

ANNALES DE L'I. H. P., SECTION A

R. GERLACH

V. WÜNSCH

Contributions to polynomial conformal tensors

Annales de l'I. H. P., section A, tome 70, n° 3 (1999), p. 313-340

<http://www.numdam.org/item?id=AIHPA_1999__70_3_313_0>

© Gauthier-Villars, 1999, tous droits réservés.

L'accès aux archives de la revue « Annales de l'I. H. P., section A » implique l'accord avec les conditions générales d'utilisation (<http://www.numdam.org/conditions>). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

NUMDAM

*Article numérisé dans le cadre du programme
Numérisation de documents anciens mathématiques
<http://www.numdam.org/>*

Contributions to polynomial conformal tensors

by

R. GERLACH and V. WÜNSCH

FSU Jena, Mathematisches Institut,
Ernst-Abbe-Platz 4, D-07743 Jena, Germany

ABSTRACT. – By means of a certain conformal covariant differentiation process we construct generating systems for conformally invariant tensors in a pseudo-Riemannian manifold and use these invariants to derive necessary conditions for the validity of Huygens' principle of some conformally invariant field equations as well as for a space-time to be conformally related to an Einstein space-time. © Elsevier, Paris

Key words: pseudo-Riemannian manifold, conformally invariant tensor, conformal covariant derivative, moment, Huygens' principle, conformal Einstein space-time.

RÉSUMÉ. – Grâce à un certain processus de différentiation covariante conforme, nous construisons des systèmes générateurs de tenseurs invariants conformes dans une variété pseudo-Riemannienne et nous utilisons ces invariants pour en déduire les conditions nécessaires à la validité du principe d'Huyghens pour quelques équations invariantes conformes de champ et pour l'application conforme d'un espace-temps à l'espace-temps d'Einstein. © Elsevier, Paris

A.M.S. subject classification: 53 B, 83 C

1. INTRODUCTION

Let (M, g) be a pseudo-Riemannian C^∞ manifold of dimension n ($n \geq 3$) and g_{ab} , g^{ab} , ∇_a , R_{abcd} , R_{ab} , R , C_{abcd} the local components of the covariant and contravariant metric tensor, the Levi-Civita connection, the curvature tensor, the Ricci tensor, the scalar curvature and the Weyl curvature tensor, respectively¹. We consider polynomial tensors, *i.e.* tensors whose components are polynomials in g^{ab} and the partial derivatives of g_{ab} . These tensors are just the elements of the tensor algebra \mathcal{R} generated by the tensors

$$g_{ab}, g^{ab}, \nabla_{(i_1} \dots \nabla_{i_r} R_{i_{r+1}|ab|i_{r+2})}, \quad r = 0, 1, 2, \dots \quad (1.1)$$

by means of the usual tensor operations [Scho; dP; GüW1].

A tensor $T[g] \in \mathcal{R}$ is said to be *conformally invariant* (or more briefly a *conformal tensor*) with weight ω if $T[g]$ has under a conformal transformation

$$\bar{g}_{ab} = e^{2\Phi} g_{ab}, \quad \Phi \in C^\infty(M) \quad (1.2)$$

the transformation law [Scho; GüW1]

$$T[\bar{g}] = e^{2\omega\Phi} T[g]. \quad (1.3)$$

Example. – The Bach tensor

$$B_{i_1 i_2} := \nabla_a \nabla_b C^a_{.i_1 i_2.}{}^b + \frac{1}{n-2} C^a_{.i_1 i_2.}{}^b R_{ab} \quad (1.4)$$

is a polynomial conformal tensor with $\omega = -1$ if $n = 4$ [Scho; GüW1].

General algebraic and differential properties of conformal tensors were investigated by Thomas [Tho], Szekeres [Sz], du Plessis [dP] and Günther/Wünsch [GüW1,2]. In [Sz; dP] particular sequences of tensors satisfying (1.3) are given which generate all tensors with this property. However, in general the elements of these sequences are not polynomial. Conformal transformations and in particular the polynomial conformal tensors have a great variety of applications, *e.g.* for a Lagrangian formulation with locality principle of both general relativity and conformal field theories and in the propagation theory of conformally invariant field equations [AMLW; BØ; Gü; McL; Ø; PR; Wü1-5]. It is an important problem to give a survey of all polynomial conformal tensors or, with less pretension, to give a method for constructing special classes of such tensors. Such a method was developed in [GüW1,2; Schi] using the *infinitesimal generator* and the *conformal covariant derivative* of a tensor $T \in \mathcal{R}$.

¹ All investigations in this paper are of local nature.

The paper is organized as follows. In Section 2, the basic ideas and results of [GüW1] are given with the help of which we derive further classes of generating systems for conformal tensors in Section 3-5. In Section 4 we transfer this method for the physically important case $n = 4$ to an extended tensor algebra by involving the Levi-Civita pseudo tensor. In Section 5 we consider the pseudo Riemannian space of dimension 3. In this case one has to replace the Weyl tensor by the conformal tensor $S_{abc} := \nabla_{[c} L_{b]} a$. Finally, in Section 6 we use generating systems for the conformal tensors for the derivation of further explicit moment equations for conformally invariant field equations for space-times and necessary conditions for a space-time to be conformally related to an Einstein space-time.

REMARK 1.1. – The Sections 3 and 4 are an English version of results which have already been published in [GeW1,2] in German, however, the journal "Wissenschaftliche Zeitschrift der Pädagogischen Hochschule Erfurt/Mühlhausen" has ceased publication and can neither be found at the Libraries and report organs nor be ordered. On the other hand some colleagues working on this field recommended the publication of these invariants in an international renowned journal.

2. SOME KNOWN RESULTS

In Section 3-5 we use the following basic ideas and results (see [GüW1,2]):

PROPOSITION 2.1. – (i) $T \in \mathcal{R}$ has, under the conformal transformation (1.2), a transformation law of the form

$$T[\bar{g}] = e^{2\omega\Phi} \left(T[g] + \sum_{k=1}^m P_k[g, \Phi] \right),$$

where the $P_k[g, \Phi]$ are tensor-valued homogeneous polynomials of degree k in the derivatives of Φ up to a certain order.

(ii) $T \in \mathcal{R}$ is conformally invariant iff this is true under infinitesimal conformal transformations, i.e. iff

$$P_1[g, \Phi] = \lim_{\varepsilon \rightarrow 0} \frac{1}{\varepsilon} (T[e^{2\varepsilon\Phi} g] - e^{2\varepsilon\omega\Phi} T[g]) = 0.$$

REMARK 2.1. – Because of this proposition it follows from $P_1[g, \Phi] = 0$ that $P_k[g, \Phi] = 0$ ($k = 2, \dots, m$), i.e. for the construction of conformal tensors it is sufficient to calculate only "to the first order in the derivatives of Φ "

Example. – For the *Schouten tensor* L_{ab} defined by

$$L_{ab} = -R_{ab} + \frac{1}{2(n-1)} R g_{ab} \quad (2.1)$$

we have

$$P_1[g, \phi]_{ab} = \frac{n-2}{2} \nabla_a \nabla_b \Phi, \quad P_2[g, \phi]_{ab} = \frac{n-2}{8} (g_{ab} g^{lk} - 2\delta_a^l \delta_b^k) \nabla_l \Phi \nabla_k \Phi.$$

Let τ be the subalgebra of those elements of \mathcal{R} which contain only *first* order derivatives of Φ in their transformation law and let i be the ideal of \mathcal{R} which is generated by the tensors

$$\nabla(i_1 \dots \nabla i_k L_{i_{k+1} i_{k+2}}), \quad k = 0, 1, 2, \dots \quad (2.2)$$

Now we define two linear operators in τ : The first one comes from infinitesimal conformal transformations, the second one comes from differentiation.

It follows from Proposition 2.1 that $T[g] \in \tau$ iff $P_1[g, \Phi]$ has the form

$$P_1[g, \Phi] = X^\gamma(T) \nabla_\gamma \Phi. \quad (2.3)$$

DEFINITION 2.1. – The linear operator X^γ defined on τ by (2.3) is called the infinitesimal generator of T .

For X^γ the Leibniz rule holds and one has $X^\gamma(g_{ab}) = 0$.

COROLLARY 2.1. – $T[g] \in \mathcal{R}$ is conformally invariant iff $T[g] \in \tau$ and $X^\gamma(T) = 0$.

Examples.

$$X^\gamma(C_{abcd}) = 0, \quad X^\gamma(\nabla_u C^u{}_{abc}) = (n-3)C^\gamma{}_{abc}, \quad (2.4)$$

$$X^\gamma(\nabla_e C_{abcd}) = 2[-\delta_e{}^\gamma C_{abcd} + \delta_{[a}{}^\gamma C_{b]ecd} + \delta_{[c}{}^\gamma C_{ab|d]e} - g_{e[a} C_{b].cd}{}^\gamma - g_{e[c} C_{ab|d].}{}^\gamma] \quad (2.5)$$

If $T \in \tau$, then we have in general $\nabla_a T \notin \tau$.

DEFINITION 2.2. – For $T \in \tau$ the tensor

$$\overset{c}{\nabla}_a T := \nabla_a T - \frac{1}{n-2} L_{a\gamma} X^\gamma(T) \quad (2.6)$$

is called the conformal covariant derivative of T ².

PROPOSITION 2.2. – (i) The conformal covariant derivative $\overset{c}{\nabla}_a$ is linear, obeys Leibniz's rule and commutes with contractions.

(ii) $\overset{c}{\nabla}_a : \tau \rightarrow \tau$

(iii) If $n \geq 4$, then τ is generated by the tensors

$$g_{ab}, g^{ab}, \overset{c}{\nabla}_{(i_1} \cdots \overset{c}{\nabla}_{i_r} C^a_{i_{r+1} i_{r+2})}{}^b, \quad r = 0, 1, 2, \dots \quad (2.7)$$

(iv) \mathcal{R} is generated by the tensors (2.7) and (2.2).

(v) For every $T \in \mathcal{R}$ there exists one and only one element $T_1 \in \tau$ with $(T - T_1) \in i$. It is $\tau \cap i = \{0\}$, and \mathcal{R}/i is isomorphic to τ .

(vi) If $T \in \tau$ has the weight ω , then

$$X^k (\overset{c}{\nabla}_a T) - \overset{c}{\nabla}_a (X^k T) = 2\omega \delta_a^k T + \triangle_a^k T, \quad (2.8)$$

where

$$\triangle_a^k (T_{ij\dots}{}^{lm\dots}) := \triangle_{as}^{kl} T_{ij\dots}{}^{sm\dots} + \triangle_{as}^{km} T_{ij\dots}{}^{ls\dots} \cdots - \triangle_{ai}^{ks} T_{sj\dots}{}^{lm\dots} - \triangle_{aj}^{ks} T_{is\dots}{}^{lm\dots} \dots \quad (2.9)$$

and

$$\triangle_{ai}^{ks} := \delta_a^k \delta_i^s + \delta_i^k \delta_a^s - g_{ai} g^{ks}.$$

(vii) The Ricci-identity for $\overset{c}{\nabla}_a$ has the form

$$\overset{c}{\nabla}_{[a} \overset{c}{\nabla}_{b]} T = (C, T)_{\alpha\beta} - \frac{1}{n-2} \nabla_{[\alpha} L_{\beta]\gamma} X^\gamma(T), \quad (2.10)$$

where $(C, T)_{\alpha\beta}$ is the term one obtains from the right hand side of the usual Ricci identity

$$\nabla_{[a} \nabla_{b]} T = (R, T)_{\alpha\beta}$$

by substitution of R by C .

Examples. – (i) $\overset{c}{\nabla}_k C_{abcd} = \nabla_k C_{abcd}$

$$(ii) \quad B_{ab} = \overset{c}{\nabla}_k \overset{c}{\nabla}_l C^k_{ab}{}^l \quad (\text{Bach tensor}) \quad (2.11)$$

Using (2.5) and (2.8), we get $X^k (B_{ab}) = 2(n-4) \overset{c}{\nabla}_k C^k_{ab}{}^\gamma$.

² This definition is different from the definition of the conformal covariant derivative introduced by Weyl and du Plessis [dP; Scho].

