

ANNALES DE L'I. H. P., SECTION A

HANS TILGNER

**A spectrum generating nilpotent group for
the relativistic free particle**

Annales de l'I. H. P., section A, tome 13, n° 2 (1970), p. 129-136

http://www.numdam.org/item?id=AIHPA_1970__13_2_129_0

© Gauthier-Villars, 1970, tous droits réservés.

L'accès aux archives de la revue « Annales de l'I. H. P., section A » implique l'accord avec les conditions générales d'utilisation (<http://www.numdam.org/conditions>). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

NUMDAM

Article numérisé dans le cadre du programme
Numérisation de documents anciens mathématiques
<http://www.numdam.org/>

A spectrum generating nilpotent group for the relativistic free particle

by

Hans TILGNER

Philipps-Universität, Marburg, Deutschland.

ABSTRACT. — We give an embedding of the full pseudoorthogonal group into the symplectic and antisymplectic matrices on the vector space of position, time, momentum and energy variables. By means of the relativistic mass operator we define a nilpotent Lie group which is to describe after unitary representation the physical concepts of mass, position, time, momentum, energy, space inversion and time reflection of a relativistic free particle. The relation to the canonical Weyl algebra of relativistic quantum mechanics is given.

RÉSUMÉ. — On donne une injection du groupe pseudo-orthogonal complet dans les matrices symplectiques et antisymplectiques d'un espace vectoriel des variables de position, temps, impulsion et énergie. Avec l'aide de l'opérateur de masse relativiste, on définit un groupe de Lie nilpotent qui, après la représentation unitaire, doit décrire les conceptions physiques de masse, position, temps, impulsion, énergie, réflexion de l'espace et du temps d'une particule relativiste libre. La relation avec l'algèbre canonique de Weyl de la mécanique quantique relativiste est décrite.

§ 1: INTRODUCTION

In § 2 we describe an embedding of the full n -dimensional pseudo-orthogonal group (which consists of four connectivity components) into the full symplectic matrix group in $2n$ dimensions (which consists of two connecti-

vity components, the symplectic and the antisymplectic matrices). Among the various possibilities for this embedding into the symplectic only and into the full symplectic group we choose one which gives no inversion of the energy components. Momentum inversion is then found only in the two group components of space inversion and time reflection. If we take for Hamilton operator the mass operator of relativistic physics (i. e. the first Casimir operator of the Poincaré group) we get a nilpotent Lie group which is a special case of the solvable groups described in [I]. Inclusion of space inversion and time reflection (which was not done in [I]) then defines an enlarged nonconnected group of four connectivity components which we suggest to describe mass, position, time, momentum, energy space inversion and time reflection of a relativistic free particle after representation into the operators on a (necessarily infinite-dimensional) Hilbert space. This group gives no concept of angular momentum, orbital or spin, contrary to the Poincaré group which conversely gives no satisfactory concept of position. Perhaps this difficulty can be overcome by inclusion of the pseudoorthogonal group as was done for the nonrelativistic Galilei group in [I; § 10]. We easily get a matrix representation of that enlarged group (which contains the full Poincaré group as a subgroup) from (17) below. To describe the relativistic free particle satisfactorily the nilpotent group in question must have (at least) two sorts of Hilbert space representations corresponding to zero mass and positive mass either. If there are additional representations with masses not yet experimentally observed (like imaginary masses), we will be confronted to the problem of « unphysical » representations as in the representation theory of the Poincaré group.

In § 4 we describe the embedding of the Lie algebra of the discussed group into the canonical Weyl algebra which is constructed over the Heisenberg subalgebra of position, time, momentum and energy variables. Since position and time, momentum and energy are treated on equal footing we will get time and energy operators after representation, which must lead to a satisfactory interpretation of time and energy for the relativistic free particle, as well as of position and momentum.

