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Mathematical analysis/Potential theory

## Harmonic measure is rectifiable if it is absolutely continuous with respect to the co-dimension-one Hausdorff measure <sup>☆</sup>



*La mesure harmonique est toujours rectifiable dès qu'elle est absolument continue par rapport à la mesure de Hausdorff de codimension 1*

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### ARTICLE INFO

#### Article history:

Received 22 September 2015

Accepted after revision 21 January 2016

Available online 3 March 2016

Presented by Gilles Pisier

### ABSTRACT

In the present paper, we sketch the proof of the fact that for any open connected set  $\Omega \subset \mathbb{R}^{n+1}$ ,  $n \geq 1$ , and any  $E \subset \partial\Omega$  with  $0 < \mathcal{H}^n(E) < \infty$ , absolute continuity of the harmonic measure  $\omega$  with respect to the Hausdorff measure on  $E$  implies that  $\omega|_E$  is rectifiable.

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### R É S U M É

Dans cet article, nous présentons les grandes lignes de la démonstration prouvant que, pour tout ensemble connexe ouvert  $\Omega \subset \mathbb{R}^{n+1}$ ,  $n \geq 1$ , et pour tout  $E \subset \partial\Omega$  avec  $0 < \mathcal{H}^n(E) < \infty$ , la continuité absolue de la mesure harmonique  $\omega$  par rapport à la mesure de Hausdorff sur  $E$  implique la rectifiabilité de  $\omega|_E$ .

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<sup>☆</sup> We acknowledge the support of the following grants: the second author – NSF DMS 1361701, the third author – SEV-2011-0087 and (FP7/2007-2013)/ERC agreement no. 615112 HAPDEGMT, the fourth author – Sloan Fellowship, NSF DMS 1344235, NSF DMS 1220089, NSF DMR 0212302, the sixth author – ERC 320501 FP7/2007-2013 (which also funded the first and fifth authors), 2014-SGR-75, MTM2013-44304-P, ITN MAnET (FP7-607647), the last author – NSF DMS 1265549. The results of this paper were obtained at the *Institut Henri Poincaré*, and at the 2015 ICMAT program *Analysis and Geometry in Metric Spaces*. All authors would like to express their gratitude to these institutions for their support and nice working environments.

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<http://dx.doi.org/10.1016/j.crma.2016.01.012>

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## Version française abrégée

Le résultat principal de cet article est le théorème suivant, qui révèle la structure fine de toutes les mesures harmoniques.

Si une mesure harmonique d'un certain domaine dans  $\mathbb{R}^{n+1}$  est absolument continue par rapport à la mesure de Hausdorff  $\mathcal{H}^n$ , alors cette mesure harmonique est  $n$ -rectifiable. En particulier, il existe un graphe (pivoté) d'une fonction lipschitzienne qui contient une partie positive de la mesure harmonique.

La preuve peut être divisée en trois parties principales. Tout d'abord, on considère le noyau de Riesz de singularité  $n$  (c'est-à-dire le gradient du noyau de Newton) et on l'applique à cette mesure harmonique. On veut estimer le potentiel que l'on vient d'obtenir. Pour cela, on utilise les cellules de David–Mattila et quelques considérations combinatoires pour trouver une partie positive de mesure harmonique sur laquelle ce potentiel est borné.

La deuxième partie de la preuve consiste à utiliser un théorème du type  $T1$  pour les mesures non homogènes obtenu par Nazarov–Treil–Volberg au début des années 2000. On obtient donc une partie positive de la mesure harmonique telle que l'opérateur de Riesz est borné par rapport à cette partie de la mesure harmonique.

Enfin, on utilise le résultat récent de Nazarov–Tolsa–Volberg (la solution d'un problème de David–Semmes) pour établir que ladite partie de la mesure harmonique est rectifiable.

## 1. Introduction and notation

In what follows,  $\Omega$  will denote a connected open set. As usual, using the Perron method and Riesz Representation Theorem, we can define the harmonic measure of  $\Omega$  with a pole at  $p \in \Omega$  (denote it by  $\omega^p$ ) for any bounded domain and use a limiting procedure to extend this definition to unbounded domains (see, e.g., [7, Section 3]).

