{1}{h_{r}} d \mathrm{~N}_{r}-\int_{t}^{1} \mathrm{D} h_{r}(.,[t, r]) d r .
$$

Proof. - First we note that if $\mathrm{H}\left(\left(1-\mathrm{T}_{1}\right) \vee 0, \ldots,\left(1-\mathrm{T}_{n}\right) \vee 0, \ldots\right)$ is a square integrable functional and its first partial derivatives are all bounded by a constant, then, using a similar argument as in [6], we have the integration by parts formula

$$
\begin{equation*}
\mathrm{E}\left[\left(\int_{0}^{1} u_{s} d \mathrm{Q}_{s}\right) \mathrm{H}\right]=-\mathrm{E}[\mathrm{DH}(., \mathrm{U})] \tag{4.12}
\end{equation*}
$$

where $\mathrm{DH}(., \mathrm{U})$ is the derivative in direction U of Definition 4.1.
A direct consequence is the product rule

$$
\begin{equation*}
\mathrm{E}\left[\mathrm{FH}\left(\int_{0}^{1} u_{s} d \mathrm{Q}_{s}\right)\right]=-\mathrm{E}[\mathrm{FDH}(., \mathrm{U})]-\mathrm{E}[\operatorname{HDF}(., \mathrm{U})] . \tag{4.13}
\end{equation*}
$$

Let $\mathrm{H}=\mathrm{G}_{1}$ be the Girsanov density, then (4.13) becomes

$$
\begin{equation*}
\mathrm{E}^{h}\left[\mathrm{~F} \int_{0}^{1} u_{s} d \mathrm{Q}_{s}\right]=-\mathrm{E}^{h}[\mathrm{DF}(., \mathrm{U})]-\mathrm{E}^{h}\left[\mathrm{FG}_{1}^{-1} \mathrm{DG}_{1}(., \mathrm{U})\right] . \tag{4.14}
\end{equation*}
$$

Now fix $t_{0} \in(0,1)$. Write $\mathrm{T}_{k}\left(t_{0}\right)$ for the $k$-th jump time of $\mathrm{N}_{t}$ greater than $t_{0}$. Suppose F is a bounded and $\mathrm{G}_{t_{0}}$ measurable function. Furthermore, we suppose that $F$ is a differentiable function (in the sense of Definition
4.1) of the form

$$
\mathrm{F}\left(\left(1-\mathrm{T}_{1}\left(t_{0}\right)\right) \vee 0, \ldots,\left(1-\mathrm{T}_{k}\left(t_{0}\right)\right) \vee 0, \ldots\right)
$$

and that the derivatives of F are bounded. Then the measure $\mathrm{DF}(., d t)$ is concentrated on $\left[t_{0}, 1\right]$ and (4.14) holds for such an $F$. Take $u_{s}=\mathrm{I}_{\left[t_{0}-\varepsilon, t_{0}\right]}(s)$ in (4.14). For such an F

$$
\begin{aligned}
\operatorname{DF}(., \mathrm{U}) & =\int_{t_{0}-\varepsilon}^{t_{0}} \operatorname{DF}(.,[r, 1]) d r \\
& =\int_{t_{0-\varepsilon}}^{t_{0}} \operatorname{DF}\left(.,\left[t_{0}, 1\right]\right) d r \\
& =\varepsilon \operatorname{DF}\left(.,\left[t_{0}, 1\right]\right) .
\end{aligned}
$$

Therefore, we have from (4.14)

$$
\begin{align*}
& \mathrm{E}^{h}\left[\left(\mathrm{Q}_{t_{0}}-\mathrm{Q}_{t_{0-\varepsilon}}\right) \mathrm{F}\right]=-\varepsilon \mathrm{E}^{h}\left[\mathrm{DF}\left(.,\left[\mathrm{t}_{0}, 1\right]\right)\right] \\
&+\mathrm{E}^{h}\left[\mathrm{FG}_{1}^{-1} \int_{t_{0-\varepsilon}}^{t_{0}} \sum_{i=1}^{\infty} \mathrm{I}_{s \leqq \mathrm{~T}_{i}<1} \frac{\partial \mathrm{G}_{1}}{\partial t_{i}} d s\right] . \tag{4.15}
\end{align*}
$$