§ 2: EMBEDDING OF THE FULL PSEUDORTHOGONAL GROUPS INTO THE FULL SYMPLECTIC GROUPS

Let V be a finite-dimensional vector space, $\mathfrak{a}$ a bilinear form on V and $\text{Aut}(V)$ the group of invertible endomorphisms on V . We define

- (1) $\text{Aut}(V, \mathfrak{a}) := \{ G \in \text{Aut}(V) / \mathfrak{a}(G\mathfrak{a}, Gy) = \mathfrak{a}(x, y) \text{ for all } x, y \in V \}$
 $\text{der}(V, \mathfrak{a}) := \{ M \in \text{End}(V) / \mathfrak{a}(M\mathfrak{a}, y) + \mathfrak{a}(x, My) = 0 \text{ for all } x, y \in V \}.$

We write the general element of the symplectic vector space E

$$x = \sum_{i=1}^{n_q} \xi^i q_i + \sum_{r=n_q+1}^n \xi^r t_r + \sum_{i=1}^{n_q} \xi_i p^i + \sum_{r=n_q+1}^n \xi_r h^r$$

$$E = E_q \oplus E_t \oplus E_p \oplus E_h$$

where $\dim(E_q) = \dim(E_p) = n_q$, $\dim(E_t) = \dim(E_h) = n_t$, $n_q + n_t = n$. We also write $E_q \oplus E_t =: E_1$, $E_p \oplus E_h =: E_2$ and $x = x_1 + x_2$ according to this decomposition. In the following $i, k, l, m = 1, \dots, n_q$ and $r, s, t, u = n_q + 1, \dots, n_q + n_t = n$. Hence we have $\dim(E) = 2n$. Let id_q , id_t , id_p and id_h be the identities in E_q , E_t , E_p and E_h respectively, and $\begin{pmatrix} \text{id}_q & 0 \\ 0 & -\text{id}_t \end{pmatrix} =: I_{qt}$. Then the $2n \times 2n$ matrix $\begin{pmatrix} I_{qt} & 0 \\ 0 & I_{ph} \end{pmatrix} =: -P$ is in $\text{Spl}(2n, \mathbb{R}, E)$. We define $\tau(x, y) := \sigma(Jx, -Py) =: \tau_1(x_1, y_1) + \tau_2(x_2, y_2)$, which induces on the vector spaces E , E_1 , E_2 symmetric nondegenerate indefinite bilinear forms τ , τ_1 , τ_2 . If we denote by $\mathbb{C}_4$ the commutative group of the elements id_n , $-\text{id}_n$, $-I_{qt}$, I_{qt} (the two latter are called *space inversion* and *time reflection* respectively) the full pseudoorthogonal group in E_1 is

$$\text{Aut}(E_1, \tau_1) = \text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1) \otimes \mathbb{C}_4 = \text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1) \cup$$

$$-I_{qt} \text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1) \cup I_{qt} \text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1) \cup -\text{id}_n \text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1)$$

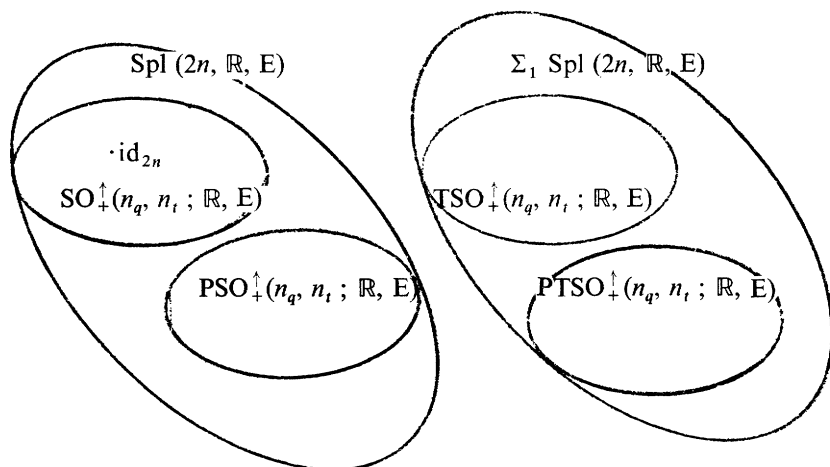
(where $\text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1)$ is its identity component). We want to embed this group into the *full symplectic group* in two dimensions, which is $\text{SpA}(2n, \mathbb{R}, E) := \text{Spl}(2n, \mathbb{R}, E) \cup \Sigma_1 \text{Spl}(2n, \mathbb{R}, E)$. This embedding is given by

$$(2) \quad G \mapsto \begin{pmatrix} G & 0 \\ 0 & (G^T)^{-1} \end{pmatrix} =: G \quad G \in \text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_q)$$