We call a Radon measure  $\sigma$  on  $\mathbb{R}^{n+1}$   $n$ -rectifiable if its (any) Borel support can be covered by countably many (rotated) graphs of scalar Lipschitz functions on  $\mathbb{R}^n$  up to zero  $\sigma$  measure (in particular any pure point measure is rectifiable in this definition, but we will apply it to harmonic measures, such measures never have point masses.)

The main result that we announce here is the following.

**Theorem 1.1.** *Let  $\Omega$  be any open connected set in  $\mathbb{R}^{n+1}$ ,  $n \geq 1$  and let  $\omega = \omega^p$ , for some  $p \in \Omega$ , be the harmonic measure on  $\Omega$ . Let  $E$  be a Borel set, and let  $\omega|_E \ll \mathcal{H}^n|_E$ . Then  $\omega|_E$  is  $n$ -rectifiable.*

A simple corollary is that if  $\omega|_E$  and  $\mathcal{H}^n|_E$  are mutually absolutely continuous then  $\mathcal{H}^n|_E$  is  $n$ -rectifiable. In Besicovitch theory, such sets are called  $n$ -rectifiable.

The structure of harmonic measures was a focus of attention of many mathematicians starting with Carleson's article [5], which proved that harmonic measure on any continuum in the plane has dimension at least  $1/2 + \epsilon$  for universal positive  $\epsilon$ . Without an  $\epsilon$  the result was known for a long time and was nothing but a re-interpretation of Kőbe distortion theorem. Carleson's result had a very difficult proof, but Makarov found a novel approach that also established that harmonic measure on any continuum on the plane has dimension exactly 1. Carleson's and Makarov's works generated a fantastic amount of very difficult and revealing results about the structure of harmonic measure. Topological restrictions (like being a continuum) and dimension played an important part from the start. Let us mention that 1) for any planar domain any harmonic measure has dimension at most 1 (Jones–Wolff, [12]); 2) there are no analogs of Makarov's or Jones–Wolff's results in dimension 3 and higher (Bourgain, [4], Wolff, [20]); 3) in dimension 2 (planar domains) if harmonic measure on a continuum dwells on a good set (meaning that its Borel support can be covered by countably many (rotated) Lipschitz graphs) then harmonic measure is absolutely continuous with respect to the surface measure  $\mathcal{H}^1$  (Bishop–Jones, [3]). It is also shown in [3] that some topological restrictions are necessary, for there is a 1-rectifiable set in a plane for which harmonic measure is positive and Hausdorff measure is zero.

Notice that Theorem 1.1 establishes the converse to Bishop–Jones result: they claim that in  $\mathbb{R}^2$  and with topological restriction (continuum) 1-rectifiable harmonic measure must be absolutely continuous with respect to measure  $\mathcal{H}^1$ . We claim that without any restrictions whatsoever and in any  $\mathbb{R}^{n+1}$  the converse is true: if harmonic measure is absolutely continuous with respect to  $\mathcal{H}^n$ , then it is  $n$ -rectifiable. Moreover, our result is very local, one can start with harmonic measure restricted to an arbitrary Borel set  $E$ , and again its absolute continuity with respect to  $\mathcal{H}^n|_E$  implies that harmonic measure restricted to  $E$  is  $n$ -rectifiable.

In this respect, let us recall that two natural directions were considered in [7,8,11]: 1) geometry-to-analysis, when some sort of (say, scalar invariant) rectifiability of the boundary implies quantitative claims on the density of harmonic measure with respect to  $\mathcal{H}^n$ , 2) analysis-to-geometry, when the absolute continuity of harmonic measure and  $\mathcal{H}^n$  in a scale invariant fashion implies the quantitative rectifiability of the boundary. As mentioned above, in geometry-to-analysis direction Bishop–Jones in [3] found that certain additional topological properties are necessary. But Theorem 1.1 shows that no such obstacles exist in the direction analysis-to-geometry.

This paper is a short announcement of the results in [1], which in turn arises from the union of two separate works, [2] and [10]. In [2], a version of Theorem 1.1 was proved under the additional assumption that the boundary of  $\Omega$  is porous in  $E$  (a certain topological restriction). In [10] the porosity assumption was removed. Both [2] and [10] exist only as preprints on ArXiv.

