## Differential geometry

# Metrics on a closed surface of genus two which maximize the first eigenvalue of the Laplacian ${ }^{*}$ 

# Métriques sur une surface fermée de genre deux qui maximisent la première valeur propre du laplacien 

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#### Abstract

In this paper, we settle in the affirmative the Jakobson-Levitin-Nadirashvili-NigamPolterovich conjecture, stating that a certain singular metric on the Bolza surface, with area normalized, should maximize the first eigenvalue of the Laplacian.


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## R É S U M É

Dans cette Note, nous donnons une réponse positive à la conjecture de Jakobson-Levitin-Nadirashvili-Nigam-Polterovich, en montrant qu'une certaine métrique singulière sur la surface de Bolza, d'aire normalisée, maximise la première valeur propre du laplacien.
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## 1. Introduction

Let $M$ be a closed surface, that is, a compact surface without boundary. Throughout this paper, we assume that $M$ is orientable. For a Riemannian metric $\mathrm{ds}^{2}$ on $M$, let

$$
\Lambda\left(\mathrm{d} s^{2}\right):=\lambda_{1}\left(\mathrm{~d} s^{2}\right) \cdot \operatorname{Area}\left(\mathrm{d} s^{2}\right)
$$

where $\lambda_{1}\left(\mathrm{ds} s^{2}\right)$ is the first positive eigenvalue of the Laplacian and $\operatorname{Area}\left(\mathrm{d} s^{2}\right)$ is the area of $M$, both with respect to $\mathrm{ds} s^{2}$. Regarding the upper bound of the quantity $\Lambda\left(\mathrm{ds}^{2}\right)$, the following results are well known.

[^0]Fact. (i) (Hersch [6]) For any metric $\mathrm{ds}^{2}$ on the sphere $S^{2}, \Lambda\left(\mathrm{ds} s^{2}\right) \leq 8 \pi$ holds.
(ii) (Yang-Yau [16]) If $M$ admits a nonconstant meromorphic function $\left(M, d s^{2}\right) \rightarrow \overline{\mathbb{C}}=\mathbb{C} \cup\{\infty\}$ of degree $d$, then $\Lambda\left(d s^{2}\right) \leq 8 \pi \cdot d$ holds. In particular, if $\gamma$ is the genus of $M$, then for any metric $\mathrm{ds}^{2}$ on $M$, we have

$$
\begin{equation*}
\Lambda\left(\mathrm{ds}^{2}\right) \leq 8 \pi \cdot\left[\frac{\gamma+3}{2}\right] \tag{1}
\end{equation*}
$$

The inequality in the statement (i) is sharp as equality holds for the standard metric of $S^{2}$. On the other hand, Nadirashvili [9] found the sharp bound $8 \pi^{2} / \sqrt{3}$ of $\Lambda\left(\mathrm{d} s^{2}\right)$ for metrics $\mathrm{ds} s^{2}$ on the torus $T^{2}$. Thus, the inequality (1) is not sharp when $\gamma=1$.

When $\gamma=2$, the inequality (1) becomes $\Lambda\left(d s^{2}\right) \leq 16 \pi$. Jakobson-Levitin-Nadirashvili-Nigam-Polterovich [7] focused their attention on the following metric. Let $B$ be the closed Riemann surface of genus two, called the Bolza surface, defined as the smooth completion of the affine complex algebraic curve $w^{2}=z\left(z^{4}+1\right)$. Topologically, $B$ is the one-point compactification of the affine curve:

$$
B=\left\{(z, w) \in \mathbb{C}^{2} \mid w^{2}=z\left(z^{4}+1\right)\right\} \cup\{(\infty, \infty)\}
$$

Let $g_{B}: B \rightarrow \overline{\mathbb{C}}$ be the meromorphic function of degree two given by $g_{B}(z, w)=z$. If we set $\mathrm{d} s_{B}^{2}=g_{B^{*}}{ }^{*} s_{S^{2}}^{2}$, where $\mathrm{d} s_{S^{2}}^{2}$ is the standard metric of $S^{2}=\overline{\mathbb{C}}$, then $\mathrm{d} s_{B}^{2}$ is a singular Riemannian metric that degenerates exactly at the ramification points of $g_{B}$. Since the map $g_{B}: B \rightarrow \overline{\mathbb{C}}$ is a two-sheeted branched covering, we have Area $\left(\mathrm{d} s_{B}^{2}\right)=8 \pi$.

Conjecture (Jakobson et al. [7]). $\lambda_{1}\left(d s_{B}^{2}\right)=2$ should hold. Therefore, $\Lambda\left(\mathrm{ds}{ }_{B}^{2}\right)=16 \pi$.
For $0<\theta<\pi / 2$, let $B_{\theta}$ be the Riemann surface of genus two defined as the smooth completion of the affine complex algebraic curve $w^{2}=z\left(z^{4}+2 \cos 2 \theta \cdot z^{2}+1\right)$ :

$$
B_{\theta}=\left\{(z, w) \in \mathbb{C}^{2} \mid w^{2}=z\left(z^{4}+2 \cos 2 \theta \cdot z^{2}+1\right)\right\} \cup\{(\infty, \infty)\}
$$

Note that $B_{\pi / 4}=B$. Let $\mathrm{d} s_{\theta}^{2}$ denote the pull-back of the standard metric of $S^{2}=\overline{\mathbb{C}}$ by the meromorphic function $g_{\theta}: B_{\theta} \ni$ $(z, w) \mapsto z \in \overline{\mathbb{C}}$.

In this paper, we prove the following theorem, and thereby settle the above conjecture in the affirmative.
Main Theorem. There exists $\theta_{1} \approx 0.65$ so that for $\theta_{1} \leq \theta \leq \pi / 2-\theta_{1}$, we have $\lambda_{1}\left(\mathrm{ds}_{B_{\theta}}^{2}\right)=2$ and therefore $\Lambda\left(\mathrm{ds} s_{\theta}^{2}\right)=16 \pi$.
Note that $16 \pi$ is a degenerate maximum for $\Lambda$ in the genus-two case, as predicted in [7]. It is also remarked in [7] that the conjecture implies that the inequality $\Lambda\left(\mathrm{ds}^{2}\right) \leq 16 \pi$ is sharp in the class of smooth metrics, although the equality may not be attained. It is worth mentioning that the Lawson minimal surface of genus two in $S^{3}$ has $\lambda_{1}=2$ [2] and Area $\approx 21.91$ [5], and therefore $\Lambda \approx 43.82<16 \pi$.

For recent progress on the existence of $\Lambda$-maximizing metrics on a closed surface, see [10,14].
In $\S 1$, we explain the relation of the above conjecture to the problem of computing the Morse index of a minimal surface in Euclidean three-space. After that, in §2, we prove the Main Theorem, assuming two technical lemmas, whose proofs are postponed to $\S 3$ and $\S 4$. The paper concludes with two appendices.

## 2. Index and nullity of a meromorphic function

The problem of estimating and computing the Morse (instability) index of a complete minimal surface in $\mathbb{R}^{3}$ (and other flat three-spaces) has been studied by various authors. In this section, we explain that the conjecture of Jakobson et al. is closely related to this problem.

Let $M$ be an orientable complete minimal surface in $\mathbb{R}^{3} . M$ is said to be stable if the second variation of area for any compactly supported variation of $M$ is nonnegative, and the plane is the only stable one. For non-planar $M$, we define the Morse index of $M$, $\operatorname{Ind}(M)$, as follows: for a relatively compact domain $\Omega \subset M$, $\operatorname{Ind}(\Omega)$ is defined as the maximal dimension of a subspace $V \subset C_{0}^{\infty}(\Omega)$ satisfying

$$
\int_{\Omega}\left(|d u|^{2}+2 K u^{2}\right) \mathrm{d} a<0, \quad \forall u \in V \backslash\{0\}
$$

where $K$ and $\mathrm{d} a$ are the Gaussian curvature and the area element of $M$, respectively. Note that $\operatorname{Ind}(\Omega)$ is necessarily finite. We then define

$$
\operatorname{Ind}(M)=\sup _{\Omega} \operatorname{Ind}(\Omega)
$$

where the supremum is taken over all relatively compact domains $\Omega \subset M$. While $\operatorname{Ind}(M)$ so defined may become infinity, it was proved by Fischer-Colbrie [4] that

$$
\operatorname{Ind}(M)<\infty \quad \Leftrightarrow \quad \int_{M}(-K) \mathrm{d} a<\infty
$$

Therefore, in studying $\operatorname{Ind}(M)$ quantitatively, we may assume that $\int_{M}(-K) \mathrm{d} a<\infty$. In this case, $M$ is conformally equivalent to a compact Riemann surface $\bar{M}$ with finitely many punctures and the Gauss map of $M, g: M \rightarrow \overline{\mathbb{C}}$, extends to a meromorphic function $\bar{g}: \bar{M} \rightarrow \overline{\mathbb{C}}$. (This is a classical result due to Osserman [13].)

In general, for a nonconstant meromorphic function $g: M \rightarrow \overline{\mathbb{C}}$ on a compact Riemann surface $M$, we pull back the standard metric of $\overline{\mathbb{C}}=S^{2}$ by $g$ and obtain a singular metric $d s_{g}^{2}$ (as we did to get $d s_{B}^{2}$ ). Let $\Delta_{g}$ denote the Laplacian defined with respect to $\mathrm{ds} s_{g}^{2}$, and $\operatorname{Ind}(g)$ (resp. $\left.\operatorname{Nul}(g)\right)$ the number of eigenvalues of $-\Delta_{g}$ less than 2 counted with multiplicity (resp. the multiplicity of eigenvalue 2 of $-\Delta_{g}$ ).

Proposition 1 (Fischer-Colbrie [4], Ejiri-Kotani [3], Montiel-Ros [8]). The Morse index $\operatorname{Ind}(M)$ of a complete minimal surface $M$ in $\mathbb{R}^{3}$ of finite total curvature coincides with the index $\operatorname{Ind}(\bar{g})$ of the extended Gauss map $\bar{g}$. The nullity $\operatorname{Nul}(\bar{g})$ equals the dimension of the vector space of all bounded Jacobi fields on $M$.

