Functional analysis/Mathematical physics

# The essential spectrum of $N$-body systems with asymptotically homogeneous order-zero interactions 

# Le spectre essentiel des systèmes à $N$-corps avec interactions asymptotiquement homogènes d'ordre zéro 

Vladimir Georgescu ${ }^{\text {a }}$, Victor Nistor ${ }^{\text {b,c }}$<br>${ }^{\text {a }}$ Département de mathématiques, Université de Cergy-Pontoise, 95000 Cergy-Pontoise, France<br>${ }^{\text {b }}$ Université de Lorraine, UFR MIM, île du Saulcy, CS 50128, 57045 Metz cedex 01, France<br>${ }^{\text {c }}$ Pennsylvania State University, Mathematics Department, University Park, PA 16802, USA

## A R T I C L E I N F O

## Article history:

Received 5 July 2014
Accepted after revision 15 September 2014
Available online 23 October 2014
Presented by the Editorial Board


#### Abstract

We study the essential spectrum of $N$-body Hamiltonians with potentials defined by functions that have radial limits at infinity. The results extend the HVZ theorem which describes the essential spectrum of usual $N$-body Hamiltonians. The proof is based on a careful study of algebras generated by potentials and their cross-products. We also describe the topology on the spectrum of these algebras, thus extending to our setting a result of A. Mageira. Our techniques apply to more general classes of potentials associated with translation invariant algebras of bounded uniformly continuous functions on a finitedimensional vector space $X$.


© 2014 Académie des sciences. Published by Elsevier Masson SAS. All rights reserved.

## R É S U M É

Nous étudions le spectre essentiel des hamiltoniens des systèmes à $N$ corps avec potentiels définis par des fonctions qui ont des limites radiales à l'infini. Les résultats étendent le théorème HVZ, qui décrit le spectre essentiel des hamiltoniens des systèmes à $N$ corps usuels. La preuve de notre théorème principal est basée sur une étude approfondie des algèbres générées par les potentiels avec des limites radiales à l'infini et de leurs produits croisés. Nous décrivons également la topologie sur le spectre de ces algèbres, étendant ainsi à notre cas un résultat de A. Mageira. Nos techniques s'appliquent à des classes plus générales de potentiels associées à des algèbres de fonctions uniformément continues bornées invariantes par translation.
© 2014 Académie des sciences. Published by Elsevier Masson SAS. All rights reserved.

[^0]http://dx.doi.org/10.1016/j.crma.2014.09.029
1631-073X/© 2014 Académie des sciences. Published by Elsevier Masson SAS. All rights reserved.

## Version française abrégée

Soit $X$ un espace vectoriel réel de dimension finie et $\mathbb{S}_{X}:=(X \backslash\{0\}) / \mathbb{R}_{+}$la sphère à l'infini de $X$. On dit qu'une fonction $v: X \rightarrow \mathbb{C}$ a des limites radiales uniformes à l'infini si $v(\hat{a}):=\lim _{r \rightarrow \infty} v(r a)$ existe uniformément en $\hat{a} \in \mathbb{S}_{X}$. Soit $V_{Y}: X / Y \rightarrow \mathbb{R}$ une fonction borélienne ayant des limites radiales uniformes à l'infini, pour chaque sous-espace linéaire $Y \subset X$. Nous supposons $V_{Y}=0$, sauf pour un nombre fini d'espaces $Y$. On note $\pi_{Y}$ la surjection canonique $X \rightarrow X / Y$ et on garde la notation $V_{Y}$ pour la fonction $V_{Y} \circ \pi_{Y}$. Dans cet article, nous utilisons des produits croisés de $C^{*}$-algèbres pour étudier le spectre essentiel des opérateurs de la forme $H:=h(P)+\sum_{Y} V_{Y}$. Ici, $h: X^{*} \rightarrow[0, \infty[$ est une fonction continue et propre et $P$ est l'observable moment (formellement $P=-\mathrm{i} \nabla$ ). Soit $v: X \rightarrow \mathbb{C}$ et $a \neq 0$ tel que $\lim _{r \rightarrow \infty} v(r a+x)$ existe pour tout $x \in X$. Cette limite est une fonction de $x \in X$, qui ne dépend que de la classe $\alpha=\hat{a}$ de $a$ dans $\mathbb{S}_{X}$, que nous noterons $\tau_{\alpha}(v)$. Par exemple, si $v=V_{Y}$ avec $V_{Y}$ comme plus haut, alors $\tau_{\alpha}\left(V_{Y}\right)=V_{Y}$ si $\alpha \subset Y$ et $\tau\left(V_{Y}\right)=V_{Y}\left(\pi_{Y}(\alpha)\right) \in \mathbb{R}$ si $\alpha \not \subset Y$, où $\pi(\alpha) \in \mathbb{S}_{X / Y}$ est naturellement défini. Plus tard (voir le Théorème 3.1), nous définirons $\tau_{\alpha}(S)$ pour une classe générale d'opérateurs $S$, en particulier pour $S=H$, ce qui donnera une nouvelle signification à la définition de $\tau_{\alpha}$.