When latin indices with subindices (e.g. i_1, \dots, i_r) appear in the sequel, we assume that symmetrization has been carried out over the indices. If T is any tensor with covariant rank r ($r \geq 2$), then we denote the trace-free part of T by $TS(T)$. For a symmetric tensor $T_{i_1 \dots i_r}$ with $r \geq 2$ we write

$$\overset{1}{T}_{i_1 \dots i_r} = \underset{*}{\overset{2}{T}}_{i_1 \dots i_r} \quad \text{iff} \quad TS(\overset{1}{T}_{i_1 \dots i_r} - \underset{*}{\overset{2}{T}}_{i_1 \dots i_r}) = 0.$$

In [Wü5; GüW2] the following was proved:

LEMMA 2.1. – *If $T_{i_1 \dots i_{k-s}}$ is a symmetric, conformally invariant tensor with covariant rank $(k-s)$ and weight ω , then*

$$X^\gamma(\overset{c}{\nabla}_{i_1} \dots \overset{c}{\nabla}_{i_s} T_{i_{s+1} \dots i_k}) = s(2\omega + s - 2k + 1) \delta_{i_1}^{\gamma} \overset{c}{\nabla}_{i_2} \dots \overset{c}{\nabla}_{i_s} T_{i_{s+1} \dots i_k}.$$

DEFINITION 2.3. – *A conformally invariant tensor T is called trivial if T is generated by*

$$\{g_{ab}, g^{ab}, C_{abcd}\}. \quad (2.12)$$

Let $\mathcal{S}_r(\omega, n)$ be the set of all nontrivial conformal, symmetric, trace-free tensors of \mathcal{R} with weight ω and covariant rank r .

LEMMA 2.2. – *If a monomial in $\mathcal{S}_r(\omega, n)$ contains α factors $\{\overset{c}{\nabla}_{i_1} \dots \overset{c}{\nabla}_{i_{k-2}} C_{i_{k-1}abi_k}\}$ and q operators $\overset{c}{\nabla}$, then*

$$r - 2\omega = 2\alpha + q \quad \text{and} \quad \omega \leq 0. \quad (2.13)$$

The number α is called C -order of a monomial of $T \in \mathcal{S}_r(\omega, n)$.

LEMMA 2.3. – *A tensor in $\mathcal{S}_r(\omega, n)$ contains a monomial with $\alpha = 1$, if and only if*

$$r = 2, \quad n \text{ is even and } \omega = \frac{n-2}{2}.$$

Let $\mathcal{S}_r^{(\alpha)}(\omega, n)$ be the subset of those elements of $\mathcal{S}_r(\omega, n)$ whose monomials have the C -order α . The Propositions 2.1, 2.2, Corollary 2.1, Definition 2.2 and Lemma 2.1 are very useful for the construction of nontrivial conformally invariant tensors. In [GüW2] generating systems are constructed for $\mathcal{S}_r(\omega, n)$ in the cases $(r, \omega) = (2, -1), (1, -2), (0, -3), (2, -2), (4, -1)$. In particular the following results were proved (see [Wü1, GüW2, Gü]):

PROPOSITION 2.3. – *If $S_r(-1, 4) \in \mathcal{S}_r(-1, 4)$ then one has*

$$S_r(-1, 4) = 0 \quad \text{for } r = 0, 1, 3, \quad S_2(-1, 4) = \alpha B, \quad S_4(-1, 4) = \sum_{m=1}^3 \beta_m W^{(m)},$$

where B is the Bach tensor,

$$\begin{aligned} W_{i_1 \dots i_4}^{(1)} := & TS[\nabla^a C_{.i_1 i_2.}^b {}^c \nabla_a C_{bi_3 i_4 c} + 16 \nabla_u C_{.i_1 i_2.}^u {}^a \nabla_\nu C_{.i_3 i_4 a}^\nu \\ & + 4 C_{.i_1 i_2.}^a {}^b \{2 \nabla_a \nabla_u C_{.i_3 i_4 b}^u - C_{ai_3 i_4.} {}^c L_{bc}\}], \end{aligned} \quad (2.14)$$

$$\begin{aligned} W_{i_1 \dots i_4}^{(2)} := & TS[2 \nabla_{i_1} C_{.i_2 i_3.}^a {}^b \nabla_u C_{.abi_4}^u + 2 \nabla_u C_{.i_1 i_2 a}^u \nabla_\nu C_{.i_3 i_4.}^\nu {}^a \\ & - C_{.i_1 i_2.}^a {}^b \{2 \nabla_{i_3} \nabla_u C_{.abi_4}^u - C_{.abi_3}^c L_{ci_4}\}] \end{aligned} \quad (2.15)$$

$$W_{i_1 \dots i_4}^{(3)} := TS[C^{abcd} C_{ai_1 i_2 d} C_{bi_3 i_4 c}]$$

and $\alpha, \beta_1, \beta_2, \beta_3 \in \mathbb{R}$.

3. GENERATING SYSTEMS OF $\mathcal{S}_6(-1, n)$ FOR $n \geq 4$

For a conformal tensor T_k the Ricci-identity (2.10) reduces to

$$\nabla_{[a} \nabla_{b]} T_c = -\frac{1}{2} C_{abc.} {}^k T_k. \quad (3.1)$$

Furthermore, we use in the sequel the Bianchi identity

$$\nabla_{[a} C_{bc]ij} + \frac{1}{n-3} [g_{j[a} \nabla_u C_{i|bc]}^u - g_{i[a} \nabla_u C_{j|bc]}^u] = 0. \quad (3.2)$$

Because of Lemma 2.3 and (2.13) the sets $\mathcal{S}_6^{(1)}(-1, n)$ and $\mathcal{S}_6^{(4)}(-1, n)$ are empty. Thus it follows from (3.1) that

$$\mathcal{S}_6(-1, n) = \mathcal{S}_6^{(2)}(-1, n) \cup \mathcal{S}_6^{(3)}(-1, n). \quad (3.3)$$

Firstly, consider the case $\alpha = 2$: Then on account of (2.13) $q = 4$. Because of (3.1) and (3.2) a tensor of $\mathcal{S}_6^{(2)}(-1, n)$ has to be a linear combination of the following linearly independent tensors ³

$$\begin{aligned} C_{i_1 \dots i_6}^{1(2)} &= TS[C_{abc i_1} {}^c {}^c {}^c {}^c {}^c {}_{i_6} C^{abc} {}_{i_6}] \\ C_{i_1 \dots i_6}^{2(2)} &= TS[C_{.i_1 i_2.}^a {}^b {}^c {}^c {}^c {}^c {}_{i_6} C_{.abi_6}^u] \\ C_{i_1 \dots i_6}^{3(2)} &= TS[C_{.i_1 i_2.}^a {}^b {}^c {}^c {}^c {}^c {}_{i_6} C_{.i_5 i_6 b}^u] \end{aligned}$$

³ One can show the linear independence by means of the methods developed in [GüW2].

$$\begin{aligned}
& \stackrel{4}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}^a C_{.i_1 i_2.}^b \stackrel{c}{\nabla}_{i_3}^c \stackrel{c}{\nabla}_{i_4}^d \stackrel{c}{\nabla}_a^e C_{b i_5 i_6 c}] \\
& \stackrel{5}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} C_{a b c i_2}^a \stackrel{c}{\nabla}_{i_3}^b \stackrel{c}{\nabla}_{i_4}^c \stackrel{c}{\nabla}_{i_5}^d C_{...}^{a b c} {}_{i_6}] \\
& \stackrel{6}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} C_{a i_2 i_3 b}^a \stackrel{c}{\nabla}_{i_4}^b \stackrel{c}{\nabla}_{i_5}^c \stackrel{c}{\nabla}_u^d C_{...}^{u a b} {}_{i_6}] \\
& \stackrel{7}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} C_{.i_2 i_3.}^a \stackrel{b}{\nabla}_{i_4}^c \stackrel{c}{\nabla}_a^d \stackrel{c}{\nabla}_u^e C_{.i_5 i_6 b}^u] \\
& \stackrel{8}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_u C_{...}^{u a b} {}_{i_1}^a \stackrel{c}{\nabla}_{i_2}^b \stackrel{c}{\nabla}_{i_3}^c \stackrel{c}{\nabla}_{i_4}^d C_{a i_5 i_6 b}] \\
& \stackrel{9}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_u C_{.i_1 i_2.}^u \stackrel{a}{\nabla}_{i_3}^c \stackrel{c}{\nabla}_{i_4}^d \stackrel{c}{\nabla}_v^e C_{.i_5 i_6 a}^v] \\
& \stackrel{10}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_a^b C_{.i_2 i_3.}^b \stackrel{c}{\nabla}_{i_4}^c \stackrel{c}{\nabla}_a^d C_{b i_5 i_6 c}] \\
& \stackrel{11}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_{i_2} C_{a b c i_3}^a \stackrel{c}{\nabla}_{i_4}^b \stackrel{c}{\nabla}_{i_5}^c C_{...}^{a b c} {}_{i_6}] \\
& \stackrel{12}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_{i_2} C_{.i_3 i_4.}^a \stackrel{b}{\nabla}_{i_5}^c \stackrel{c}{\nabla}_u^d C_{.a b i_6}^u] \\
& \stackrel{13}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_{i_2} C_{.i_3 i_4.}^a \stackrel{b}{\nabla}_a^c \stackrel{c}{\nabla}_u^d C_{.i_5 i_6 b}^u] \\
& \stackrel{14}{C}_{i_1 \dots i_6}^{(2)} = TS[\stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_u C_{.i_2 i_3.}^u \stackrel{a}{\nabla}_{i_4}^c \stackrel{c}{\nabla}_v^d C_{.i_5 i_6 a}^v]. \tag{3.4}
\end{aligned}$$

We put

$$\begin{aligned}
& \stackrel{1}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \delta_{i_1}^\gamma \stackrel{c}{\nabla}_{i_2} \stackrel{c}{\nabla}_a C_{b i_3 i_4 c} \stackrel{c}{\nabla}_{i_5} C_{...}^{a b c} {}_{i_6} \\
& \stackrel{3}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \delta_{i_1}^\gamma \stackrel{c}{\nabla}_{i_2} \stackrel{c}{\nabla}_u C_{.a b i_3}^u \stackrel{c}{\nabla}_{i_4}^a C_{.i_5 i_6 b}^b \\
& \stackrel{5}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \delta_{i_1}^\gamma \stackrel{c}{\nabla}_{i_2} \stackrel{c}{\nabla}_u C_{.i_3 i_4.}^u \stackrel{a}{\nabla}_v^c \stackrel{c}{\nabla}_{i_5} C_{.i_6 a}^v \\
& \stackrel{7}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_{i_2} C_{.a b i_3}^\gamma \stackrel{c}{\nabla}_{i_4}^a C_{.i_5 i_6 b}^b \\
& \stackrel{9}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_{i_2} C_{a i_3 i_4 b} \stackrel{c}{\nabla}_{i_5}^c C_{...}^{\gamma a b} \\
& \stackrel{2}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \delta_{i_1}^\gamma \stackrel{c}{\nabla}_{i_2} \stackrel{c}{\nabla}_a C_{b i_3 i_4 c} \stackrel{c}{\nabla}_u^a C_{.i_5 i_6.}^b \\
& \stackrel{4}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \delta_{i_1}^\gamma \stackrel{c}{\nabla}_a \stackrel{c}{\nabla}_u C_{.i_2 i_3 b}^u \stackrel{c}{\nabla}_{i_4}^a C_{.i_5 i_6.}^b \\
& \stackrel{6}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \delta_{i_1}^\gamma \stackrel{c}{\nabla}_{i_2} \stackrel{c}{\nabla}_{i_3} C_{a i_4 i_5 b} \stackrel{c}{\nabla}_u^a C_{...}^{u a b} {}_{i_6} \\
& \stackrel{8}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_a C_{.i_2 i_3 b}^\gamma \stackrel{c}{\nabla}_{i_4}^a C_{.i_5 i_6.}^b \\
& \stackrel{10}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_{i_2} C_{a i_3 i_4 b} \stackrel{c}{\nabla}_u^a C_{.i_5 i_6.}^\gamma \\
& \stackrel{11}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \stackrel{c}{\nabla}_{i_1} \stackrel{c}{\nabla}_{i_2} C_{.i_3 i_4 a}^\gamma \stackrel{c}{\nabla}_u^a C_{.i_5 i_6.}^b \\
& \stackrel{13}{Q}_{i_1 \dots i_6}^{(2)\gamma} = \delta_{i_1}^\gamma C_{.i_2 i_3.}^a \stackrel{b}{\nabla}_{i_4}^c \stackrel{c}{\nabla}_{i_5}^d \stackrel{c}{\nabla}_u^e C_{.a b i_6}^u \\
& \stackrel{15}{Q}_{i_1 \dots i_6}^{(2)\gamma} = C_{.i_1 i_2.}^\gamma \stackrel{a}{\nabla}_{i_3}^c \stackrel{c}{\nabla}_{i_4}^d \stackrel{c}{\nabla}_u^e C_{.i_5 i_6 a}^u \\
& \stackrel{17}{Q}_{i_1 \dots i_6}^{(2)\gamma} = C_{.a b i_1}^\gamma \stackrel{c}{\nabla}_{i_2} \stackrel{c}{\nabla}_{i_3} \stackrel{c}{\nabla}_{i_4}^a C_{.i_5 i_6.}^b \\
& \stackrel{19}{Q}_{i_1 \dots i_6}^{(2)\gamma} = C_{a i_1 i_2 b} \stackrel{c}{\nabla}_{i_3} \stackrel{c}{\nabla}_{i_4} \stackrel{c}{\nabla}_{i_5}^a C_{...}^{\gamma a b} {}_{i_6} \tag{3.5}
\end{aligned}$$