Since constant functions are necessarily eigenfunctions of the eigenvalue 0 of $-\Delta_{g}$, we have $\operatorname{Ind}(g) \geq 1$. The conjecture of Jakobson et al. asserts that when $g=g_{B}$, the second least eigenvalue of $-\Delta_{g_{B}}$ should equal 2 , and so it is equivalent to asserting that $\operatorname{Ind}\left(g_{B}\right)=1$.

## 3. Proof of the Main Theorem

In this section, we prove the Main Theorem, assuming two technical Lemmas 3 and 5 . The proofs of these lemmas are contained in $\S 3$ and $\S 4$. Note that the equation of $B_{\theta}$ can be rewritten as

$$
w^{2}=z\left(z-\mathrm{e}^{\mathrm{i}(\pi / 2-\theta)}\right)\left(z-\mathrm{e}^{\mathrm{i}(\pi / 2+\theta)}\right)\left(z-\mathrm{e}^{-\mathrm{i}(\pi / 2-\theta)}\right)\left(z-\mathrm{e}^{-\mathrm{i}(\pi / 2+\theta)}\right)
$$

Let $g_{\theta}$ and $d s_{\theta}^{2}$ be as in the introduction, and $\Delta_{\theta}$ the Laplacian corresponding to $d s_{\theta}^{2}$. The meromorphic function $g_{\theta}: B_{\theta} \rightarrow \overline{\mathbb{C}}$ gives a two-sheeted branched covering that ramifies at the six points $(0,0),\left(\mathrm{e}^{ \pm \mathrm{i}(\pi / 2 \pm \theta)}, 0\right),(\infty, \infty)$. ds $s_{\theta}^{2}$ is a singular metric that degenerates precisely at the six ramification points of $g_{\theta}$. Define three great circular arcs $C_{1}, C_{2}, C_{3}$ on $S^{2}=\overline{\mathbb{C}}$ by

$$
\begin{aligned}
& C_{1}=\{t \mid t \geq 0\} \cup\{\infty\}, \quad C_{2}=\left\{\mathrm{e}^{\mathrm{i}(\pi / 2+t)} \mid-\theta \leq t \leq \theta\right\}, \\
& C_{3}=\left\{\mathrm{e}^{-\mathrm{i}(\pi / 2+t)} \mid-\theta \leq t \leq \theta\right\} .
\end{aligned}
$$

Then ( $B_{\theta}, \mathrm{ds} s_{\theta}^{2}$ ) can be represented as the gluing of two copies of $\left(S^{2}, d s_{S^{2}}^{2}\right)$ along $C_{1}, C_{2}, C_{3}$. As $\theta \rightarrow 0$, the two arcs $C_{2}, C_{3}$ collapse to points, and by neglecting the contact at these two points, we obtain the metric that is the gluing of two copies of ( $S^{2}, \mathrm{~d} s_{S^{2}}^{2}$ ) along $C_{1}$. The last metric, denoted by $\mathrm{d} s_{0}^{2}$, is nothing but the pull-back of $\mathrm{d} s_{S^{2}}^{2}$ by the degree two rational function $g_{0}: \overline{\mathbb{C}} \ni z \mapsto z^{2} \in \overline{\mathbb{C}}$. Let $\Delta_{0}$ be the Laplacian defined with respect to $\mathrm{d} s_{0}^{2}$. Then we have the following lemma regarding the eigenvalues of $-\Delta_{\theta}$ and $-\Delta_{0}$.

Lemma 2. For every positive integer $k$, the $k$-th eigenvalue $\lambda_{k}\left(\mathrm{ds} s_{\theta}^{2}\right.$ ) of $-\Delta_{\theta}$ is continuous in $\theta$, and as $\theta \rightarrow 0$ it converges to the $k$-th eigenvalue $\lambda_{k}\left(\mathrm{ds}_{0}^{2}\right)$ of $-\Delta_{0}$.

This lemma may be proved by arguments similar to those in the proof of [12, Theorem 1].
In [11], by computing all the eigenvalues of $-\Delta_{0}$ explicitly, it is shown that $\operatorname{Ind}\left(g_{0}\right)=3$ and $\operatorname{Nul}\left(g_{0}\right)=3$. On the other hand, it is known that $\operatorname{Nul}(g) \geq 3$ for any nonconstant meromorphic function $g$. In fact, the pull-back of three independent eigenfunctions of the eigenvalue 2 of $-\Delta_{S^{2}}$, the Laplacian with respect to $d s_{S^{2}}^{2}$, by $g$ give eigenfunctions of the eigenvalue 2 of $-\Delta_{g}$. From these facts and from Lemma 2, it follows that $\operatorname{Ind}\left(g_{\theta}\right)=3$ and $\operatorname{Nul}\left(g_{\theta}\right)=3$ for $\theta$ sufficiently close to 0 .

We now observe the change of $\operatorname{Nul}\left(g_{\theta}\right)$ as $\theta$ increases up to $\pi / 4$. To do this, we use the work of Ejiri-Kotani [3] and Montiel-Ros [8]. If $g$ is a nonconstant meromorphic function such that $\operatorname{Nul}(g)>3$, then there exists an extra eigenfunction, that is, an eigenfunction of the eigenvalue 2 of $-\Delta_{g}$ that is not the pull-back of an eigenfunction of the eigenvalue 2 of $-\Delta_{S^{2}}$ by $g$. As shown in [3,8], any extra eigenfunction can be written as the support function (that is, the inner product of the position vector field and the unit normal vector field) of a complete branched minimal surface of finite total curvature whose extended Gauss map is $g$ and whose ends are contained in the ramification locus of $g$ and are all planar. By using the Weierstrass representation, we can express such a minimal surface as follows. Let $P$ and $B=\sum_{j=1}^{l} e_{j} p_{j}$ be the polar
and ramification divisors of $g$, respectively, where $e_{j}$ is the multiplicity with which $g$ takes its value at $p_{j}$. Set $D=B-2 P$. Suppose that there exists a non-zero $\omega \in H^{0}\left(M, K_{M} \otimes D\right)$ satisfying

$$
\begin{equation*}
\operatorname{Res}_{p_{j}} \omega=0, \quad 1 \leq \forall j \leq l, \tag{2}
\end{equation*}
$$

and

$$
\begin{equation*}
\Re \int_{\ell} t\left(1-g^{2}, \mathrm{i}\left(1+g^{2}\right), 2 g\right) \omega=\mathbf{o}, \quad \forall \ell \in H_{1}(M, \mathbb{Z}) \tag{3}
\end{equation*}
$$

where $K_{M}$ is the canonical divisor of $M$. Then, for any such $\omega$,

$$
X_{\omega}(p)=\Re \int_{p_{0}}^{p} t\left(1-g^{2}, \mathrm{i}\left(1+g^{2}\right), 2 g\right) \omega
$$

gives a minimal surface with the above properties.
We now apply the general result as above to ( $B_{\theta}, g_{\theta}$ ). We can determine the values of $\theta$ for which there exists a non-zero $\omega \in H^{0}\left(B_{\theta}, K_{B_{\theta}} \otimes D\right)$ satisfying (2) and (3). In fact, we have

## Lemma 3. Set

$$
\begin{aligned}
& A=\int_{0}^{\infty} \frac{\mathrm{d} t}{\sqrt{t\left(t^{4}+2 \cos 2 \theta \cdot t^{2}+1\right)}}, \quad B=\int_{0}^{\infty} \frac{\mathrm{d} t}{\sqrt{t\left(t^{4}-2 \cos 2 \theta \cdot t^{2}+1\right)}}, \\
& C=\int_{0}^{\infty} \frac{t^{3} \mathrm{~d} t}{{\sqrt{t\left(t^{4}+2 \cos 2 \theta \cdot t^{2}+1\right)}}^{3}}, \quad D=\int_{0}^{\infty} \frac{t^{3} \mathrm{~d} t}{{\sqrt{t\left(t^{4}-2 \cos 2 \theta \cdot t^{2}+1\right)}}^{3}} .
\end{aligned}
$$

Let $\theta_{1}(\approx 0.65)$ be the unique solution to

$$
A\left(B^{2}+16 D^{2} \sin ^{2} 2 \theta\right)+8(A D+B C)\left(B \cos 2 \theta-4 D \sin ^{2} 2 \theta\right)=0
$$

and set $\theta_{2}=\pi / 2-\theta_{1}(\approx 0.91)$. Then there exists a non-zero $\omega \in H^{0}\left(B_{\theta}, K_{B_{\theta}} \otimes D\right)$ satisfying (2) and (3) if and only if $\theta=\theta_{1}, \theta_{2}$. If $\theta=\theta_{1}$, then any such $\omega$ is given by a real linear combination of

$$
\begin{aligned}
\omega_{1} & :=-\frac{A D+3 B C}{4(A D+B C)} \frac{\mathrm{d} z}{w}-\frac{A D+3 B C}{4(A D+B C)} \frac{\mathrm{d} z}{w^{3}}+\frac{z}{w^{3}} \mathrm{~d} z \\
& +\frac{A B+(A D-B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{2}}{w^{3}} \mathrm{~d} z+\frac{A B+2(A D+B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{3}}{w^{3}} \mathrm{~d} z \\
& +\frac{3 A D+B C}{4(A D+B C)} \frac{z^{4}}{w^{3}} \mathrm{~d} z, \\
\omega_{2} & :=i\left(-\frac{A D+3 B C}{4(A D+B C)} \frac{\mathrm{d} z}{w}+\frac{A D+3 B C}{4(A D+B C)} \frac{\mathrm{d} z}{w^{3}}+\frac{z}{w^{3}} \mathrm{~d} z\right. \\
& -\frac{A B+(A D-B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{2}}{w^{3}} \mathrm{~d} z+\frac{A B+2(A D+B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{3}}{w^{3}} \mathrm{~d} z \\
& \left.-\frac{3 A D+B C}{4(A D+B C)} \frac{z^{4}}{w^{3}} \mathrm{~d} z\right) .
\end{aligned}
$$

(We can obtain a similar assertion for $\theta=\theta_{2}$.)