Nous énonçons maintenant un cas particulier de notre résultat principal : si les fonctions $V_{Y}: X / Y \rightarrow \mathbb{R}$ sont bornées et ont des limites radiales uniformes à l'infini et si, pour chaque $\alpha \in \mathbb{S}_{X}$, on pose $\tau_{\alpha}(H)=h(P)+\sum_{Y \supset \alpha} V_{Y}+\sum_{Y \not \supset \alpha} V_{Y}\left(\pi_{Y}(\alpha)\right)$, alors le spectre essentiel de $H$ est $\sigma_{\text {ess }}(H)=\bar{\bigcup}_{\alpha \in \mathbb{S}_{X}} \sigma\left(\tau_{\alpha}(H)\right)$.

## 1. Introduction

Let $X$ be a finite dimensional real vector space and, for each linear subspace $Y$ of $X$, let $V_{Y}: X / Y \rightarrow \mathbb{R}$ be a Borel function. We assume $V_{Y}=0$, except for a finite number of $Y$. We keep the notation $V_{Y}$ for the function on $X$ given by $V_{Y} \circ \pi_{Y}$, where $\pi_{Y}: X \rightarrow X / Y$ is the natural map. In this paper, we use crossed-products of $C^{*}$-algebras to study the essential spectrum of Hamiltonians of the form $H:=h(P)+\sum_{Y} V_{Y}$, under certain conditions on the potentials $V_{Y}$. Here $h: X^{*} \rightarrow[0, \infty[$ is a continuous, proper function and $P$ is the momentum observable (recall that proper means that $\left.\lim _{|k| \rightarrow \infty} h(k)=+\infty\right)$. More precisely, $h(P)=\mathcal{F}^{-1} M_{h} \mathcal{F}$, where $\mathcal{F}: L^{2}(X) \rightarrow L^{2}\left(X^{*}\right)$ is the Fourier transform and $M_{h}$ is the operator of multiplication by $h$ (formally $P=-i \nabla$ ). Operators of this form cover the Hamiltonians that are currently the most interesting (from a physical point of view) Hamiltonians of $N$-body systems. Here are two main examples. In a generalized version of the non-relativistic case, a scalar product is given on $X$, so, by taking $h(\xi)=|\xi|^{2}$, we get $h(P)=\Delta$, the positive Laplacian. In the simplest relativistic case, $X=\left(\mathbb{R}^{3}\right)^{N}$ and, writing the momentum $P=\left(P_{1}, \ldots, P_{N}\right)$, we have $h(P)=\sum_{k=1}^{N}\left(P_{k}^{2}+m_{k}^{2}\right)^{1 / 2}$ for some real numbers $m_{k}$. We refer to [3] for a thorough introduction to the subject and study of these systems.

Let $\mathbb{S}_{X}:=(X \backslash\{0\}) / \mathbb{R}_{+}$be the sphere at infinity of $X$, i.e. the set of all half-lines $\hat{a}:=\mathbb{R}_{+} a$. A function $v: X \rightarrow \mathbb{C}$ is said to have uniform radial limits at infinity if $v(\hat{a}):=\lim _{r \rightarrow \infty} v(r a)$ exists uniformly in $\hat{a} \in \mathbb{S}_{X}$. From the definition of the topology on $\mathbb{S}_{X}$, we get $v(\hat{a})=\lim _{r \rightarrow \infty} v(r a+x), \forall x \in X$. More generally, we are interested in functions $v$ such that $\lim _{r \rightarrow \infty} v(r a+x)$ exists for all $x \in X$. The limit may depend on $x$ and defines a function $\tau_{\alpha}(v): X \rightarrow \mathbb{C}$, where $\alpha:=\hat{a}$. For example, let us consider $v=V_{Y}$. Then $\tau_{\alpha}\left(V_{Y}\right)(x)=\lim _{r \rightarrow \infty} V_{Y}\left(r \pi_{Y}(a)+\pi_{Y}(x)\right)$. In particular, $\tau_{\alpha}\left(V_{Y}\right)=V_{Y}$ whenever $\alpha:=\hat{a} \subset Y$ (i.e. $a \in Y$ ). On the other hand, if $V_{Y}: X / Y \rightarrow \mathbb{C}$ has uniform radial limits at infinity and $\hat{a}=\alpha \not \subset Y$, then $\pi_{Y}(\alpha):=\mathbb{R}_{+} \pi_{Y}(a) \in \mathbb{S}_{X / Y}$ is well defined and $\tau_{\alpha}\left(V_{Y}\right)(x)=V_{Y}\left(\pi_{Y}(\alpha)\right)$ turns out to be a constant.

Theorem 1.1. Let $V_{Y}: X / Y \rightarrow \mathbb{R}$ be bounded with uniform radial limits at infinity. If $\alpha \in \mathbb{S}_{X}$ set

$$
\begin{equation*}
\tau_{\alpha}(H)=h(P)+\sum_{Y} \tau_{\alpha}\left(V_{Y}\right)=h(P)+\sum_{Y \supset \alpha} V_{Y}+\sum_{Y \not \supset \alpha} V_{Y}\left(\pi_{Y}(\alpha)\right) \tag{1}
\end{equation*}
$$

Then $\sigma\left(\tau_{\alpha}(H)\right)=\left[c_{\alpha}, \infty\right)$ for some real $c_{\alpha}$ and $\sigma_{\text {ess }}(H)=\bigcup_{\alpha \in \mathbb{S}_{X}} \sigma\left(\tau_{\alpha}(H)\right)=\left[\inf _{\alpha} c_{\alpha}, \infty\right)$.

