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F. BRICKELL R.S. CLARK

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## CONFORMALLY RIEMANNIAN STRUCTURES, I;

BY

### F. BRICKELL AND R. S. CLARK

(Southampton).

Introduction. — We define a conformally Riemannian structure on a differentiable (1) manifold M of dimension n to be a differentiable subordinate structure of the tangent bundle to M whose group G consists of the non-zero scalar multiples of the orthogonal  $n \times n$  matrices. The method of equivalence of E. Cartan [1], as described by S. Chern [3], associates with a given conformal structure a certain principal fibre bundle on which a set of linear differential forms is defined globally. We obtain such a bundle and set of forms explicitly and show their relation to the normal conformal connection of E. Cartan [2].

The first paragraph contains an exposition of conformal connections in the light of C. Ehresmann's general theory of Cartan connections [4]. In the second paragraph we show how this leads to the normal conformal connection on a manifold admitting a conformally Riemannian structure. The third paragraph summarises the method of Cartan-Chern and we apply this, in the fourth paragraph, to the special case of a conformally Riemannian structure. In the fifth paragraph we show how these ideas are related.

1. Conformal Cartan connections. — We first collect together the information we require on conformal space and on Cartan connections.

Conformal space of dimension n is defined to be the homogeneous space K/K', where K is the linear group on n+2 variables  $\{\xi_0, \xi_1, \ldots, \xi_{n+1}\}$  leaving invariant the quadratic form

$$\sum_{i=1\ldots n} \xi_i^2 + \xi_0 \xi_{n+1}$$

<sup>(1)</sup> The word differentiable will always mean differentiable of class C  $^{\omega}$ .

and K' is the subgroup of K leaving invariant the point  $\{1, 0, \ldots, 0\}$ . Explicitly, K' consists of matrices of the form

$$\begin{bmatrix}
b & p & c \\
o & A & Aq \\
o & o & a
\end{bmatrix}$$

where A is an orthogonal  $n \times n$  matrix and the remaining elements satisfy the relations

$$(1.1) ab = 1, ap + \tilde{q} = 0, 2ac + \tilde{q}q = 0$$

 $\tilde{q}$  denoting the transpose of q.

The linear group of isotropy  $L'_n$  of the conformal space at  $\{1,0,\ldots,0\}$  is isomorphic with the group G of non-zero scalar multiples of the orthogonal  $n \times n$  matrices. We identify  $L'_n$  with G in such a way that the canonical homomorphism  $\varphi$  of K' onto  $L'_n$  is

$$\begin{bmatrix} b & p & c \\ o & A & Aq \\ o & o & a \end{bmatrix} \stackrel{\circ}{\rightarrow} aA.$$

The Lie algebra  $\mathcal{L}$  K is isomorphic with the Lie algebra of the  $(n+2) \times (n+2)$  matrices of the type

$$\begin{bmatrix} -\mu & -\tilde{\Psi} & o \\ \omega & \Omega & \psi \\ o & -\tilde{\omega} & \mu \end{bmatrix}$$

where the  $n \times n$  matrix  $\Omega$  is skew-symmetric. A representation of the subalgebra  $\mathcal{E}K'$  is obtained by imposing the condition  $\omega = 0$ . The translation operations on these Lie algebras are obtained by matrix multiplication.

- G. EHRESMANN [4] has given necessary and sufficient conditions for the existence of a Cartan connection on M of type K/K', that is, a conformal Cartan connection. These are:
- (i) that the tangent bundle of M should admit a subordinate structure with group  $L'_n$ ;
- (ii) that there should exist a principal fibre bundle  $\mathcal{H}' = H'(M, K')$  with which the homomorphism  $\varphi$  associates the subordinate structure.

Since K' is a subgroup of K,  $\mathcal{X}'$  defines canonically a principal bundle  $\mathcal{X} = H(M, K)$ . A conformal Cartan connection on M is a connection on  $\mathcal{X}$ , in the usual sense, such that no horizontal directions on H are tangent to the subspace H'.

We shall construct  $\partial \mathcal{C}'$  from a cocycle  $\mathcal{K}'_{\alpha\beta}$ , whit values in K', defined on an open covering  $\{U_{\alpha}\}$  of M. Then H' is the quotient of the sum  $\sum_{\alpha} U_{\alpha} \times K'$  by the equivalence relation

$$(m_{\alpha}, k'_{\alpha})_{\beta} \sim (m_{\beta}, k'_{\beta}) = if \quad m_{\alpha} = m_{\beta}, \quad k'_{\alpha} = (k'_{\alpha\beta}m_{\alpha}) k'_{\beta}.$$

A Cartan connection can then be obtained from local 1-forms  $\Gamma_{\alpha}$  with values in  $\mathcal{L}K$  defined on  $U_{\alpha}$ , provided that on  $U_{\alpha} \cap U_{\beta}$  they satisfy the relation.