The lemma implies that there are two independent extra eigenfunctions when $\theta=\theta_{1}, \theta_{2}$. Thus we obtain Proposition 4.

## Proposition 4.

$$
\operatorname{Nul}\left(g_{\theta}\right)= \begin{cases}5, & \theta=\theta_{1}, \theta_{2}  \tag{4}\\ 3, & \theta \neq \theta_{1}, \theta_{2}\end{cases}
$$



Fig. 1. The fundamental domain $\Omega$ for $H$.


Fig. 2. Fixed point sets of $s_{1}, j \circ s_{1}, s_{3}, j \circ s_{3}$.

To see how Ind $\left(g_{\theta}\right)$ changes as $\theta$ increases and passes $\theta_{1}$, we use symmetries of $B_{\theta}$. Let $j: B_{\theta} \rightarrow B_{\theta}$ be the hyperelliptic involution given by $j(z, w)=(z,-w)$, and $s_{1}, s_{2}, s_{3}: B_{\theta} \rightarrow B_{\theta}$ the anti-holomorphic involutions given by $s_{1}(z, w)=(\bar{z}, \bar{w})$, $s_{2}(z, w)=(-\bar{z}, i \bar{w}), s_{3}(z, w)=\left(1 / \bar{z}, \bar{w} / \bar{z}^{3}\right)$. We have

$$
s_{1} \circ s_{2}=j \circ s_{2} \circ s_{1}, \quad s_{2} \circ s_{3}=s_{3} \circ s_{2}, \quad s_{3} \circ s_{1}=s_{1} \circ s_{3}
$$

Thus the three involutions $j, s_{1}, s_{3}$ of $B_{\theta}$ commute with one another, and the group of symmetries, $H$, generated by them is an abelian group of order eight. A fundamental domain for the action of $H$ on $B_{\theta}$ is given by the intersection of the upper half plane and the unit disk, denoted by $\Omega$ (see Fig. 1).

Recall that $B_{\theta}$ is the gluing of two copies of $\overline{\mathbb{C}}$. The fixed point sets of the anti-holomorphic involutions $s_{1}, j \circ s_{1}, s_{3}$, $j \circ s_{3}$ are as follows. (See Fig. 2.)

- The fixed point set of $s_{1}$ is the red half-line on the real axis,
- The fixed point set of $j \circ s_{1}$ is the blue half-line on the real axis,
- The fixed point set of $s_{3}$ is the union of the red arcs on the unit circle,
- The fixed point set of $j \circ s_{3}$ is the union of the blue arcs on the unit circle.

For example, $s_{1}(z, w)=(z, w)$ if and only if $z(=: x), w(=: y)$ are real. Since

$$
y^{2}=x\left(x^{4}+2 \cos 2 \theta \cdot x^{2}+1\right)=x\left\{\left(x^{2}+\cos 2 \theta\right)^{2}+\sin ^{2} 2 \theta\right\} \geq 0
$$

and $\left(x^{2}+\cos 2 \theta\right)^{2}+\sin ^{2} 2 \theta>0$, one must have $x \geq 0$.
Since $H$ is abelian and preserves $d s_{\theta}^{2}$, each eigenspace of $-\Delta_{\theta}$ is invariant under the action of $H$ and spanned by simultaneous eigenvectors for all $s \in H$. Let $u_{i}, i=1,2$, be the support functions of the branched minimal immersions $X_{\omega_{i}}$, in whose definition we choose $p_{0}=(1, \sqrt{2+2 \cos 2 \theta})$ as the base point. They are extra eigenfunctions for $\theta=\theta_{1}$. The following lemma shows how $H$ acts on $u_{1}, u_{2}$.

## Lemma 5.

$$
\begin{aligned}
& s_{1}^{*} u_{1}=u_{1}, \quad s_{3}^{*} u_{1}=u_{1}, \quad j^{*} u_{1}=-u_{1}+\left\langle c_{1}, N\right\rangle, \\
& s_{1}^{*} u_{2}=-u_{2}, \quad s_{3}^{*} u_{2}=u_{2}, \quad j^{*} u_{2}=-u_{2}+\left\langle c_{2}, N\right\rangle,
\end{aligned}
$$

where $c_{i} \in \mathbb{R}^{3}, i=1,2$, and $N$ is the unit normal vector field of $X_{\omega_{i}}$.


Fig. 3. Fixed point sets in $\partial \Omega$.
In order to get extra eigenfunctions that behave properly with respect to the actions of $j \circ s_{1}$ and $j \circ s_{3}$, we set

$$
\begin{aligned}
& v_{1}=u_{1}-\left(j \circ s_{1}\right)^{*} u_{1}-\left(j \circ s_{3}\right)^{*} u_{1}+\left(j \circ s_{1}\right)^{*} \circ\left(j \circ s_{3}\right)^{*} u_{1}, \\
& v_{2}=u_{2}+\left(j \circ s_{1}\right)^{*} u_{2}-\left(j \circ s_{3}\right)^{*} u_{2}-\left(j \circ s_{1}\right)^{*} \circ\left(j \circ s_{3}\right)^{*} u_{2}
\end{aligned}
$$

By Lemma 5, we have

$$
\begin{array}{ll}
s_{1}^{*} v_{1}=v_{1}, & \left(j \circ s_{1}\right)^{*} v_{1}=-v_{1},
\end{array} \quad s_{3}^{*} v_{1}=v_{1}, \quad\left(j \circ s_{3}\right)^{*} v_{1}=-v_{1}, ~ 子 v_{1}^{*}, \quad\left(j \circ s_{1}\right)^{*} v_{2}=v_{2}, \quad s_{3}^{*} v_{2}=v_{2}, \quad\left(j \circ s_{3}\right)^{*} v_{2}=-v_{2} .
$$

Henceforth, we regard $v_{1}$ and $v_{2}$ as functions on $\Omega$. (See Fig. 3.) Then the preceding observations mean that $v_{1}$ satisfies the Dirichlet (resp. Neumann) condition on the blue (resp. red) segments in the unit circle and on the blue (resp. red) segment in the real axis. As $\theta$ increases, the blue (resp. red) segment in the unit circle becomes longer (resp. shorter). Hence, by the variational characterization of eigenvalues, the eigenvalues of the Laplacian in $\Omega$ under the boundary conditions as above monotonically increase. Similarly, $v_{2}$ satisfies the Dirichlet (resp. Neumann) condition on the blue (resp. red) segment in the unit circle and on the red (resp. blue) segment in the real axis, and therefore the eigenvalues of the Laplacian in $\Omega$ under the boundary conditions of $v_{2}$ also monotonically increase.

The two assertions we just made mean that there exist two independent eigenfunctions of $-\Delta_{\theta}$ with the same type of symmetry as $v_{1}$ and $v_{2}$, respectively, such that the corresponding eigenvalues increase monotonically and continuously. On the other hand, for $0<\theta<\theta_{2}$, extra eigenfunctions with the other types of symmetry do not occur. Hence, the number of the eigenvalues of $-\Delta_{\theta}$ less than 2 , whose eigenfunctions have the other types of symmetry, remains unchanged throughout ( $0, \theta_{2}$ ). (Here we use the continuity of eigenvalues in $\theta$ again.)

We may now conclude that as $\theta$ increases and passes $\theta_{1}$, two eigenvalues of $-\Delta_{\theta}$ will monotonically increase and pass 2 upward, and thus the number of eigenvalues less than 2 decreases by two. One can also verify that if $\theta$ increases further and passes $\theta_{2}$, then two eigenvalues of $-\Delta_{\theta}$ will decrease and pass 2 downward, and the number of eigenvalues less than 2 increases by two. To summarize, we have proved the following

## Theorem 6.

$$
\operatorname{Ind}\left(g_{\theta}\right)= \begin{cases}3, & 0<\theta<\theta_{1}  \tag{5}\\ 1, & \theta_{1} \leq \theta \leq \theta_{2} \\ 3, & \theta_{2}<\theta<\pi / 2\end{cases}
$$

This theorem implies the Main Theorem.

## 4. Proof of Lemma 3

This section is devoted to the proof of Lemma 3.
Recall that $K_{B_{\theta}}$ is the canonical divisor of $B_{\theta}$ and $D=B-2 P$, where $P$ and $B=\sum_{j=1}^{l} e_{j} p_{j}$ are the polar and ramification divisors of $g_{\theta}$, respectively. Let $\widehat{H}\left(g_{\theta}\right)$ denote the set of all $\omega \in H^{0}\left(B_{\theta}, K_{B_{\theta}} \otimes D\right)$ satisfying

$$
\begin{equation*}
\operatorname{Res}_{p_{j}} \omega=0, \quad 1 \leq \forall j \leq l \tag{6}
\end{equation*}
$$

and $H\left(g_{\theta}\right)$ the set of all $\omega \in \widehat{H}\left(g_{\theta}\right)$ satisfying

$$
\begin{equation*}
\Re \int_{\ell} t\left(1-g_{\theta}^{2}, \mathrm{i}\left(1+g_{\theta}^{2}\right), 2 g_{\theta}\right) \omega=\mathbf{0}, \quad \forall \ell \in H_{1}\left(B_{\theta}, \mathbb{Z}\right) \tag{7}
\end{equation*}
$$

Note that $\widehat{H}\left(g_{\theta}\right)$ is a complex vector space. We should determine the values of $\theta$ for which $H\left(g_{\theta}\right) \neq\{0\}$.