We finish this section with some notation and preliminaries. Given a signed Radon measure  $\nu$  in  $\mathbb{R}^{n+1}$  we consider the  $n$ -dimensional Riesz transform

$$\mathcal{R}\nu(x) = \int \frac{x - y}{|x - y|^{n+1}} d\nu(y),$$

whenever the integral makes sense. For  $\varepsilon > 0$ , its  $\varepsilon$ -truncated version is given by

$$\mathcal{R}_\varepsilon \nu(x) = \int_{|x-y>\varepsilon} \frac{x - y}{|x - y|^{n+1}} d\nu(y), \quad \mathcal{R}_{*,\delta} \nu(x) = \sup_{\varepsilon>\delta} |\mathcal{R}_\varepsilon \nu(x)|.$$

We also consider the maximal operator

$$\mathcal{M}_\delta^n \nu(x) = \sup_{r>\delta} \frac{|\nu|(B(x, r))}{r^n},$$

In the case  $\delta = 0$  we write  $\mathcal{R}_* \nu(x) := \mathcal{R}_{*,0} \nu(x)$  and  $\mathcal{M}^n \nu(x) := \mathcal{M}_0^n \nu(x)$ .

Next, let  $\mathcal{H}_\infty^s$  denote the Hausdorff content of order  $s$ . We recall a result of Bourgain from [4].

**Lemma 1.2** (Bourgain’s lemma). *There is  $\delta_0 > 0$  depending only on  $n \geq 1$  so that the following holds for  $\delta \in (0, \delta_0)$ . Let  $\Omega \subsetneq \mathbb{R}^{n+1}$  be a bounded domain,  $n + 1 \geq s > n - 1$ ,  $\xi \in \partial\Omega$ ,  $r > 0$ , and  $B = B(\xi, r)$ . Then*

$$\omega^x(B) \gtrsim_{n,s} \frac{\mathcal{H}_\infty^s(\partial\Omega \cap \delta B)}{(\delta r)^s} \quad \text{for all } x \in \delta B \cap \Omega.$$

Let us point out that it is not difficult to prove that if Theorem 1.1 holds on bounded domains, then it remains valid for any open set  $\Omega \subset \mathbb{R}^{n+1}$ . We shall restrict our discussion here to the case when  $\Omega$  is bounded. Also, for brevity, we consider the case  $n \geq 2$  only. For details on these and other issues we refer the reader to [1].

## 2. David–Mattila cells, doubling cells

What follows is the sketch of the proof of our main result. The details should be found in [1]. To start, we consider the dyadic lattice of cubes with small boundaries of David–Mattila associated with  $\omega^p$ . It has been constructed in [6, Theorem 3.2] (with  $\omega^p$  replaced by a general Radon measure).

**Lemma 2.1.** (See David and Mattila [6].) *Consider two constants  $C_0 > 1$  and  $A_0 > 5000 C_0$  and denote  $W = \text{supp } \omega^p$ . Then there exists a sequence of partitions of  $W$  into Borel subsets  $Q$ ,  $Q \in \mathcal{D}_k$ , with the following properties: 1) for each integer  $k \geq 0$ ,  $W$  is the disjoint union of the “cubes”  $Q$ ,  $Q \in \mathcal{D}_k$ , and if  $k < l$ ,  $Q \in \mathcal{D}_l$ , and  $R \in \mathcal{D}_k$ , then either  $Q \cap R = \emptyset$  or else  $Q \subset R$ . 2) For each  $k \geq 0$  and each cube  $Q \in \mathcal{D}_k$ , there is a ball  $B(Q) = B(z_Q, r(Q))$  such that*

$$z_Q \in W, \quad A_0^{-k} \leq r(Q) \leq C_0 A_0^{-k}, \quad W \cap B(Q) \subset Q \subset W \cap 28 B(Q) = W \cap B(z_Q, 28r(Q)),$$

and the balls  $5B(Q)$ ,  $Q \in \mathcal{D}_k$ , are disjoint. 3) The cubes  $Q \in \mathcal{D}_k$  have “small boundaries” – see [6] for a precise definition. 4) Denote by  $\mathcal{D}_k^{db}$  the family of cubes  $Q \in \mathcal{D}_k$  for which

$$\omega^p(100B(Q)) \leq C_0 \omega^p(B(Q)). \tag{1}$$

We have that  $r(Q) = A_0^{-k}$  when  $Q \in \mathcal{D}_k \setminus \mathcal{D}_k^{db}$  and  $\omega^p(100B(Q)) \leq C_0^{-l} \omega^p(100^{l+1}B(Q))$  for all  $l \geq 1$  such that  $100^l \leq C_0$  and  $Q \in \mathcal{D}_k \setminus \mathcal{D}_k^{db}$ .