Using (3.1), (3.2) and the transformation formulae (2.4), (2.5), (2.8), we obtain

$$\begin{bmatrix} X^\gamma(C_{i_1 \dots i_6}^{(2)}) \\ \vdots \\ \vdots \\ X^\gamma(C_{i_1 \dots i_6}^{(14)}) \end{bmatrix} = A_6^{(2)}(-1, n) \begin{bmatrix} Q_{i_1 \dots i_6}^{(2)\gamma} \\ \vdots \\ \vdots \\ Q_{i_1 \dots i_6}^{(19)\gamma} \end{bmatrix},$$

where

$$A_6^{(2)}(-1, n) =$$

$$\begin{vmatrix} 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -8 & -24 & 8 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -18 & 0 & 3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & -14 & -6 & a & 0 & 0 \\ -2 & -12 & 4/a & 0 & 0 & -\frac{6}{a} & 4 & 2 & -4 & -6 & 2\frac{c}{a} & \frac{2}{a} & -\frac{4}{a} & 0 & 2\frac{b}{a} & -4 & 2 & -1 & -2 \\ -30 & 0 & 30/a & 0 & 0 & 0 & 6 & 0 & -6 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 2 & -3 & -2 \\ 0 & 0 & -10 & 0 & 0 & 0 & a & 0 & 0 & 0 & 0 & 2 & -4 & 0 & -1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & -6 & 0 & 0 & 0 & a & 0 & 0 & 0 & 0 & -5 & 0 & -4 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -18 & 0 & 0 & 0 & 0 & 0 & -3 & 0 & 0 & 0 & 0 & 0 & a & 0 & 0 \\ 0 & 0 & 0 & 0 & -12 & 0 & 0 & 0 & 0 & 0 & a & 0 & 0 & 0 & 0 & a & 0 & 0 & 0 \\ -4 & -10 & -6/a & 0 & 0 & \frac{2}{a} & -6 & -10 & 6 & 2 & \frac{2}{a} & 2\frac{b+c}{a} & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ -32 & 0 & 32/a & 0 & 0 & 0 & -8 & 0 & 8 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -10 & 0 & 0 & -4 & 0 & 0 & a & 0 & 1 & -2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -10 & 0 & 1 & 0 & 0 & 0 & a & -4 & 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & -10 & 0 & 0 & 0 & 0 & 0 & 0 & 2a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \end{vmatrix}$$

with $a = n - 3$, $b = n - 4$, $c = n - 5$.

For $n > 4$, $\text{rank } A_6^{(2)}(-1, n) = 12$. If $n = 4$, one gets more identities concerning the tensors (3.4) and (3.5) (see [RW]; [Thi]). We obtain the following result:

PROPOSITION 3.1. – (i) If $n > 4$ then $S_6^{(2)}(-1, n)$ is generated by the tensors

$$\begin{aligned} S_6^{(2)}(-1, n) = & 7(n-3)C_{i_1 \dots i_6}^{(2)} - 228C_{i_1 \dots i_6}^{(2)} - 180C_{i_1 \dots i_6}^{(2)} \\ & - 45(n-3)C_{i_1 \dots i_6}^{(2)} - 41(n-3)C_{i_1 \dots i_6}^{(2)} \\ & + 1026C_{i_1 \dots i_6}^{(2)} + 630C_{i_1 \dots i_6}^{(2)} + 228C_{i_1 \dots i_6}^{(2)} \\ & + 90\frac{(n-4)}{(n-3)}C_{i_1 \dots i_6}^{(2)} + 54(n-3)C_{i_1 \dots i_6}^{(2)} \\ & + \frac{69}{2}(n-3)C_{i_1 \dots i_6}^{(2)} - 1026C_{i_1 \dots i_6}^{(2)} - 378C_{i_1 \dots i_6}^{(2)} \\ & - 108\frac{(n-4)}{(n-3)}C_{i_1 \dots i_6}^{(2)} \end{aligned}$$

$$\begin{aligned}
S_6^{(2)}(-1, n) = & -2(n-3)C_{i_1 \dots i_6}^{(2)} + 48C_{i_1 \dots i_6}^{(2)} + 16(n-3)C_{i_1 \dots i_6}^{(2)} \\
& - 216C_{i_1 \dots i_6}^{(2)} - 48C_{i_1 \dots i_6}^{(2)} - \frac{360}{(n-3)}C_{i_1 \dots i_6}^{(2)} \\
& - 15(n-3)C_{i_1 \dots i_6}^{(2)} + 216C_{i_1 \dots i_6}^{(2)} + \frac{432}{(n-3)}C_{i_1 \dots i_6}^{(2)}.
\end{aligned}$$

(ii) $S_6^{(2)}(-1, 4)$ is generated by the tensors

$$\begin{aligned}
S_6^{(2)}(-1, 4) = & -228C_{i_1 \dots i_6}^{(2)} - 180C_{i_1 \dots i_6}^{(2)} - 45C_{i_1 \dots i_6}^{(2)} + 1026C_{i_1 \dots i_6}^{(2)} \\
& + 630C_{i_1 \dots i_6}^{(2)} + 228C_{i_1 \dots i_6}^{(2)} + 54C_{i_1 \dots i_6}^{(2)} - 1026C_{i_1 \dots i_6}^{(2)} \\
& - 378C_{i_1 \dots i_6}^{(2)} \\
S_6^{(2)}(-1, 4) = & 48C_{i_1 \dots i_6}^{(2)} - 216C_{i_1 \dots i_6}^{(2)} - 48C_{i_1 \dots i_6}^{(2)} - 360C_{i_1 \dots i_6}^{(2)} \\
& + 216C_{i_1 \dots i_6}^{(2)} + 432C_{i_1 \dots i_6}^{(2)}
\end{aligned}$$

Now let be $\alpha = 3$: Then on account of (2.13) it is $q = 2$. Because of (3.1) and (3.2) a tensor of $S_6^{(3)}(-1, n)$ has to be a linear combination of the following linearly independent tensors ³⁾

$$\begin{aligned}
C_{i_1 \dots i_6}^{(3)} &= TS[C_{acbd} C_{.i_1 i_2}^a {}^b {}^c {}^c {}^c {}^d] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{acdi_1} C_{b..i_2}^c {}^c {}^c {}^c {}^c {}^b] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{acdi_1} C_{b..i_2}^c {}^c {}^c {}^c {}^c {}^b] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{acdi_3} {}^c {}^c {}^c {}^c {}^d] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{acdi_3} {}^c {}^c {}^c {}^c {}^d] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{acdi_3} {}^c {}^c {}^c {}^c {}^d] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{.i_3 i_4}^c {}^d {}^c {}^c {}^c {}^d] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{.i_3 i_4}^c {}^d {}^c {}^c {}^c {}^d] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{ai_3 i_4} {}^c {}^c {}^c {}^c {}^d] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{cab i_3} {}^c {}^c {}^c {}^u {}^c] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{.i_1 i_2}^a {}^b C_{ai_3 i_4 c} {}^c {}^c {}^c {}^u {}^c] \\
C_{i_1 \dots i_6}^{(3)} &= TS[C_{acbd} {}^c {}^c {}^c {}^c {}^c {}^d]
\end{aligned}$$

$$\begin{aligned}
& \overset{13}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{acd i_1} \overset{c}{\nabla}_{i_2} C_{.i_3 i_4}^a \overset{b}{\nabla}_{i_5} C_{b..i_6}^{cd}] \\
& \overset{14}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{acd i_1} \overset{c}{\nabla}_{i_2} C_{.i_3 i_4}^a \overset{b}{\nabla}_{i_5} C_{b..i_6}^{dc}] \\
& \overset{15}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{acd i_1} \overset{c}{\nabla}_{i_2} C_{.i_3 i_4}^a \overset{b}{\nabla}_b C_{.i_5 i_6}^c \overset{d}{\nabla}] \\
& \overset{16}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_{i_3} C_{acd i_4} \overset{c}{\nabla}_{i_5} C_{b..i_6}^{cd}] \\
& \overset{17}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_{i_3} C_{acd i_4} \overset{c}{\nabla}_{i_5} C_{b..i_6}^{dc}] \\
& \overset{18}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_{i_3} C_{acd i_4} \overset{c}{\nabla}_b C_{.i_5 i_6}^c \overset{d}{\nabla}] \\
\\
& \overset{19}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_a C_{ci_3 i_4 d} \overset{c}{\nabla}_b C_{.i_5 i_6}^c \overset{d}{\nabla}] \\
& \overset{20}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_{i_3} C_{.i_4 i_5}^c \overset{d}{\nabla}_{i_6} C_{acbd}] \\
& \overset{21}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_{i_3} C_{.i_4 i_5}^c \overset{d}{\nabla}_b C_{cad i_6}] \\
& \overset{22}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{cab i_1} \overset{c}{\nabla}_{i_2} C_{.i_3 i_4}^a \overset{b}{\nabla}_u C_{.i_5 i_6}^u \overset{c}{\nabla}] \\
& \overset{23}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_u C_{.i_3 i_4}^u \overset{c}{\nabla}_{i_5} C_{cab i_6}] \\
& \overset{24}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_u C_{.i_3 i_4}^u \overset{c}{\nabla}_a C_{bi_5 i_6 c}] \\
& \overset{25}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_u C_{.aci_3}^u \overset{c}{\nabla}_{i_4} C_{bi_5 i_6}^c] \\
& \overset{26}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_u C_{.ca i_3}^u \overset{c}{\nabla}_{i_4} C_{bi_5 i_6}^c] \\
& \overset{27}{C}_{i_1 \dots i_6}^{(3)} = TS[C_{.i_1 i_2}^a \overset{b}{\nabla}_u C_{.i_3 i_4 a}^u \overset{c}{\nabla}_{i_5} C_{.i_5 i_6 b}^v]. \tag{3.6}
\end{aligned}$$

We set

$$\begin{aligned}
& \overset{1}{Q}_{i_1 \dots i_6}^{(3)\gamma} = \delta_{i_1}^\gamma C_{.i_2 i_3}^a \overset{b}{\nabla}_{a i_4 i_5} C_{.b c i_6}^u \quad \overset{2}{Q}_{i_1 \dots i_6}^{(3)\gamma} = \delta_{i_1}^\gamma C_{.i_2 i_3}^a \overset{b}{\nabla}_{c a i_4} C_{.b c i_6}^u \overset{c}{\nabla} \\
& \overset{3}{Q}_{i_1 \dots i_6}^{(3)\gamma} = C_{.i_1 i_2}^a \overset{b}{\nabla}_{i_3 i_4 a} C_{.i_5 i_6 b}^u \quad \overset{4}{Q}_{i_1 \dots i_6}^{(3)\gamma} = C_{.i_1 i_2}^a \overset{b}{\nabla}_{a i_3 i_4 b} C_{.i_5 i_6}^u \overset{c}{\nabla} \\
& \overset{5}{Q}_{i_1 \dots i_6}^{(3)\gamma} = \delta_{i_1}^\gamma C_{a..i_2}^{cd} C_{bcd i_3} \overset{c}{\nabla}_{i_4} C_{.i_5 i_6}^a \overset{b}{\nabla} \quad \overset{6}{Q}_{i_1 \dots i_6}^{(3)\gamma} = \delta_{i_1}^\gamma C_{a..i_2}^{cd} C_{bi_3 cd} \overset{c}{\nabla}_{i_4} C_{.i_5 i_6}^a \overset{b}{\nabla} \\
& \overset{7}{Q}_{i_1 \dots i_6}^{(3)\gamma} = C_{.i_1 i_2}^a \overset{b}{\nabla}_{i_3 a} C_{.i_4}^{\gamma} \overset{c}{\nabla}_{i_4} C_{bi_5 i_6 c} \quad \overset{8}{Q}_{i_1 \dots i_6}^{(3)\gamma} = C_{.i_1 i_2}^a \overset{b}{\nabla}_{aci_3}^{\gamma} \overset{c}{\nabla}_{i_4} C_{bi_5 i_6}^c
\end{aligned}$$