We first find a basis for $\widehat{H}\left(g_{\theta}\right)$. The polar and ramification divisors of $g_{\theta}$ are given by

$$
P=2(\infty, \infty), \quad B=2(0,0)+2\left(\mathrm{e}^{ \pm \mathrm{i}(\pi / 2 \pm \theta)}, 0\right)+2(\infty, \infty)
$$

and therefore

$$
D=2(0,0)+2\left(\mathrm{e}^{ \pm \mathrm{i}(\pi / 2 \pm \theta)}, 0\right)-2(\infty, \infty)
$$

By the Riemann-Roch theorem, $H^{0}\left(B_{\theta}, K_{B_{\theta}} \otimes D\right)$ has dimension nine, and

$$
\left\{\frac{\mathrm{d} z}{w}, \frac{\mathrm{~d} z}{w^{2}}, \frac{z}{w^{2}} \mathrm{~d} z, \frac{z^{2}}{w^{2}} \mathrm{~d} z, \frac{\mathrm{~d} z}{w^{3}}, \frac{z}{w^{3}} \mathrm{~d} z, \frac{z^{2}}{w^{3}} \mathrm{~d} z, \frac{z^{3}}{w^{3}} \mathrm{~d} z, \frac{z^{4}}{w^{3}} \mathrm{~d} z\right\}
$$

is a basis for it. It is easy to verify that

$$
\left\{\frac{\mathrm{d} z}{w}, \frac{\mathrm{~d} z}{w^{3}}, \frac{z}{w^{3}} \mathrm{~d} z, \frac{z^{2}}{w^{3}} \mathrm{~d} z, \frac{z^{3}}{w^{3}} \mathrm{~d} z, \frac{z^{4}}{w^{3}} \mathrm{~d} z\right\}
$$

is a basis for $\widehat{H}\left(g_{\theta}\right)$. Therefore, $\omega \in \widehat{H}\left(g_{\theta}\right)$ has the form

$$
\begin{equation*}
\omega=\alpha_{1} \frac{\mathrm{~d} z}{w}+\alpha_{2} \frac{\mathrm{~d} z}{w^{3}}+\alpha_{3} \frac{z}{w^{3}} \mathrm{~d} z+\alpha_{4} \frac{z^{2}}{w^{3}} \mathrm{~d} z+\alpha_{5} \frac{z^{3}}{w^{3}} \mathrm{~d} z+\alpha_{6} \frac{z^{4}}{w^{3}} \mathrm{~d} z \tag{8}
\end{equation*}
$$

where $\alpha_{1}, \ldots, \alpha_{6}$ are complex numbers.
We now consider the period condition (7). First, we express the above basis elements of $\widehat{H}\left(g_{\theta}\right)$ as linear combinations of the abelian differentials of the second kind $\mathrm{d} z / w, z \mathrm{~d} z / w, z^{3} \mathrm{~d} z / w^{3}, z^{4} \mathrm{~d} z / w^{3}$ up to exact forms. It is easy to show that

$$
\begin{align*}
d\left(z^{p} w^{q}\right) & =\frac{1}{2} z^{p-1} w^{q-2}\left\{(2 p+5 q) w^{2}-4 q \cos 2 \theta \cdot z^{3}-4 q z\right\} \mathrm{d} z  \tag{9}\\
& =\frac{1}{2} z^{p} w^{q-2}\left\{(2 p+5 q) z^{4}+2 \cos 2 \theta \cdot(2 p+3 q) z^{2}+2 p+q\right\} \mathrm{d} z \tag{10}
\end{align*}
$$

For two meromorphic one-forms $\eta_{1}, \eta_{2}$ on $B_{\theta}$, we write $\eta_{1} \sim \eta_{2}$ if there exists a meromorphic function $f$ on $B_{\theta}$ such that $\eta_{1}=\eta_{2}+d f$. By using (9), (10) we deduce the following relations:

$$
\begin{align*}
& \frac{z}{w^{3}} \mathrm{~d} z \sim \frac{3}{4} \frac{\mathrm{~d} z}{w}-\cos 2 \theta \cdot \frac{z^{3}}{w^{3}} \mathrm{~d} z  \tag{11}\\
& \frac{z^{2}}{w^{3}} \mathrm{~d} z \sim \frac{1}{4} \frac{z}{w} \mathrm{~d} z-\cos 2 \theta \cdot \frac{z^{4}}{w^{3}} \mathrm{~d} z  \tag{12}\\
& \frac{\mathrm{~d} z}{w^{3}} \sim-\frac{3}{2} \cos 2 \theta \cdot \frac{z}{w} \mathrm{~d} z+\left(-5+6 \cos ^{2} 2 \theta\right) \frac{z^{4}}{w^{3}} \mathrm{~d} z  \tag{13}\\
& \frac{z^{5}}{w^{3}} \mathrm{~d} z \sim \frac{1}{4} \frac{\mathrm{~d} z}{w}-\cos 2 \theta \cdot \frac{z^{3}}{w^{3}} \mathrm{~d} z  \tag{14}\\
& \frac{z^{6}}{w^{3}} \mathrm{~d} z \sim \frac{3}{4} \frac{z}{w} \mathrm{~d} z-\cos 2 \theta \cdot \frac{z^{4}}{w^{3}} \mathrm{~d} z  \tag{15}\\
& \frac{z^{2}}{w} \mathrm{~d} z \sim-\cos 2 \theta \cdot \frac{\mathrm{~d} z}{w}-4 \sin ^{2} 2 \theta \cdot \frac{z^{3}}{w^{3}} \mathrm{~d} z \tag{16}
\end{align*}
$$

In fact, (11) and (12) follow immediately from (9) with choices $(p, q)=(1,-1)$ and $(p, q)=(2,-1)$, respectively. (13) follows by using (10) with $(p, q)=(0,-1)$ and then applying (12). (14) and (15) follow by substituting $z^{5}=$ $w^{2}-2 \cos 2 \theta \cdot z^{3}-z$ and then applying (11) and (12), respectively. Finally, (16) follows by using (9) with $(p, q)=(3,-1)$ and then applying (14).

For

$$
\omega=\alpha_{1} \frac{\mathrm{~d} z}{w}+\alpha_{2} \frac{\mathrm{~d} z}{w^{3}}+\alpha_{3} \frac{z}{w^{3}} \mathrm{~d} z+\alpha_{4} \frac{z^{2}}{w^{3}} \mathrm{~d} z+\alpha_{5} \frac{z^{3}}{w^{3}} \mathrm{~d} z+\alpha_{6} \frac{z^{4}}{w^{3}} \mathrm{~d} z
$$

as in (8), we find by using the above relations

$$
\begin{align*}
\omega \underbrace{\sim}_{(11),(12),(13)} & \left(\alpha_{1}+\frac{3}{4} \alpha_{3}\right) \frac{\mathrm{d} z}{w}+\left(-\frac{3}{2} \cos 2 \theta \cdot \alpha_{2}+\frac{\alpha_{4}}{4}\right) \frac{z}{w} \mathrm{~d} z  \tag{17}\\
& +\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right) \frac{z^{3}}{w^{3}} \mathrm{~d} z \\
& +\left(\left(-5+6 \cos ^{2} 2 \theta\right) \alpha_{2}-\cos 2 \theta \cdot \alpha_{4}+\alpha_{6}\right) \frac{z^{4}}{w^{3}} \mathrm{~d} z \\
z \omega \underbrace{\sim}_{(11),(12),(14)} & \left(\frac{3}{4} \alpha_{2}+\frac{\alpha_{6}}{4}\right) \frac{\mathrm{d} z}{w}+\left(\alpha_{1}+\frac{\alpha_{3}}{4}\right) \frac{z}{w} \mathrm{~d} z  \tag{18}\\
& +\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right) \frac{z^{3}}{w^{3}} \mathrm{~d} z \\
& +\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right) \frac{z^{4}}{w^{3}} \mathrm{~d} z
\end{align*}
$$

and

$$
\begin{align*}
z^{2} \omega \underbrace{\sim}_{(12),(14),(15),(16)} & \left(-\cos 2 \theta \cdot \alpha_{1}+\frac{\alpha_{5}}{4}\right) \frac{\mathrm{d} z}{w}+\left(\frac{\alpha_{2}}{4}+\frac{3}{4} \alpha_{6}\right) \frac{z}{w} \mathrm{~d} z  \tag{19}\\
& +\left(-4 \sin ^{2} 2 \theta \cdot \alpha_{1}+\alpha_{3}-\cos 2 \theta \cdot \alpha_{5}\right) \frac{z^{3}}{w^{3}} \mathrm{~d} z \\
& +\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right) \frac{z^{4}}{w^{3}} \mathrm{~d} z
\end{align*}
$$

Let $\varphi: B_{\theta} \rightarrow B_{\theta}$ be the automorphism given by $\varphi(z, w)=(-z, i w)$. Note that $\varphi^{2}=j$, the hyperelliptic involution of $B_{\theta}$. Define paths $C_{4}, C_{5}$ on $B_{\theta}$ by

$$
\begin{aligned}
& C_{4}=\left\{(z, w)=\left(t, \sqrt{t\left(t^{4}+2 \cos 2 \theta \cdot t^{2}+1\right)}\right) \mid 0 \leq t \leq \infty\right\} \\
& C_{5}=\left\{(z, w)=\left(i t, \mathrm{e}^{\mathrm{i} \pi / 4} \sqrt{t\left(t^{4}-2 \cos 2 \theta \cdot t^{2}+1\right)}\right) \mid 0 \leq t \leq \infty\right\}
\end{aligned}
$$

