# $L^{p}$-theory for vector potentials and Sobolev's inequalities for vector fields 

# Théorie $L^{p}$ pour les potentiels vecteurs et inégalités de Sobolev pour des champs de vecteurs 

Chérif Amrouche ${ }^{\text {a }}$, Nour El Houda Seloula ${ }^{\text {a,b }}$<br>${ }^{\text {a }}$ Laboratoire de mathématiques appliquées, CNRS UMR 5142, université de Pau et des Pays de l'Adour, IPRA, avenue de l'université, 64000 Pau, France<br>${ }^{\text {b }}$ EPI Concha, LMA UMR CNRS 5142, INRIA Bordeaux-Sud-Ouest, 64000 Pau, France

## A R TICLE I N F O

## Article history:

Received 7 July 2010
Accepted 31 March 2011
Available online 4 May 2011
Presented by Philippe G. Ciarlet


#### Abstract

In a three-dimensional bounded possibly multiply-connected domain, we prove the existence and uniqueness of vector potentials in $L^{p}$-theory, associated with a divergencefree function and satisfying some boundary conditions. We also present some results concerning scalar potentials and weak vector potentials. Furthermore, various Sobolev-type inequalities are given.


© 2011 Académie des sciences. Published by Elsevier Masson SAS. All rights reserved.

## R É S U M É

Dans un ouvert borné tridimensionnel, éventuellement multiplement connexe, nous prouvons l'existence et l'unicité des potentiels vecteurs en théorie $L^{p}$, associés à des champs de vecteurs à divergence nulle et vérifiant plusieurs conditions aux limites. On présente également des résultats concernant les potentiels scalaires et les potentiels vecteurs faibles. De plus, plusieurs inégalités de Sobolev sont données.
© 2011 Académie des sciences. Published by Elsevier Masson SAS. All rights reserved.

## Version française abrégée

Dans cette Note on s'intéresse à la théorie des potentiels vecteurs dans un ouvert $\Omega$ borné tridimensionnel éventuellement non simplement connexe à bord $\Gamma$ de classe $\mathcal{C}^{1,1}$. Le cadre hilbertien est déjà traité par $C$. Amrouche, C. Bernardi, M. Dauge et V. Girault [1]. L’originalité de notre travail est de développer des résultats similaires en théorie $L^{p}$ lorsque $1<p<\infty$. Le résultat de base concernant l'existence d'un potentiel vecteur sans conditions aux limites est donné dans le Théorème 3.1. En particulier dans les Théorèmes 3.2, 3.3 et 3.4 , plusieurs conditions aux limites sont proposées. Les autres résultats concernent la régularité de tels potentiels vecteurs. On s'intéresse ensuite au cas des potentiels scalaires et potentiels vecteurs très faibles.

## 1. Introduction

In this work, we assume that $\Omega$ is a bounded open connected set of $\mathbb{R}^{3}$ of class $\mathcal{C}^{1,1}$ with boundary $\Gamma$. Let $\Gamma_{i}, 0 \leqslant i \leqslant I$, denote the connected components of the boundary $\Gamma, \Gamma_{0}$ being the exterior boundary of $\Omega$. We do not assume that $\Omega$ is simply-connected but we suppose that there exist $J$ connected open surfaces $\Sigma_{j}, 1 \leqslant j \leqslant J$, called 'cuts', contained in

[^0]$\Omega$, such that each surface $\Sigma_{j}$ is an open subset of a smooth manifold. The boundary of each $\Sigma_{j}$ is contained in $\Gamma$. The intersection $\overline{\Sigma_{i}} \cap \overline{\Sigma_{j}}$ is empty for $i \neq j$, and finally the open set $\Omega^{\circ}=\Omega \backslash \bigcup_{j=1}^{J} \Sigma_{j}$ is simply-connected. We denote by $[\cdot]_{j}$ the jump of a function over $\Sigma_{j}$, for $1 \leqslant j \leqslant J$. The pair $\langle\cdot, \cdot\rangle_{X, X^{\prime}}$ denotes the duality product between the spaces $X$ and $X^{\prime}$. We then define the spaces:
\[

$$
\begin{aligned}
& \boldsymbol{H}^{p}(\operatorname{curl}, \Omega)=\left\{\boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega) ; \operatorname{curl} \boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega)\right\}, \quad \boldsymbol{H}^{p}(\operatorname{div}, \Omega)=\left\{\boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega) ; \operatorname{div} \boldsymbol{v} \in L^{p}(\Omega)\right\}, \\
& \boldsymbol{X}^{p}(\Omega)=\boldsymbol{H}^{p}(\operatorname{curl}, \Omega) \cap \boldsymbol{H}^{p}(\operatorname{div}, \Omega),
\end{aligned}
$$
\]

equipped with the graph norm. We also define their subspaces:

$$
\begin{aligned}
& \boldsymbol{H}_{0}^{p}(\text { curl, } \Omega)=\left\{\boldsymbol{v} \in \boldsymbol{H}^{p}(\text { curl, } \Omega) ; \boldsymbol{v} \times \boldsymbol{n}=\mathbf{0} \text { on } \Gamma\right\}, \\
& \boldsymbol{H}_{0}^{p}(\operatorname{div}, \Omega)=\left\{\boldsymbol{v} \in \boldsymbol{H}^{p}(\operatorname{div}, \Omega) ; \boldsymbol{v} \cdot \boldsymbol{n}=0 \text { on } \Gamma\right\}, \\
& \boldsymbol{X}_{N}^{p}(\Omega)=\left\{\boldsymbol{v} \in \boldsymbol{X}^{p}(\Omega) ; \boldsymbol{v} \times \boldsymbol{n}=\mathbf{0} \text { on } \Gamma\right\}, \quad \boldsymbol{X}_{T}^{p}(\Omega)=\left\{\boldsymbol{v} \in \boldsymbol{X}^{p}(\Omega) ; \boldsymbol{v} \cdot \boldsymbol{n}=0 \text { on } \Gamma\right\}
\end{aligned}
$$

and $\boldsymbol{X}_{0}^{p}(\Omega)=\boldsymbol{X}_{N}^{p}(\Omega) \cap \boldsymbol{X}_{T}^{p}(\Omega)$. We also define the space

$$
\boldsymbol{W}_{\sigma}^{1, p}(\Omega)=\left\{\boldsymbol{v} \in \boldsymbol{W}^{1, p}(\Omega), \operatorname{div} \boldsymbol{v}=0 \text { in } \Omega\right\}
$$

As in the Hilbertian case, we can prove that the space $\boldsymbol{X}_{0}^{p}(\Omega)$ coincides with $\boldsymbol{W}_{0}^{1, p}(\Omega)$ for $1<p<\infty$. We can also prove that $\mathcal{D}(\bar{\Omega})$ is dense in $\boldsymbol{H}^{p}(\mathbf{c u r l}, \Omega), \boldsymbol{H}^{p}(\operatorname{div}, \Omega)$ and $\boldsymbol{X}^{p}(\Omega)$. Also $\mathcal{D}(\Omega)$ is dense in $\boldsymbol{H}_{0}^{p}(\mathbf{c u r l}, \Omega)$ and $\boldsymbol{H}_{0}^{p}(\operatorname{div}, \Omega)$. For any function $q$ in $W^{1, p}\left(\Omega^{\circ}\right), \operatorname{grad} q$ can be extended to $\boldsymbol{L}^{p}(\Omega)$. We denote this extension by $\widetilde{\operatorname{grad} q}$. We finally define the spaces:

$$
\begin{aligned}
& \boldsymbol{K}_{T}^{p}(\Omega)=\left\{\boldsymbol{v} \in \boldsymbol{X}_{T}^{p}(\Omega) ; \text { curl } \boldsymbol{v}=\mathbf{0}, \operatorname{div} \boldsymbol{v}=0 \text { in } \Omega\right\}, \\
& \boldsymbol{K}_{N}^{p}(\Omega)=\left\{\boldsymbol{v} \in \boldsymbol{X}_{N}^{p}(\Omega) ; \text { curl } \boldsymbol{v}=\mathbf{0}, \operatorname{div} \boldsymbol{v}=0 \text { in } \Omega\right\} .
\end{aligned}
$$

As shown in [1], Proposition 3.14, for the case $p=2$, we can prove that the space $\boldsymbol{K}_{T}^{p}(\Omega)$ is of dimension $J$ and is spanned by the functions $\widetilde{\operatorname{grad}} q_{j}^{T}, 1 \leqslant j \leqslant J$, where each $q_{j}^{T} \in W^{1, p}\left(\Omega^{\circ}\right)$ is unique up to an additive constant and satisfies $\Delta q_{j}^{T}=0$ in $\Omega^{\circ}, \partial_{n} q_{j}^{T}=0$ on $\Gamma,\left[q_{j}^{T}\right]_{k}=$ constant, $\left[\partial_{n} q_{j}^{T}\right]_{k}=0,1 \leqslant k \leqslant J$, and $\left\langle\partial_{n} q_{j}^{T}, 1\right\rangle_{\Sigma_{k}}=\delta_{j k}, 1 \leqslant k \leqslant J$. We note that $\boldsymbol{K}_{T}^{p}(\Omega)=\{0\}$ if $J=0$, where $J$ is the second Betti number. Similarly, we can prove that the dimension of the space $\boldsymbol{K}_{N}^{p}(\Omega)$ is $I$ and that it is spanned by the functions grad $q_{i}^{N}, 1 \leqslant i \leqslant I$, where each $q_{i}^{N} \in W^{1, p}(\Omega)$ is the unique solution to the problem $\Delta q_{i}^{N}=0$ in $\Omega, q_{i}^{N}=0$ in $\Gamma_{0}, q_{i}^{N}=$ constant in $\Gamma_{k},\left\langle\partial_{n} q_{i}^{N}, 1\right\rangle_{\Gamma_{0}}=-1$ and $\left\langle\partial_{n} q_{i}^{N}, 1\right\rangle_{\Gamma_{k}}=\delta_{i k}, 1 \leqslant k \leqslant I$. We note that $I$ is the first Betti number and if $\Gamma=\Gamma_{0}$, then $\boldsymbol{K}_{N}^{p}(\Omega)=\{0\}$. In the sequel, the letter $C$ denotes a constant that is not necessarily the same at its various occurrences. The detailed proofs of the results announced in this Note are given in [3].