$$\begin{aligned}
Q_{i_1 \dots i_6}^{(3)\gamma} &= C_{.i_1 i_2.}^a {}^b C_{cab i_3} {}^c {}_{\nabla i_4} C_{.i_5 i_6.}^{\gamma} & Q_{i_1 \dots i_6}^{(3)\gamma} &= C_{abc i_1} C_{.i_2 i_3.}^{\gamma} {}^a {}_{\nabla i_4} C_{.i_5 i_6.}^b \\
Q_{i_1 \dots i_6}^{(3)\gamma} &= \delta_{i_1}^{\gamma} C_{.i_2 i_3.}^a {}^b C_{acbd} {}^c {}_{\nabla i_4} C_{.i_5 i_6.}^d & Q_{i_1 \dots i_6}^{(3)\gamma} &= \delta_{i_1}^{\gamma} C_{.i_2 i_3.}^a {}^b C_{.i_4 i_5.}^c {}^d {}_{\nabla c} C_{dabi_6} \\
Q_{i_1 \dots i_6}^{(3)\gamma} &= C_{.i_1 i_2.}^a {}^b C_{ai_3 i_4} {}^c {}_{\nabla i_5} C_{.b. i_6}^{\gamma c} & Q_{i_1 \dots i_6}^{(3)\gamma} &= C_{.i_1 i_2.}^a {}^b C_{ai_3 i_4} {}^c {}_{\nabla b} C_{.i_5 i_6.}^{\gamma c} \\
Q_{i_1 \dots i_6}^{(3)\gamma} &= C_{.i_1 i_2.}^a {}^b C_{.i_3 i_4.}^{\gamma} {}^c {}_{\nabla c} C_{ai_5 i_6 b} & Q_{i_1 \dots i_6}^{(3)\gamma} &= C_{.i_1 i_2.}^a {}^b C_{.i_3 i_4.}^{\gamma} {}^c {}_{\nabla a} C_{bi_5 i_6 c} \\
Q_{i_1 \dots i_6}^{(3)\gamma} &= \delta_{i_1}^{\gamma} C_{.i_2 i_3.}^a {}^b C_{a..i_4} {}^c {}_{\nabla b} C_{ci_5 i_6 d} & Q_{i_1 \dots i_6}^{(3)\gamma} &= \delta_{i_1}^{\gamma} C_{.i_2 i_3.}^a {}^b C_{a..i_4} {}^c {}_{\nabla i_5} C_{bcd i_6} \\
Q_{i_1 \dots i_6}^{(3)\gamma} &= \delta_{i_1}^{\gamma} C_{.i_2 i_3.}^a {}^b C_{a..i_4} {}^c {}_{\nabla i_5} C_{bcd i_6}. & & (3.7)
\end{aligned}$$

Using (3.1), (3.2) and the transformation formulae (2.4), (2.5), (2.8), we obtain

$$\begin{bmatrix} X^{\gamma}(C_{i_1 \dots i_6}^{(3)}) \\ \vdots \\ X^{\gamma}(C_{i_1 \dots i_6}^{27}) \end{bmatrix}_* = A_6^{(3)}(-1, n) \begin{bmatrix} Q_{i_1 \dots i_6}^{(3)\gamma} \\ \vdots \\ Q_{i_1 \dots i_6}^{19} \end{bmatrix}.$$

where

$$A_6^{(3)}(-1, n) =$$

$$\left[\begin{array}{ccccccccccccccccccccccccc}
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -4 & 0 & -10 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & -10 & 0 & 0 & 0 & 4 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & -10 & -10 & 0 & 0 & 0 & 4 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & -2 & 0 & -2 & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -8 \\
0 & 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & -1 & -3 & 0 & 0 & 0 & -1 & 0 & 0 & -5 & -1 & -1 & -1 & -1 & -1 & -1 & -1 & -1 \\
12/a & 0 & -8/a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 12 & 0 & 0 & 8 & -8 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -4 & 0 & 4 & -6 & -4 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
-4 & 0 & 1 & 0 \\
0 & -5 & 0 & 1 & 0 & 0 & 0 & 0 & a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
1 & 0 & -4 & 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -4 & -8 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & -1/a & 0 & -3 & 0 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 1 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & -4 & 0 \\
0 & 0 & -1/a & 0 & -3 & -3 & 0 & 0 & -1 & 1 & 0 & 0 & 0 & 0 & 0 & 1 & -1 & 0 & -4 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & -2 & -1 & 2 & 1 & -2 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & -4 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 2/a & 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 & -2 & 0 & -2 & 2 & 0 & 0 & 0 & -6 \\
0 & 0 & 2/a & 0 & 0 & 0 & -2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -2 & 2 & 0 & -6 & 0 \\
0 & 0 & 3/a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 1 & 1 & -4 & 3 & -3 & -1 & -1 \\
0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & -4 & -4 & 4 & -4 & 0 & 0 \\
8/a & 0 & 2/a & 0 & 0 & 0 & 0 & 2 & 2 & -2 & 8 & 0 & 0 & -2 & 2 & 0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0 & 0 & 0 & -1 & 1 & -1 & 3 & -1 & 4 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\
0 & -4 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & -3 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & a & -a & 0 & 0 & 0 \\
0 & 1 & -2 & -1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & a & 0 & 0 & 0 \\
-4 & 0 & 0 & 0 & 0 & 0 & a & a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
-4 & 0 & -1 & 0 & 0 & 0 & a & a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\
0 & 0 & 2a & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0 & 0
\end{array} \right]$$

with $a = n - 3$.

It is rank $A_6^{(3)}(-1, n) = 17$ if $n > 4$. For $n = 4$ there are more identities concerning the tensors (3.6) and (3.7) (see [Thi]). So we get

PROPOSITION 3.2. – (i) If $n > 4$, then $S_6^{(3)}(-1, n)$ is generated by the tensors

$$\begin{aligned} S_6^{(3)}(-1, n) = & -346C_{i_1 \dots i_6}^{(3)} + 36C_{i_1 \dots i_6}^{(3)} - 234C_{i_1 \dots i_6}^{(3)} + 140C_{i_1 \dots i_6}^{(3)} \\ & - 310C_{i_1 \dots i_6}^{(3)} - 240C_{i_1 \dots i_6}^{(3)} - 56C_{i_1 \dots i_6}^{(3)} + 52C_{i_1 \dots i_6}^{(3)} - \frac{280}{n-3}C_{i_1 \dots i_6}^{(3)} \\ & + \frac{448}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{448}{n-3}C_{i_1 \dots i_6}^{(3)} + 405C_{i_1 \dots i_6}^{(3)} - 220C_{i_1 \dots i_6}^{(3)} \\ & + 680C_{i_1 \dots i_6}^{(3)} + 300C_{i_1 \dots i_6}^{(3)} + 220C_{i_1 \dots i_6}^{(3)} \\ & - \frac{560}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{1030}{(n-3)^2}C_{i_1 \dots i_6}^{(3)} \\ S_6^{(3)}(-1, n) = & -64C_{i_1 \dots i_6}^{(3)} - 4C_{i_1 \dots i_6}^{(3)} - 40C_{i_1 \dots i_6}^{(3)} + 4C_{i_1 \dots i_6}^{(3)} - 56C_{i_1 \dots i_6}^{(3)} \\ & - 32C_{i_1 \dots i_6}^{(3)} - 12C_{i_1 \dots i_6}^{(3)} + 4C_{i_1 \dots i_6}^{(3)} - \frac{8}{n-3}C_{i_1 \dots i_6}^{(3)} \\ & + \frac{48}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{48}{n-3}C_{i_1 \dots i_6}^{(3)} + 75C_{i_1 \dots i_6}^{(3)} + 120C_{i_1 \dots i_6}^{(3)} \\ & + 40C_{i_1 \dots i_6}^{(3)} + 20C_{i_1 \dots i_6}^{(3)} - \frac{60}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{130}{(n-3)^2}C_{i_1 \dots i_6}^{(3)} \\ S_6^{(3)}(-1, n) = & 3(n-3)C_{i_1 \dots i_6}^{(3)} - 3(n-3)C_{i_1 \dots i_6}^{(3)} + 5(n-3)C_{i_1 \dots i_6}^{(3)} \\ & - 5(n-3)C_{i_1 \dots i_6}^{(3)} - 10C_{i_1 \dots i_6}^{(3)} - 10(n-3)C_{i_1 \dots i_6}^{(3)} \\ & + 10(n-3)C_{i_1 \dots i_6}^{(3)} + 10C_{i_1 \dots i_6}^{(3)} + \frac{10}{n-3}C_{i_1 \dots i_6}^{(3)} \\ S_6^{(3)}(-1, n) = & + 20(n-3)C_{i_1 \dots i_6}^{(3)} + 120(n-3)C_{i_1 \dots i_6}^{(3)} - 120(n-3)C_{i_1 \dots i_6}^{(3)} \\ & + 232(n-3)C_{i_1 \dots i_6}^{(3)} - 168(n-3)C_{i_1 \dots i_6}^{(3)} - 96(n-3)C_{i_1 \dots i_6}^{(3)} \\ & + 4(n-3)C_{i_1 \dots i_6}^{(3)} + 12(n-3)C_{i_1 \dots i_6}^{(3)} - 464C_{i_1 \dots i_6}^{(3)} \\ & + 144C_{i_1 \dots i_6}^{(3)} - 144C_{i_1 \dots i_6}^{(3)} - 25(n-3)C_{i_1 \dots i_6}^{(3)} \\ & - 440(n-3)C_{i_1 \dots i_6}^{(3)} + 360(n-3)C_{i_1 \dots i_6}^{(3)} \\ & + 120(n-3)C_{i_1 \dots i_6}^{(3)} - 180C_{i_1 \dots i_6}^{(3)} + 440C_{i_1 \dots i_6}^{(3)} - \frac{170}{n-3}C_{i_1 \dots i_6}^{(3)} \end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, n) = & -596C_{i_1 \dots i_6}^{(3)} - 78C_{i_1 \dots i_6}^{(3)} - 186C_{i_1 \dots i_6}^{(3)} - 76C_{i_1 \dots i_6}^{(3)} \\
& - 256C_{i_1 \dots i_6}^{(3)} - 272C_{i_1 \dots i_6}^{(3)} + 4C_{i_1 \dots i_6}^{(3)} + 12C_{i_1 \dots i_6}^{(3)} \\
& + \frac{152}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{560}{n-3}C_{i_1 \dots i_6}^{(3)} + \frac{560}{n-3}C_{i_1 \dots i_6}^{(3)} \\
& + 745C_{i_1 \dots i_6}^{(3)} + 220C_{i_1 \dots i_6}^{(3)} + 580C_{i_1 \dots i_6}^{(3)} + 120C_{i_1 \dots i_6}^{(3)} \\
& + 220C_{i_1 \dots i_6}^{(3)} + \frac{700}{n-3}C_{i_1 \dots i_6}^{(3)} + \frac{1810}{(n-3)^2}C_{i_1 \dots i_6}^{(3)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, n) = & -636C_{i_1 \dots i_6}^{(3)} - 98C_{i_1 \dots i_6}^{(3)} - 386C_{i_1 \dots i_6}^{(3)} - 100C_{i_1 \dots i_6}^{(3)} \\
& - 580C_{i_1 \dots i_6}^{(3)} - 520C_{i_1 \dots i_6}^{(3)} - 4C_{i_1 \dots i_6}^{(3)} - 12C_{i_1 \dots i_6}^{(3)} \\
& - \frac{240}{n-3}C_{i_1 \dots i_6}^{(3)} + \frac{912}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{912}{n-3}C_{i_1 \dots i_6}^{(3)} \\
& + 795C_{i_1 \dots i_6}^{(3)} + 220C_{i_1 \dots i_6}^{(3)} + 1180C_{i_1 \dots i_6}^{(3)} + 320C_{i_1 \dots i_6}^{(3)} \\
& + 440C_{i_1 \dots i_6}^{(3)} - \frac{1140}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{2250}{(n-3)^2}C_{i_1 \dots i_6}^{(3)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, n) = & -68C_{i_1 \dots i_6}^{(3)} - 12C_{i_1 \dots i_6}^{(3)} - 36C_{i_1 \dots i_6}^{(3)} + 12C_{i_1 \dots i_6}^{(3)} - 88C_{i_1 \dots i_6}^{(3)} \\
& - 96C_{i_1 \dots i_6}^{(3)} + 4C_{i_1 \dots i_6}^{(3)} + 12C_{i_1 \dots i_6}^{(3)} - \frac{24}{n-3}C_{i_1 \dots i_6}^{(3)} \\
& + \frac{144}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{144}{n-3}C_{i_1 \dots i_6}^{(3)} + 85C_{i_1 \dots i_6}^{(3)} + 80C_{i_1 \dots i_6}^{(3)} \\
& + 120C_{i_1 \dots i_6}^{(3)} + 80C_{i_1 \dots i_6}^{(3)} - \frac{180}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{390}{(n-3)^2}C_{i_1 \dots i_6}^{(3)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, n) = & -68C_{i_1 \dots i_6}^{(3)} + 72C_{i_1 \dots i_6}^{(3)} - 120C_{i_1 \dots i_6}^{(3)} + 92C_{i_1 \dots i_6}^{(3)} \\
& - 168C_{i_1 \dots i_6}^{(3)} - 96C_{i_1 \dots i_6}^{(3)} + 4C_{i_1 \dots i_6}^{(3)} + 12C_{i_1 \dots i_6}^{(3)} \\
& - \frac{24}{n-3}C_{i_1 \dots i_6}^{(3)} + \frac{144}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{144}{n-3}C_{i_1 \dots i_6}^{(3)} \\
& + 85C_{i_1 \dots i_6}^{(3)} - 280C_{i_1 \dots i_6}^{(3)} + 360C_{i_1 \dots i_6}^{(3)} + 120C_{i_1 \dots i_6}^{(3)} \\
& + 80C_{i_1 \dots i_6}^{(3)} - \frac{180}{n-3}C_{i_1 \dots i_6}^{(3)} - \frac{390}{(n-3)^2}C_{i_1 \dots i_6}^{(3)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, n) = & 20(n-3)C_{i_1 \dots i_6}^{(3)} - (n-3)C_{i_1 \dots i_6}^{(2)} - 10(n-3)C_{i_1 \dots i_6}^{(3)} \\
& + (n-3)C_{i_1 \dots i_6}^{(4)} - 14(n-3)C_{i_1 \dots i_6}^{(5)} - 8(n-3)C_{i_1 \dots i_6}^{(6)} \\
& + 4(n-3)C_{i_1 \dots i_6}^{(7)} + 12(n-3)C_{i_1 \dots i_6}^{(8)} - 2C_{i_1 \dots i_6}^{(9)} + 166C_{i_1 \dots i_6}^{(10)} \\
& - 56C_{i_1 \dots i_6}^{(11)} - 25(n-3)C_{i_1 \dots i_6}^{(12)} + 30(n-3)C_{i_1 \dots i_6}^{(14)} \\
& + 10(n-3)C_{i_1 \dots i_6}^{(15)} - 180(n-3)C_{i_1 \dots i_6}^{(22)} + 110C_{i_1 \dots i_6}^{(24)} - \frac{60}{n-3}C_{i_1 \dots i_6}^{(27)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, n) = & 332(n-3)C_{i_1 \dots i_6}^{(1)} + 12(n-3)C_{i_1 \dots i_6}^{(2)} + 120(n-3)C_{i_1 \dots i_6}^{(3)} \\
& - 12(n-3)C_{i_1 \dots i_6}^{(4)} + 168(n-3)C_{i_1 \dots i_6}^{(5)} + 96(n-3)C_{i_1 \dots i_6}^{(6)} \\
& - 4(n-3)C_{i_1 \dots i_6}^{(7)} - 12(n-3)C_{i_1 \dots i_6}^{(8)} + 24C_{i_1 \dots i_6}^{(9)} \\
& + 296C_{i_1 \dots i_6}^{(10)} + 144C_{i_1 \dots i_6}^{(11)} - 415(n-3)C_{i_1 \dots i_6}^{(12)} \\
& - 360(n-3)C_{i_1 \dots i_6}^{(14)} - 120(n-3)C_{i_1 \dots i_6}^{(15)} - 700(n-3)C_{i_1 \dots i_6}^{(22)} \\
& + 440C_{i_1 \dots i_6}^{(23)} - \frac{270}{n-3}C_{i_1 \dots i_6}^{(27)}.
\end{aligned}$$