(1.2) 
$$\Gamma_{\beta} = (k'_{\alpha\beta})^{-1} \{ \Gamma_{\alpha} k'_{\alpha\beta} + dk'_{\alpha\beta} \}$$

and possess the further property that  $\Gamma_{\alpha} \overrightarrow{m} \in \mathcal{L}K'$  if and only if the tangent vector  $\overrightarrow{m}$  of  $U_{\alpha}$  is zero.

Denote by h' the projection  $H' \to M$  and by  $h'^*$  the dual mapping on the differential forms in M. From the local product representation, we have functions  $k'_{\alpha}$  with values in K' on  $(h')^{-1}U_{\alpha}$ . The connection form  $\Gamma$  is defined locally in H' by

(1.3) 
$$\Gamma = (k'_{\alpha})^{-1} \left\{ (h'^{*}\Gamma_{\alpha}) k'_{\alpha} + dk'_{\alpha} \right\}$$

and this extends uniquely to H.

2. The normal conformal connection. — Suppose now that a conformally Riemannian structure is given on M, so that the tangent bundle of M admits a given subordinate structure with group G. We shall construct a particular Cartan connection on M called the normal conformal connection.

The first condition of EHRESMANN is satisfied since the linear isotropic group  $L'_n$  is isomorphic to G. We have to construct a bundle  $\mathcal{H}' = H'(M, K')$  which gives rise to the above subordinate structure, using the homomorphism  $\varphi: K' \to G$ .

We are given a covering of M by open sets,  $U_{\alpha}$ , each admitting a coordinate system  $x_{\alpha} = \{x_{\alpha}^1, \ldots, x_{\alpha}^n\}$  and a function  $X_{\alpha}$  with values in the general linear group GL(n, R), such that on  $U_{\alpha} \cap U_{\beta}$  the function

$$g_{\alpha\beta} = X_{\alpha} M_{\alpha\beta} X_{\overline{\beta}}^{-1}$$

where  $M_{\alpha\beta} = [\partial x_{\alpha}^{i}/\partial x_{\beta}^{i}]$ , has values in G. If  $dx_{\alpha}$  is the natural coframe on  $U_{\alpha}$ , then the coframe

$$\omega_{\alpha} = X_{\alpha} dx_{\alpha}$$

is adapted to the G-structure, since on  $U_{\alpha} \cap U_{\beta}$ :

$$\omega_{\alpha} = X_{\alpha} dx_{\alpha} = X_{\alpha} M_{\alpha\beta} dx_{\beta} = g_{\alpha\beta} \omega_{\beta}.$$

From this adapted coframe, we define a local Riemannian metric  $\mathfrak{G}_{\alpha}\mathfrak{G}_{\alpha}$  on  $U_{\alpha}$ .

To construct a cocycle on M which will define a bundle  $\partial \mathcal{C}$ , we remark that any matrix of G can be expressed uniquely as aA, where A is an orthogonal  $n \times n$  matrix and the real number a is positive. If we split up the functions  $g_{\alpha\beta}$  in this way

$$(2.1) g_{\alpha\beta} = a_{\alpha\beta} A_{\alpha\beta},$$

the following cocycle relations are satisfied

$$a_{\alpha\gamma} = a_{\alpha\beta} a_{\beta\gamma}, \quad A_{\alpha\gamma} = A_{\alpha\beta} A_{\beta\gamma}.$$

We use these functions to define

$$K_{\alpha\beta} = \begin{bmatrix} b_{\alpha\beta} & p_{\alpha\beta} & c_{\alpha\beta} \\ o & A_{\alpha\beta} & A_{\alpha\beta}q_{\alpha\beta} \\ o & o & a_{\alpha\beta} \end{bmatrix},$$

where  $q_{\alpha\beta}$  is defined by the relation

$$(2.2) q_{\alpha\beta}\omega_{\beta} = d(\log a_{\alpha\beta})$$

and the remaining components are determined by the relations (1.1). These new functions  $k'_{\alpha\beta}$  on  $U_{\alpha} \cap U_{\beta}$  have values in K' and it can be shown that they satisfy the cocycle relations, consequently they define a bundle  $\mathcal{H}' = H'(M, K')$ . Since the cocycle  $g_{\alpha\beta}$  is the image of the cocycle  $k'_{\alpha\beta}$  under the homomorphism  $\varphi$ , this bundle  $\mathcal{H}'$  satisfies Ehresmann's second condition. In fact,  $k'_{\alpha\beta}$  has values in the subgroup K'' of K' defined by  $\alpha > 0$ . We denote by  $\mathcal{H}'' = H''(M, K'')$  the principal bundle with group K'' defined by the cocycle  $k'_{\alpha\beta}$ . It is a sub-bundle of  $\mathcal{H}'$ .