for the identity component of $\text{Aut}(E_1, \tau_1)$. We have $G \in \text{Aut}(E_1, \tau_1)$ iff $G \in \text{Aut}(E, \sigma) \cap \text{Aut}(E, \tau)$. This $2n$ -dimensional representation of $\text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1)$ is denoted by $\text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E)$ in the future. To continue this representation to the full pseudoorthogonal group we consider among the various representations of $\mathbb{C}_4$ in $\text{Spl}(2n, \mathbb{R}, E)$ and $\text{SpA}(2n, \mathbb{R}, E)$ the « physical » one

$$(3) \quad \begin{aligned} \text{id}_n &\mapsto \text{id}_{2n} & -I_{qt} &\mapsto P \\ I_{qt} &\mapsto \begin{pmatrix} I_{qt} & 0 \\ 0 & -I_{ph} \end{pmatrix} =: T & -\text{id} &\mapsto \begin{pmatrix} -\text{id}_n & 0 \\ 0 & \text{id}_n \end{pmatrix} = PT \end{aligned}$$

with id_{2n} , $P \in \text{Spl}(2n, \mathbb{R}, E)$ and $T, PT \in \Sigma_1 \text{Spl}(2n, \mathbb{R}, E)$. We choose the terminus physical for this representation because it has the physical transformation properties of the *momentum* components x_p and no reflection of the *energy* components x_h of $x \in E$. If we want to embed projective Hilbert space representations of the full pseudoorthogonal group into those of the full symplectic group (for instance when we use the latter as a so-called noninvariance group of the relativistic free particle) we have to represent id_{2n} and P by unitary and T and PT by antiunitary operators [2; p. 7].



The two components of $\text{Aut}(E_1, \tau_1)$ in $\text{Spl}(2n, \mathbb{R}, E)$ may be called *causal* subgroup of $\text{Aut}(E_1, \tau_1)$ [2; p. 2]. The Lie algebra of $\text{SO}_+^\uparrow(n_q, n_t; \mathbb{R}, E_1)$ is $\text{der}(E_1, \tau_1)$. The condition for M to be in $\text{der}(E_1, \tau_1)$ is in matrix form $M^T I_{qt} + I_{qt} M = 0$. The embedding of $\text{der}(E_1, \tau_1)$ in $\text{spl}(2n, \mathbb{R}, E)$ follows then from [I; p. 77]

$$(4) \quad M \mapsto \begin{pmatrix} M & 0 \\ 0 & I_{qt} M I_{qt} \end{pmatrix}.$$

§ 3: THE NILPOTENT GROUP OF THE RELATIVISTIC FREE PARTICLE

For the nilpotent $2n \times 2n$ matrix $\begin{pmatrix} 0 & -I_{ph} \\ 0 & 0 \end{pmatrix} =: F \in \text{spl}(2n, \mathbb{R}, E)$ the one-parameter subgroup $\exp(\alpha F) = \text{id}_n + \alpha F$ of $\text{Spl}(2n, \mathbb{R}, E)$ consists of unipotent matrices. This group is isomorphic as a Lie group to the

straight line by $\exp(\mathbb{R}F) \mapsto \mathbb{R}F$, if we write the latter as the additive group $\mathbb{R}F$. For the relativistic free particle we consider the following nilpotent groups on the analytic manifolds $\exp(\mathbb{R}F) \times E \times \mathbb{R}$, $\exp(\mathbb{R}F) \times E \times e^{i\mathbb{R}}$, $\mathbb{R}F \times E \times \mathbb{R}$ with compositions

$$(5) \quad (e^{\alpha F}, x, \beta)(e^{\alpha' F}, y, \beta') = (e^{(\alpha + \alpha')F}, x + e^{\alpha F}y, \sigma(x, e^{\alpha F}y) + \beta + \beta')$$

$$(6) \quad (e^{\alpha F}, x, e^{i\beta})(e^{\alpha' F}, y, e^{i\beta'}) = (e^{(\alpha + \alpha')F}, x + e^{\alpha F}y, e^{i(\sigma(x, e^{\alpha F}y) + \beta + \beta')})$$

$$(7) \quad (\alpha F, x, \beta)(\alpha' F, x, \beta') = ((\alpha + \alpha')F, x + e^{\alpha F}y, \sigma(x, e^{\alpha F}y) + \beta + \beta')$$