(ii) $S_6^{(3)}(-1, 4)$ is generated by the tensors

$$\begin{aligned}
S_6^{(3)}(-1, 4) = & 8C_{i_1 \dots i_6}^{(2)} - 7C_{i_1 \dots i_6}^{(16)} + 104C_{i_1 \dots i_6}^{(23)} + 10C_{i_1 \dots i_6}^{(24)} - 218C_{i_1 \dots i_6}^{(25)} \\
& + 72C_{i_1 \dots i_6}^{(27)} - 10C_{i_1 \dots i_6}^{(28)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, 4) = & 4C_{i_1 \dots i_6}^{(5)} - 5C_{i_1 \dots i_6}^{(16)} - 40C_{i_1 \dots i_6}^{(23)} - 10C_{i_1 \dots i_6}^{(24)} + 90C_{i_1 \dots i_6}^{(25)} \\
& - 40C_{i_1 \dots i_6}^{(27)} + 10C_{i_1 \dots i_6}^{(28)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, 4) = & 32C_{i_1 \dots i_6}^{(9)} - 3C_{i_1 \dots i_6}^{(16)} - 24C_{i_1 \dots i_6}^{(23)} - 14C_{i_1 \dots i_6}^{(24)} + 30C_{i_1 \dots i_6}^{(25)} \\
& - 56C_{i_1 \dots i_6}^{(27)} + 14C_{i_1 \dots i_6}^{(28)}
\end{aligned}$$

$$\begin{aligned}
S_6^{(3)}(-1, 4) = & 8C_{i_1 \dots i_6}^{(8)} - 7C_{i_1 \dots i_6}^{(16)} + 24C_{i_1 \dots i_6}^{(23)} + 82C_{i_1 \dots i_6}^{(24)} - 34C_{i_1 \dots i_6}^{(25)} \\
& + 184C_{i_1 \dots i_6}^{(27)} - 2C_{i_1 \dots i_6}^{(28)}
\end{aligned}$$

$$S_6^{(3)}(-1, 4) = C_{i_1 \dots i_6}^{14(3)} + 4C_{i_1 \dots i_6}^{23(3)} + 2C_{i_1 \dots i_6}^{24(3)} - 10C_{i_1 \dots i_6}^{25(3)} + 6C_{i_1 \dots i_6}^{27(3)} - 2C_{i_1 \dots i_6}^{28(3)}$$

$$S_6^{(3)}(-1, 4) = C_{.i_1 i_2.}^a {}^b C_{a i_3 i_4 b} {}^c {}^c \nabla_c \nabla_u C_{.i_5 i_6.}^u {}^c.$$

We summarize our results in the following

THEOREM 3.1. – (i) If $n > 4$, $\mathcal{S}_6(-1, n)$ is generated by

$$\{S_6^{(2)}(-1, n), S_6^{(2)}(-1, n), S_6^{(3)}(-1, n), \dots, S_6^{(10)}(-1, n)\}.$$

(ii) $\mathcal{S}_6(-1, 4)$ is generated by ⁴

$$\{S_6^{(2)}(-1, 4), S_6^{(2)}(-1, 4), S_6^{(3)}(-1, 4), \dots, S_6^{(6)}(-1, 4)\}.$$

4. ADDITION CONFORMAL TENSORS FOR $n = 4$

In a four dimensional pseudo Riemannian manifold (M, g) with a smooth metric of the signature (+---) the *moments* of some conformally invariant field equations are symmetric, trace-free conformal tensors (see Section 6 and [Gü; Wü5]). In case of odd order these moments also depend on the Levi-Civita pseudo tensor e_{abcd} . Thus for a general construction of such conformal tensors we have to extend the tensor algebra \mathcal{R} (see 1.1) by the pseudo tensor e_{abcd} . Let \mathcal{R}^* be that tensor algebra which is generated by the tensors (1.1) and e_{abcd} . Then all the results of Section 2 remain valid with exception of Proposition 2.2 (γ) where the tensors (2.7) have to be extended by e_{abcd} .

The Bianchi identity for the dual Weyl tensor

$${}^*C_{abcd} := \frac{1}{2} e_{ab}{}^{ef} C_{efcd} \quad (4.1)$$

has the form

$$\nabla_{[a} {}^*C_{bc]ij} + g_{j[a|} \nabla_u {}^*C_{.i|bc]}^u - g_{i[a|} \nabla_u {}^*C_{.j|bc]}^u = 0. \quad (4.2)$$

Let $\mathcal{S}_r^*(\omega)$ be the set of all nontrivial conformal, symmetric, trace-free tensors of \mathcal{R}^* with weight ω and covariant rank r . Let $\mathcal{S}_r^{*(\alpha)}(\omega)$ be the subset of those elements of $\mathcal{S}_r^*(\omega)$ whose monomials have the order α .

⁴ This generating system for $\mathcal{S}_6(-1, 4)$ is the same as the one derived in [GeW2] and has already been used in [Wü5], in detail $S_6^{k(\nu)}(-1, 4)$ is equal to $S_6^{(\nu, k)}(-1, 4)$ of [Wü5], where $(k, \nu) \in \{(2, 1), (2, 2), (3, 1), (3, 2), (3, 3), (3, 4), (3, 5), (3, 6)\}$.

In [Wü1] it was proved

PROPOSITION 4.1. – *The sets $\mathcal{S}_1^*(-1)$, $\mathcal{S}_3^*(-1)$ are empty.*

Now we consider *the case $r = 5$:*

Because of $\overset{c}{\nabla}_a \overset{c}{\nabla}_b {}^*C_{.i_1 i_2.}^{a b} = 0$ (see [KNT]) and Lemma (2.3) $\mathcal{S}_r^*(1)(-1)$ is empty and we have

$$\mathcal{S}_5^*(-1) = \mathcal{S}_5^{*(2)}(-1) \cup \mathcal{S}_5^{*(3)}(-1). \quad (4.3)$$

Firstly, let $\alpha = 2$: Then on account of (2.13) it is $q = 3$.

Under consideration of (4.2),

$$\overset{c}{\nabla}_{i_1} \dots \overset{c}{\nabla}_{i_{r-1}} {}^*C_{...i_r}^{abc} \overset{c}{\nabla}_{i_{r+1}} \dots \overset{c}{\nabla}_{i_{s-1}} C_{abci_5} = 0 \quad (r, s = 1, 2, 3, \dots) \quad (4.4)$$

and of some more identities with respect to the derivatives of $C...$ and ${}^*C...$ which one can show by means of the spinor formalism (see [Thi; GeW1]) a tensor of $\mathcal{S}_5^{*(2)}(-1)$ has to be a linear combination of the following linearly independent³⁾ monomials ([GeW1; Ge])

$$\overset{1}{C}_{i_1 \dots i_5}^{(2)} = TS[C_{.i_1 i_2.}^a \overset{b}{\nabla}_{i_3} \overset{c}{\nabla}_{i_4} \overset{c}{\nabla}_u C_{.abi_5}^u]$$

$$\overset{2}{C}_{i_1 \dots i_5}^{(2)} = TS[C_{.i_1 i_2.}^a \overset{b}{\nabla}_{i_3} \overset{c}{\nabla}_a \overset{c}{\nabla}_u C_{.i_4 i_5 b}^u]$$

$$\overset{3}{C}_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_{i_1} C_{.i_2 i_3.}^a \overset{b}{\nabla}_{i_4} \overset{c}{\nabla}_u C_{.abi_5}^u]$$

$$\overset{4}{C}_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_{i_1} C_{.i_2 i_3.}^a \overset{b}{\nabla}_a \overset{c}{\nabla}_u C_{.i_4 i_5 b}^u]$$

$$\overset{5}{C}_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_u C_{.abi_1}^u \overset{c}{\nabla}_{i_2} \overset{c}{\nabla}_{i_3} C_{.i_4 i_5.}^{a b}]$$

$$\overset{6}{C}_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_u C_{.i_1 i_2 a}^u \overset{c}{\nabla}_{i_3} \overset{c}{\nabla}_v C_{.i_4 i_5.}^{a b}]$$

$$\overset{7}{C}_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_d C_{.i_1 i_2.}^a \overset{b}{\nabla}_{i_3} \overset{c}{\nabla}_d C_{a i_4 i_5 b}^u]$$

$${}^*C_{i_1 \dots i_5}^{(2)} = TS[{}^*C_{.i_1 i_2.}^a \overset{b}{\nabla}_{i_3} \overset{c}{\nabla}_{i_4} \overset{c}{\nabla}_u C_{.abi_5}^u]$$

$${}^*C_{i_1 \dots i_5}^{(2)} = TS[{}^*C_{.i_1 i_2.}^a \overset{b}{\nabla}_{i_3} \overset{c}{\nabla}_a \overset{c}{\nabla}_u C_{.i_4 i_5 b}^u]$$

$${}^*C_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_{i_1} {}^*C_{.i_2 i_3.}^a \overset{b}{\nabla}_{i_4} \overset{c}{\nabla}_u C_{.abi_5}^u]$$

$${}^*C_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_u {}^*C_{.abi_1}^u \overset{c}{\nabla}_{i_2} \overset{c}{\nabla}_{i_3} C_{.i_4 i_5.}^{a b}]$$

$${}^*C_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_u {}^*C_{.i_1 i_2 a}^u \overset{c}{\nabla}_{i_3} \overset{c}{\nabla}_v C_{.i_4 i_5.}^{a b}]$$

$${}^*C_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_d {}^*C_{.i_1 i_2.}^a \overset{b}{\nabla}_{i_3} \overset{c}{\nabla}_d C_{a i_4 i_5 b}^u]$$

$${}^*C_{i_1 \dots i_5}^{(2)} = TS[\overset{c}{\nabla}_d {}^*C_{.i_1 i_2.}^a \overset{b}{\nabla}_{i_3} \overset{c}{\nabla}_d C_{a i_4 i_5 b}^u]$$