We are now ready to construct on  $U_{\alpha}$  the local 1-form  $\Gamma_{\alpha}$  with values in  $\mathcal{L}$  K which will define the Cartan connection. We shall take this to be

$$\Gamma_{\alpha} = \begin{bmatrix} \mathrm{o} & -\widetilde{\psi}_{\alpha} & \mathrm{o} \\ \mathrm{o}_{\alpha} & \Omega_{\alpha} & \psi_{\alpha} \\ \mathrm{o} & -\widetilde{\mathrm{o}}_{\alpha} & \mathrm{o} \end{bmatrix},$$

where the 1-forms  $\Omega_{\alpha}$  and  $\psi_{\alpha}$  are still to be determined. We have, of course, to verify that the choices for these remaining components are such that  $\Gamma_{\alpha}$  satisfies the relation (1.2); the further condition on  $\Gamma_{\alpha}$  is satisfied already since the forms  $\omega_{\alpha}^{l}$  are linearly independent. Cartan determines  $\Omega_{\alpha}$  and  $\psi_{\alpha}$  in terms of the local Riemannian metric  $\widetilde{\omega}_{\alpha}\omega_{\alpha}$  on  $U_{\alpha}$  by imposing certain conditions on the curvature of the Cartan connection and this will be done by imposing conditions on the local curvature form

$$d\Gamma_{\alpha} + \Gamma_{\alpha} \wedge \Gamma_{\alpha}$$

consistent with relation (1.2).

This local curvature form has values in  $\mathcal{L}K$  and so it has components

$$egin{bmatrix} -B_{oldsymbollpha} & - ilde{D}_{oldsymbollpha} & \mathrm{o} \ T_{oldsymbollpha} & C_{oldsymbollpha} & D_{oldsymbollpha} \ \mathrm{o} & - ilde{T}_{oldsymbollpha} & B_{oldsymbollpha} \ \end{bmatrix}$$

where the values of the 2-form  $C_{\alpha}$  are skew-symmetric. The first condition  $T_{\alpha} = 0$  is consistent with (1.2); since

$$T_{\alpha} = d\omega_{\alpha} + \Omega_{\alpha} \wedge \omega_{\alpha}$$

it implies that  $\Omega_x$  is the connection form of the local Riemannian metric (calculated relative to the coframe  $\omega_x$ ). It now follows that on  $U_x \cap U_\beta$ :

$$C_{\beta} = g_{\beta\alpha} C_{\alpha} g_{\alpha\beta}.$$

Consequently if

$$C_{oldsymbollpha}\!=\!rac{{ extbf{I}}}{2}\,C^{i}_{jh\,k}\omega^{h}_{oldsymbollpha}\wedge\omega^{k}_{oldsymbollpha},$$

the second condition  $C^i_{jhi} = 0$  is consistent with (1.2) and, if  $n \ge 3$ , it can be shown to determine the form  $\psi_{\alpha}$  uniquely. Thus a Cartan connection has been determined from the conformal structure of M; it is the normal conformal connection of E. Cartan.

We shall need to calculate  $\psi_{\alpha}$  explictly and we suppose that  $\psi_{\alpha} = \psi_{ih} \omega_{\alpha}^{h}$ . Since

$$C_{\alpha} = R_{\alpha} - \omega_{\alpha} \wedge \tilde{\psi}_{\alpha} - \psi_{\alpha} \wedge \tilde{\omega}_{\alpha},$$

where  $R_{\alpha} = d\Omega_{\alpha} + \Omega_{\alpha} \wedge \Omega_{\alpha}$  is the curvature form of the local Riemannian metric then, if

$$R_{\alpha} = \frac{1}{2} R^{i}_{jhk} \omega^{h}_{\alpha} \wedge \omega^{k}_{\alpha},$$

it follows that

$$C_{jhk}^{i} = R_{jhk}^{i} + \delta_{ik}\psi_{jh} - \delta_{ih}\psi_{jk} + \delta_{jh}\psi_{ik} - \delta_{jk}\psi_{ih}.$$

The condition  $C_{jhi}^i = 0$  then shows that, for  $n \ge 3$ ,

$$\psi_{\mathbf{z}} = \frac{1}{n-2} \left\langle \frac{R}{2(n-1)} \, \hat{o}_{ih} - R_{ih} \right\rangle \omega_{\mathbf{z}}^{h}$$

where  $R_{jh} = R^i_{jhi}$  and  $R = R_{jh} \delta^{jh}$ . Consequently  $C_{\alpha}$  is the Weyl conformal curvature form for the local Riemannian metric.