Here $\bar{\bigcup}_{\alpha}$ is the closure of the union. Sometimes the union is already closed [11]. Unbounded potentials are considered in Theorem 3.2. If all the radial limits are zero, which is the case of the usual $N$-body potentials, then the terms corresponding to $\alpha \not \subset Y$ are dropped in Eq. (1). Consequently, if $h(P)=\Delta$ is the non-relativistic kinetic energy, we recover the Hunziker, van Winter, Zhislin (HVZ) theorem. Descriptions of the essential spectrum of various classes of self-adjoint operators in terms of limits at infinity of translates of the operators have already been obtained before, see for example [7,12,8] (in historical order). Our approach is based on the "localization at infinity" technique developed in [5,6] in the context of crossed-product $C^{*}$-algebras.

Let us sketch the main idea of this approach. Let $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ be the algebra of bounded uniformly continuous functions, $\mathcal{C}_{0}(X)$ the ideal of functions vanishing at infinity, and $\mathcal{C}\left(X^{+}\right)=\mathbb{C}+\mathcal{C}_{0}(X)$. Consider a translation invariant $C^{*}$-subalgebra $\mathcal{A} \subset \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ containing $\mathcal{C}\left(X^{+}\right)$and let $\hat{\mathcal{A}}$ be its character space. Note that $\hat{\mathcal{A}}$ is a compact topological space that naturally contains $X$ as an open dense subset and $\delta(\mathcal{A})=\hat{\mathcal{A}} \backslash X$ can be thought of as a boundary of $X$ at infinity. Recall that a self-adjoint operator $H$ on a Hilbert space $\mathcal{H}$ is said to be affiliated to a $C^{*}$-algebra $\mathscr{A}$ of operators on $\mathcal{H}$ if one has $(H+i)^{-1} \in \mathscr{A}$. Then with each self-adjoint operator $H$ affiliated to the crossed product $\mathcal{A} \rtimes X$ of $\mathcal{A}$ by the action of $X$, one may associate a family of self-adjoint operators $H_{\varkappa}$ affiliated to $\mathcal{A} \rtimes X$ indexed by the characters $\varkappa \in \delta(\mathcal{A})$. This family
completely describes the image of $H$ (in the sense of affiliated operators) in the quotient of $\mathcal{A} \rtimes X$ with respect to the ideal of compact operators. In particular, the essential spectrum of $H$ is the closure of the union of the spectra of the operators $H_{\varkappa}$. These operators are the localizations at infinity of $H$, more precisely, $H_{\varkappa}$ is the localization of $H$ at point $\varkappa$.

Once chosen the algebra $\mathcal{A}$, in order to use these techniques of this paper, we also need: (1) to have a good description of the character space of the Abelian algebra $\mathcal{A}$, and (2) to have an efficient criterion for affiliation to the crossed product $\mathcal{A} \rtimes X$. We also indicate how to achieve (1) and (2).

## 2. Crossed products and localizations at infinity

For $p \in X^{*}$ and $q \in X$ let $\left(S_{p} f\right)(x)=\mathrm{e}^{\mathrm{i}\langle x \mid p\rangle} f(x)$ and $\left(T_{q} f\right)(q)=f(x+q)$. We say that $A \in \mathcal{B}\left(L^{2}(X)\right)$ has the positionmomentum limit property if $\lim _{p \rightarrow 0}\left\|\left[S_{p}, A\right]\right\|=0$ and $\lim _{q \rightarrow 0}\left\|\left(T_{q}-1\right) A^{(*)}\right\|=0$ (where $A^{(*)}$ means that the relation holds for $A$ and $A^{*}$ ). The set of such operators is a $C^{*}$-algebra equal to the crossed product $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X) \rtimes X$ [5]. Note that if $\mathcal{A}$ is a translation invariant $C^{*}$-subalgebra of $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$, then there is a natural realization of the abstract crossed product $\mathcal{A} \rtimes X$ as a $C^{*}$-algebra of operators on $L^{2}(X)$ and we do not distinguish the two algebras. We describe this concrete version of $\mathcal{A} \rtimes X$ below.

If $\varphi: X \rightarrow \mathbb{C}$ and $\psi: X^{*} \rightarrow \mathbb{C}$ are measurable functions, then $\varphi(Q)$ and $\psi(P)$ are the operators on $L^{2}(X)$ defined as follows: $\varphi(Q):=M_{\phi}$ acts as multiplication by $\varphi$ and $\psi(P)=\mathcal{F}^{-1} M_{\psi} \mathcal{F}$, where $\mathcal{F}$ is the Fourier transform $L^{2}(X) \rightarrow$ $L^{2}\left(X^{*}\right)$ and $M_{\psi}$ is the operator of multiplication by $\psi$. Then $\psi \mapsto \psi(P)$ is an isomorphism between $\mathcal{C}_{0}\left(X^{*}\right)$ and the group $C^{*}$-algebra $C^{*}(X)$ and $\mathcal{A} \rtimes X$ is the norm closed linear space of bounded operators on $L^{2}(X)$ generated by the products $\varphi(Q) \psi(P)$ with $\varphi \in \mathcal{A}$ and $\psi \in \mathcal{C}_{0}\left(X^{*}\right)$. In particular, $\mathcal{A} \rtimes X$ consists of operators that have the position-momentum limit property.