## 2. Sobolev's inequalities and compactness properties

We introduce the following two operators:

$$
T \lambda(\boldsymbol{x})=-\frac{1}{2 \pi} \int_{\Gamma} \lambda(\boldsymbol{\xi}) \frac{\partial}{\partial \boldsymbol{n}}|\boldsymbol{x}-\boldsymbol{\xi}|^{-1} \mathrm{~d} \sigma_{\xi}, \quad R \lambda(\boldsymbol{x})=\frac{1}{2 \pi} \int_{\Gamma} \operatorname{curl}\left(\frac{\lambda(\boldsymbol{\xi})}{|\boldsymbol{x}-\boldsymbol{\xi}|}\right) \times \boldsymbol{n} \mathrm{d} \sigma_{\xi},
$$

where $T$ is compact from $L^{p}(\Gamma)$ into $L^{p}(\Gamma)$ and $R$ is compact from $\boldsymbol{L}^{p}(\Gamma)$ into $\boldsymbol{L}^{p}(\Gamma)$ (see [6]). Using the Fredholm alternative, we can check that the null spaces $\operatorname{Ker}(I d+T)$ and $\operatorname{Ker}(I d+R)$ are of finite dimension and are respectively spanned by the traces of the functions $\operatorname{grad} q_{i}^{N} \cdot \boldsymbol{n}$ on $\Gamma$ for any $1 \leqslant i \leqslant I$ and the traces of the functions $\widetilde{\operatorname{grad}} q_{j}^{T} \times \boldsymbol{n}$ on $\Gamma$ for any $1 \leqslant j \leqslant J$. The next lemma is a generalization of the one in [6] to the case $I \geqslant 1$ and $J \geqslant 1$. We expect that to estimate $\nabla \boldsymbol{v}$, in addition to $\operatorname{div} \boldsymbol{v}$ and $\operatorname{curl} \boldsymbol{v}$, the quantity $\sum_{j=1}^{J}\langle\boldsymbol{v} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}$ is necessarily in the case where $\boldsymbol{v} \cdot \boldsymbol{n}$ vanishes on $\Gamma$ (respectively the quantity $\sum_{i=1}^{I}\langle\boldsymbol{v} \cdot \boldsymbol{n}, 1\rangle_{\Gamma_{i}}$ is necessarily in the case where $\boldsymbol{v} \times \boldsymbol{n}$ vanishes on $\Gamma$ ).

## Lemma 2.1.

(i) Any function $\boldsymbol{v} \in \boldsymbol{W}^{1, p}(\Omega) \cap \boldsymbol{X}_{N}^{p}(\Omega)$ satisfies:

$$
\|\nabla \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} \leqslant C\left(\|\operatorname{div} \boldsymbol{v}\|_{L^{p}(\Omega)}+\|\operatorname{curl} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)}+\sum_{i=1}^{I}\left|\langle\boldsymbol{v} \cdot \boldsymbol{n}, 1\rangle_{\Gamma_{i}}\right|\right) .
$$

(ii) Any function $\boldsymbol{v} \in \boldsymbol{W}^{1, p}(\Omega) \cap \boldsymbol{X}_{T}^{p}(\Omega)$ satisfies:

$$
\|\nabla \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)} \leqslant C\left(\|\operatorname{div} \boldsymbol{v}\|_{L^{p}(\Omega)}+\|\operatorname{curl} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)}+\sum_{j=1}^{J}\left|\langle\boldsymbol{v} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}\right|\right) .
$$

Sketch of the proof. We use the integral representations, the properties of the operators $T$ and $R$ and the CalderonZygmund inequalities.