We put

$$\begin{aligned}
 Q_{i_1 \dots i_5}^{(2)\gamma} &:= \delta_{i_1}^\gamma C_{.i_2 i_3.}^a b \nabla_{i_4}^c \nabla_u^c C_{.abi_5}^u & Q_{i_1 \dots i_5}^{(2)\gamma} &:= \delta_{i_1}^\gamma C_{.i_2 i_3.}^a b \nabla_a^c \nabla_u^c C_{.i_4 i_5 b}^u \\
 Q_{i_1 \dots i_5}^{(2)\gamma} &:= C_{.i_1 i_2.}^a b \nabla_{i_3}^c \nabla_a^c C_{.i_4 i_5 b}^\gamma & Q_{i_1 \dots i_5}^{(2)\gamma} &:= C_{.abi_1}^\gamma c \nabla_{i_2}^c \nabla_{i_3}^c C_{.i_4 i_5.}^a \\
 Q_{i_1 \dots i_5}^{(2)\gamma} &:= C_{.i_1 i_2.}^\gamma a \nabla_{i_3}^c \nabla_u^c C_{.i_4 i_5 a}^c & Q_{i_1 \dots i_5}^{(2)\gamma} &:= \delta_{i_1}^\gamma c d C_{.i_2 i_3.}^a b \nabla_d^c C_{.a i_4 i_5 b}^c \\
 Q_{i_1 \dots i_5}^{(2)\gamma} &:= \delta_{i_1}^\gamma c \nabla_{i_2}^c C_{.i_3 i_4.}^a b \nabla_u^c C_{.abi_5}^u & Q_{i_1 \dots i_5}^{(2)\gamma} &:= \delta_{i_1}^\gamma c \nabla_u^c C_{.i_2 i_3 a}^u c \nabla_v^v C_{.i_4 i_5.}^a \\
 Q_{i_1 \dots i_5}^{(2)\gamma} &:= c \nabla_{i_1}^c C_{.i_2 i_3.}^a b \nabla_{i_4}^c C_{.abi_5}^\gamma & Q_{i_1 \dots i_5}^{(2)\gamma} &:= c \nabla_{i_1}^c C_{.i_2 i_3.}^a b \nabla_a^c C_{.i_4 i_5 b}^\gamma \\
 Q_{i_1 \dots i_5}^{(2)\gamma} &:= c \nabla_{i_1}^c C_{.i_2 i_3.}^\gamma a \nabla_u^c C_{.i_4 i_5 a}^u
 \end{aligned}$$

$$\begin{aligned}
 *Q_{i_1 \dots i_5}^{(2)\gamma} &= \delta_{i_1}^\gamma *C_{.i_2 i_3.}^a b \nabla_{i_4}^c \nabla_u^c C_{.abi_5}^u & *Q_{i_1 \dots i_5}^{(2)\gamma} &= \delta_{i_1}^\gamma *C_{.i_2 i_3.}^a b \nabla_a^c \nabla_u^c C_{.i_4 i_5 b}^u \\
 *Q_{i_1 \dots i_5}^{(2)\gamma} &= *C_{.i_1 i_2.}^a b \nabla_{i_3}^c \nabla_a^c C_{.i_4 i_5 b}^\gamma & *Q_{i_1 \dots i_5}^{(2)\gamma} &= *C_{.i_1 i_2.}^a b \nabla_{i_3}^c \nabla_{i_4}^c C_{.abi_5}^\gamma \\
 *Q_{i_1 \dots i_5}^{(2)\gamma} &= *C_{.i_1 i_2.}^\gamma a \nabla_{i_3}^c \nabla_u^c C_{.i_4 i_5 a}^u & *Q_{i_1 \dots i_5}^{(2)\gamma} &= \delta_{i_1}^\gamma c \nabla_{i_2}^c *C_{.i_3 i_4.}^a b \nabla_u^c C_{.abi_5}^u \\
 *Q_{i_1 \dots i_5}^{(2)\gamma} &= c \nabla_{i_1}^c *C_{.i_2 i_3.}^a b \nabla_{i_4}^c C_{.i_4 i_5 b}^\gamma & *Q_{i_1 \dots i_5}^{(2)\gamma} &= c \nabla_{i_1}^c *C_{.i_2 i_3.}^\gamma a \nabla_u^c C_{.i_4 i_5 a}^u.
 \end{aligned}$$

Using the transformation formulae (2.4), (2.5), (2.8), one gets

$$\begin{bmatrix} X^\gamma(C_{i_1 \dots i_5}^{(2)}) \\ \vdots \\ X^\gamma(C_{i_1 \dots i_5}^{(7)}) \\ X^\gamma(*C_{i_1 \dots i_5}^{(1)}) \\ \vdots \\ \vdots \\ X^\gamma(*C_{i_1 \dots i_5}^{(7)})^\gamma \end{bmatrix} = A_5^{*(2)}(-1) \begin{bmatrix} Q_{i_1 \dots i_5}^{(2)\gamma} \\ \vdots \\ Q_{i_1 \dots i_5}^{(11)\gamma} \\ *Q_{i_1 \dots i_5}^{(1)\gamma} \\ \vdots \\ \vdots \\ *Q_{i_1 \dots i_5}^{(8)\gamma} \end{bmatrix}$$

with

$$A_5^{*(2)}(-1) = \begin{bmatrix} A & 0 \\ 0 & B \end{bmatrix}$$

and

$$A = \begin{bmatrix} -10 & 0 & 0 & 1 & 2 & 0 & 0 & 0 & 0 & 0 & 0 \\ 1 & -6 & 1 & 0 & -5 & 0 & 0 & 0 & 0 & 0 & 0 \\ -4 & 0 & 0 & 0 & -1 & 0 & -4 & 0 & 1 & 0 & 1 \\ 0 & -4 & 0 & 0 & 1 & 0 & 1 & 0 & 0 & 1 & -4 \\ 0 & 0 & 0 & 1 & 0 & 0 & -10 & 0 & 0 & 0 & -2 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & -5 & 0 & 0 & 1 \\ -4 & 0 & -4 & 0 & 0 & -5 & -4 & 0 & 0 & -4 & 0 \end{bmatrix},$$

$$B = \begin{bmatrix} -10 & 0 & 0 & 1 & 2 & 0 & 0 & 0 \\ 1 & -6 & 1 & 0 & -5 & 0 & 0 & 0 \\ -4 & 0 & 0 & 0 & -1 & -4 & 0 & 1 \\ 0 & -4 & 0 & 0 & 1 & 1 & 1 & -4 \\ 0 & 0 & 0 & 1 & 0 & 10 & 0 & 2 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & -1 \\ -4 & 0 & -4 & 0 & 0 & 6 & 6 & 2 \end{bmatrix}$$

Because of rank $A = 7$, rank $B = 6$ we obtain

PROPOSITION 4.2. $-S_5^{*(2)}(-1)$ is generated by the tensor

$$S_5^{(2)}(-1) = 4^*C_{i_1 \dots i_5}^{(2)} - 6^*C_{i_1 \dots i_5}^{(4)} + 26^*C_{i_1 \dots i_5}^{(6)} + {}^*C_{i_1 \dots i_5}^{(7)}. \quad (4.5)$$

Finally, in the case $\alpha = 3$ we have $q = 1$ because of (2.13). A tensor of $S_5^{*(3)}(-1)$ has to be a linear combination of the following linearly independent³⁾ monomials

$$C_{i_1 \dots i_5}^{(3)} = TS[C_{.i_1 i_2.}^a {}^b C_{a i_3 i_4.} {}^d {}^c \nabla_u C_{.b d i_5}^u]$$

$$C_{i_1 \dots i_5}^{(3)} = TS[C_{a..i_1}^c {}^d C_{b c d i_2} {}^c \nabla_{i_3} C_{.i_4 i_5.}^a {}^b]$$

$$C_{i_1 \dots i_5}^{(3)} = TS[C_{a..i_1}^c {}^d C_{b i_2 c d} {}^c \nabla_{i_3} C_{.i_4 i_5.}^a {}^b]$$

$${}^*C_{i_1 \dots i_5}^{(3)} = TS[{}^*C_{.i_1 i_2.}^a {}^b C_{a i_3 i_4.} {}^d {}^c \nabla_u C_{.b d i_5}^u]$$

$${}^*C_{i_1 \dots i_5}^{(3)} = TS[C_{a..i_1}^c {}^d C_{b c d i_2} {}^c \nabla_{i_3} {}^*C_{.i_4 i_5.}^a {}^b]$$

$${}^*C_{i_1 \dots i_5}^{(3)} = TS[{}^*C_{a..i_1}^c {}^d C_{b i_2 c d} {}^c \nabla_{i_3} C_{.i_4 i_5.}^a {}^b]$$

After using the transformation formulae (2.4), (2.5), (2.8) and consideration of further identities [Thi], we obtain

$$\begin{aligned} X^\gamma(C_{i_1 \dots i_5}^{(3)})_* &= C_{.i_1 i_2.}^a {}^b C_{ai_3 i_4.} {}^c C_{.bci_5}^\gamma =: Q_{i_1 \dots i_5}^{(3)\gamma} \\ X^\gamma(C_{i_1 \dots i_5}^{(3)})_* &= -12 Q_{i_1 \dots i_5}^{(3)\gamma} \\ X^\gamma(C_{i_1 \dots i_5}^{(3)})_* &= 8 Q_{i_1 \dots i_5}^{(3)\gamma} \\ X^\gamma({}^*C_{i_1 \dots i_5}^{(3)}) &= {}^*C_{.i_1 i_2.}^a {}^b C_{ai_3 i_4.} {}^c C_{.bci_5}^\gamma =: {}^*Q_{i_1 \dots i_5}^{(3)\gamma} \\ X^\gamma({}^*C_{i_1 \dots i_5}^{(3)})_* &= -12 {}^*Q_{i_1 \dots i_5}^{(3)\gamma} \\ X^\gamma({}^*C_{i_1 \dots i_5}^{(3)})_* &= 8 {}^*Q_{i_1 \dots i_5}^{(3)\gamma}. \end{aligned}$$

Thus we get

PROPOSITION 4.3. – $\mathcal{S}_5^{*(3)}(-1)$ is generated by the tensors

$$S_5^{(3)}(-1) = 12 {}^*C_{i_1 \dots i_5}^{(3)} + {}^*C_{i_1 \dots i_5}^{(3)} \quad (4.6)$$

$$S_5^{(3)}(-1) = 8 {}^*C_{i_1 \dots i_5}^{(3)} - {}^*C_{i_1 \dots i_5}^{(3)} \quad (4.7)$$

$$S_5^{(3)}(-1) = 12 C_{i_1 \dots i_5}^{(3)} + C_{i_1 \dots i_5}^{(3)} \quad (4.8)$$

$$S_5^{(3)}(-1) = 8 C_{i_1 \dots i_5}^{(3)} - C_{i_1 \dots i_5}^{(3)}. \quad (4.9)$$

It follows from the Propositions 4.1, 4.2, 4.3.