We recall the definition of localizations at infinity for such operators. Assume $\mathcal{C}\left(X^{+}\right) \subset \mathcal{A}$, so $\hat{\mathcal{A}}$ is a compactification of $X$ and $\delta(\mathcal{A})=\hat{\mathcal{A}} \backslash X$ is a compact. If $q \in X$ and $\varphi$ is a function on $X$ then $T_{q} \varphi$ is its translation by $q$. We extend this definition of $T_{q}$ by replacing $q \in X$ with $x \in \hat{\mathcal{A}}:\left(T_{\varkappa} \varphi\right)(x)=\varkappa\left(T_{x} \varphi\right)$, for any $\varphi \in \mathcal{A}, x \in \hat{\mathcal{A}}$, and, $x \in X$. It is clear that $T_{\varkappa} \varphi \in \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ and that its definition coincides with the previous one if $x=q \in X$. Moreover, we also get "translations at infinity" of $\varphi \in \mathcal{A}$ by elements $x \in \delta(\mathcal{A})$; note however that such a translation does not belong to $\mathcal{A}$ in general. Also, the function $\varkappa \mapsto T_{\varkappa} \varphi \in \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ defined on $\hat{\mathcal{A}}$ is continuous if $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ is equipped with the topology of local uniform convergence, hence $T_{\varkappa} \varphi=\lim _{q \rightarrow \varkappa} T_{q} \varphi$ in this topology for any $\varkappa \in \delta(\mathcal{A})$. If $A$ is an operator on $L^{2}(X)$, let $\tau_{q}(A)=T_{q} A T_{q}^{*}$ be its translation by $q \in X$. Clearly $\tau_{q}(\varphi(Q))=\left(T_{q} \varphi\right)(Q)$. If $A \in \mathcal{A} \rtimes X$, then we may also consider "translations at infinity" by elements of the boundary $\delta(\mathcal{A})$ of $X$ in $\hat{\mathcal{A}}$ and we get a useful characterization of the compact operators. The following are mainly consequences of [6, Theorem 1.15]:
(i) For each $\varkappa \in \hat{\mathcal{A}}$, there is a unique morphism $\tau_{\varkappa}: \mathcal{A} \rtimes X \rightarrow \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X) \rtimes X$ such that $\tau_{\varkappa}(\varphi(Q) \psi(P))=\left(T_{\varkappa} \varphi\right)(Q) \psi(P)$, $\varphi \in \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X), \psi \in \mathcal{C}_{0}(X)$. (ii) $\bigcap_{\varkappa \in \delta(\mathcal{A})} \operatorname{ker} \tau_{\varkappa}=\mathcal{C}_{0}(X) \rtimes X \equiv \mathscr{K}(X)=$ ideal of compact operators on $L^{2}(X)$. (iii) If $H$ is a self-adjoint operator on $L^{2}(X)$ affiliated to $\mathcal{A}$ then for each $x \in \delta(\mathcal{A})$ the limit $\tau_{\varkappa}(H):=\mathrm{s}-\lim _{q \rightarrow \varkappa} T_{q} H T_{q}^{*}$ exists and $\sigma_{\text {ess }}(H)=\bar{\bigcup}_{\varkappa \in \delta(\mathcal{A})} \sigma\left(\tau_{\varkappa}(H)\right)$.

To be precise, the last strong limit means: $\tau_{\varkappa}(H)$ is a self-adjoint operator (not necessarily densely defined) on $L^{2}(X)$ and s- $\lim _{q \rightarrow \varkappa} \theta\left(T_{q} H T_{q}^{*}\right)=\theta\left(\tau_{\varkappa}(H)\right)$ for all $\theta \in \mathcal{C}_{0}(\mathbb{R})$. It is clear that in the last three statements above one may replace $\delta(\mathcal{A})$ by a subset $\pi$ if for each $A \in \mathcal{A} \rtimes X$ we have: $\tau_{\varkappa}(A)=0 \forall \varkappa \in \pi \Rightarrow \tau_{\varkappa}(A)=0 \forall \varkappa \in \delta(\mathcal{A})$. In the case of groupoid (pseudo)differential algebras (that is, when $\hat{\mathcal{A}}$ is a manifold with corners), the morphisms $\tau_{\varkappa}$ can be defined using restrictions to fibers, as in [9], and the last three statements above (i)-(iii) remain valid.