Using the previous result, the density of $\boldsymbol{W}^{1, p}(\Omega) \cap \boldsymbol{X}_{N}^{p}(\Omega)$ in $\boldsymbol{X}_{N}^{p}(\Omega)$ and the density of $\boldsymbol{W}^{1, p}(\Omega) \cap \boldsymbol{X}_{T}^{p}(\Omega)$ in $\boldsymbol{X}_{T}^{p}(\Omega)$, we obtain the following continuous embeddings:

$$
\boldsymbol{X}_{N}^{p}(\Omega) \hookrightarrow \boldsymbol{W}^{1, p}(\Omega) \quad \text { and } \quad \boldsymbol{X}_{T}^{p}(\Omega) \hookrightarrow \boldsymbol{W}^{1, p}(\Omega)
$$

We now introduce the following spaces for $s \in \mathbb{R}, s \geqslant 1$ :

$$
\begin{aligned}
& \boldsymbol{X}^{s, p}(\Omega)=\left\{\boldsymbol{v} \in \mathbf{L}^{p}(\Omega) ; \operatorname{div} \boldsymbol{v} \in W^{s-1, p}(\Omega), \text { curl } \boldsymbol{v} \in \boldsymbol{W}^{s-1, p}(\Omega) \text { and } \boldsymbol{v} \cdot \boldsymbol{n} \in \boldsymbol{W}^{s-\frac{1}{p}, p}(\Gamma)\right\} \\
& \boldsymbol{Y}^{s, p}(\Omega)=\left\{\boldsymbol{v} \in \boldsymbol{L}^{p}(\Omega) ; \operatorname{div} \boldsymbol{v} \in W^{s-1, p}(\Omega), \operatorname{curl} \boldsymbol{v} \in \boldsymbol{W}^{s-1, p}(\Omega) \text { and } \boldsymbol{v} \times \boldsymbol{n} \in \boldsymbol{W}^{s-\frac{1}{p}, p}(\Omega)\right\} .
\end{aligned}
$$

The following result is an extension to the case where the boundary conditions $\boldsymbol{v} \cdot \boldsymbol{n}=0$ and $\boldsymbol{v} \times \boldsymbol{n}=0$ on $\Gamma$ are replaced by inhomogeneous ones:

Theorem 2.2. The spaces $\boldsymbol{X}^{1, p}(\Omega)$ and $\boldsymbol{Y}^{1, p}(\Omega)$ are both continuously imbedded in $\boldsymbol{W}^{1, p}(\Omega)$ :
(i) Any $\boldsymbol{v}$ in $\boldsymbol{X}^{1, p}(\Omega)$ satisfies

$$
\|\boldsymbol{v}\|_{\boldsymbol{W}^{1, p}(\Omega)} \leqslant C\left(\|\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{curl} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{div} \boldsymbol{v}\|_{L^{p}(\Omega)}+\|\boldsymbol{v} \cdot \boldsymbol{n}\|_{\boldsymbol{w}^{1-\frac{1}{p}, p}(\Gamma)}\right) .
$$

(ii) Any $\boldsymbol{v}$ in $\boldsymbol{Y}^{1, p}(\Omega)$ satisfies

$$
\|\boldsymbol{v}\|_{\boldsymbol{W}^{1, p}(\Omega)} \leqslant C\left(\|\boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{curl} \boldsymbol{v}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{div} \boldsymbol{v}\|_{L^{p}(\Omega)}+\|\boldsymbol{v} \times \boldsymbol{n}\|_{\boldsymbol{W}^{1-\frac{1}{p}, p}(\Gamma)}\right) .
$$

(iii) Let $m \in \mathbb{N}^{*}$ and $\Omega$ of class $\mathcal{C}^{m, 1}$. Then, the spaces $\boldsymbol{X}^{m, p}(\Omega)$ and $\boldsymbol{Y}^{m, p}(\Omega)$ are both continuously imbedded in $\boldsymbol{W}^{m, p}(\Omega)$.
(iv) Let $s=m+\sigma, m \in \mathbb{N}^{*}$ and $0<\sigma \leqslant 1$. Assume that $\Omega$ is of class $\mathcal{C}^{m+1,1}$. Then, the spaces $\boldsymbol{X}^{s, p}(\Omega)$ and $\boldsymbol{Y}^{s, p}(\Omega)$ are both continuously imbedded in $\boldsymbol{W}^{s, p}(\Omega)$.

The first point of the following result is proven by Costabel [4] for the case $p=2$ in a bounded simply connected domain. We extend this result to the case of a multiply connected domain $\Omega$ and for any $1<p<\infty$.

Theorem 2.3. Let $\boldsymbol{u} \in \boldsymbol{X}^{p}(\Omega)$ with $\boldsymbol{u} \cdot \boldsymbol{n} \in L^{p}(\Gamma)$ (respectively with $\boldsymbol{u} \times \boldsymbol{n} \in \boldsymbol{L}^{p}(\Gamma)$ ). Then, $\boldsymbol{u} \in \boldsymbol{W}^{\frac{1}{p}, p}(\Omega)$ and satisfies the inequality

$$
\|\boldsymbol{u}\|_{\boldsymbol{w}^{\frac{1}{p}, p}(\Omega)} \leqslant C\left(\|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{curl} \boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{div} \boldsymbol{u}\|_{L^{p}(\Omega)}+\|\boldsymbol{u} \cdot \boldsymbol{n}\|_{L^{p}(\Gamma)}\right)
$$