From (4.15), for almost all $t$
$\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon}\left[\left(\mathrm{Q}_{t_{0}}-\mathrm{Q}_{t_{0}-\varepsilon}\right) \mathrm{F}\right]=-\mathrm{E}^{h}\left[\mathrm{DF}\left(.,\left[t_{0}, 1\right]\right)\right]$

$$
\begin{equation*}
+\mathrm{E}^{h}\left[\mathrm{FG}_{1}^{-1} \sum_{i=1}^{\infty} \mathrm{I}_{t_{0} \leqq \mathrm{~T}_{i}<1} \frac{\partial \mathrm{G}_{1}}{\partial t_{i}}\right] \tag{4.16}
\end{equation*}
$$

Using (4.15) again with $\varepsilon=t_{0}=t$, we have
$-\mathrm{E}^{h}[\mathrm{DF}(.,[t, 1])]=\frac{1}{t} \mathrm{E}^{h}\left[\mathrm{Q}_{t} \mathrm{~F}\right]$

$$
\begin{equation*}
-\frac{1}{t} \mathrm{E}^{h}\left[\mathrm{FG}_{1}^{-1} \int_{0}^{t} \sum_{i=1}^{\infty} \mathrm{I}_{s \leqq \mathrm{~T}_{i}<1} \frac{\partial \mathrm{G}_{1}}{\partial t_{i}} d s\right] \tag{4.17}
\end{equation*}
$$

Now let $u_{s}=\mathrm{I}_{[0, t]}(s)$ in Theorem 4.2 to obtain

$$
\begin{aligned}
&-\int_{0}^{t}\left(\sum_{i=1}^{\infty} \mathrm{I}_{s \leqq \mathrm{~T}_{i}<1} \frac{\partial \mathrm{G}_{1}}{\partial t_{i}}\right) \mathrm{G}_{1}^{-1} d s \\
&=\int_{0}^{t} \int_{s}^{1}\left[\frac{\partial h_{r}}{\partial r}+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<r\right\}} \frac{\partial h_{r}}{\partial t_{j}}+\mathrm{D} h_{r}(.,[s, r])\right] \frac{1}{h_{r}} d \mathrm{~N}_{r} d s \\
& \quad-\int_{0}^{t} \int_{s}^{1} \mathrm{D} h_{r}(.,[s, r]) d r d s .
\end{aligned}
$$

Hence (4.17) becomes

$$
\begin{equation*}
-\mathrm{E}^{h}[\mathrm{DF}(.,[t, 1])]=\frac{1}{t} \mathrm{E}^{h}\left[\mathrm{Q}_{t} \mathrm{~F}\right]+\frac{1}{t} \mathrm{E}^{h}\left[a_{t} \mathrm{~F}\right] \tag{4.18}
\end{equation*}
$$

Now take $u_{s}=\mathrm{I}_{[t-\varepsilon, t]}(s)$ in Theorem 4.2 to obtain

$$
\begin{align*}
&-\int_{t-\varepsilon}^{t}\left(\sum_{i=1}^{\infty} \mathrm{I}_{s \leq \mathrm{T}_{i}<1} \frac{\partial \mathrm{G}_{1}}{\partial t_{i}}\right) \mathrm{G}_{1}^{-1} d s \\
&=\int_{t-\varepsilon}^{t} \int_{s}^{1}\left[\frac{\partial h_{r}}{\partial r}+\sum_{j=1}^{\infty} \mathrm{I}_{\left\{\mathrm{T}_{j}<r\right\}} \frac{\partial h_{r}}{\partial t_{j}}\right.\left.+\mathrm{D} h_{r}(.,[s, r])\right] \frac{1}{h_{r}} d \mathrm{~N}_{r} d s \\
& \quad-\int_{t-\varepsilon}^{t} \int_{s}^{1} \mathrm{D} h_{r}(.,[s, r]) d r d s . \tag{4.19}
\end{align*}
$$

Multiply both sides of (4.19) by F, and then take expectations

$$
\begin{equation*}
-\mathrm{E}^{h}\left[\mathrm{~F} \int_{t-\varepsilon}^{t}\left(\sum_{i=1}^{\infty} \mathrm{I}_{s \leqq \mathrm{~T}_{i}<1} \frac{\partial \mathrm{G}_{1}}{\partial t_{i}}\right) \mathrm{G}_{1}^{-1} d s\right]=\mathrm{E}^{h}\left[\mathrm{~F} \int_{t-\varepsilon}^{t} b_{s} d s\right] . \tag{4.20}
\end{equation*}
$$