## 3. Main results

As a warm-up and in order to introduce some general notation, we treat first the two-body case, where complete results may be obtained by direct arguments. The algebra of interactions in the standard two-body case is $\mathcal{C}\left(X^{+}\right)$, and hence the Hamiltonian algebra is

$$
\begin{equation*}
\mathcal{C}\left(X^{+}\right) \rtimes X=\mathbb{C} \rtimes X+\mathcal{C}_{0}(X) \rtimes X=C^{*}(X)+\mathscr{K}(X) \tag{2}
\end{equation*}
$$

where the sums are direct. Thus $\mathcal{C}\left(X^{+}\right) \rtimes X / \mathscr{K}(X)=C^{*}(X)$, which finishes the theory. Another elementary case, which has been considered as an example in [5], is $X=\mathbb{R}$ with $\mathcal{C}\left(\mathbb{R}^{+}\right)$replaced by the algebra $\mathcal{C}(\overline{\mathbb{R}})$ of continuous functions that have limits (distinct in general) at $\pm \infty$. Then there is no natural direct sum decomposition of $\mathcal{C}(\overline{\mathbb{R}}) \rtimes \mathbb{R}$ as in (2), but one has, by standard arguments, $\mathcal{C}(\overline{\mathbb{R}}) \rtimes \mathbb{R} / \mathscr{K}(\mathbb{R}) \simeq C^{*}(\mathbb{R}) \oplus C^{*}(\mathbb{R})$. Our purpose in this section is to extend this equation to arbitrary $X$.

Let $\mathcal{C}(\bar{X})$ be the closure in $\mathcal{C}_{\mathrm{b}}(X)$ of the subalgebra of functions homogeneous of degree zero outside a compact set. Then $\mathcal{C}(\bar{X})=\left\{u \in \mathcal{C}(X) \mid \lim _{\lambda \rightarrow+\infty} u(\lambda a)\right.$ exists uniformly in $\left.\hat{a} \in \mathbb{S}_{X}\right\}$, where, we recall, $\hat{a}:=\mathbb{R}_{+} a$ and $\mathbb{S}_{X}:=(X \backslash\{0\}) / \mathbb{R}_{+}$, so $\hat{a} \in \mathbb{S}_{X}$. As a set, the character space of $\mathcal{C}(\bar{X})$ can be identified with the disjoint union $\bar{X}=X \cup \mathbb{S}_{X}$. The topology induced by the character space on $X$ is the usual one and the intersections with $X$ of the neighborhoods of some $\alpha \in \mathbb{S}_{X}$ are the sets that contain a truncated cone $C$ such that there is $a \in \alpha$ such $\lambda a \in C$ if $\lambda \geq 1$. The set of such subsets is a filter $\tilde{\alpha}$ on $X$ and, if $Y$ is a Hausdorff space and $u: X \rightarrow Y$, then $\lim _{\tilde{\alpha}} u=y$ means that $u^{-1}(V) \in \tilde{\alpha}$ for any neighborhood $V$ of $u$. We shall
write $\lim _{x \rightarrow \alpha} u(x)$ instead of $\lim _{\tilde{\alpha}} u$. We have that $\mathcal{C}(\bar{X})$ is a translation invariant $C^{*}$-subalgebra of $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ and so the crossed product $\mathcal{C}(\bar{X}) \rtimes X$ is well defined. We have the following explicit description of this algebra.

Proposition 3.1. The algebra $\mathcal{C}(\bar{X}) \rtimes X$ acting on $L^{2}(X)$ consists of bounded operators $A$ that have the position-momentum limit property and are such that the limit $\tau_{\alpha}(A)=s-\lim _{a \rightarrow \alpha} T_{a} A T_{a}^{*}$ exists for each $\alpha=\hat{a} \in \mathbb{S}_{X}$. If $A \in \mathcal{C}(\bar{X}) \rtimes X$ and $\alpha \in \mathbb{S}_{X}$, then $\tau_{\alpha}(A) \in$ $C^{*}(X)$ and $\tau(A): \alpha \mapsto \tau_{\alpha}(A)$ is norm continuous. The map $\tau: \mathcal{C}(\bar{X}) \rtimes X \rightarrow C\left(\mathbb{S}_{X}\right) \otimes C^{*}(X)$ is a surjective morphism whose kernel is the set of compact operators on $L^{2}(X)$, which gives $\mathcal{C}(\bar{X}) \rtimes X / \mathscr{K}(X) \cong C\left(\mathbb{S}_{X}\right) \otimes C^{*}(X)$. If $H$ is a self-adjoint operator affiliated to $\mathcal{C}(\bar{X}) \rtimes X$ then $\tau_{\alpha}(H)=s-\lim _{a \rightarrow \alpha} T_{a} H T_{a}^{*}$ exists for all $\alpha \in \mathbb{S}_{X}$ and $\sigma_{\text {ess }}(H)=\bigcup_{\alpha} \sigma\left(\tau_{\alpha}(H)\right)$.

In the next two examples $H=h(P)+V$ with $h: X^{*} \rightarrow[0, \infty[$ continuous and proper. We denote by $|\cdot|$ a fixed norm on $X^{*}$ and by $\mathcal{H}^{s}$ we denote the usual Sobolev spaces on $X(s \in \mathbb{R})$.