THEOREM 4.1. – $\mathcal{S}_5^*(-1)$ is generated by the tensors⁵

$$S_5^{(2)}(-1), S_5^{(3)}(-1), \dots, S_5^{(4)}(-1).$$

⁵ This generating system for $\mathcal{S}_5^*(-1)$ is the same as the one derived in [GeW1] and has already been used in [Wü5], in detail we have

$$\{S_5^{(2)}(-1), S_5^{(3)}(-1), \dots, S_5^{(4)}(-1)\} \equiv \{S_5^{(2)}(-1), S_5^{(3,1)}(-1), \dots, S_5^{(3,4)}(-1)\}_{[Wü5]}.$$

5. THE CASE $n = 3$

If $n = 3$, the Weyl tensor vanishes. However, the tensor

$$S_{abc} := \overset{c}{\nabla}_{[c} L_{b]a} \quad (5.1)$$

is a conformal tensor of the weight 0 [Sz]. It is $S_{abc} = 0$ iff (M, g) is locally conformal flat. One can show easily the following properties of S_{abc} [Scho,Ge]:

$$S_{.ab}^a = S_{.ba}^a = S_{ab.}^b = 0 \quad (5.2)$$

$$\nabla_{[i} S_{a|bc]} = 0, \quad \nabla^u S_{uab} = 0, \quad \nabla^u S_{[ab]u} = 0 \quad (5.3)$$

$$S_{i[ab} g_{c]j} g_{d]k} = 0. \quad (5.4)$$

LEMMA 5.1. – For the tensor S_{abc} it holds the identity

$$T_{.abcd}^j := \delta_{[a}^j S_{b]cd} - g_{b[c} S_{a|d].}^j + g_{a[c} S_{b|d].}^j \equiv 0. \quad (5.5)$$

Proof. – The tensor $T_{.abcd}^j$ satisfies the conditions

$$T_{.abcd}^j = -T_{.bacd}^j = -T_{.abdc}^j$$

and $T_{.akc.}^j k = 0$. Consequently we have $T_{.abcd}^j = 0$. (see [Lo]).

COROLLARY 5.1. – It is

$$\nabla_{[a} S_{b]cd} = g_{b[c} \nabla_u S_{a|d].}^u - g_{a[c} \nabla_u S_{b|d].}^u. \quad (5.6)$$

All results of Section 2 with exception of Proposition 2.2 where one has to replace the tensors (2.7) by

$$g^{ab}, \quad g_{ab}, \quad \overset{c}{\nabla}_{i_1} \dots \overset{c}{\nabla}_{i_r} S_{i_{r+1}ab} \quad r = 0, 1, 2, \dots \quad (5.7)$$

remain valid. In particular we have

$$X^\gamma(S_{abc}) = 0, \quad \overset{c}{\nabla}_i S_{abc} = \nabla_i S_{abc} \quad (5.8)$$

$$X^\gamma(\overset{c}{\nabla}_i S_{abc}) = -3\delta_i^\gamma S_{abc} - \delta_a^\gamma S_{ibc} - \delta_b^\gamma S_{aic} - \delta_c^\gamma S_{abi} + g_{ia} S_{.bc}^\gamma + g_{ib} S_{a.c}^\gamma + g_{ic} S_{ab.}^\gamma \quad (5.9)$$

$$X^\gamma(\overset{c}{\nabla}^u S_{abu}) = -2S_{(ab)}^\gamma. \quad (5.10)$$

$$\begin{aligned} X^\gamma(\overset{c}{\nabla}_i \overset{c}{\nabla}_j S_{abc}) &= -8\delta_{(i}^\gamma \nabla_{j)} S_{abc} + g_{ij} \nabla^\gamma S_{abc} - 2\delta_{(i}^\gamma \nabla_{(j} S_{j)b} c \\ &\quad - 2\delta_b^\gamma \nabla_{(i} S_{a|j)c} - 2\delta_c^\gamma \nabla_{(i} S_{ab|j)} \\ &\quad + 2g_{a(i} \nabla_{j)} S_{.bc}^\gamma + g_{b(i} \nabla_{j)} S_{a.c}^\gamma + g_{c(i} \nabla_{j)} S_{ab.}^\gamma; \end{aligned} \quad (5.11)$$

The tensors of $\mathcal{S}_r^{(\alpha)}(\omega, 3)$ are representable as a sum of monomials containing α factors $\{\overset{c}{\nabla}_{i_1} \dots \overset{c}{\nabla}_{i_k} S_{i_{k+1}ab}\}$ and $q := r - 2\omega - 2\alpha - 2$ operators $\overset{c}{\nabla}$ (with respect to S_{abc}).

In the following we construct generating systems for $\mathcal{S}_r^{(2)}(\omega, 3)$ in the cases $(r, \omega) = (0, -4), (1, -3), (2, -3), (3, -2), (4, -2)$.

Firstly, let be $r = 0$:

Because of (5.2), ..., (5.4) and (5.6) a tensor of $\mathcal{S}_0^{(2)}(-4, 3)$ has to be a linear combination of the following monomials:

$$\overset{1}{S} = \nabla_a S_{bcd} \nabla^a S^{bcd} \dots$$

$$\overset{2}{S} = \nabla_a S_{bcd} \nabla^b S^{acd} \dots$$

$$\overset{3}{S} = S_{\dots}^{abc} \overset{c}{\nabla} d \overset{c}{\nabla} d S_{abc}.$$

Using the transformation formulae (5.8), ..., (5.10) and the identities (5.2), ..., (5.4), (5.6) one obtains

$$\begin{bmatrix} X^\gamma(\overset{1}{S}) \\ X^\gamma(\overset{2}{S}) \\ X^\gamma(\overset{3}{S}) \end{bmatrix} = \begin{bmatrix} -6 & -4 \\ -2 & -8 \\ -5 & 4 \end{bmatrix} \begin{bmatrix} \overset{1}{T}^\gamma \\ \overset{2}{T}^\gamma \end{bmatrix}.$$

with the transformation terms

$$\overset{1}{T}^\gamma = S_{abc} \nabla^\gamma S^{abc}$$

$$\overset{2}{T}^\gamma = S_{abc} \nabla^a S^{\gamma bc}.$$

The transformation matrix has the rank 2 from where we get

PROPOSITION 5.1. — $\mathcal{S}_0^{(2)}(-4, 3)$ is generated by the tensor

$$S_0^{(2)}(-4, 3) = -12\overset{1}{S} + 11\overset{2}{S} + 10\overset{3}{S}.$$

Now, let be $r = 2$:

Considering (5.2), ..., (5.6) the following linear independent terms are left

$$\begin{aligned} \overset{1}{S}_{i_1 i_2} &= TS[\nabla_{i_1} S_{\dots}^{abc} \nabla_{i_2} S_{abc}] & \overset{2}{S}_{i_1 i_2} &= TS[\nabla^a S_{i_1 i_2} \overset{b}{\nabla} u S_{abu}] \\ \overset{3}{S}_{i_1 i_2} &= TS[\nabla_{i_1} S_{\dots}^{abc} \nabla_a S_{i_2 bc}] & \overset{4}{S}_{i_1 i_2} &= TS[S_{\dots}^{abc} \overset{c}{\nabla}_{i_1} \overset{c}{\nabla}_{i_2} S_{abc}] \\ \overset{5}{S}_{i_1 i_2} &= TS[S_{\dots}^{abc} \overset{c}{\nabla}_{i_1} \overset{c}{\nabla}_a S_{i_2 bc}] & \overset{6}{S}_{i_1 i_2} &= TS[S_{i_1 \dots}^{bc} \overset{c}{\nabla} a \overset{c}{\nabla} a S_{i_2 bc}]. \end{aligned}$$

With due regard to the identities (5.2), ..., (5.6), to those following out of them and the transformation formulae (5.8), ..., (5.11) we obtain after a

conformal transformation

$$\begin{bmatrix} X^\gamma(S_{i_1 i_2}) \\ \vdots \\ X^\gamma(S_{i_1 i_2}) \end{bmatrix} = * \begin{bmatrix} -2 & -4 & -8 & 0 & 0 & 0 \\ -\frac{1}{2} & \frac{5}{2} & 1 & -\frac{3}{2} & -1 & -4 \\ 1 & -3 & -6 & -3 & -2 & 0 \\ -12 & 4 & 8 & 0 & 0 & 0 \\ -6 & 0 & 12 & 3 & -6 & 0 \\ 0 & 2 & 0 & -5 & 2 & 4 \end{bmatrix} \begin{bmatrix} T_{i_1 i_2}^\gamma \\ \vdots \\ T_{i_1 i_2}^\gamma \end{bmatrix}$$

with the transformation terms

$$\begin{aligned} T_{i_1 i_2}^1 &= \delta_{i_1}^\gamma S_{abc} \nabla_{i_2} S^{abc} & T_{i_1 i_2}^2 &= S_{i_1 ab} \nabla_{i_2} S^{\gamma ab} \\ T_{i_1 i_2}^3 &= S_{abi_1} \nabla_{i_2} S^{ab\gamma} & T_{i_1 i_2}^4 &= S_{i_1 ab} \nabla^\gamma S_{i_2}^{ab} \\ T_{i_1 i_2}^5 &= \delta_{i_1}^\gamma S_{abc} \nabla^a S_{i_2}^{bc} & T_{i_1 i_2}^6 &= S_{ab} \nabla^a S_{i_1 i_2}^b. \end{aligned}$$

The rank of the transformation matrix is 5 from where we get

PROPOSITION 5.2. – $\mathcal{S}_2^{(2)}(-3, 3)$ is generated by the tensor

$$\mathcal{S}_2^{(2)}(-3, 3) = 18S_{i_1 i_2}^1 + 12S_{i_1 i_2}^2 - 18S_{i_1 i_2}^3 - 9S_{i_1 i_2}^4 + 8S_{i_1 i_2}^5 + 12S_{i_1 i_2}^6.$$

In an analogous manner we obtain

PROPOSITION 5.3. – $\mathcal{S}_4^{(2)}(-2, 3)$ is generated by the tensor

$$\begin{aligned} S_4^1(-2, 3) &= -8S_{i_1 \dots i_4}^1 - 2S_{i_1 \dots i_4}^2 - 5S_{i_1 \dots i_4}^3 + 2S_{i_1 \dots i_4}^4 + 4S_{i_1 \dots i_4}^6 \\ S_4^2(-2, 3) &= -12S_{i_1 \dots i_4}^1 - 7S_{i_1 \dots i_4}^2 - 12S_{i_1 \dots i_4}^3 + 2S_{i_1 \dots i_4}^5 + 12S_{i_1 \dots i_4}^6, \end{aligned}$$

where

$$\begin{aligned} S_{i_1 \dots i_4}^1 &= TS[\nabla_{i_1} S_{i_2 i_3}^a \nabla^u S_{i_4 au}] & S_{i_1 i_4}^2 &= TS[\nabla_{i_1} S_{i_2}^{ab} \nabla_{i_3} S_{i_4 ab}] \\ S_{i_1 \dots i_4}^3 &= TS[\nabla^a S_{i_1 i_2}^b \nabla_a S_{i_3 i_4 b}] & S_{i_1 \dots i_4}^4 &= TS[S_{i_1 i_2}^a \nabla_{i_3}^c \nabla_u^b S_{i_4 au}] \\ S_{i_1 \dots i_4}^5 &= TS[S_{i_1}^{ab} \nabla_{i_2}^c \nabla_{i_3}^d S_{i_4 ab}] & S_{i_1 \dots i_4}^6 &= TS[S_{i_1 i_2}^a \nabla^b \nabla^c S_{i_3 i_4 a}]. \end{aligned}$$

PROPOSITION 5.4. – The sets $\mathcal{S}_1^{(2)}(-3, 3)$ and $\mathcal{S}_3^{(2)}(-2, 3)$ are empty.

Remark 5.1. – One can show the linear independence of the in this section occurring tensors either by means of the spinor formalism for three dimensional manifolds [II] or with the help of suitably choosen test metrics.

6. APPLICATIONS ON SPACE-TIMES

1. Huygen's principle for conformally invariant field equations

Let (M, g) be a space-time, i.e. a 4-manifold with a smooth metric of Lorentzian signature. The following conformally invariant field equations are considered ([Gü; Wü5; CMcL; McL]):

$$\text{Scalar wave equation} \quad g^{ab} \nabla_a \nabla_b u - \frac{1}{6} R u = 0 \quad E_1$$

$$\text{Maxwell's equations} \quad \nabla_{[a} F_{bc]} = 0, \quad \nabla_a F^a_{\cdot b} = 0 \quad E_2$$

$$\text{Weyl's neutrino equation} \quad \nabla^A \dot{\chi}_A = 0, \quad E_3$$

where $F_{ab} dx^a dx^b$ is the Maxwell 2-form, φ a valence 1-spinor and $\nabla_{A\dot{X}}$ the covariant derivative on spinors. For any of the equations E_1 - E_3 Huygens' principle (in the sense of Hadamard's "minor premise") is valid if and only if the tail term with respect to E_σ , $\sigma = 1, 2, 3$ vanishes ([Gü; Wü; CMcL]). Since the functional relationship between the tail terms and the metric is very complicated, the problem of the determination of all metrics for which any equation E_σ satisfies Huygens' principle is not yet completely solved (see [Gü; Wü1-5; CMcL; McL; McLS; AMLW; RW]).