(respectively $\|\boldsymbol{u}\|_{\boldsymbol{W}^{\frac{1}{p}, p}(\Omega)} \leqslant C\left(\|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{curl} \boldsymbol{u}\|_{\mathbf{L}^{p}(\Omega)}+\|\operatorname{div} \boldsymbol{u}\|_{L^{p}(\Omega)}+\|\boldsymbol{u} \times \boldsymbol{n}\|_{\boldsymbol{L}^{p}(\Gamma)}\right)$ ). If in addition $\boldsymbol{u} \cdot \boldsymbol{n} \in W^{s-1 / p, p}(\Gamma)$ (respectively $\left.\boldsymbol{u} \times \boldsymbol{n} \in \boldsymbol{W}^{s-1 / p, p}(\Gamma)\right)$ with $1 / p<s \leqslant 1$, then $\boldsymbol{u} \in \boldsymbol{W}^{s, p}(\Omega)$.

As for the case $p=2$ (see [1]), we can prove that the imbedding of $\boldsymbol{X}^{p}(\Omega)$ into $\boldsymbol{L}^{p}(\Omega)$ is not compact and that the homogeneous normal or tangential boundary conditions are sufficient to insure compactness. More precisely we have the following result:

Theorem 2.4. The imbeddings of $\boldsymbol{X}_{N}^{p}(\Omega)$ and $\boldsymbol{X}_{T}^{p}(\Omega)$ into $\boldsymbol{L}^{p}(\Omega)$ are compact.

## 3. Vector potentials

We begin with a first result concerning vector potentials without boundary conditions. The result is known for $p=2$ (see [1]) and we can give a different proof for $1<p<\infty$ by using the fundamental solution of the Laplacian.

Theorem 3.1. A vector field $\boldsymbol{u}$ in $\boldsymbol{H}^{p}(\operatorname{div}, \Omega)$ satisfies

$$
\begin{equation*}
\operatorname{div} \boldsymbol{u}=0 \quad \text { in } \Omega \quad \text { and } \quad\langle\boldsymbol{u} \cdot \boldsymbol{n}, 1\rangle_{\Gamma_{i}}=0, \quad 0 \leqslant i \leqslant I, \tag{1}
\end{equation*}
$$

if and only if there exists a vector potential $\boldsymbol{\psi}_{0}$ in $\boldsymbol{W}^{1, p}(\Omega)$ such that

$$
\boldsymbol{u}=\operatorname{curl} \psi_{0}
$$

Moreover, we can choose $\psi_{0}$ such that $\operatorname{div} \psi_{0}=0$ and we have the estimate

$$
\left\|\boldsymbol{\psi}_{0}\right\|_{\boldsymbol{W}^{1, p}(\Omega)} \leqslant C\|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)},
$$

where $C>0$ depends only on $p$ and $\Omega$.
Theorem 3.2. A function $\boldsymbol{u}$ in $\boldsymbol{H}^{p}(\operatorname{div}, \Omega)$ satisfies (1) if and only if there exists a vector potential $\boldsymbol{\psi}$ in $\boldsymbol{W}^{1, p}(\Omega)$ such that

$$
\begin{equation*}
\boldsymbol{u}=\operatorname{curl} \psi \quad \text { and } \quad \operatorname{div} \psi=0 \quad \text { in } \Omega, \quad \psi \cdot \boldsymbol{n}=0 \quad \text { on } \Gamma, \quad\langle\boldsymbol{\psi} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}=0, \quad 1 \leqslant j \leqslant J . \tag{2}
\end{equation*}
$$

This function $\psi$ is unique and we have the estimate:

$$
\begin{equation*}
\|\boldsymbol{\psi}\|_{\boldsymbol{W}^{1, p}(\Omega)} \leqslant C\|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)} \tag{3}
\end{equation*}
$$

Sketch of the proof. Let $\psi_{0} \in \boldsymbol{W}^{1, p}(\Omega)$ be the function associated with $\boldsymbol{u}$ by Theorem 3.1 and $\chi \in W^{1, p}(\Omega)$ be a unique solution up to an additive constant, of the problem: $-\Delta \chi=0$ in $\Omega$ and $\partial_{n} \chi=\boldsymbol{\psi}_{0} \cdot \boldsymbol{n}$ on $\Gamma$. Then, the vector function

$$
\boldsymbol{\psi}=\psi_{0}-\operatorname{grad} \chi-\sum_{j=1}^{J}\left\langle\left(\psi_{0}-\operatorname{grad} \chi\right) \cdot \boldsymbol{n}, 1\right\rangle_{\Sigma_{j}} \widetilde{\operatorname{grad}} q_{j}^{T}
$$

satisfies the properties in (2) and the estimate (3).
Applying the Peetre-Tartar Lemma (cf. Ref. [5, Chapter I, Theorem 2.1]), we can prove the following Poincaré-type inequality:

Corollary 3.3. On the space $\boldsymbol{X}_{T}^{p}(\Omega)$, the seminorm

$$
\boldsymbol{w} \mapsto\|\operatorname{curl} \boldsymbol{w}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{div} \boldsymbol{w}\|_{L^{p}(\Omega)}+\sum_{j=1}^{J}\left|\langle\boldsymbol{w} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}\right|
$$

is equivalent to the norm $\|\cdot\|_{\boldsymbol{X}^{p}(\Omega)}$.
Theorem 3.4. A function $\boldsymbol{u}$ in $\boldsymbol{H}^{p}(\operatorname{div}, \Omega)$ satisfies:

$$
\begin{equation*}
\operatorname{div} \boldsymbol{u}=0 \quad \text { in } \Omega, \quad \boldsymbol{u} \cdot \boldsymbol{n}=0 \quad \text { on } \Gamma \quad \text { and } \quad\langle\boldsymbol{u} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}=0, \quad 1 \leqslant j \leqslant J \tag{4}
\end{equation*}
$$

if and only if there exists a vector potential $\boldsymbol{\psi}$ in $\boldsymbol{W}^{1, p}(\Omega)$ such that

$$
\begin{equation*}
\boldsymbol{u}=\operatorname{curl} \psi \quad \text { and } \quad \operatorname{div} \psi=0 \quad \text { in } \Omega, \quad \psi \times \boldsymbol{n}=\mathbf{0} \quad \text { on } \Gamma, \quad\langle\boldsymbol{\psi} \cdot \boldsymbol{n}, 1\rangle_{\Gamma_{i}}=0, \quad \text { for any } 0 \leqslant i \leqslant I . \tag{5}
\end{equation*}
$$

This function $\psi$ is unique and we have the estimate:

$$
\begin{equation*}
\|\boldsymbol{\psi}\|_{\boldsymbol{W}^{1, p}(\Omega)} \leqslant C\|\boldsymbol{u}\|_{\boldsymbol{L}^{p}(\Omega)} \tag{6}
\end{equation*}
$$

Sketch of the proof. We solve the problem

$$
-\Delta \boldsymbol{\xi}=0 \quad \text { and } \quad \operatorname{div} \boldsymbol{\xi}=0 \quad \text { in } \Omega, \quad \boldsymbol{\xi} \cdot \boldsymbol{n}=0, \quad \operatorname{curl} \boldsymbol{\xi} \times \boldsymbol{n}=\boldsymbol{\psi}_{0} \times \boldsymbol{n} \quad \text { on } \Gamma, \quad\langle\boldsymbol{\xi} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}=0, \quad 1 \leqslant j \leqslant J
$$

where $\boldsymbol{\psi}_{0}$ is the function associated with $\boldsymbol{u}$ by Theorem 3.1. This problem is equivalent to the following: find $\boldsymbol{\xi} \in \boldsymbol{V}_{T}^{p}(\Omega)$ such that

$$
\begin{equation*}
\forall \varphi \in \boldsymbol{V}_{T}^{p^{\prime}}(\Omega), \quad \int_{\Omega} \operatorname{curl} \xi \cdot \operatorname{curl} \varphi \mathrm{d} \boldsymbol{x}=\int_{\Omega} \psi_{0} \cdot \operatorname{curl} \varphi \mathrm{~d} \boldsymbol{x}-\int_{\Omega} \varphi \cdot \operatorname{curl} \psi_{0} \mathrm{~d} \boldsymbol{x} \tag{7}
\end{equation*}
$$

where

$$
\boldsymbol{V}_{T}^{p}(\Omega)=\left\{\boldsymbol{w} \in \boldsymbol{X}_{T}^{p}(\Omega) ; \operatorname{div} \boldsymbol{w}=0 \text { in } \Omega \text { and }\langle\boldsymbol{w} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}=0,1 \leqslant j \leqslant J\right\}
$$

We can prove that (7) satisfies the following Inf-Sup condition:

$$
\begin{equation*}
\inf _{\boldsymbol{\varphi} \in \boldsymbol{V}_{T}^{p^{\prime}}(\Omega)} \sup _{\boldsymbol{\xi} \in \boldsymbol{V}_{T}^{p}(\Omega)} \frac{\int_{\Omega} \operatorname{curl} \boldsymbol{\xi} \cdot \operatorname{curl} \boldsymbol{\varphi} \mathrm{d} \boldsymbol{x}}{\|\boldsymbol{\xi}\|_{\boldsymbol{X}_{T}^{p}(\Omega)}\|\boldsymbol{\varphi}\|_{\boldsymbol{X}_{T}^{p^{\prime}}(\Omega)}}>0 \tag{8}
\end{equation*}
$$

and that the solution $\boldsymbol{\xi}$ belongs to $\boldsymbol{W}^{2, p}(\Omega)$. Then, the vector function

$$
\boldsymbol{\psi}=\psi_{0}-\operatorname{curl} \xi-\sum_{i=1}^{I}\left\langle\left(\psi_{0}-\operatorname{curl} \xi\right) \cdot \boldsymbol{n}, 1\right\rangle_{\Gamma_{i}} \operatorname{grad} q_{i}^{N},
$$

satisfies (5) and the estimate (6).