Example 1. Let $V$ be a bounded symmetric operator satisfying: (1) $\lim _{p \rightarrow 0}\left\|\left[S_{p}, V\right]\right\|=0$ and (2) the limit $\tau_{\alpha}(V)=$ $s$ - $\lim _{a \rightarrow \alpha} T_{a} V T_{a}^{*}$ exists for each $\alpha \in \mathbb{S}_{X}$. Then $H$ is affiliated to $\mathcal{C}(\bar{X}) \rtimes X$ and $\tau_{\alpha}(H)=h(P)+\tau_{\alpha}(V)$. Moreover, if $V$ is a function, then $\tau_{\alpha}(V)$ is a number, but in general we have $\tau_{\alpha}(V)=v_{\alpha}(P)$ for some function $v_{\alpha} \in \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}\left(X^{*}\right)$.

Example 2. Assume that $h$ is locally Lipschitz and that there exist $c, s>0$ such that, for all $p$ with $|p|>1,|\nabla h(p)| \leq c(1+$ $h(p))$ and $c^{-1}|p|^{s} \leq(1+h(p))^{1 / 2} \leq c|p|^{s}$. Let $V: \mathcal{H}^{s} \rightarrow \mathcal{H}^{-s}$ such that $\pm V \leq \mu h(P)+v$ for some numbers $\mu, v$ with $\mu<1$ and satisfying the next two conditions: (1) $\lim _{p \rightarrow 0}\left\|\left[S_{p}, V\right]\right\|_{\mathcal{H}^{s} \rightarrow \mathcal{H}^{-s}}=0$, (2) $\forall \alpha \in \mathbb{S}_{X}$ the limit $\tau_{\alpha}(V)=s-\lim _{a \rightarrow \alpha} T_{a} V T_{a}^{*}$ exists strongly in $\mathcal{B}\left(\mathcal{H}^{s}, \mathcal{H}^{-s}\right)$. Then $h(P)+V$ and $h(P)+\tau_{\alpha}(V)$ are symmetric operators $\mathcal{H}^{s} \rightarrow \mathcal{H}^{-s}$ that induce self-adjoint operators $H$ and $\tau_{\alpha}(H)$ in $L^{2}(X)$ affiliated to $\mathcal{C}(\bar{X}) \rtimes X$ and $\sigma_{\text {ess }}(H)=\bigcup_{\alpha} \sigma\left(\tau_{\alpha}(H)\right)$.

We now treat the $N$-body case. We first indicate a general way of constructing $N$-body Hamiltonians. For each linear subspace $Y \subset X$, let $\mathcal{A}(X / Y) \subset \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X / Y)$ be a translation invariant $C^{*}$-subalgebra containing $\mathcal{C}_{0}(X / Y)$ with $\mathcal{A}(X / X)=\mathcal{A}(0)=\mathbb{C}$. We embed $\mathcal{A}(X / Y) \subset \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ as usual by identifying $v$ with $v \circ \pi_{Y}$. Then the $C^{*}$-algebra $\mathcal{A}$ generated by these algebras is a translation invariant $C^{*}$-subalgebra of $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ containing $\mathcal{C}\left(X^{+}\right)$and thus we may consider the crossed product $\mathcal{A} \rtimes X$ which is equal to the $C^{*}$-algebra generated by the crossed products $\mathcal{A}(X / Y) \rtimes X$. The operators affiliated to $\mathcal{A} \rtimes X$ are $N$-body Hamiltonians. The standard $N$-body algebra corresponds to the minimal choice $\mathcal{A}(X / Y)=\mathcal{C}_{0}(X / Y)$ and has remarkable properties, which makes its study relatively easy (it is graded by the lattice of subspaces of $X$ ). Our purpose in this paper is to study what could arguably be considered to be the simplest extension of the classical $N$-body obtained by choosing $\mathcal{A}(X / Y)=\mathcal{C}(\overline{X / Y})$ for all $Y$. The next more general case would correspond to the choice $\mathcal{A}(X / Y)=\mathcal{V}(X / Y)$ (slowly oscillating functions, i.e. the closure in sup norm of the set of bounded functions of class $C^{1}$ with derivatives tending to zero at infinity).

Definition 3.2. Let $\mathcal{E}(X)$ be the $C^{*}$-subalgebra of $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ generated by $\bigcup_{Y} \mathcal{C}(\overline{X / Y})$.
Clearly $\mathcal{E}(X)$ is a translation invariant $C^{*}$-subalgebra of $\mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X)$ containing $\mathcal{C}\left(X^{+}\right):=\mathcal{C}_{0}(X)+\mathbb{C}$. If $Y$ is a linear subspace of $X$ then the $C^{*}$-algebra $\mathcal{E}(X / Y) \subset \mathcal{C}_{\mathrm{b}}^{\mathrm{u}}(X / Y)$ is well defined and naturally embedded in $\mathcal{E}(X)$ : it is the $C^{*}$-algebra generated by $\cup_{Z \supset Y} C(\overline{X / Z})$. We have $\mathbb{C}=\mathcal{E}(0)=\mathcal{E}(X / X) \subset \mathcal{E}(X / Y) \subset \mathcal{E}(X / Z) \subset \mathcal{E}(X)$. If $\alpha \in \mathbb{S}_{X}$, we shall denote by abuse of notation $X / \alpha$ be the quotient $X /[\alpha]$ of $X$ by the subspace $[\alpha]:=\mathbb{R} \alpha$ generated by $\alpha$ and let us set $\pi_{\alpha}=\pi_{[\alpha]}$. It is clear that $\tau_{\alpha}(u)(x)=\lim _{r \rightarrow+\infty} u(r a+x)$ exists $\forall u \in \mathcal{E}(X)$ and that the resulting function $\tau_{\alpha}(u)$ belongs to $\mathcal{E}(X)$. The map $\tau_{\alpha}$ is an endomorphism of $\mathcal{E}(X)$ and a linear projection of $\mathcal{E}(X)$ onto the $C^{*}$-subalgebra $\mathcal{E}(X / \alpha)$.