Finally, we obtain a local formula for the connection form  $\Gamma$  on H''. From the local product structure of H'' we have functions  $k''_{\alpha}$  with values in K'' on  $(h'')^{-1}U_{\alpha}$  and we put

(2.3) 
$$k_{\alpha}'' = \begin{bmatrix} b_{\alpha} & p_{\alpha} & c_{\alpha} \\ o & A_{\alpha} & A_{\alpha}q_{\alpha} \\ o & o & a_{\alpha} \end{bmatrix},$$

where  $A_{\alpha}$  is orthogonal and

$$a_{\alpha} > 0$$
,  $a_{\alpha}b_{\alpha} = 1$ ,  $a_{\alpha}p_{\alpha} + \tilde{q}_{\alpha} = 0$ ,  $2a_{\alpha}c_{\alpha} + \tilde{q}_{\alpha}q_{\alpha} = 0$ .

Since

$$(k'_{\alpha})^{-1} = \begin{bmatrix} a_{\alpha} & \widetilde{q}_{\alpha}\widetilde{A}_{\alpha} & c_{\alpha} \\ o & \widetilde{A}_{\alpha} & \widetilde{p}_{\alpha} \\ o & o & b_{\alpha} \end{bmatrix},$$

the formula (1.3) applied to H'' shows that

$$\Gamma \!=\! \begin{bmatrix} -\mu & -\tilde{\psi} & \mathbf{0} \\ \mathbf{\omega} & \Omega & \psi \\ \mathbf{0} & -\tilde{\mathbf{\omega}} & \mu \end{bmatrix}\!,$$

where the global forms are defined locally by

$$(2.4) \begin{cases} \omega = \frac{1}{a_{\alpha}} \tilde{A}_{\alpha}(h''^{*}\omega_{\alpha}), \\ \mu = \frac{da_{\alpha}}{a_{\alpha}} - \tilde{\omega} q_{\alpha}, \\ \Omega = \tilde{A}_{\alpha} \{ (h''^{*}\Omega_{\alpha}) A_{\alpha} + d\tilde{A}_{\alpha} \} - \omega \tilde{q}_{\alpha} + q_{\alpha} \tilde{\omega}, \\ \psi = dq_{\alpha} + \Omega q_{\alpha} - \mu q_{\alpha} + a_{\alpha} \tilde{A}_{\alpha}(h''^{*}\psi_{\alpha}) - (\tilde{q}_{\alpha} \omega) q_{\alpha} + \frac{1}{2} (\tilde{q}_{\alpha} q_{\alpha}) \omega. \end{cases}$$

3. The method of equivalence of E. Cartan and S. Chern. — In this paragraph we shall suppose that G is any closed subgroup of the linear group and that the tangent bundle of a manifold M admits a subordinate structure with group G. In the nomenclature of S. Chern, M admits a G-structure. In [3], Chern gives a procedure for constructing a sequence of fibre bundles and differential forms for a G-structure. We give a short account of his work.

From the definition of a subordinate structure, there exists an open covering of M by coordinate neighbourhoods  $U_{\alpha}$  on which are defined functions  $X_{\alpha}$ , with values in the linear group, such that on  $U_{\alpha} \cap U_{\beta}$  the functions

$$g_{\alpha\beta} = X_{\alpha} M_{\alpha\beta} X_{\overline{\beta}}^{-1}$$

have values in G. The coframe

$$\omega_{\alpha} = X_{\alpha} dx_{\alpha}$$

on  $U_{\alpha}$  is adapted to the G-structure, since on  $U_{\alpha} \cap U_{\beta}$ ,

$$\omega_{\alpha} = g_{\alpha\beta}\omega_{\beta}$$
.

The first fibre bundle in the sequence is the principal bundle  $\mathfrak{G} = B(M, G)$  associated with the reduced structure and it is defined by the cocycle  $g_{\alpha\beta}$ . As usual, we shall denote by b the projection  $B \to M$  and by  $b^*$  the dual mapping on the forms in M. Let  $g_{\alpha}$  denote the local functions with values in G on  $V_{\alpha} = b^{-1}U_{\alpha}$  defined by the local product structure, so that on  $V_{\alpha} \cap V_{\beta}$ ,

$$g_{\alpha} = (b^{\star}g_{\alpha\beta})g_{\beta}.$$

Using the local 1-forms  $\omega_{\alpha}$  on  $U_{\alpha}$ , we construct on B a global 1-form  $\theta$  with values in  $R^n$ . It is defined on  $V_{\alpha}$  by

$$(3.1) \theta = g_{\alpha}^{-1}(b^*\omega_{\alpha}),$$

and its exterior derivative is given on  $V_{\alpha}$  by

$$d\theta = g_{\mathbf{x}}^{-1} b^{\star} (d\omega_{\mathbf{x}}) - g_{\mathbf{x}}^{-1} dg_{\mathbf{x}} \wedge \theta.$$