respectively. These $2n + 2$ -dimensional groups are connected, noncompact, and locally isomorphic with Lie algebras

$$(8) \quad [(\alpha F, x, \beta), (\alpha' F, y, \beta')]_- = (0, \alpha Fy - \alpha' Fx, \sigma(x, y))$$

on the vector spaces $\mathbb{R}F \oplus E \oplus \mathbb{R}$ (after a suitable change of basis). The groups (5) and (7) are simply connected and therefore global isomorphic. The group (6) is infinitely connected since the kernel of the covering homomorphism

$$(9) \quad (e^{\alpha F}, x, \beta) \mapsto (e^{\alpha F}, x, e^{i\beta})$$

is $2\pi\mathbb{Z}$. (5) and (6) are central extensions of the Lie groups $e^{\mathbb{R}F} \otimes E$ cf. [I; theorem 12] by $\mathbb{R}$ and $e^{i\mathbb{R}}$ respectively. The proofs of these statements are given in [I; § 2, 3, 10]. For $S = e^{\alpha F}$ we get a faithful $2n + 2$ -dimensional representation of (5) from [I; (42)]. We don't know similar representations for (6). We study the group (6) besides the group (5) since it is a central extension of $\exp(\mathbb{R}F) \otimes E$ by the torus, which is the standard connection we need for the ray representation theory of the latter [3].

A $2n + 2$ -dimensional representation of the Lie algebras (8) is given in [I; (44)] for $V = \alpha F$. The proof of [I; § 10] for $n \geq 2$ is applicable to our special case with only minor modifications. Hence we get for the derivation Lie algebra of (8) the set of matrices

$$(10) \quad \begin{pmatrix} v & 0 & 0 \\ -F\beta p & W + \frac{1}{2}v\Sigma_3 + \beta id_{2n} & 0 \\ \rho & p^T & \beta \end{pmatrix}$$

$$\text{with } v, \rho, \beta \in \mathbb{R}, \Sigma_3 := \begin{pmatrix} id_n & 0 \\ 0 & -id_n \end{pmatrix} \text{ and } W = \begin{pmatrix} L & B \\ 0 & I_{qt}LI_{qt} \end{pmatrix},$$

B a symmetric $n \times n$ matrix, $L \in \text{der}(E_1, \tau_1)$, i. e. $L^T I_{qt} + I_{qt} L = 0$. It is

straightforward to show that because of [I ; theorem 38] the automorphism group of (8) is given by the set of $2n + 2 \times 2n + 2$ matrices

$$(11) \quad \begin{pmatrix} \delta & 0 & 0 \\ -\frac{1}{\alpha} \mathbf{G} \mathbf{F} \mathbf{J} \mathbf{b} & \mathbf{G} & 0 \\ \tau & \mathbf{b}^T & \alpha \end{pmatrix} \quad \begin{array}{l} 0 \neq \alpha \in \mathbb{R} \\ 0 \neq \delta \in \mathbb{R} \\ \tau \in \mathbb{R} \end{array}$$

where the $2n \times 2n$ matrix $\mathbf{G}$ has the form

$$\mathbf{G} = \sqrt{|\alpha|} \left\{ \frac{1 + \text{sign } \alpha}{2} \begin{pmatrix} \mathbf{A} & \mathbf{B} \\ 0 & \frac{1}{\delta} \mathbf{I}_{q_t} \mathbf{A} \mathbf{I}_{q_t} \end{pmatrix} + \frac{1 - \text{sign } \alpha}{2} \begin{pmatrix} \mathbf{C} & \mathbf{D} \\ 0 & \frac{1}{\delta} \mathbf{I}_{q_t} \mathbf{C} \mathbf{I}_{q_t} \end{pmatrix} \right\}$$

and the $n \times n$ matrices $\mathbf{B}^T \mathbf{I}_{q_t} \mathbf{A} \mathbf{I}_{q_t}$ and $\mathbf{D}^T \mathbf{I}_{q_t} \mathbf{C} \mathbf{I}_{q_t}$ are symmetric, where as $\mathbf{A}$ and $\mathbf{C}$ are subject to $\mathbf{A}^T \mathbf{I}_{q_t} \mathbf{A} = \delta \mathbf{I}_{q_t}$ and $\mathbf{C}^T \mathbf{I}_{q_t} \mathbf{C} = -\delta \mathbf{I}_{q_t}$. The adjoint representations of the groups (5) and (6), which are the same for both, are given by [I ; (47)], those of the Lie algebras (8) are given by [I ; (46)].