If $\alpha \in \mathbb{S}_{X}$ and $\beta \in \mathbb{S}_{X / \alpha}$, then $\beta$ generates a one-dimensional linear subspace $[\beta]:=\mathbb{R} \beta \subset X / \alpha$, as above, and hence $\pi_{\alpha}^{-1}([\beta])$ is a two-dimensional subspace of $X$ that we shall denote $[\alpha, \beta]$. We shall identify $(X / \alpha) / \beta$ with $X /[\alpha, \beta]$. Then we have two idempotent morphisms $\tau_{\alpha}: \mathcal{E}(X) \rightarrow \mathcal{E}(X / \alpha)$ and $\tau_{\beta}: \mathcal{E}(X / \alpha) \rightarrow \mathcal{E}(X /[\alpha, \beta])$. Thus $\tau_{\beta} \tau_{\alpha}: \mathcal{E}(X) \rightarrow \mathcal{E}(X /[\alpha, \beta])$ is an idempotent morphism. This construction extends in an obvious way to families $\vec{\alpha}=\left(\alpha_{1}, \ldots, \alpha_{n}\right)$ with $n \leq \operatorname{dim} X$ and $\alpha_{1} \in \mathbb{S}_{X}, \alpha_{2} \in \mathbb{S}_{X / \alpha_{1}}, \alpha_{3} \in \mathbb{S}_{X /\left[\alpha_{1}, \alpha_{2}\right]}, \ldots$ (we allow $n=0$ and denote $A$ the set of all such families). The endomorphism $\tau_{\vec{\alpha}}$ of $\mathcal{E}(X)$ is defined by induction: $\tau_{\vec{\alpha}}=\tau_{\alpha_{n}} \ldots \tau_{\alpha_{1}}$. We also define $\left[\alpha_{1}, \alpha_{2}, \ldots, \alpha_{n}\right.$ ] by induction, so this is an $n$-dimensional subspace of $X$ associated with $\left(\alpha_{1}, \alpha_{2}, \ldots, \alpha_{n}\right)$ and we denote $X / \vec{\alpha}$ the quotient of $X$ with respect to it. Thus $\tau_{\vec{\alpha}}$ is an endomorphism of $\mathcal{E}(X)$ and a projection of $\mathcal{E}(X)$ onto $\mathcal{E}(X / \vec{\alpha})$.

Proposition 3.3. If $\vec{\alpha} \in A$ and $a \in X / \vec{\alpha}$, then $\varkappa(u)=\left(\tau_{\vec{\alpha}} u\right)(a)$ defines a character of $\mathcal{E}(X)$. Conversely, each character of $\mathcal{E}(X)$ is of this form.

Remark 1. A natural Abelian $C^{*}$-algebra in the present context is the set $\mathcal{R}(X)$ of all bounded uniformly continuous functions $v: X \rightarrow \mathbb{C}$ such that $\lim _{r \rightarrow \infty} v(r a+x)$ exists locally uniformly in $x \in X$ for each $a \in X$. It would be interesting to find an explicit description of its spectrum.

This description of the spectrum of $\mathcal{E}(X)$ extends [10]. We now state our main results.

Theorem 3.1. Let $H$ be a self-adjoint operator on $L^{2}(X)$ affiliated to $\mathcal{E}(X) \rtimes X$. Then for any $a \in X \backslash\{0\}$ the limit $s-\lim _{r \rightarrow+\infty} T_{r a} H T_{r a}^{*}=$ : $\tau_{\hat{a}}(H)$ exists and $\sigma_{\text {ess }}(H)=\bar{\bigcup}_{\alpha \in \mathbb{S}_{X}} \sigma\left(\tau_{\alpha}(H)\right)$.