We can express  $g_{\mathbf{x}}^{-1} b^{\star}(d\omega_{\mathbf{x}})$  as  $\frac{1}{2} C^{i}_{hk} \theta^{h} \wedge \theta^{k}$  and so, if we put

$$II_{\alpha} = g_{\alpha}^{-1} dg_{\alpha} + \varepsilon_{\alpha},$$

where  $\varepsilon_{\alpha} = \varepsilon_{jh}^{i} \theta^{h}$  is a 1-form on  $V_{\alpha}$  with values in the Lie algebra  $\mathcal{L}G$  whose coefficients are to be determined, the above formula for  $d\theta$  becomes

(3.2) 
$$d\theta + \Pi_{\alpha} \wedge \theta = \frac{1}{2} \left( \varepsilon_{kh}^i - \varepsilon_{hk}^i + C_{hk}^i \right) \theta^h \wedge \theta^k.$$

We impose as many linear relations with constant coefficients between the quantities  $\frac{1}{2}(\varepsilon^i_{hh}-\varepsilon^i_{hk}+C^i_{hk})$  as possible. These quantities are then determined uniquely. This implies that if the coefficients of the form  $\eta \wedge 0$  satisfy the same linear relations, where  $\eta$  is any 1-form  $\eta^i_{jh}\theta^h$  with values in  $\mathcal{E}G$ , then  $\eta \wedge 0 = 0$ . The relations may, or may not, determine the coefficients  $\varepsilon^i_{jh}$ . If they do and if the coefficients of  $\eta \wedge 0$  satisfy the same relations, then  $\eta = 0$ .

Thus on  $V_{\alpha}$  we have the formula

$$d\theta + \Pi_{\alpha} \wedge \theta = \tau_{\alpha}$$

and on  $V_{\mathfrak{B}}$ 

$$d\theta + \mathbf{H}_3 \wedge \theta = \tau_3$$

where the coefficients of  $\tau_{\alpha}$  and  $\tau_{\beta}$  are determined by the imposed linear relations. Since on  $V_{\alpha} \cap V_{\beta}$ ,

$$\tau_{\alpha} - \tau_{\beta} = (\Pi_{\alpha} - \Pi_{\beta}) \wedge \theta,$$

the coefficients of the form  $(\Pi_{\alpha} - \Pi_{\beta}) \wedge \theta$  also satisfy these linear relations. But the form  $\Pi_{\alpha} - \Pi_{\beta}$  has values in  $\mathcal{E}G$  and, since

$$(3.3) g_{\overline{\alpha}}^{-1} dg_{\alpha} - g_{\overline{\beta}}^{-1} dg_{\beta} = g_{\overline{\beta}}^{-1} b^{*} (g_{\alpha\beta}^{-1} dg_{\alpha\beta}) g_{\beta},$$

it is linear in  $\theta^i$ . Consequently

$$(\Pi_{\alpha} - \Pi_{\beta}) \wedge \theta = 0$$

and we have a global 2-form  $\tau$  on B defined on  $V_{\alpha}$  by  $\tau = \tau_{\alpha}$ . If the imposed relations determine the coefficients  $z_{jh}^{l}$ , then  $\Pi_{\alpha} = \Pi_{\beta}$  and we have a global 1-form  $\Pi$  on B defined on  $V_{\alpha}$  by  $\Pi = \Pi_{\alpha}$ .

But in the general case,

$$\Pi_{\alpha} - \Pi_{\beta} = \lambda_{\alpha\beta}^{\nu} \Lambda_{\nu} \quad (\nu = 1, \ldots, d_1)$$

where  $\Lambda_{\gamma}$  are a basis for the  $d_1$ -dimensional vector space of 1-forms on B, with values in  $\mathcal{L}G$ , which satisfy the equation  $\eta \wedge \theta = 0$  and whose components are linear in  $\theta^i$  with constant coefficients. The functions  $\lambda_{\alpha\beta}$  on  $V_{\alpha} \cap V_{\beta}$  form a cocycle on B with values in the additive group  $R^{d_1}$  and so they define a principal bundle

$$\mathfrak{B}^{\scriptscriptstyle 1} = B^{\scriptscriptstyle 1}(B, R^{d_{\scriptscriptstyle 1}}).$$