Again, using the Peetre-Tartar Lemma, we have the following Poincaré-type inequality:
Corollary 3.5. On the space $\boldsymbol{X}_{N}^{p}(\Omega)$, the seminorm

$$
\boldsymbol{w} \mapsto\|\operatorname{curl} \boldsymbol{w}\|_{\boldsymbol{L}^{p}(\Omega)}+\|\operatorname{div} \boldsymbol{w}\|_{L^{p}(\Omega)}+\sum_{i=1}^{I}\left|\langle\boldsymbol{w} \cdot \boldsymbol{n}, 1\rangle_{\Gamma_{i}}\right|
$$

is equivalent to the norm $\|\cdot\|_{\boldsymbol{X}^{p}(\Omega)}$.

In a similar fashion and by characterizing the kernel space

$$
\boldsymbol{K}_{0}^{p}(\Omega)=\left\{\boldsymbol{w} \in \boldsymbol{W}_{0}^{1, p}(\Omega) ; \text { curl } \boldsymbol{w}=\mathbf{0} \text { and } \operatorname{div}(\Delta \boldsymbol{w})=0 \text { in } \Omega\right\}
$$

we can give another type of vector potential:

Theorem 3.6. A function $\boldsymbol{u}$ in $\boldsymbol{H}^{p}(\operatorname{div}, \Omega)$ satisfies:

$$
\operatorname{div} \boldsymbol{u}=0 \quad \text { in } \Omega, \quad \boldsymbol{u} \cdot \boldsymbol{n}=0 \quad \text { on } \Gamma \quad \text { and } \quad\langle\boldsymbol{u} \cdot \boldsymbol{n}, 1\rangle_{\Sigma_{j}}=0, \quad 1 \leqslant j \leqslant J
$$

if and only if there exists a vector potential $\boldsymbol{\psi}$ in $\boldsymbol{W}^{1, p}(\Omega)$ such that

$$
\boldsymbol{u}=\operatorname{curl} \boldsymbol{\psi} \quad \text { and } \quad \operatorname{div}(\Delta \psi)=0 \quad \text { in } \Omega, \quad \boldsymbol{\psi}=0 \quad \text { on } \Gamma, \quad\left\langle\partial_{n}(\operatorname{div} \boldsymbol{\psi}), 1\right\rangle_{\Gamma_{i}}=0, \quad \text { for any } 0 \leqslant i \leqslant I
$$

This function $\boldsymbol{\psi}$ is unique.

## 4. Scalar potentials

In this section we present several results concerning scalar potentials. We begin with the following fundamental result:

Theorem 4.1. For any function $\boldsymbol{f} \in \boldsymbol{L}^{p}(\Omega)$ that satisfies

$$
\operatorname{curl} \boldsymbol{f}=\mathbf{0} \quad \text { and } \quad \int_{\Omega} \boldsymbol{f} \cdot \boldsymbol{v} \mathrm{d} \boldsymbol{x}=0 \quad \text { for all } \boldsymbol{v} \in \boldsymbol{K}_{T}^{p^{\prime}}(\Omega)
$$

there exists a unique scalar potential $\chi \in W^{1, p}(\Omega) / \mathbb{R}$ such that $\boldsymbol{f}=\operatorname{grad} \chi$ and the following estimate holds:

$$
\|\chi\|_{W^{1, p}(\Omega) / \mathbb{R}} \leqslant C\|\boldsymbol{f}\|_{\mathbf{L}^{p}(\Omega)}
$$

Sketch of the proof. It suffices to prove that for any $\boldsymbol{v} \in \boldsymbol{H}_{0}^{p^{\prime}}(\operatorname{div}, \Omega)$ such that $\operatorname{div} \boldsymbol{v}=0$ in $\Omega$, we have $\int_{\Omega} \boldsymbol{f} \cdot \boldsymbol{v} \mathrm{d} \boldsymbol{x}=0$. Next, we apply [2, Lemma 2.7].

We are now interested in the case of singular data.

Theorem 4.2. For any $\boldsymbol{f}$ in the dual space of $\boldsymbol{H}_{0}^{p^{\prime}}(\operatorname{div}, \Omega)$ with curl $\boldsymbol{f}=0$ in $\Omega$ and $\boldsymbol{f}$ satisfying:

$$
\forall \boldsymbol{v} \in \boldsymbol{K}_{T}^{p^{\prime}}(\Omega), \quad\langle\boldsymbol{f}, \boldsymbol{v}\rangle_{\left[\boldsymbol{H}_{0}^{p^{\prime}}(\operatorname{div}, \Omega)\right]^{\prime} \times \boldsymbol{H}_{0}^{p^{\prime}}(\operatorname{div}, \Omega)}=0,
$$

there exists a unique scalar potential $\chi$ in $L^{p}(\Omega) / \mathbb{R}$ such that $\boldsymbol{f}=\mathbf{g r a d} \chi$ and the following estimate holds:

$$
\|\chi\|_{L^{p}(\Omega) / \mathbb{R}} \leqslant C\|\boldsymbol{f}\|_{\boldsymbol{H}_{0}^{p^{\prime}}(\operatorname{div}, \Omega)^{\prime}} .
$$