We use the notation  $\mathcal{D} = \bigcup_{k \geq 0} \mathcal{D}_k$ . Observe that the families  $\mathcal{D}_k$  are only defined for  $k \geq 0$ . So the diameters of the cubes from  $\mathcal{D}$  are uniformly bounded from above. Given  $Q \in \mathcal{D}_k$ , we denote  $J(Q) = k$ . We call  $\ell(Q) = 56 C_0 A_0^{-k}$  the side length of  $Q$ . Observe that  $r(Q) \sim \text{diam}(B(Q)) \sim \ell(Q)$ . Also, we call  $z_Q$  the center of  $Q$ , and the cube  $Q' \in \mathcal{D}_{k-1}$  such that  $Q' \supset Q$  the parent of  $Q$ . We set  $B_Q = 28 B(Q) = B(z_Q, 28r(Q))$ .

We denote  $\mathcal{D}^{db} = \bigcup_{k \geq 0} \mathcal{D}_k^{db}$ . Note that, in particular, from (1) it follows that

$$\omega^p(3B_Q) \leq \omega^p(100B(Q)) \leq C_0 \omega^p(Q) \quad \text{if } Q \in \mathcal{D}^{db}. \tag{2}$$

For this reason we will refer to the cubes from  $\mathcal{D}^{db}$  as doubling. As shown in [6, Lemma 5.28], every cube  $R \in \mathcal{D}$  can be covered  $\omega^p$ -a.e. by a family of doubling cubes.

### 3. Good and bad cells. The estimate of $\mathcal{R}$ on good doubling cells

We introduce the notions of bad and good David–Mattila cells. First we need the  $n$ -dimensional Frostman measure. Let  $g \in L^1(\omega^p)$  be such that  $\omega^p|_E = g \mathcal{H}^n|_{\partial\Omega}$ . Given  $M > 0$ , let

$$E_M = \{x \in \partial\Omega : M^{-1} \leq g(x) \leq M\}.$$

Take  $M$  big enough so that  $\omega^p(E_M) \geq \omega^p(E)/2$ , say. Consider an arbitrary compact set  $F_M \subset E_M$  with  $\omega^p(F_M) > 0$ . We will show that there exists  $G_0 \subset F_M$  with  $\omega^p(G_0) > 0$  which is  $n$ -rectifiable. Clearly, this suffices to prove that  $\omega^p|_{E_M}$  is  $n$ -rectifiable, and letting  $M \rightarrow \infty$  we get the full  $n$ -rectifiability of  $\omega^p|_E$ .

Let  $\mu$  be an  $n$ -dimensional Frostman measure for  $F_M$ . That is,  $\mu$  is a non-zero Radon measure supported on  $F_M$  such that  $\mu(B(x, r)) \leq Cr^n$  for all  $x \in \mathbb{R}^{n+1}$ . Further, by renormalizing  $\mu$ , we can assume that  $\|\mu\| = 1$ . Of course the constant  $C$  above will depend on  $\mathcal{H}^n_\infty(F_M)$ , and the same may happen for all the constants  $C$  to appear, but this will not bother us. Notice that  $\mu \ll \mathcal{H}^n|_{F_M} \ll \omega^p$ . In fact, for any set  $H \subset F_M$ ,

$$\mu(H) \leq C \mathcal{H}^n_\infty(H) \leq C \mathcal{H}^n(H) \leq CM \omega^p(H). \tag{3}$$

The cell  $Q \in \mathcal{D}$  is called bad if it is a maximal cube satisfying one of the conditions below:

- high density (HD):  $\omega^p(3B_Q) \geq A\ell(Q)^n$ , where  $A$  is a suitably large number,
- or low measure  $\mu$  (LM):  $\mu(B(Q)) \leq \tau\omega^p(Q)$ , where  $\tau$  is a suitably small number.