Denote by  $b^{\dagger}$  the projection  $B^{\dagger} \rightarrow B$  and by  $\lambda_{\alpha}$  the local functions with values in  $R^{d_1}$  on  $V_{\alpha}^{\dagger} = (b^{\dagger})^{-1} V_{\alpha}$ . Since on  $V_{\alpha}^{\dagger} \cap V_{\beta}^{\dagger}$ ,

$$\lambda_{\alpha} - \lambda_{\beta} = b^{1*} \lambda_{\alpha\beta},$$

we have global t-forms  $\theta^{\dagger}$ ,  $\mathbf{H}^{\dagger}$  on  $B^{\dagger}$  defined by

$$\begin{array}{c} \theta^{\scriptscriptstyle \rm I} = b^{\scriptscriptstyle \rm I}{}^{\star}\theta, \\ \Pi^{\scriptscriptstyle \rm I} = b^{\scriptscriptstyle \rm I}{}^{\star}\Pi_{\rm x} - \lambda_{\rm x}^{\scriptscriptstyle y}(b^{\scriptscriptstyle \rm I}{}^{\star}\Lambda_{\rm y}). \end{array}$$

We now use the same procedure to construct a decomposition for  $d\theta^{\dagger}$  and  $d\Pi^{\dagger}$  and thus obtain further local forms  $\chi_{\alpha}$  on  $V_{\alpha}^{\dagger}$ . Defining a third bundle

$$\mathfrak{G}^2 = B^2(B^1, R^{d_2}),$$

we then construct global forms  $\theta^2$ ,  $\Pi^2$ ,  $\chi^2$  on  $B^2$ . And so on. If the new forms are defined globally at any stage, the process terminates. The final bundle space  $B^r$  then carries a structure whose group is the identity. This solves the problem of local equivalence in the sense now to be explained.

Suppose that M' is a second manifold carrying a G-structure and denote quantities arising from M', corresponding to those already defined for M, by an accent. The two G-structures on M and M' are locally equivalent at points m and m' if there exists a local diffeomorphism of some neighbourhood  $U_{\alpha}$  of m onto a neighbourhood  $U_{\alpha}$  of m' such that

$$(\omega'_{\alpha'})^* = g\omega_{\alpha}$$

where \* denotes the dual mapping defined by the diffeomorphism and g is some differentiable function on  $U_{\alpha}$  with values in G. Two such diffeomorphisms are said to give the same local equivalence of the structures at m, m' if they coı̈ncide in some neighbourhood of m. It follows from the work of E. Cartan [1] that the local equivalences for the G-structures on M, M' can be obtained from the local equivalences for the identity-structures on  $B^r$ ,  $B^{\prime r}$ . Cartan gives a finite algorithm for finding the latter.

4. Application of the method of Cartan-Chern to conformal structure.

— We now return to our original notation and suppose that G is the group of non-zero scalar multiples of the orthogonal  $n \times n$  matrices. Its Lie algebra  $\mathcal{L}G$  is isomorphic with the algebra of  $n \times n$  matrices A such that

$$A + \tilde{A} \equiv \rho I$$

where p is any scalar.

We first construct the bundle  $\mathfrak{B} = B(M, G)$  and the form  $\theta$  on B as in the preceding paragraph. We can then find local forms  $\Pi_{\alpha}$  on  $V_{\alpha}$  in many ways so that the equation (3.2) becomes

$$d\theta + \Pi_{\alpha} \wedge \theta = 0.$$

In order to make a definite choice, we put

(4.1) 
$$\Pi_{\alpha} = g_{\alpha}^{-1} dg_{\alpha} + g_{\alpha}^{-1} (b^{*}\Omega_{\alpha}) g_{\alpha}$$

where, as in paragraph 2,  $\Omega_{\alpha}$  is the connection form of the local Riemannian metric  $\mathfrak{S}_{\alpha}\omega_{\alpha}$  on  $U_{\alpha}$ .  $\Pi_{\alpha}$  is then the corresponding local connection form on  $V_{\alpha}$ .

Suppose that  $\eta = \eta^i_{jh} \theta^h$  is any local 1-form with values in  $\mathcal{L}G$  and such that  $\eta \wedge \theta = 0$ . Then

$$\eta^l_{jh} + \eta^j_{ih} = 2 \lambda^h \delta_{ij}, \qquad \eta^l_{jh} - \eta^l_{hj} = 0.$$

These equations show that

$$egin{aligned} \eta^i_{jh} &= rac{1}{2} \left( \eta^h_{ij} + \eta^i_{hj} - \eta^h_{ji} - \eta^j_{hl} + \eta^i_{jh} + \eta^j_{ih} 
ight) \ &= \lambda^j \, \delta_{ih} - \lambda^l \delta_{jh} + \lambda^h \, \delta_{ij} \end{aligned}$$

and so it follows that

$$\eta = \theta \tilde{\lambda} - \lambda \tilde{\theta} + (\tilde{\lambda}\theta)I.$$

Thus any such form is determined by a function  $\lambda$  with values in  $R^n$ . In particular,  $\mathbf{H}_{\alpha} - \mathbf{H}_{\beta}$  will be determined by functions  $\lambda_{\alpha\beta}$  on  $V_{\alpha} \cap V_{\beta}$ ,

We downot calculate these functions explicitly at present. They form a cocycle on B and this defines a principal bundle  $\mathfrak{G}^1 = B^1(B, \mathbb{R}^n)$ .