Divide both sides of (4.20) by $\varepsilon$, and then let $\varepsilon \downarrow 0$, to obtain for almost all $t$

$$
\begin{equation*}
-\mathrm{E}^{h}\left[\mathrm{~F}\left(\sum_{i=1}^{\infty} \mathrm{I}_{t \leqq \mathrm{~T}_{i}<1} \frac{\partial \mathrm{G}_{1}}{\partial t_{i}}\right) \mathrm{G}_{1}^{-1}\right]=\mathrm{E}^{h}\left[b_{t} \mathrm{~F}\right] . \tag{4.21}
\end{equation*}
$$

Combining (4.16), (4.18) and (4.21), we have

$$
\lim _{\varepsilon \downarrow 0} \frac{1}{\varepsilon} \mathrm{E}^{h}\left[\left(\mathrm{Q}_{t}-\mathrm{Q}_{t-\varepsilon}\right) \mathrm{F}\right]=\frac{1}{t} \mathrm{E}^{h}\left[\left(a_{t}+\mathrm{Q}_{t}\right) \mathrm{F}\right]-\mathrm{E}^{h}\left[b_{t} \mathrm{~F}\right] .
$$

Thus we have proved (4.11).
As a consequence of Theorem $4.3, \mathrm{Z}_{t}$ is a reverse time $\mathrm{G}_{t}$-quasimartingale having the decomposition given by (4.10). It follows immediately that the integrand $\alpha_{t}$ in (4.10) is given by

$$
\alpha_{t}=\mathrm{E}^{h}\left[b_{t}+h_{t}-1 \mid \mathrm{G}_{t}\right]-\frac{1}{t} \mathrm{E}^{h}\left[a_{t}+\mathrm{Q}_{t} \mid \mathrm{G}_{t}\right] .
$$

## REFERENCES

[1] K. Bichteler, J.-B. Gravereaux and J. Jacod, Malliavin Calculus for Processes with Jumps, Stochastics Monographs, Vol. 2, Gordon and Breach, New York, London, 1987.
[2] R. J. Elliott, Stochastic Calculus and Applications, Applications of Mathematics, Vol. 18, Springer-Verlag, Berlin, Heidelberg, New York, 1982.
[3] R. J. Elliott, Filtering and control for point process observations, Research Report No. 09.86.1, Department of Statistics and Applied Probability, University of Alberta, 1986.
[4] R. J. Elliott and B. D. O. Anderson, Reverse time diffusions, Stoch. Proc. Appl., Vol. 19, 1985, pp. 327-339.
[5] R. J. Elliott and P. E. Kopp, Eauivalent Martingale measures for bridge processes, Technical Report No. 89.17, Department of Statistics and Applied Probability, University of Alberta, 1989.
[6] R. J. Elliott and A. H. TsoI, Integration by parts for Poisson processes, Technical Report No. 89.09, Department of Statistics and Applied Probability, University of Alberta, 1989.
[7] H. Föllmer, Random fields and diffusion processes, École d'Été de Probabilités de Saint-Flour XV-XVII, 1985-87, Lect. Notes Math., Vol. 1362, Springer-Verlag, Berlin.
[8] U. G. Haussman and E. Pardoux, Time reversal of diffusions, Ann. Prob., Vol. 14, 1986, pp. 1188-1205.
[9] J. Jacod and P. Protter, Time reversal on Lévy processes, to appear.
[10] T. Jeulin and M. Yor, Inégalité de Hardy, semi-martingales et faux-amis, Séminaire de Probabilités XIII, Lect. Notes Math., Vol. 721, 1979, pp. 332-359, Springer-Verlag, Berlin.
[11] E. Nelson, Dynamical theories of Brownian motion, Princeton University Press, 1967.
[12] C. Stricker, Une Characterization des quasimartingales, Séminaire de Probabilités IX, Lect. Notes Math., Vol. 465, 1975, pp. 420-424, Springer-Verlag, Berlin.
(Manuscript received September 4, 1989)
(In revised form December 11, 1989.)


[^0]:    Classification A.M.S. : 60 G 55, 60 J 75.
    ( ${ }^{1}$ ) Research partially supported by N.S.E.R.C. Grant A-7964, the Air Force Office of Scientific Research, United States Air Force, under contract AFOSR-86-0332, and the U.S. Army Research Office under contract DAAL 03-87-0102.
    ( ${ }^{2}$ ) Research partially supported by N.S.E.R.C. Grant A-7964 and a University of Alberta Dissertation Fellowship.