Any cube which is not contained in a bad cell will be called good.

Notice that  $\sum_{Q \text{ bad}} \mu(Q) \leq \tau \sum_{Q \text{ bad}} \omega^p(Q) \leq \tau \|\omega\| = \tau = \tau \mu(F_M)$ . Therefore, taking into account that  $\tau$  is small, that  $\mu$  is dominated by  $\mathcal{H}^n$  and that  $\omega^p|_{F_M}$  and  $\mathcal{H}^n|_{F_M}$  are boundedly equivalent, we conclude that  $\omega^p(F_M \setminus \bigcup_{Q \text{ bad}} Q) > 0$ . Also, taking into account that upper density of  $\omega^p$ ,  $\theta_{\omega^p}^*(x) = \limsup_{r \rightarrow 0} \frac{\omega^p(B(x,r))}{r^n}$ , is finite  $\omega^p|_E$ -a.e. (again because  $\omega^p|_E \ll \mathcal{H}^n|_E$ ), for  $A$  large enough, we can dispose of (HD) cubes as well. Ultimately, a little more careful consideration shows that  $\omega^p(F_M \cap \bigcup_{Q \in \mathcal{D}_0^{db}} Q \setminus \bigcup_{Q \text{ bad}} Q) > 0$ .<sup>1</sup> Notice also that the (HD) condition on bad cubes implies that

$$\mathcal{M}^n \omega^p(x) \lesssim A \quad \text{for } \omega^p\text{-a.e. } x \in F_M \setminus \bigcup_{Q \text{ bad}} Q. \tag{4}$$

It remains to concentrate on good cells contained in some cube from  $\mathcal{D}_0^{db}$ . We start with

**Lemma 3.1.** *For any good doubling cell  $R$  one has the estimate*

$$|\mathcal{R}_{r(B_R)} \omega^p(x)| \lesssim C(A, M, \tau, \text{dist}(p, \partial\Omega)) \quad \text{for any } x \in R.$$

The idea is to reduce the desired bound to certain estimates on the Green function of  $\Omega$ , which can be written as

$$G(x, y) = \mathcal{E}(x - y) - \int_{\partial\Omega} \mathcal{E}(x - z) d\omega^y(z) \quad \text{for } m\text{-a.e. } x \in \mathbb{R}^{n+1},$$

where  $\mathcal{E}$  is the fundamental solution for the Laplacian. Since the kernel of the Riesz transform is  $K(x) = c_n \nabla \mathcal{E}(x)$ , essentially differentiating the expression above and using a trivial estimate on  $\nabla \mathcal{E}(x)$  in terms of distance from  $p$  to  $\partial\Omega$ , we are left with considering the gradient of the Green function.<sup>2</sup> In fact, all we need is the bound  $\frac{1}{r} |G(y, p)| \lesssim 1$  for all  $y \in B(x, r) \cap \Omega$ ,  $r \approx r(B_Q)$ .

To this end, take  $\delta > 0$  from Bourgain’s Lemma and fix a point  $x^R$  such that  $\mu(B(x^R, \delta r(R)))$  is maximal (or almost maximal) possible. Denote  $B^R := B(x^R, \delta r(R))$ . Then automatically,  $\delta^n \mu(R) \lesssim \mu(B^R)$ . Thus, by Bourgain’s lemma and by (3) we have  $\inf_{z \in \partial(2B^R)} \omega^z(3B_R) \geq c \frac{\mu(R)}{r(R)^n}$ . This very easily implies (for  $n \geq 2$  at least) that for every  $z \in \partial(2B^R)$  and every  $y \in B^R$  one has  $\frac{1}{r(R)} G(z, y) \lesssim \frac{\omega^z(3B_R)}{\mu(R)}$ . Now the maximal principle yields (we also use that  $R$  is good and doubling in the last inequality)

$$\forall y \in B^R, \quad \frac{1}{r(R)} G(y, p) \lesssim \frac{\omega^p(3B_R)}{\mu(R)} \lesssim \tau^{-1}, \tag{5}$$

as desired. The similar inequality can be proved for  $n = 1$  but requires more work.