On  $B^1$  we define global forms  $\theta^1$ ,  $\Pi^1$  where

$$( \begin{array}{l} \theta^{_{1}} = b^{_{1}\star}\theta \\ ( \mathbf{\Pi}^{_{1}} = b^{_{1}\star}\mathbf{\Pi}_{\alpha} - \theta^{_{1}}\tilde{\lambda}_{\alpha} + \lambda_{\alpha}\tilde{\theta}^{_{1}} - (\tilde{\lambda}_{\alpha}\theta^{_{1}})I. \end{array}$$

A calculation of their exterior derivatives gives

$$\begin{cases} d\theta^{1} = -\mathbf{H}^{1} \wedge \theta^{1}, \\ d\mathbf{H}^{1} = d\lambda_{x} \wedge \tilde{\theta}^{1} + \theta^{1} \wedge d\tilde{\lambda}_{x} - (d\lambda_{x} \wedge \tilde{\theta}^{1})\mathbf{I} - \mathbf{H}^{1} \wedge \mathbf{H}^{1} + \mathbf{a} + \mathbf{b}, \end{cases}$$

where (a) involves mixed products of components from  $\Pi^1$  and  $\theta^1$  and (b) involves products of components of  $\theta^1$ . Following the general method, we put

$$\chi_{\alpha} = d\lambda_{\alpha} + \chi'_{\alpha} + \chi''_{\alpha}$$

where  $\chi'_{\alpha}$  and  $\chi''_{\alpha}$  are 1-forms on  $V'_{\alpha}$  with values in  $R^n$  which are linear in the components of  $\theta^1$  and  $\Pi^1$  respectively. We can show that  $\chi'_{\alpha}$  and  $\chi''_{\alpha}$  are uniquely determined by requiring that

$$d\mathbf{H}^{1} = \chi_{\alpha} \wedge \mathbf{\tilde{\theta}}^{1} + \mathbf{\theta}^{1} \wedge \mathbf{\tilde{\chi}}_{\alpha} - (\mathbf{\tilde{\chi}}_{\alpha} \wedge \mathbf{\theta}^{1})I - \mathbf{H}^{1} \wedge \mathbf{H}^{1} + \mathbf{\Phi},$$

where the form  $\Phi = \frac{1}{2} \Phi^{l}_{jhk} \theta^{1h} \wedge \theta^{1k}$  satisfies the relations

$$\Phi^{i}_{ihi} = 0$$
.

Explicitly, we find that

$$\begin{cases} \chi_{\mathbf{x}}' = b^{1*}(\tilde{\mathbf{g}}_{\mathbf{x}}(b^*\psi_{\mathbf{x}})) = (\tilde{\lambda}_{\mathbf{x}}\theta^1)\lambda_{\mathbf{x}} + \frac{1}{2}(\tilde{\lambda}_{\mathbf{x}}\lambda_{\mathbf{x}})\theta^1, \\ \chi_{\mathbf{x}}'' = -\tilde{\mathbf{H}}^1\lambda_{\mathbf{x}} \end{cases}$$

and that  $\Phi = b^{**}(g_{\alpha}^{-1}(b^*C_{\alpha})g_{\alpha})$ . The local forms  $\psi_{\alpha}$  and  $C_{\alpha}$ , which arise from the Riemannian metric on  $U_{\alpha}$ , have been defined in paragraph 2.

From the general theory of paragraph 3, the local forms  $\chi_2$  define a global form  $\chi^1$  on  $B^1$  and hence  $\Phi$  is also defined globally. The forms  $\theta^1$ ,  $\Pi^1$ ,  $\chi^1$  contain  $n + \frac{1}{2}n(n-1) + 1 + n$  linearly independent components and so they define an identity-structure on  $B^1$ . This structure, as explained in paragraph 3, solves the problem of local equivalence.

5. The relation between the two theories. — Starting from a given conformally Riemannian structure on M, we constructed, in paragraph 2, global forms  $\omega$ ,  $\mu$ ,  $\Omega$  and  $\psi$  on H'' which defined a normal conformal connection. In paragraph h, we carried out the Chern process for the conformal structure and obtained global forms  $\theta^1$ ,  $\Pi^1$  and  $\chi^1$  on  $B^1$ . We shall set up a diffeomorphism mapping H'' onto  $B^1$  and then find the relation between these two sets of forms.

We must first calculate the functions  $\lambda_{\alpha\beta}$  on  $V_{\alpha} \cap V_{\beta}$  explicitly. From (4.2), we have

(5.1) trace 
$$(\mathbf{II}_{\alpha} - \mathbf{II}_{\beta}) = n \tilde{\lambda}_{\alpha\beta} \theta$$
.