§ 4: CONNECTION WITH THE CANONICAL FORMALISM

We discuss the relation of the above groups and their Lie algebras to the canonical formalism, i. e. to the associative infinite-dimensional Weyl algebra of [I] which we get from the sub algebra $\text{heis}(2n)$ in (8). For Hamilton operator we take the *mass operator* of the relativistic free particle

$$(12) \quad \mathbf{H}_F := \frac{1}{2} \left(\sum_{i=1}^{n_q} p^i p^i - \sum_{r=n_q+1}^n h^r h^r \right), \quad \mathbf{R}(\mathbf{H}_F) = \mathbf{F}$$

($\mathbf{R}$ being the isomorphism [I ; (74)]). The sub algebra $\mathbb{R}\mathbf{F} \rtimes \text{heis}(2n)$ of the Lie algebra $\text{spl}(2n, \mathbb{R}) \rtimes \text{heis}(2n)$, which is formed by $\Lambda \mathbf{W}_2 \oplus \mathbf{E} \oplus \mathbb{R} \mathbf{1}$ in $\text{weyl}(\mathbf{E}, \sigma)$, then is isomorphic to the Lie algebra (8), and the Lie algebra of linear transformations on $\mathbf{E} \subset \text{weyl}(\mathbf{E}, \sigma)$ given by $\text{ad } (\mathbf{E} \oplus \Lambda \mathbf{W}_2) |_{\mathbf{E}}$ is isomorphic to the adjoint Lie algebra of (8), describes by [I; (18), (46)]. The centralizer of $\mathbf{H}_F$ in $\mathbf{E} \oplus \Lambda \mathbf{W}_2$, which is supposed to be the smallest Lie algebra the universal enveloping algebra of which after embedding in

$\text{weyl}(E, \sigma)$ is the centralizer of H_F in $\text{weyl}(E, \sigma)$, is given by the $\frac{1}{2}n(n-1)$ elements

$$(13) \quad M_{ik} := \Lambda(q_i p^k - q_k p^i), \quad M_{ir} := \Lambda(q_i h^r + t_r p^i), \quad M_{rs} := \Lambda(t_r h^s - t_s h^r)$$

which fulfill the well known commutation relations of the pseudoorthogonal Lie algebra, together with the n elements p^i and h^r . The R-image [I; (74)] of the Lie algebra of the M 's is given by the matrices (4), the exponentials of those matrices by (2). One easily proves that

$$\text{ad}(M)|_E \in \text{der}(E, \sigma) \cap \text{der}(E, \tau)$$

for M a linear combination of the M 's above. For the realization of some outer automorphisms of $\mathbb{R}F \rhd \text{heis}(2n)$ in $\text{weyl}(E, \sigma)$ we get the corresponding results to [I; § 10].

§ 5: INCLUSION OF SPACE INVERSION AND TIME REFLECTION

Consider the one-parameter nonconnected subgroup

$$(14) \quad \mathbb{C}_4 \otimes e^{\mathbb{R}F} = e^{\mathbb{R}F} \cup P e^{\mathbb{R}F} \cup T e^{\mathbb{R}F} \cup P T e^{\mathbb{R}F} \\ = (\text{id}_{2n} + \mathbb{R}F) \cup (P + \mathbb{R}Q) \cup (T - \mathbb{R}Q) \cup (PT - \mathbb{R}F)$$

of $\text{SpA}(2n, \mathbb{R}, E)$, where $Q := \begin{pmatrix} 0 & \text{id}_n \\ 0 & 0 \end{pmatrix}$, which has four connectivity components and consists no longer of unipotent matrices only. Its composition on the manifold $\mathbb{C}_4 \times e^{\mathbb{R}F}$ is given by

$$(15) \quad (Z, e^{\alpha F})(Z', e^{\alpha' F}) := (ZZ', e^{((-1)^v \alpha + \alpha') F})$$

where v depends on Z and is zero for $Z \in \text{Spl}(2n, \mathbb{R}, E)$ and one for $Z \in \Sigma_1 \text{ Spl}(2n, \mathbb{R}, E)$. We have $(Z, e^{\alpha F})^{-1} = (Z, e^{(-1)^{v+1} \alpha F})$; for the commutant of two elements we get $(\text{id}_{2n}, \exp((-1)^v)((-1)^v - 1)((-1)^v \alpha + \alpha') F)$, which shows that the second commutant on $\mathbb{C}_4 \otimes e^{\mathbb{R}F}$ is the identity element $(\text{id}_{2n}, \text{id}_{2n})$; hence the group is nilpotent.