Then the four closed paths

$$
C_{4} \cup\left(-j\left(C_{4}\right)\right), \varphi\left(C_{4} \cup\left(-j\left(C_{4}\right)\right)\right), C_{5} \cup\left(-j\left(C_{5}\right)\right), \varphi\left(C_{5} \cup\left(-j\left(C_{5}\right)\right)\right)
$$

form a homology basis, as verified by integrating the holomorphic differentials $\mathrm{d} z / w, z \mathrm{~d} z / w$ over them.
Straightforward calculations yield

$$
\begin{gathered}
\int_{C_{4} \cup\left\{-j\left(C_{4}\right)\right\}} \frac{\mathrm{d} z}{w}=2 \int_{0}^{\infty} \frac{\mathrm{d} t}{\sqrt{t\left(t^{4}+2 \cos 2 \theta \cdot t^{2}+1\right)}}=2 A, \\
\int_{C_{4} \cup\left\{-j\left(C_{4}\right)\right\}} \frac{z}{w} \mathrm{~d} z=2 \int_{0}^{\infty} \frac{t \mathrm{~d} t}{\sqrt{t\left(t^{4}+2 \cos 2 \theta \cdot t^{2}+1\right)}} \underbrace{}_{s=1 / t}=2 A, \\
\int_{\varphi\left(C_{4} \cup\left\{-j\left(C_{4}\right)\right\}\right)} \frac{\mathrm{d} z}{w}=2 \mathrm{i} A, \quad \int_{C^{2}} \frac{z}{w} \mathrm{~d} z=-2 \mathrm{i} A, \\
\int_{C_{5} \cup\left\{-j\left(C_{5}\right)\right\}} \frac{\mathrm{d} z}{w}=2 \mathrm{e}^{\frac{\pi}{4} \mathrm{i}} B, \quad \int_{\left.C_{4} \cup\left\{-j\left(C_{4}\right)\right\}\right)} \frac{z}{w} \mathrm{~d} z=-2 \mathrm{e}^{-\frac{\pi}{4} \mathrm{i}} B, \\
\left.\int_{C_{5} \cup\left\{-j\left(C_{5}\right)\right\}} \frac{\mathrm{d} z}{w}=-2 \mathrm{e}^{-\frac{\pi}{4} \mathrm{i}} B, \quad \int_{5} \cup\left\{-j\left(C_{5}\right)\right\}\right) \\
\int_{C_{4}} \frac{z^{2}}{w} \mathrm{~d} z=2 \mathrm{e}^{\frac{\pi}{4} \mathrm{i}} B, \\
\left.C_{4} \cup-j\left(C_{4}\right)\right\} \\
w^{3} \\
\mathrm{~d} z=2 \int_{0}^{\infty} \frac{\left.\left.t^{3} \mathrm{~d} t\left(C_{5}\right)\right\}\right)}{\sqrt{t\left(t^{4}+2 \cos 2 \theta \cdot t^{2}+1\right)}}{ }^{3}=2 C,
\end{gathered}
$$

$$
\begin{gathered}
\int_{C_{4} \cup\left\{-j\left(C_{4}\right)\right\}} \frac{z^{4}}{w^{3}} \mathrm{~d} z=2 \int_{0}^{\infty} \frac{t^{4} \mathrm{~d} t}{\sqrt{t\left(t^{4}+2 \cos 2 \theta \cdot t^{2}+1\right)^{3}}} \underbrace{=}_{s=1 / t} 2 C, \\
\int_{\varphi\left(C_{4} \cup\left\{-j\left(C_{4}\right)\right\}\right)} \frac{z^{3}}{w^{3}} \mathrm{~d} z=2 \mathrm{i} C, \quad \int_{\varphi\left(C_{4} \cup\left\{\left(-j\left(C_{4}\right)\right\}\right)\right.} \frac{z^{4}}{w^{3}} \mathrm{~d} z=-2 \mathrm{i} C, \\
\int_{C_{5} \cup\left\{-j\left(C_{5}\right)\right\}} \frac{z^{3}}{w^{3}} \mathrm{~d} z=-2 \mathrm{e}^{\frac{\pi}{4} \mathrm{i}} D, \quad \int_{C_{5} \cup\left\{-j\left(C_{5}\right)\right\}} \frac{z^{4}}{w^{3}} \mathrm{~d} z=2 \mathrm{e}^{-\frac{\pi}{4} \mathrm{i}} D, \\
\int_{\varphi\left(C_{5} \cup\left\{-j\left(C_{5}\right)\right\}\right)} \frac{z^{3}}{w^{3}} \mathrm{~d} z=2 \mathrm{e}^{-\frac{\pi}{4} \mathrm{i}} D, \quad \int_{\varphi\left(C_{5} \cup\left\{-j\left(C_{5}\right)\right\}\right)} \frac{z^{4}}{w^{3}} \mathrm{~d} z=-2 \mathrm{e}^{\frac{\pi}{4} \mathrm{i}} D .
\end{gathered}
$$

Note that the period condition (7) can be rewritten as

$$
\begin{equation*}
\int_{\ell} \omega=\overline{\int_{\ell} g_{\theta}^{2} \omega}, \quad \Re \int_{\ell} g_{\theta} \omega=\mathbf{0}, \quad \forall \ell \in H_{1}\left(B_{\theta}, \mathbb{Z}\right) \tag{20}
\end{equation*}
$$

By using (17)-(19) and the calculation we have just made, one can express the former relation of (20) for the above homology basis as

$$
\begin{align*}
& \left(\alpha_{1}+\frac{3}{4} \alpha_{3}\right) A+\left(-\frac{3}{2} \cos 2 \theta \cdot \alpha_{2}+\frac{\alpha_{4}}{4}\right) A  \tag{21}\\
& +\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right) C+\left(\left(-5+6 \cos ^{2} 2 \theta\right) \alpha_{2}-\cos 2 \theta \cdot \alpha_{4}+\alpha_{6}\right) C \\
& =\overline{\left(-\cos 2 \theta \cdot \alpha_{1}+\frac{\alpha_{5}}{4}\right) A+\left(\frac{\alpha_{2}}{4}+\frac{3}{4} \alpha_{6}\right) A} \\
& \overline{+\left(-4 \sin ^{2} 2 \theta \cdot \alpha_{1}+\alpha_{3}-\cos 2 \theta \cdot \alpha_{5}\right) C+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right) C} \text {, } \\
& \left(\alpha_{1}+\frac{3}{4} \alpha_{3}\right) \text { i } A+\left(-\frac{3}{2} \cos 2 \theta \cdot \alpha_{2}+\frac{\alpha_{4}}{4}\right)(-\mathrm{i} A)  \tag{22}\\
& +\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right) \mathrm{i} C+\left(\left(-5+6 \cos ^{2} 2 \theta\right) \alpha_{2}-\cos 2 \theta \cdot \alpha_{4}+\alpha_{6}\right)(-\mathrm{i} C) \\
& =\overline{\left(-\cos 2 \theta \cdot \alpha_{1}+\frac{\alpha_{5}}{4}\right) \mathrm{i} A+\left(\frac{\alpha_{2}}{4}+\frac{3}{4} \alpha_{6}\right)(-\mathrm{i} A)} \\
& \overline{+\left(-4 \sin ^{2} 2 \theta \cdot \alpha_{1}+\alpha_{3}-\cos 2 \theta \cdot \alpha_{5}\right) \mathrm{i} C+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right)(-\mathrm{i} C)}, \\
& \left(\alpha_{1}+\frac{3}{4} \alpha_{3}\right)(1+\mathrm{i}) B+\left(-\frac{3}{2} \cos 2 \theta \cdot \alpha_{2}+\frac{\alpha_{4}}{4}\right)(-1+\mathrm{i}) B  \tag{23}\\
& +\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right)(-1-\mathrm{i}) D \\
& +\left(\left(-5+6 \cos ^{2} 2 \theta\right) \alpha_{2}-\cos 2 \theta \cdot \alpha_{4}+\alpha_{6}\right)(1-\mathrm{i}) D \\
& =\left(-\cos 2 \theta \cdot \alpha_{1}+\frac{\alpha_{5}}{4}\right)(1+\mathrm{i}) B+\left(\frac{\alpha_{2}}{4}+\frac{3}{4} \alpha_{6}\right)(-1+\mathrm{i}) B \\
& \overline{+\left(-4 \sin ^{2} 2 \theta \cdot \alpha_{1}+\alpha_{3}-\cos 2 \theta \cdot \alpha_{5}\right)(-1-\mathrm{i}) D} \\
& \overline{+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right)(1-\mathrm{i}) D} \text {, } \\
& \left(\alpha_{1}+\frac{3}{4} \alpha_{3}\right)(-1+\mathrm{i}) B+\left(-\frac{3}{2} \cos 2 \theta \cdot \alpha_{2}+\frac{\alpha_{4}}{4}\right)(1+\mathrm{i}) B  \tag{24}\\
& +\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right)(1-\mathrm{i}) D \\
& +\left(\left(-5+6 \cos ^{2} 2 \theta\right) \alpha_{2}-\cos 2 \theta \cdot \alpha_{4}+\alpha_{6}\right)(-1-\mathrm{i}) D \\
& =\overline{\left(-\cos 2 \theta \cdot \alpha_{1}+\frac{\alpha_{5}}{4}\right)(-1+\mathrm{i}) B+\left(\frac{\alpha_{2}}{4}+\frac{3}{4} \alpha_{6}\right)(1+\mathrm{i}) B} \\
& \overline{+\left(-4 \sin ^{2} 2 \theta \cdot \alpha_{1}+\alpha_{3}-\cos 2 \theta \cdot \alpha_{5}\right)(1-\mathrm{i}) D} \\
& \overline{+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right)(-1-\mathrm{i}) D} \text {. }
\end{align*}
$$

Likewise, one expresses the latter relation from (20) for the homology basis as

$$
\begin{align*}
& \mathfrak{R}\left[\left(\frac{3}{4} \alpha_{2}+\frac{\alpha_{6}}{4}\right) A+\left(\alpha_{1}+\frac{\alpha_{3}}{4}\right) A\right.  \tag{25}\\
& \left.+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right) C+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right) C\right]=0, \\
& \mathfrak{R}\left[\left(\frac{3}{4} \alpha_{2}+\frac{\alpha_{6}}{4}\right) \mathrm{i} A+\left(\alpha_{1}+\frac{\alpha_{3}}{4}\right)(-\mathrm{i} A)\right.  \tag{26}\\
& \left.+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right) \mathrm{i} C+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right)(-\mathrm{i} C)\right]=0, \\
& \mathfrak{R}\left[\left(\frac{3}{4} \alpha_{2}+\frac{\alpha_{6}}{4}\right)(1+\mathrm{i}) B+\left(\alpha_{1}+\frac{\alpha_{3}}{4}\right)(-1+\mathrm{i}) B\right.  \tag{27}\\
& +\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right)(-1-\mathrm{i}) D \\
& \left.+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right)(1-\mathrm{i}) D\right]=0, \\
& \mathfrak{R}\left[\left(\frac{3}{4} \alpha_{2}+\frac{\alpha_{6}}{4}\right)(-1+\mathrm{i}) B+\left(\alpha_{1}+\frac{\alpha_{3}}{4}\right)(1+\mathrm{i}) B\right.  \tag{28}\\
& +\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right)(1-\mathrm{i}) D \\
& \left.+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right)(-1-\mathrm{i}) D\right]=0 \text {. }
\end{align*}
$$