Since the values of  $\Omega_{\alpha}$  are skew-symmetric matrices, it follows from (4.1), (3.3) and (2.1) that

trace 
$$(\mathbf{II}_{\alpha} - \mathbf{II}_{\beta}) \equiv \text{trace } (g_{\alpha}^{-1} dg_{\alpha} - g_{\beta}^{-1} dg_{\beta})$$
  
 $\equiv \text{trace } b^{*}(g_{\alpha\beta}^{-1} dg_{\alpha\beta})$   
 $\equiv \text{trace } b^{*}(d(\log a_{\alpha\beta})I + \tilde{A}_{\alpha\beta} dA_{\alpha\beta})$   
 $\equiv nb^{*}(d(\log a_{\alpha\beta})).$ 

Then using (2.2) and (3.1), we find that

trace 
$$(\mathbf{II}_{\alpha} - \mathbf{II}_{\beta}) = nb^{\star}(\mathbf{\tilde{q}}_{\alpha\beta}\omega_{\beta}) = n(b^{\star}\mathbf{\tilde{q}}_{\alpha\beta})g_{\beta}\theta.$$

Comparing this result with (5.1), it follows that

$$\lambda_{\alpha\beta} = \tilde{g}_{\beta}(b^{\star}q_{\alpha\beta}),$$

We recall from paragraph 2 that the bundle  $\partial \mathcal{E}''$  is defined by means of the cocycle  $K_{\alpha\beta}$  on M. Consequently H'' is the quotient of the sum  $\sum U_{\alpha} \times K''$  by the equivalence relation

$$(m_{\alpha}, k''_{\alpha}) \sim (m_{\beta}, k''_{\beta})$$
 if  $m_{\alpha} = m_{\beta}$ ,  $k''_{\alpha} = k'_{\alpha\beta} k''_{\beta}$ .

In paragraph 4 we defined  $\mathcal{B}$  by means of the cocycle  $g_{\alpha\beta}$  on M and  $\mathcal{B}^1$  by means of the cocycle  $\lambda_{\alpha\beta}$  on B. Combining these definitions and using (5.2), it follows that  $B^1$  is the quotient of the sum  $\sum U_{\alpha} \times G \times R^n$  by the equivalence relation

$$(m_{\alpha}, g_{\alpha}, \lambda_{\alpha}) \sim (m_{\beta}, g_{\beta}, \lambda_{\beta})$$

if  $m_{\alpha} = m_{\beta}$ ,  $g_{\alpha} = g_{\alpha\beta}g_{\beta}$ ,  $\lambda_{\alpha} = \lambda_{\beta} + \tilde{g}_{\beta}q_{\alpha\beta}$ . The functions  $k'_{\alpha\beta}$ ,  $g_{\alpha\beta}$  and  $q_{\alpha\beta}$  are all to be evaluated at  $m_{\alpha} = m_{\beta}$ .

We now set up a local diffeomorphism of  $U_{\alpha} \times K''$  onto  $U_{\alpha} \times G \times R''$ .

$$(m_{\alpha}, k''_{\alpha}) \rightarrow (m_{\alpha}, a_{\alpha} A_{\alpha}, q_{\alpha})$$

where  $a_{\alpha}$ ,  $A_{\alpha}$  and  $q_{\alpha}$  are obtained from the decomposition (2.3) for any element  $k_{\alpha}^{"}$  of  $K^{"}$ . It can be shown that these local diffeomorphisms commute with the above equivalence relations and so they define a global diffeomorphism of  $H^{"}$  onto  $B^{1}$ . Denoting the dual mapping on the forms in  $B^{1}$  by  $\bigstar$ , it follows that

$$\begin{cases} \bigstar \lambda_{\alpha} = q_{\alpha}, \\ \bigstar (b^{1\star} g_{\alpha}) = a_{\alpha} A_{\alpha}, \\ \bigstar (b^{1\star} b^{\star} \zeta_{\alpha}) = h''^{\star} \zeta_{\alpha} & \text{for any form } \zeta_{\alpha} \text{ on } U_{\alpha}. \end{cases}$$

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Using the definitions of the forms  $0^1$ ,  $\Pi^1$ ,  $\chi^1$  on  $B^1$  from paragraph 4 and the definitions of the forms  $\omega$ ,  $\mu$ ,  $\Omega$ ,  $\psi$  on H'' from paragraph 2, it is then easy to see that

$$\star \theta^{\scriptscriptstyle 1} = \omega, \quad \star \Pi^{\scriptscriptstyle 1} = \Omega + \mu I, \quad \star \chi^{\scriptscriptstyle 1} = \psi.$$

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F. BRICKELL and R. S. CLARK.
Department of Mathematics,
University of Southampton,
Southampton (Grande-Bretagne).