The usual method for solving this problem is the derivation and the exploitation of the moment equations ([Gü; Wü5])

$$I_{i_1 \dots i_r}^\sigma = 0 \quad \sigma = 1, 2, 3 \quad r = 0, 1, 2, \dots, \quad \text{ME}_r^\sigma$$

where the moments $I_r^\sigma = I_{i_1 \dots i_r}^\sigma$ are symmetric, trace-free, conformally invariant tensors of the weight -1 . They are derived from the tail terms with respect to E_σ , $\sigma = 1, 2, 3$ by means of the conformal covariant derivative (2.6) ([Gü; Wü5]). If g is analytic, we have the following relationship between the moments and the validity of Huygens' principle: The equation E_σ , $\sigma = 1, 2, 3$ satisfies Huygens' principle if and only if all corresponding moments vanish on M ([Gü; Wü5]). Using the results on the theory of conformal tensors, in particular the results on generating systems of Sections 3 and 4, one obtains information about the general algebraic structure of the moments for $0 \leq r \leq 6$ ([Gü; Wü5]).

The following proposition was proved in ([Gü; Wü2,5]):

PROPOSITION 6.1. – *One has*

- (i) $I_r^\sigma \equiv 0$ if $(r, \sigma) \in \mathcal{M} := \{(k, 1) : k \text{ odd}\} \cup \{(m, \sigma) : m \in \{0, 1, 3\}, \sigma \in \{1, 2, 3\}\}$,
- (ii) $I_r^\sigma \in \mathcal{S}_r(-1, 4)$ if r is even,
- (iii) $I_r^\sigma \in \mathcal{S}_r^*(-1)$ if r is odd.

The propositions 2.3, 6.1 and the Theorems 3.1 (ii), 4.1 imply

PROPOSITION 6.2. – There are real coefficients

$$\alpha^{(\sigma)}, \beta_k^{(\sigma)}, \gamma^{(2,\sigma)}, \gamma_l^{(3,\sigma)}, \delta_m^{(2,\sigma)}, \delta_p^{(3,\sigma)}$$

with

$$\begin{aligned} I_2^\sigma &= \alpha^{(\sigma)} B, & I_4^\sigma &= \sum_{k=1}^3 \beta_k^{(\sigma)} W^{(k)}, \\ I_5^\sigma &= \gamma^{(2,\sigma)} S_5^{(2)}(-1) + \sum_{l=1}^2 \gamma_l^{(3,\sigma)} {}_5^l S_5^{(3)}(-1) \\ I_6^\sigma &= \sum_{m=1}^2 \delta_m^{(2,\sigma)} {}_6^m S_6^{(2)}(-1, 4) + \sum_{p=1}^6 \delta_p^{(3,\sigma)} {}_6^p S_6^{(3)}(-1, 4) \quad (\sigma = 1, 2, 3). \end{aligned}$$

Detailed information about the coefficients of Proposition 6.2 are given in [Wü5]. The following proposition was proved in [Wü2,3,5; McL; CMcL; McLS; Gü]:

PROPOSITION 6.3. – If g is an Einstein metric, a central symmetric metric, a metric of Petrov type N or D, then it follows from the moment equations

$$I_r^\sigma = 0, \quad \sigma \in \{1, 2, 3\}, 0 \leq r \leq 6$$

that g is conformally equivalent to a plane wave metric or to a flat metric ⁶.

2. Conformal Einstein space-times

The conformal tensors B (see (2.11)), W^2 (see (2.15)) and $S_6^{(2)}(-1, 4)$ (see Proposition 3.1(ii)) vanish in a special Einstein space-time, i.e. a space-time with $R_{ab} = 0$, if $n = 4$. Thus we have:

PROPOSITION 6.4. – In the case $n = 4$ the conditions

$$B = 0, \quad W^{(2)} = 0, \quad S_6^{(2)}(-1, 4) = 0$$

⁶ A conjecture is that the moment equations $I_r^\sigma = 0$ ($r = 2, 4, 5, 6$) are also sufficient for the validity of Huygens' principle for E_σ , $\sigma = 1, 2, 3$ and that these equations are fulfilled if and only if g is conformally equivalent to a plane wave metric or to a flat metric [W2,5; CMcL].

are necessary for (M, g) to be conformally related to a special Einstein space-time.

Remark 6.1. – H.W. Brinkmann [Bri; Scho] studied necessary and sufficient conditions for Riemannian spaces to be conformally related to Einstein spaces. However, since his arguments involved the existence and compatibility of differential equations, a constructive set of necessary and sufficient conditions is very difficult to infer. Kozamek, Newman, Tod [KNT] and the second author [Wü4] solved the problem in the physically interesting case $n = 4$ for all space-times excluding space-times of Petrov type N , using the conditions $B = 0$, $W^{(2)} = 0$. In the case of N -type the derivation of a constructive set of necessary and sufficient conditions for (M, g) to be conformal to an Einstein space is much more difficult [KNT; Wü4].

REFERENCES

- [AW] M. ALVAREZ and V. WÜNSCH, Zur Gültigkeit des Huygenschen Prinzips bei der Weyl-Gleichung und den homogenen Maxwellschen Gleichungen für Metriken vom Petrov-Typ N, *Wiss. Zeitschr. der Päd. Hochschule Erfurt/Mühlhausen, Math.-Naturwissenschaft. Reihe*, **27**, 1991, H.2., pp. 77-91.
- [AMLW] W. G. ANDERSON, R. G. MCLEAGHAN and T. F. WALTON, An explicit determination of the non-self-adjoint wave equations that satisfy Huygens' principle on Petrov type III background space-times, *Z.f. Anal. u. ihre Anw.*, **16**, 1997, 1, pp. 37-58.
- [Ba] J. R. BASTON, Verma Modules and Differential Conformal Invariants, *J. Differential Geometrie*, **32**, 1990, pp. 851-898.
- [BE] J. R. BASTON, and M. G. EASTWOOD, Invariant operators, Twistors in Mathematics and Physics, LMS Lecture Notes 156, (eds. BAILY and BASTON) Cambridge University Press, 1990.
- [BØ] T. BRANSON and B. ØRSTED B., Conformal indices of Riemannian manifolds, *Composito Mathematica*, **60**, 1986, pp. 261-293.
- [Bra] T. BRANSON, Differential Operators Canonically Associated to a Conformal Structure, *Mathematica Scandinavica*, **57**, 1985, pp. 293-345.
- [Bri] H. W. BRINKMAN, Riemann spaces conformal to Einstein spaces, *Math. Ann.*, **91**, 1924, pp. 269-287.
- [CMcL] J. CARMINATI and R. G. MCLEAGHAN, An explicit determination of the space-times on which the conformally invariant scalar wave equation satisfies Huygens' principle, *Ann. Inst. Henri Poincaré, Phys. théor.*, Vol. **44**, 1986, pp. 115-153; Part II, Vol. **47**, 1987, pp. 337-354; Part III, Vol. **48**, 1988, pp. 77-96; Vol. **54**, 1991, p. 9.
- [dP] J. C. DU PLESSIS, Polynomial conformal tensors, *Proc. Camb. Phil. Soc.*, **68**, 1970, pp. 329-344.
- [E] M. G. EASTWOOD, The Fefferman Graham conformal invariant, *Twistor Newsletter*, **20**, 1985, p. 46.
- [FG] C. FEFFERMAN and C. R. GRAHAM, Conformal invariants, *Élie Cartan et les Mathématiques d'Aujourd'hui*, Astérisque, 1985, pp. 95-116.

- [Ge] R. GERLACH, Beiträge zur konformen Geometrie in drei- und vierdimensionalen pseudo-Riemannschen Räumen, *Dissertation A, Pädagogische Hochschule Erfurt*, 1996.
- [GeW1] R. GERLACH and V. WÜNSCH, Über konforminvariante Tensoren ungerader Stufe in gekrümmten Raum-Zeiten, *Wiss. Zeitschr. d. Päd. Hochschule Erfurt/Mühlhausen, Math.-Naturwissensch. Reihe*, **26**, 1990, H.1, 20-32.
- [GeW2] GERLACH R. and V. WÜNSCH, Über konforminvariante Tensoren der Stufe sechs in gekrümmten Raum-Zeiten, *Wiss. Zeitschr. d. Päd. Hochschule Erfurt/Mühlhausen, Math.-Naturwissensch. Reihe*, **27**, 1991, H.2, 117-143.
- [Gü] P. GÜNTHER, Huygens' Principle and Hyperbolic Equations, *Academic Press, Boston*, 1988.
- [GüW1] P. GÜNTHER and V. WÜNSCH, Contributions to a theory of polynomial conformal tensors, *Math. Nachr.*, **126**, 1986, pp. 83-100.
- [GüW2] P. GÜNTHER and V. WÜNSCH, On some polynomial conformal tensors, *Math. Nach.*, **124**, 1985, pp. 217-238.
- [Il] R. ILLGE, Zur Gültigkeit des Huygensschen Prinzips bei hyperbolischen Differentialgleichungssystemen in statischen Raum-Zeiten, *Zs. für Anal. und Anwendungen*, **6**, 1987, pp. 385-407.
- [KNT] C. N. KOZAMEK, E. T. NEWMAN and K. P. TOD, Conformal Einstein-spaces, *Gen. Rel. and Grav.*, **17**, 4, 1985, pp. 343-352.
- [Lo] D. LOVELOCK, Dimensionally dependent identities, *Proc. Camb. Phil. Soc.*, **68**, 1970, pp. 345-350.
- [McL] R. G. MCLENAGHAN, An explicit determination of the empty space-times on which the wave equation satisfies Huygens' principle, *Proc. Cambridge Philos. Soc.*, **65**, 1969, pp. 139-155.
- [McLS] R. G. MCLENAGHAN and F. D. SASSE, Nonexistence of Petrov-type III space-times on which Weyl's neutrino equation or Maxwell's equations satisfy Huygens' principle, *Ann. Inst. Henri Poincaré, Phys. Théor.*, **65**, 1996, pp. 253-271.
- [Ø] B. ØRSTED, The Conformal Invariance of Huygens' Principle, *Diff. Geom.*, **16**, 1981, pp. 1-9.
- [PR] R. PENROSE and W. RINDLER, Spinors and Space-Time I, II, *Cambridge University Press*, 1984, 1986.
- [RW] B. RINKE and V. WÜNSCH, Zum Huygensschen Prinzip bei der skalaren Wellengleichung, *Beiträge zur Analysis*, **18**, 1981, pp. 43-75.
- [Schi] R. SCHIMMING, Konforminvarianten vom Gewicht -1 eines Zusammenhangs oder Eichfeldes, *Zeitschrift für Analysis und ihre Anwendungen*, **3** (5), 1984, pp. 401-412.
- [Scho] J. A. SCHOUTEN and Ricci-Calculus, *Springer-Verlag*, Berlin, 1954.
- [Sz] P. SZEKERES, Conformal tensors, *Proc. Roy. Soc. A.*, **304**, 1968, pp. 113-122.
- [Tho] T. Y. THOMAS, Conformal tensors, *Proc. of the Nat. Acad. of Sci, Washington*, **12** (152), 1926, **18** (103u. 189), 1932.
- [Thi] R. THIELKEN, Herleitung einiger neuer Identitäten für den Konformkrümmungstensor mit Hilfe des Spinorkalküls, *Diplomarbeit, Universität Leipzig, Pädagogische Hochschule Erfurt*, 1989.
- [Wü1] V. WÜNSCH, Über eine Klasse konforminvarianter Tensoren, *Math. Nachr.*, **73**, 1976, pp. 37-58.
- [Wü2] V. WÜNSCH, Cauchy-Problem und Huygenssches Prinzip bei einigen Klassen spinorieller Feldgleichungen I, II, *Beitr. zu Analysis*, **12**, 1978, pp. 47-76; **13**, 1979, pp. 147-177.
- [Wü3] V. WÜNSCH, Huygens' principle on Petrov type D Space-Times, *Ann. d. Physik, 7.Folge*, **46** (8), 1989, pp. 593-597.

- [Wü4] V. WÜNSCH, Conformal C- and Einstein Spaces, *Math. Nach.*, **146**, 1990, pp. 237-245.
- [Wü5] V. WÜNSCH, Moments and Huygens' principle for conformally invariant field equation in curved space-times, *Ann. Inst. Henri Poincaré, Physique théorique*, **60**, (4), 1994, pp. 433-455.

(*Manuscript received September 9th, 1997;
revised February 9th, 1998.*)