(21), (22) are equivalent to

$$
\begin{align*}
& \left(\alpha_{1}+\frac{3}{4} \alpha_{3}\right) A+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right) C  \tag{29}\\
& =\overline{\left(\frac{\alpha_{2}}{4}+\frac{3}{4} \alpha_{6}\right) A+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right) C} \\
& \left(-\frac{3}{2} \cos 2 \theta \cdot \alpha_{2}+\frac{\alpha_{4}}{4}\right) A+\left(\left(-5+6 \cos ^{2} 2 \theta\right) \alpha_{2}-\cos 2 \theta \cdot \alpha_{4}+\alpha_{6}\right) C  \tag{30}\\
& =\overline{\left(-\cos 2 \theta \cdot \alpha_{1}+\frac{\alpha_{5}}{4}\right) A+\left(-4 \sin ^{2} 2 \theta \cdot \alpha_{1}+\alpha_{3}-\cos 2 \theta \cdot \alpha_{5}\right) C}
\end{align*}
$$

(23), (24) are equivalent to
$\left(\alpha_{1}+\frac{3}{4} \alpha_{3}\right) B+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right)(-D)$
$=-\overline{\left\{\left(\frac{\alpha_{2}}{4}+\frac{3}{4} \alpha_{6}\right) B+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right)(-D)\right\}}$,
$\left(-\frac{3}{2} \cos 2 \theta \cdot \alpha_{2}+\frac{\alpha_{4}}{4}\right) B+\left(\left(-5+6 \cos ^{2} 2 \theta\right) \alpha_{2}-\cos 2 \theta \cdot \alpha_{4}+\alpha_{6}\right)(-D)$
$=-\overline{\left\{\left(-\cos 2 \theta \cdot \alpha_{1}+\frac{\alpha_{5}}{4}\right) B+\left(-4 \sin ^{2} 2 \theta \cdot \alpha_{1}+\alpha_{3}-\cos 2 \theta \cdot \alpha_{5}\right)(-D)\right\}}$.
(25), (26) are equivalent to

$$
\begin{align*}
& \left(\alpha_{1}+\frac{\alpha_{3}}{4}\right) A+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right) C  \tag{33}\\
& =-\left\{\left(\frac{3}{4} \alpha_{2}+\frac{\alpha_{6}}{4}\right) A+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right) C\right\}
\end{align*}
$$

(27), (28) are equivalent to

$$
\begin{align*}
& \left(\alpha_{1}+\frac{\alpha_{3}}{4}\right) B+\left(-\cos 2 \theta \cdot \alpha_{3}+\alpha_{5}\right)(-D)  \tag{34}\\
& =\overline{\left(\frac{3}{4} \alpha_{2}+\frac{\alpha_{6}}{4}\right) B+\left(-\cos 2 \theta \cdot \alpha_{2}+\alpha_{4}-\cos 2 \theta \cdot \alpha_{6}\right)(-D)}
\end{align*}
$$

The equations (29)-(34) are summarized as

$$
\left(\begin{array}{ll}
X_{1} & X_{2}
\end{array}\right)\left(\begin{array}{l}
\alpha_{1}  \tag{35}\\
\alpha_{5} \\
\overline{\alpha_{2}} \\
\overline{\alpha_{4}} \\
\frac{\alpha_{3}}{\alpha_{6}}
\end{array}\right)=\left(\begin{array}{l}
0 \\
0 \\
0 \\
0 \\
0 \\
0
\end{array}\right),
$$

where

$$
\begin{aligned}
& X_{1}=\left(\begin{array}{ccc}
A & C & \frac{3}{4} A-C \cos 2 \theta \\
B & -D & -\frac{3}{4} B-D \cos 2 \theta \\
B & -D & \frac{B}{4}+D \cos 2 \theta \\
A & C & -\frac{A}{4}+C \cos 2 \theta \\
-\left(A \cos 2 \theta+4 C \sin ^{2} 2 \theta\right) & \frac{A}{4}-C \cos 2 \theta & \frac{3}{2} A \cos 2 \theta+\left(5-6 \cos ^{2} 2 \theta\right) C \\
B \cos 2 \theta-4 D \sin ^{2} 2 \theta & -\left(\frac{B}{4}+D \cos 2 \theta\right) & \frac{3}{2} B \cos 2 \theta+\left(-5+6 \cos ^{2} 2 \theta\right) D
\end{array}\right), \\
& X_{2}=\left(\begin{array}{ccc}
C & \frac{A}{4}-C \cos 2 \theta & \frac{A}{4}-C \cos 2 \theta \\
D & \frac{B}{4}+D \cos 2 \theta & -\frac{B}{4}-D \cos 2 \theta \\
-D & \frac{3}{4} B+D \cos 2 \theta & \frac{3}{4} B+D \cos 2 \theta \\
-C & \frac{3}{4} A-C \cos 2 \theta & -\frac{3}{4} A+C \cos 2 \theta \\
-\frac{A}{4}+C \cos 2 \theta & C & -C \\
-\frac{B}{4}-D \cos 2 \theta & D & D
\end{array}\right) .
\end{aligned}
$$

By applying elementary transformations as listed in Appendix A, it can be verified that the above system of linear equations is equivalent to

$$
\left(\begin{array}{lll}
Y_{1} & Y_{2} & Y_{3}
\end{array}\right)\left(\begin{array}{l}
\alpha_{1}  \tag{36}\\
\alpha_{5} \\
\overline{\alpha_{2}} \\
\overline{\alpha_{4}} \\
\alpha_{3} \\
\overline{\alpha_{6}}
\end{array}\right)=\left(\begin{array}{l}
0 \\
0 \\
0 \\
0 \\
0 \\
0
\end{array}\right),
$$

where

$$
\begin{aligned}
& Y_{1}=\left(\begin{array}{cccc}
A(A D+B C)^{2} & 0 & 0 & 0 \\
0 & -(A D+B C)^{2} & 0 & 0 \\
0 & 0 & A D+B C & 0 \\
0 & 0 & 0 & -2 C(A D+B C) \\
0 & 0 & 0 & 0 \\
0 & 0 & 0 & 0
\end{array}\right), \\
& Y_{2}=\left(\begin{array}{c}
\frac{1}{2} A(A D+B C)^{2}+\frac{1}{8} A(-A D+B C)^{2} \\
(A D+B C)^{2} \cos 2 \theta+\frac{1}{4} A B(-A D+B C) \\
\frac{1}{2}(-A D+B C) \\
(A B+(A D-B C) \cos 2 \theta) C \\
\frac{1}{16} A C(3 A D+B C)\left[-B\left(A^{2}+16 C^{2} \sin ^{2} 2 \theta\right)+8(A D+B C)\left(A \cos 2 \theta+4 C \sin ^{2} 2 \theta\right)\right] \\
-\frac{1}{16} B D(A D+3 B C)\left[A\left(B^{2}+16 D^{2} \sin ^{2} 2 \theta\right)+8(A D+B C)\left(B \cos 2 \theta-4 D \sin ^{2} 2 \theta\right)\right]
\end{array}\right), \\
& Y_{3}=\left(\begin{array}{c}
\frac{1}{2} A(A D+B C)(-A D+B C) \\
A B(A D+B C) \\
A D+B C \\
0 \\
\frac{1}{4} A C(A D+B C)\left[B\left(A^{2}+16 C^{2} \sin ^{2} 2 \theta\right)-8(A D+B C)\left(A \cos 2 \theta+4 C \sin ^{2} 2 \theta\right)\right] \\
-\frac{1}{4} B D(A D+B C)\left[A\left(B^{2}+16 D^{2} \sin ^{2} 2 \theta\right)+8(A D+B C)\left(B \cos 2 \theta-4 D \sin ^{2} 2 \theta\right)\right]
\end{array}\right) .
\end{aligned}
$$

It is easy to see that this system has a nontrivial solution if and only if the matrix

$$
\left(\begin{array}{ll}
\left(Y_{2}\right)_{5} & \left(Y_{3}\right)_{5} \\
\left(Y_{2}\right)_{6} & \left(Y_{3}\right)_{6}
\end{array}\right)
$$

is not invertible, where $\left(Y_{i}\right)_{j}$ is the $j$-th component of $Y_{i}$. In conclusion, the necessary and sufficient condition that (35) has a nontrivial solution is that either

$$
\begin{equation*}
A\left(B^{2}+16 D^{2} \sin ^{2} 2 \theta\right)+8(A D+B C)\left(B \cos 2 \theta-4 D \sin ^{2} 2 \theta\right)=0 \tag{37}
\end{equation*}
$$

or

$$
\begin{equation*}
B\left(A^{2}+16 C^{2} \sin ^{2} 2 \theta\right)-8(A D+B C)\left(A \cos 2 \theta+4 C \sin ^{2} 2 \theta\right)=0 \tag{38}
\end{equation*}
$$

holds.
One can verify that Eq. (37) has a unique solution $\theta_{1} \approx 0.65<\pi / 4$ in the range $0<\theta<\pi / 2$. We shall give a proof of this fact in Appendix B. Note that the change of variable $\theta \mapsto \pi / 2-\theta$ transforms (37) to (38) and vice versa. Therefore, $\theta_{2}:=\pi / 2-\theta_{1} \approx 0.91>\pi / 4$ gives a unique solution to Eq. (38) in the range $0<\theta<\pi / 2$.