Theorem 3.2. Let $h$ be as in Example 2 and $V=\sum V_{Y}$ with $V_{Y}: \mathcal{H}^{s} \rightarrow \mathcal{H}^{-s}$ symmetric operators such that $V_{Y}=0$ but for a finite number of $Y$ and satisfying: (i) $\exists \mu_{Y}, \nu_{Y} \geq 0$ with $\sum_{Y} \mu_{Y}<1$ such that $\pm V_{Y} \leq \mu_{Y} h(P)+\nu_{Y}$, (ii) $\lim _{p \rightarrow 0}\left\|\left[S_{p}, V_{Y}\right]\right\|_{\mathcal{H}^{s} \rightarrow \mathcal{H}^{-s}=0 \text {, }}$ (iii) $\left[T_{y}, V_{Y}\right]=0$ for all $y \in Y$, (iv) $\tau_{\alpha}\left(V_{Y}\right):=s-\lim _{a \rightarrow \alpha} T_{a} V_{Y} T_{a}^{*}$ exists in $B\left(\mathcal{H}^{s}, \mathcal{H}^{-s}\right)$ for all $\alpha \in \mathbb{S}_{X / Y}$. Then the maps $\mathcal{H}^{s} \rightarrow \mathcal{H}^{-s}$ given by $h(P)+V$ and $h(P)+\sum_{Y} \tau_{\alpha}\left(V_{Y}\right)$ induce self-adjoint operators $H$ and $\tau_{\alpha}(H)$ in $L^{2}(X)$ affiliated to $\mathscr{E}(X)$ and $\sigma_{\text {ess }}(H)=$ $\bar{\bigcup}_{\alpha \in \mathbb{S}_{X}} \sigma\left(\tau_{\alpha}(H)\right)$.

Example 3. Using [2], we also obtain that Theorem 3.2 covers uniformly elliptic operators of the form $H=\sum_{|k|,|\ell| \leq s} P^{k} a_{k \ell} P^{\ell}$, where $a_{k \ell}$ are finite sums of functions of the form $v_{Y} \circ \pi_{Y}$ with $v_{Y}: X / Y \rightarrow \mathbb{R}$ bounded measurable such that $\lim _{z \rightarrow \alpha} v_{Y}(z)$ exists uniformly in $\alpha \in \mathbb{S}_{X / Y}$. The fact that we allow $a_{k \ell}$ to be only bounded measurable for $|k|=|\ell|=s$ is not trivial.

In addition to the above-mentioned results, we also use general results on cross-product $C^{*}$-algebras, their ideals, and their representations [4,13]. The maximal ideal spectrum of the algebra $\mathcal{E}(X)$ is of independent interest and can be used to study the regularity properties of the eigenvalues of the $N$-body Hamiltonian [1]. Its relation to the constructions of Vasy in [14] will be studied elsewhere.

## Acknowledgements

We thank Bernd Ammann for several useful discussions. Victor Nistor was partially supported by ANR SINGSTAR 2014-18 and by NSF Grant DMS-1016556.

## References

[1] B. Ammann, C. Carvalho, V. Nistor, Regularity for eigenfunctions of Schrödinger operators, Lett. Math. Phys. 101 (2012) 49-84.
[2] M. Damak, V. Georgescu, Self-adjoint operators affiliated to C*-algebras, Rev. Math. Phys. 16 (2) (2004) 257-280.
[3] J. Dereziński, C. Gérard, Scattering Theory of Classical and Quantum N-Particle Systems, Texts and Monographs in Physics, Springer-Verlag, Berlin, 1997.
[4] T. Fack, G. Skandalis, Sur les représentations et idéaux de la $C^{*}$-algèbre d'un feuilletage, J. Oper. Theory 8 (1) (1982) 95-129.
[5] V. Georgescu, A. Iftimovici, Crossed products of $C^{*}$-algebras and spectral analysis of quantum Hamiltonians, Commun. Math. Phys. 228 (3) (2002) 519-560.
[6] V. Georgescu, A. Iftimovici, Localizations at infinity and essential spectrum of quantum Hamiltonians. I. General theory, Rev. Math. Phys. 18 (4) (2006) 417-483.
[7] B. Helffer, A. Mohamed, Caractérisation du spectre essentiel de l'opérateur de Schrödinger avec un champ magnétique, Ann. Inst. Fourier (Grenoble) 38 (2) (1988) 95-112.
[8] Y. Last, B. Simon, The essential spectrum of Schrödinger, Jacobi, and CMV operators, J. Anal. Math. 98 (2006) 183-220.
[9] R. Lauter, B. Monthubert, V. Nistor, Pseudodifferential analysis on continuous family groupoids, Doc. Math. 5 (2000) 625-655 (electronic).
[10] A. Mageira, Graded C*-algebras, J. Funct. Anal. 254 (6) (2008) 1683-1701.
[11] V. Nistor, N. Prudhon, Exhausting families of representations and spectra of pseudodifferential operators (in final preparation).
[12] V. Rabinovich, S. Roch, B. Silbermann, Limit Operators and Their Applications in Operator Theory, Operator Theory: Advances and Applications, vol. 150, Birkhäuser Verlag, Basel, Switzerland, 2004.
[13] J. Renault, A Groupoid Approach to C*-Algebras, Lecture Notes in Mathematics, vol. 793, Springer, Berlin, 1980.
[14] A. Vasy, Propagation of singularities in many-body scattering, Ann. Sci. Éc. Norm. Super. (4) 34 (3) (2001) 313-402.


[^0]:    E-mail addresses: vladimir.georgescu@math.cnrs.fr (V. Georgescu), victor.nistor@univ-lorraine.fr (V. Nistor).