<sup>1</sup> Here we start working with a slightly enlarged collection of doubling cubes, but let us disregard these details in the present brief sketch.  
<sup>2</sup> In a careful argument, one has to utilize a smooth truncation of the kernel of the Riesz transform. The difference between the latter and  $\mathcal{R}_{r(B_R)} \omega^p(x^R)$  is fairly directly controlled by the  $\mathcal{M}^n \omega^p$ , which is in turn controlled by (4).

Consider now the case when  $Q$  is good and non-doubling,  $Q \notin \mathcal{D}^{db}$ , such that, in addition,  $Q \subset R_0$ , where  $R_0 \in \mathcal{D}_0^{db}$ . Let  $R \supset Q$  be the cube from  $\mathcal{D}^{db}$  with minimal side length containing  $Q$ . Then for all  $y \in B_Q$  we have

$$|\mathcal{R}_{\tau(B_Q)}\omega^p(y)| \leq |\mathcal{R}_{\tau(B_R)}\omega^p(y)| + C \sum_{P \in \mathcal{D}: Q \subset P \subset R} \theta_{\omega^p}(2B_P),$$

where  $\theta_{\omega^p}(B) := \frac{\omega^p(B)}{\tau(B)^n}$ . The first term is bounded by some constant as above since  $R$  is good and doubling. The sum is bounded by  $\theta_{\omega^p}(4B_R)$  as for all such cubes  $\omega^p(100B(P)) \leq A_0^{-10n(J(P)-J(R)-1)}\omega^p(100B(R))$  – see [6, Lemma 5.31]. Finally,  $\theta_{\omega^p}(4B_R)$  is bounded by  $A$  as  $R$  is not high density – see (HD). This combines with Lemma 3.1 to yield the key theorem:

**Theorem 3.2.** *Let  $G_M := F_M \cap \bigcup_{Q \in \mathcal{D}_0^{db}} Q \setminus \bigcup_{Q\text{-bad}} Q$ . Then  $\omega^p(G_M) > 0$  and*

$$\mathcal{R}_*\omega^p(x) \lesssim C(A, M, \tau, \text{dist}(p, \partial\Omega)) \quad \text{for } \omega^p\text{-a.e. } x \in G_M.$$

#### 4. Harmonic measure: its singular integral and its rectifiability

**Theorem 4.1.** *Let  $\sigma$  be a Radon measure with compact support on  $\mathbb{R}^{n+1}$  and consider a  $\sigma$ -measurable set  $G$  with  $\sigma(G) > 0$  such that*

$$G \subset \{x \in \mathbb{R}^d : \mathcal{M}^n\sigma(x) < \infty \text{ and } \mathcal{R}_*\sigma(x) < \infty\}.$$

*Then there exists a Borel subset  $G_0 \subset G$  with  $\sigma(G_0) > 0$  such that  $\sup_{x \in G_0} \mathcal{M}^n\sigma|_{G_0}(x) < \infty$  and  $\mathcal{R}_{\sigma|_{G_0}}$  is bounded in  $L^2(\sigma|_{G_0})$ .*

This result follows from the non-homogeneous  $Tb$  theorem of Nazarov, Treil and Volberg in [16] (see also [19]) in combination with the methods in [17]. For the detailed proof in the case of the Cauchy transform, see [18, Theorem 8.13]. The same arguments with very minor modifications work for the Riesz transform. This result applies to quite general antisymmetric Calderón–Zygmund operators. However, the next theorem uses very prominently that we work with the Riesz transform, and that its singularity is precisely  $n$ . It is the main result in [15], which is based in its turn on the solution to the (co-dimension-1) David–Semmes conjecture in [14] (see also [9] in the context of uniform domains).

**Theorem 4.2.** *Let  $\sigma$  be Radon measure on  $\mathbb{R}^{n+1}$  and  $\sigma \ll \mathcal{H}^n$ . If the vector  $n$ -dimensional Riesz transform  $\mathcal{R}$  is a bounded operator on  $L^2(\sigma)$ , then  $\sigma$  is rectifiable.*

By the latter theorem (or the David–Léger theorem [13] for  $n = 1$ ), we deduce that  $\omega^p|_{G_0}$  is  $n$ -rectifiable.

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