If $\theta=\theta_{1}$, then it is easy to verify that the corresponding nontrivial solutions are given by real linear combinations of $\omega_{1}$ and $\omega_{2}$ as in the statement of Lemma 3.

## 5. Proof of Lemma 5

In this section, we shall prove Lemma 5.
Note that $u_{i}=\left\langle X_{\omega_{i}}, N\right\rangle$, where $N$ is the unit normal vector field of $X_{\omega_{i}}$, related to $g_{\theta_{1}}$ by

$$
N={ }^{t}\left(\frac{2 \Re g_{\theta_{1}}}{\left|g_{\theta_{1}}\right|^{2}+1}, \frac{2 \Im g_{\theta_{1}}}{\left|g_{\theta_{1}}\right|^{2}+1}, \frac{\left|g_{\theta_{1}}\right|^{2}-1}{\left|g_{\theta_{1}}\right|^{2}+1}\right)
$$

We have $s_{1}^{*} \omega_{1}=\overline{\omega_{1}}, s_{1}^{*} \omega_{2}=-\overline{\omega_{2}}$,

$$
s_{1}^{*}\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
\mathrm{i}\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right)=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & -1 & 0 \\
0 & 0 & 1
\end{array}\right) \overline{\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
\mathrm{i}\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right)}, \quad s_{1}^{*} N=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & -1 & 0 \\
0 & 0 & 1
\end{array}\right) N .
$$

Since $s_{1}\left(p_{0}\right)=p_{0}$, it follows from these formulae that $s_{1}^{*} u_{1}=u_{1}$ and $s_{1}^{*} u_{2}=-u_{2}$.
Let $\psi(z, w)=\left(1 / z, w / z^{3}\right)$. By straightforward calculation, we get

$$
\begin{aligned}
& \psi^{*}\left(\frac{\mathrm{~d} z}{w}\right)=-\frac{z}{w} \mathrm{~d} z, \quad \psi^{*}\left(\frac{\mathrm{~d} z}{w^{3}}\right)=-\frac{z^{7}}{w^{3}} \mathrm{~d} z, \quad \psi^{*}\left(\frac{z}{w^{3}}\right) \mathrm{d} z=-\frac{z^{6}}{w^{3}} \mathrm{~d} z \\
& \psi^{*}\left(\frac{z^{2}}{w^{3}}\right) \mathrm{d} z=-\frac{z^{5}}{w^{3}} \mathrm{~d} z, \quad \psi^{*}\left(\frac{z^{3}}{w^{3}}\right) \mathrm{d} z=-\frac{z^{4}}{w^{3}} \mathrm{~d} z, \quad \psi^{*}\left(\frac{z^{4}}{w^{3}}\right) \mathrm{d} z=-\frac{z^{3}}{w^{3}} \mathrm{~d} z
\end{aligned}
$$

Therefore,

$$
\begin{aligned}
\psi^{*} \omega_{1} & =z^{2}\left(\frac{A D+3 B C}{4(A D+B C)} \frac{z^{4}+2 \cos 2 \theta \cdot z^{2}+1}{w^{3}} \mathrm{~d} z+\frac{A D+3 B C}{4(A D+B C)} \frac{z^{5}}{w^{3}} \mathrm{~d} z-\frac{z^{4}}{w^{3}} \mathrm{~d} z\right. \\
& -\frac{A B+(A D-B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{3}}{w^{3}} \mathrm{~d} z-\frac{A B+2(A D+B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{2}}{w^{3}} \mathrm{~d} z \\
& \left.-\frac{3 A D+B C}{4(A D+B C)} \frac{z}{w^{3}} \mathrm{~d} z\right) \\
& =z^{2}\left(\frac{-3 A D-B C}{4(A D+B C)} \frac{z^{4}}{w^{3}} \mathrm{~d} z+\frac{-A B+(-A D+B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{2}}{w^{3}} \mathrm{~d} z\right. \\
& +\frac{A D+3 B C}{4(A D+B C)} \frac{\mathrm{d} z}{w^{3}}-\frac{A B+2(A D+B C) \cos 2 \theta}{2(A D+B C)} \frac{z^{3}}{w^{3}} \mathrm{~d} z-\frac{z}{w^{3}} \mathrm{~d} z \\
& \left.+\frac{A D+3 B C}{4(A D+B C)} \frac{z^{5}+2 \cos 2 \theta \cdot z^{3}+z}{w^{3}} \mathrm{~d} z\right) \\
& =-z^{2} \omega_{1}
\end{aligned}
$$

Likewise, we obtain $\psi^{*} \omega_{2}=z^{2} \omega_{2}$. Since we also have

$$
\psi^{*}\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
\mathrm{i}\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right)=\frac{1}{z^{2}}\left(\begin{array}{ccc}
-1 & 0 & 0 \\
0 & 1 & 0 \\
0 & 0 & 1
\end{array}\right)\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
i\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right), \quad \psi^{*} N=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & -1 & 0 \\
0 & 0 & -1
\end{array}\right) N
$$

and $\psi\left(p_{0}\right)=p_{0}$, we find $\psi^{*} u_{1}=u_{1}$ and $\psi^{*} u_{2}=-u_{2}$. Since $s_{3}=\psi \circ s_{1}$, we conclude that $s_{3}^{*} u_{1}=u_{1}$ and $s_{3}^{*} u_{2}=u_{2}$.
We have $j^{*} \omega_{1}=-\omega_{1}, j^{*} \omega_{2}=-\omega_{2}$,

$$
j^{*}\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
\mathrm{i}\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right)=\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
\mathrm{i}\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right), \quad j^{*} N=N
$$

from which it follows that

$$
\begin{aligned}
j^{*} u_{1}(p) & =\left\langle\Re \int_{j\left(p_{0}\right)}^{j(p)}\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
\mathrm{i}\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right) \omega_{1}, N(j(p))\right\rangle+\left\langle\Re \int_{p_{0}}^{j\left(p_{0}\right)}\left(\begin{array}{c}
1-g_{\theta_{1}}^{2} \\
\mathrm{i}\left(1+g_{\theta_{1}}^{2}\right) \\
2 g_{\theta_{1}}
\end{array}\right) \omega_{1}, N(j(p))\right\rangle \\
& =-u_{1}(p)+\left\langle c_{1}, N(p)\right\rangle
\end{aligned}
$$

where $c_{1}=\Re \int_{p_{0}}^{j\left(p_{0}\right) t}\left(1-g_{\theta_{1}}^{2}, \mathrm{i}\left(1+g_{\theta_{1}}^{2}\right), 2 g_{\theta_{1}}\right) \omega_{1}$, and $j^{*} u_{2}=-u_{2}+\left\langle c_{2}, N\right\rangle$.

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## Appendix A

As mentioned in the proof of Lemma 3, the system (35) of linear equations can be reduced to an equivalent one of simpler form by applying elementary transformations. For the reader's convenience, we shall list all the elementary transformations explicitly.

We apply the following operations, where $\mathrm{Rj}(1 \leq j \leq 6)$ denotes the $j$-th row, to the matrix $\left(\begin{array}{ll}X_{1} & X_{2}\end{array}\right)$.
(i) $\mathrm{R} 4 \longrightarrow \mathrm{R} 4-\mathrm{R} 1$.
(ii) R5 $\longrightarrow \mathrm{R} 5+\mathrm{R} 1 \times \cos 2 \theta$.
(iii) R3 $\longrightarrow$ R3 - R2.
(iv) $\mathrm{R} 6 \longrightarrow \mathrm{R} 6-\mathrm{R} 2 \times \cos 2 \theta$.
(v) R5 $\longrightarrow \mathrm{R} 5 \times A+\mathrm{R} 1 \times 4 C \sin ^{2} 2 \theta$.
(vi) $\mathrm{R} 6 \longrightarrow \mathrm{R} 6 \times B+\mathrm{R} 2 \times 4 D \sin ^{2} 2 \theta$.
(vii) $R 2 \longrightarrow R 2+R 3 \times 1 / 2$.
(viii) $\mathrm{R} 6 \longrightarrow \mathrm{R} 6+\mathrm{R} 3 \times\left(-B \cos 2 \theta+2 D \sin ^{2} 2 \theta\right)$.
(ix) $\mathrm{R} 1 \longrightarrow \mathrm{R} 1+\mathrm{R} 4 \times 1 / 2$.
(x) $\mathrm{R} 5 \longrightarrow \mathrm{R} 5+\mathrm{R} 4 \times\left(A \cos 2 \theta+2 C \sin ^{2} 2 \theta\right)$.
(xi) $\mathrm{R} 2 \longrightarrow \mathrm{R} 2 \times A+\mathrm{R} 1 \times(-B)$.
(xii) $\mathrm{R} 5 \longrightarrow \mathrm{R} 5 \times C+\mathrm{R} 4 \times\left(-A^{2} / 8\right)$.
(xiii) $\mathrm{R} 6 \longrightarrow \mathrm{R} 6 \times D+\mathrm{R} 3 \times\left(-B^{2} / 8\right)$.
(xiv) $\mathrm{R} 1 \longrightarrow \mathrm{R} 1 \times(A D+B C)+\mathrm{R} 2 \times C$.
$(\mathrm{xv}) \mathrm{R} 3 \longrightarrow \mathrm{R} 3 \times \mathrm{C}+\mathrm{R} 4 \times(-\mathrm{D})$.
(xvi) $\mathrm{R} 5 \longrightarrow \mathrm{R} 5 \times(A D+B C)+\mathrm{R} 2 \times\left(A^{2} C / 4+4 C^{3} \sin ^{2} 2 \theta\right)$.
(xvii) $\mathrm{R} 6 \longrightarrow \mathrm{R} 6 \times(A D+B C)+\mathrm{R} 2 \times\left(-B^{2} D / 4-4 D^{3} \sin ^{2} 2 \theta\right)$.
(xviii) $\mathrm{R} 1 \longrightarrow \mathrm{R} 1 \times(A D+B C)+\mathrm{R} 3 \times A(-A D+B C) / 4$.
(xix) $\mathrm{R} 2 \longrightarrow \mathrm{R} 2 \times(A D+B C)+\mathrm{R} 3 \times A B / 2$.
(xx) $\mathrm{R} 4 \longrightarrow \mathrm{R} 4 \times(A D+B C)+\mathrm{R} 3 \times(A-2 C \cos 2 \theta)$.
$(x x i) \mathrm{R} 5 \longrightarrow \mathrm{R} 5 \times(A D+B C)+\mathrm{R} 3 \times\left[-A^{2}\left(A^{2} D / 8+(A D+B C) C \cos 2 \theta+6 C^{2} D \sin ^{2} 2 \theta\right)-4 A B C^{3} \sin ^{2} 2 \theta\right]$.
(xxii) $\mathrm{R} 6 \longrightarrow \mathrm{R} 6 \times(A D+B C)+\mathrm{R} 3 \times\left[-B^{2}\left(-B^{2} C / 8+(A D+B C) D \cos 2 \theta-6 C D^{2} \sin ^{2} 2 \theta\right)+4 A B D^{3} \sin ^{2} 2 \theta\right]$.