On the nonconnected manifold $\mathbb{C}_4 \times e^{\mathbb{R}F} \times E \times \mathbb{R}$ of four connectivity components we define a group composition by

$$(16) \quad (Z, e^{\alpha F}, x, \beta)(Z', e^{\alpha' F}, y, \beta') := \\ (ZZ', e^{((-1)^v \alpha + \alpha') F}, x + i^v Z e^{\alpha F} y, i^v \sigma(x, Z e^{\alpha F} y) + \beta + \beta').$$

This group

$$\mathbb{C}_4 \otimes (e^{\mathbb{R}^F} \otimes \text{Heis}(2n)) = (\mathbb{C}_4 \otimes e^{\mathbb{R}^F}) \otimes (\text{Heis}(2n)) = e^{\mathbb{R}^F} \otimes (\mathbb{C}_4 \otimes \text{Heis}(2n))$$

is again nilpotent. A faithful $2n + 2$ -dimensional representation is given by the matrices

$$(17) \quad (Z, e^{\alpha_F}, x, \beta) \mapsto \begin{pmatrix} 1 & 0 & 0 \\ x & i^v Z e^{\alpha_F} & 0 \\ \beta & i^v x^T \tilde{J} Z e^{\alpha_F} & 1 \end{pmatrix}$$

By substituting $\mathbb{R} \rightarrow e^{i\mathbb{R}}$ into the group (16) we get the corresponding results for the group (6), and by dropping that last factor we get the corresponding results for the *adjoint group* $\mathbb{C}_4 \otimes (e^{\mathbb{R}^F} \otimes E)$. A faithful $2n + 2$ -dimensional representation of this adjoint group is given by the matrices [7; (47)] if we substitute $e^{\alpha_F} \mapsto Z e^{\alpha_F}$ and use $ZF = (-1)^v FZ$. We summarize the relation between the various groups of the relativistic case by the commutative diagram of short exact sequences

$$\begin{array}{ccccc}
 & e^{i\mathbb{R}} & & \mathbb{R} & \\
 & \searrow & & \swarrow & \\
 & e^{\mathbb{R}^F} \times E \times e^{i\mathbb{R}} & \xleftarrow{\varphi} & e^{\mathbb{R}^F} \times E \times \mathbb{R} & \xleftarrow{2\pi\mathbb{Z}} \\
 & \searrow & \swarrow & \searrow & \swarrow \\
 e^{i\mathbb{R}} & & e^{\mathbb{R}^F} \times E & & \mathbb{R} \\
 & \searrow & \downarrow & \searrow & \swarrow \\
 \mathbb{C}_4 \times e^{\mathbb{R}^F} \times E \times e^{i\mathbb{R}} & \xleftarrow{\quad} & \mathbb{C}_4 \times e^{\mathbb{R}^F} \times E \times \mathbb{R} & \xleftarrow{2\pi\mathbb{Z}} & \\
 & \searrow & \downarrow & \searrow & \swarrow \\
 & \mathbb{C}_4 & \mathbb{C}_4 \times e^{\mathbb{R}^F} \times E & \mathbb{C}_4 & \\
 & \updownarrow & \updownarrow & \updownarrow & \\
 & \mathbb{C}_4 & \mathbb{C}_4 & \mathbb{C}_4 &
 \end{array}$$

(we dropped the trivial parts). φ is the covering homomorphism (9).

This work was supported in part by the Deutsche Forschungsgemeinschaft.

REFERENCES

- [1] H. TILGNER, A Class of Solvable Lie Groups and their Relation to the Canonical Formalism.
- [2] D. SIMMS, *Lie Groups in Quantum Mechanics*. Springer. Lecture Notes in Mathematics, 52, Berlin, 1968.
- [3] C.-M. EDWARDS et J. LEWIS, Twisted Group Algebras. I. *Commun. math. Phys.*, t. 13, 1969, p. 119-130.

(Manuscript reçu le 10 octobre 1969).