Sketch of the proof. We can prove that for any $\boldsymbol{f}$ in $\left[\boldsymbol{H}_{0}^{p^{\prime}}(\operatorname{div}, \Omega)\right]^{\prime}$, there exists $\boldsymbol{\psi} \in \boldsymbol{L}^{p}(\Omega)$ and $\chi_{0} \in L^{p}(\Omega)$ such that $\boldsymbol{f}=\boldsymbol{\psi}+\boldsymbol{g r a d} \chi_{0}$. We then apply Theorem 4.1 to $\boldsymbol{\psi}$.

## 5. Weak vector potentials

The aim of this section is the proof of the existence of a new type of vector potential called weak vector potentials.

Theorem 5.1. For any $\boldsymbol{f}$ in the dual space of $\boldsymbol{H}_{0}^{p^{\prime}}(\operatorname{curl}, \Omega)$ with div $\boldsymbol{f}=0$ in $\Omega$ and $\boldsymbol{f}$ satisfying:

$$
\forall \boldsymbol{v} \in \boldsymbol{K}_{N}^{p^{\prime}}(\Omega), \quad\langle\boldsymbol{f}, \boldsymbol{v}\rangle_{\left[\boldsymbol{H}_{0}^{p^{\prime}}(\mathbf{c u r l}, \Omega)\right]^{\prime} \times \boldsymbol{H}_{0}^{p^{\prime}}(\mathbf{c u r l}, \Omega)}=0,
$$

there exists a vector potential $\boldsymbol{\xi}$ in $\mathbf{L}^{p}(\Omega)$ such that

$$
\boldsymbol{f}=\operatorname{curl} \boldsymbol{\xi}, \quad \text { with } \operatorname{div} \boldsymbol{\xi}=0 \quad \text { in } \Omega \quad \text { and } \quad \boldsymbol{\xi} \cdot \boldsymbol{n}=0 \quad \text { on } \Gamma,
$$

and such that the following estimate holds:

$$
\|\boldsymbol{\xi}\|_{\boldsymbol{L}^{p}(\Omega)} \leqslant C\|\boldsymbol{f}\|_{\boldsymbol{H}_{0}^{p^{\prime}}(\mathbf{c u r l}, \Omega)^{\prime}} .
$$

Sketch of the proof. We can prove that for any $\boldsymbol{f}$ in $\left[\boldsymbol{H}_{0}^{p^{\prime}}(\mathbf{c u r l}, \Omega)\right]^{\prime}$, there exists $\boldsymbol{\psi} \in \boldsymbol{L}^{p}(\Omega)$ and $\boldsymbol{\xi}_{0} \in \boldsymbol{L}^{p}(\Omega)$ such that $\boldsymbol{f}=\boldsymbol{\psi}+\operatorname{curl} \xi_{0}$ with $\operatorname{div} \boldsymbol{\xi}_{0}=0$ in $\Omega$ and $\xi_{0} \cdot \boldsymbol{n}=0$ on $\Gamma$. We then apply Theorem 3.2 to $\psi$.

## References

[1] C. Amrouche, C. Bernardi, M. Dauge, V. Girault, Vector potentials in three-dimensional nonsmooth domains, Math. Methods Appl. Sci. 21 (1998) $823-864$.
[2] C. Amrouche, V. Girault, Decomposition of vector spaces and application to the Stokes problem in arbitrary dimension, Czechoslovak Math. J. 119 (44) (1994) 109-140.
[3] C. Amrouche, N. Seloula, $L^{p}$-theory for vector potentials and Sobolev's inequalities for vector fields. Application to the Stokes problem's with pressure boundary conditions, submitted for publication.
[4] M. Costabel, A remark on the regularity of solutions of Maxwell's equations on Lipschitz domains, Math. Methods Appl. Sci. Theory 12 (1990) $365-368$.
[5] V. Girault, P.A. Raviart, Finite Element Methods for the Navier-Stokes Equations, Theory and Algorithms, Springer, Berlin, 1986.
[6] W. von Wahl, Estimating $\nabla u$ by $\operatorname{div} u$, curl $u$, Math. Methods Appl. Sci. 15 (1992) 123-143.


[^0]:    E-mail addresses: cherif.amrouche@univ-pau.fr (C. Amrouche), nourelhouda.seloula@etud.univ-pau.fr (N. Seloula).
    1631-073X/\$ - see front matter © 2011 Académie des sciences. Published by Elsevier Masson SAS. All rights reserved.
    doi:10.1016/j.crma.2011.04.008