Then we finally obtain the matrix $\left(\begin{array}{lll}Y_{1} & Y_{2} & Y_{3}\end{array}\right)$ as in the proof of Lemma 3.

## Appendix B

In this appendix, we prove that Eq. (37) has a unique solution $\theta_{1}<\pi / 4$ in the range $0<\theta<\pi / 2$.
We first prove that (37) has a unique solution in the range $0<\theta<\pi / 4$. Though it is possible to verify this fact by a direct elementary argument, here we present an indirect one, assuming that (37) has no solutions in the range $\pi / 4 \leq \theta<\pi / 2$, which we will prove afterwards. Since the left-hand side of (37) is positive near $\theta=0$ and negative at $\theta=\pi / 4$, (37) has at least one solution by the intermediate value theorem. On the other hand, (38) has no solutions in the range $0<\theta<\pi / 4$ by the remark at the end of the proof of Lemma 3. Suppose that there is more than one solution to (37), and let $\varphi_{1}<\varphi_{2}$ be the first and second smallest ones. Then, since the argument for proving Theorem 6 depends only on the fact that $\theta_{1}$ is a solution to (37), we deduce that the number of eigenvalues of $-\Delta_{\theta}$ less than 2 decreases by two each time when $\theta$ passes $\varphi_{1}$ and $\varphi_{2}$. But this is impossible because there are exactly three such eigenvalues for $\theta<\varphi_{1}$. Thus, the solutions to (37) must be unique.

We now proceed to prove that Eq. (37) has no solutions in the range $\pi / 4 \leq \theta<\pi / 2$. We start by rewriting the integrals $A, B, C, D$ using the complete elliptic integrals

$$
K(\mathbf{k})=\int_{0}^{\frac{\pi}{2}} \frac{\mathrm{~d} \theta}{\sqrt{1-\mathbf{k}^{2} \sin ^{2} \theta}}, \quad E(\mathbf{k})=\int_{0}^{\frac{\pi}{2}} \sqrt{1-\mathbf{k}^{2} \sin ^{2} \theta} \mathrm{~d} \theta
$$

defined for $0<\mathbf{k}<1$. Clearly, $K(\mathbf{k})$ (resp. $E(\mathbf{k})$ ) is a monotone increasing (resp. decreasing) function of $\mathbf{k}$. Computing with the change of variable $u=\sqrt{t}-1 / \sqrt{t}$ and using 222 of [1], we obtain

$$
\begin{aligned}
& A=\frac{2}{\sqrt{2(1+\sin \theta)}} K(k), \quad B=\frac{2}{\sqrt{2(1+\cos \theta)}} K(l), \\
& C=\frac{1}{4 \sqrt{2(1+\sin \theta)} \sin ^{2} \theta(1-\sin \theta)}(E(k)-(1-\sin \theta) K(k)), \\
& D=\frac{1}{4 \sqrt{2(1+\cos \theta)} \cos ^{2} \theta(1-\cos \theta)}(E(l)-(1-\cos \theta) K(l)),
\end{aligned}
$$

where $k=\sqrt{2 \sin \theta /(1+\sin \theta)}$ and $l=\sqrt{2 \cos \theta /(1+\cos \theta)}$.
The left-hand side of (37) can be rewritten as

$$
\cos 2 \theta\left(A B^{2} \cos 2 \theta+8 A B D+8 B^{2} C\right)+\sin ^{2} 2 \theta\left(A B^{2}-16 A D^{2}-32 B C D\right)
$$

Therefore, it suffices to verify that both

$$
\begin{equation*}
A B^{2} \cos 2 \theta+8 A B D+8 B^{2} C>0, \quad A B^{2}-16 A D^{2}-32 B C D<0 \tag{39}
\end{equation*}
$$

hold in the range $\pi / 4 \leq \theta<\pi / 2$.
We first reduce these inequalities to several simpler ones, with details discussed later on. The former inequality of (39) follows from

$$
\begin{equation*}
A \cos 2 \theta+8 C>0, \quad \pi / 4 \leq \theta<\pi / 2 \tag{40}
\end{equation*}
$$

For the latter inequality of (39), since

$$
\begin{aligned}
& A B^{2}-16 A D^{2}-32 B C D \\
& \quad=\left\{\begin{array}{l}
A\left(B-\frac{192}{25} D\right)\left(B+\frac{25}{12} D\right)+B D\left(\frac{1679}{300} A-32 C\right), \quad \pi / 4 \leq \theta \leq 5 \pi / 16 \\
A(B-10 D)\left(B+\frac{8}{5} D\right)+B D\left(\frac{42}{5} A-32 C\right), \quad 5 \pi / 16 \leq \theta \leq 3 \pi / 8 \\
A(B-16 D)(B+D)+B D(15 A-32 C), \quad 3 \pi / 8 \leq \theta<\pi / 2
\end{array}\right.
\end{aligned}
$$

it suffices to show

$$
\begin{align*}
& 25 B-192 D<0, \quad \pi / 4 \leq \theta \leq 5 \pi / 16,  \tag{41}\\
& 1679 A-9600 C<0, \quad \pi / 4 \leq \theta \leq 5 \pi / 16,  \tag{42}\\
& B-10 D<0, \quad 5 \pi / 16 \leq \theta \leq 3 \pi / 8,  \tag{43}\\
& 21 A-80 C<0, \quad 5 \pi / 16 \leq \theta \leq 3 \pi / 8,  \tag{44}\\
& B-16 D<0, \quad 3 \pi / 8 \leq \theta<\pi / 2,  \tag{45}\\
& 15 A-32 C<0, \quad 3 \pi / 8 \leq \theta<\pi / 2 . \tag{46}
\end{align*}
$$

We now present a detailed proof of (41). Since the proofs of (40) and (42)-(46) are similar, they are left to the reader. We have

$$
25 B-192 D=f(l)\left[\left(49 l^{4}-96 l^{2}+96\right)\left(1-l^{2}\right) K(l)-12\left(2-l^{2}\right)^{3} E(l)\right]
$$

where $f(l)$ is a positive function of $l$. Therefore, one must show that

$$
\left(49 l^{4}-96 l^{2}+96\right)\left(1-l^{2}\right) K(l)-12\left(2-l^{2}\right)^{3} E(l)<0
$$

in the range

$$
0.7142 \cdots=\frac{2 \cos \frac{5}{16} \pi}{1+\cos \frac{5}{16} \pi} \leq l^{2} \leq \frac{2 \cos \frac{\pi}{4}}{1+\cos \frac{\pi}{4}}=0.8284 \cdots
$$

Using

$$
\frac{\mathrm{d}}{\mathrm{~d} \mathbf{k}}\left[\left(1-\mathbf{k}^{2}\right) K(\mathbf{k})\right]=\frac{E(\mathbf{k})}{\mathbf{k}}-\frac{1+\mathbf{k}^{2}}{\mathbf{k}} K(\mathbf{k}), \quad \frac{\mathrm{d}}{\mathrm{~d} \mathbf{k}} E(\mathbf{k})=\frac{E(\mathbf{k})-K(\mathbf{k})}{\mathbf{k}}
$$

(cf. [1, 710]), we obtain

$$
\begin{align*}
& \frac{\mathrm{d}}{\mathrm{~d} l}\left[\left(49 l^{4}-96 l^{2}+96\right)\left(1-l^{2}\right) K(l)-12\left(2-l^{2}\right)^{3} E(l)\right]  \tag{47}\\
& \quad=l\left[-\left(257 l^{4}-507 l^{2}+336\right) K(l)+\left(84 l^{4}-311 l^{2}+336\right) E(l)\right]
\end{align*}
$$

Observe that $257 l^{4}-507 l^{2}+336$ and $84 l^{4}-311 l^{2}+336$ are positive and monotone decreasing in the range $0.71<l^{2}<$ 0.83. Then we can show that the right-hand side of (47) is negative in the range $0.71<l^{2}<0.83$ by estimating it in $0.71<l^{2} \leq 0.81$ and $0.81 \leq l^{2}<0.83$ separately. E.g., in $0.71<l^{2} \leq 0.81$,

$$
\begin{aligned}
- & \left(257 l^{4}-507 l^{2}+336\right) K(l)+\left(84 l^{4}-311 l^{2}+336\right) E(l) \\
\leq & -\left(257 \cdot 0.81^{2}-507 \cdot 0.81+336\right) K(\sqrt{0.71}) \\
& +\left(84 \cdot 0.71^{2}-311 \cdot 0.71+336\right) E(\sqrt{0.71}) \\
= & -1.723 \cdots<0 .
\end{aligned}
$$

Therefore, $\left(49 l^{4}-96 l^{2}+96\right)\left(1-l^{2}\right) K(l)-12\left(2-l^{2}\right)^{3} E(l)$ is monotone decreasing. Since its value at $l^{2}=0.714$ is $-0.033 \cdots<0$, (41) is proved.

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