Global geometry of $T^2$-symmetric spacetimes with weak regularity

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1. Introduction

We are interested in spacetimes satisfying the Einstein vacuum equations and we provide a description of the global geometry of the development of an initial data set, having only weak regularity but enjoying certain symmetry properties. We consider a manifold that is diffeomorphic to the torus $T^3$ and initial data that are invariant under an action of the Lie group $T^2$. This requirement characterizes the so-called $T^2$-symmetric spacetimes on $T^3$. A large literature is available on such spacetimes when sufficiently high regularity is imposed; cf. Moncrief [13], Chruściel [4], Berger, Chruściel, Isenberg, and Moncrief [1], and Isenberg and Weaver [6] and, for further references, [14–16].

The present Note is motivated by recent work by LeFloch and Stewart [11,12] and LeFloch and Rendall [9], who treated Gowdy-symmetric matter spacetimes described by the Einstein–Euler equations. It was recognized that, due to the formation of shock waves in the fluid and by virtue of the Einstein equations, only weak regularity on the geometry should be allowed.
Here, we carry out this strategy further, and encompass (vacuum) spacetimes under less restrictive regularity and symmetry conditions. Importantly, we present a fully geometric framework within a class of weakly regular spacetimes and precisely identify the weak regularity conditions allowing us to, both, establish the existence of such spacetimes and tackle the analysis of their global geometric structure.

Even in the smooth class, there remain challenging open problems for these spacetimes, such as the question of geodesic completeness in the expanding direction or the structure of singularities in the contracting direction, and we believe that the new framework and estimates developed in the present work will be useful to tackle some of these issues.

The initial value problem in general relativity (in the vacuum case) is formulated as follows. An initial data set for the vacuum Einstein equations is a triple \((\Sigma, h, K)\) such that \((\Sigma, h)\) is a Riemannian 3-manifold and \(K\) is a symmetric 2-tensor field on \(\Sigma\), satisfying the so-called Einstein’s constraint equations. A solution to the initial value problem with initial data \((\Sigma, h, K)\) is a \((3+1)\)-dimensional Lorentzian manifold \((M, g)\) satisfying the vacuum Einstein equations (that is, the Ricci flat condition \(\text{Ric} = 0\)), together with an embedding \(\phi : \Sigma \rightarrow M\) such that \(\phi(\Sigma)\) is a Cauchy hypersurface of \((M, g)\) and the pull-back of its first and second fundamental forms coincides with \(h\) and \(K\), respectively.

Recall that the existence of a unique (up to diffeomorphism) maximal globally hyperbolic development \((M, g)\) was established in pioneering work by Choquet-Bruhat and Geroch [5,3] (reviewed in [2] together with extensive recent material). The optimal existence result as far as the regularity of the initial data is concerned is currently given in Klainerman and Rodnianski [7] in the Sobolev space \(H^{2+\epsilon}\) (with \(\epsilon > 0\)).

2. Statement of the main result

By restricting attention to the class of \(T^2\)-symmetric spacetimes, our regularity assumptions will go far below the regularity covered in previous works. We assume here that the area \(R\) of the orbits of \(T^2\)-symmetry is a \(C^1\)-function. Introducing a frame of commuting vector fields \((T, \Theta, X, Y)\) where \(X, Y\) are Killing fields, the remaining components of the metric belong to the Sobolev space \(H^1\) on every spacelike slice of a foliation, except for the cross terms between the Killing fields \(X, Y\) and the vector field \(\Theta\) which have lower regularity. (See Section 3, below.) Importantly, not all Christoffel symbols and curvature components make sense as distributions, but only those that are relevant to formulate the Einstein equations in a weak sense. To do otherwise, additional regularity would be necessary, as recognized in [8].

Under these regularity conditions, our main result establishes the existence of a weakly regular development, and provides detailed information about its global geometry. We refer to the next section and [10] for the precise terminology used in the following theorem:

**Theorem 2.1.**

- If \((\Sigma, h, K)\) is a weakly regular \(T^2\)-symmetric triple, then the Einstein constraint equations (suitably reformulated) make sense as equality between distributions. Similarly, if \((M, g)\) is a weakly regular \(T^2\)-symmetric Lorentzian manifold, then the Einstein equations (suitably reformulated) make sense as equality between distributions.
- For all weakly regular \(T^2\)-symmetric initial data set \((\Sigma, h, K)\) with topology \(T^3\) and constant area function \(R\), there exists a weakly regular vacuum spacetime with \(T^2\)-symmetry on \(T^3\), \((M, g)\), which is a development of \((\Sigma, h, K)\) and admits a global foliation by the level sets of the area \(R \in (R_\ast, \infty)\) (for some \(R_\ast > 0\)). Moreover, unless \((M, g)\) is a flat Kasner spacetime, one has \(R_\ast = 0\).

The compactness arguments developed in [10] for the existence proof are very different from the classical techniques used in the high regularity case [13,1]. In the rest of this Note, we discuss in further detail the class of initial data and present our new geometric formulation which is an essential contribution of [10].

3. Weakly regular \(T^2\)-symmetric manifolds

While all topological spacetimes under consideration are of class \(C^\infty\), the metric structures of interest here have low regularity. For instance, the Lie derivative \(\mathcal{L}_Z h\) of a measurable and locally integrable 2-tensor \(h\) is defined in the distribution sense, for any \(C^1\)-vector fields \(X, T, Z\), by

\[
(\mathcal{L}_Z h)(T, Z) := X(h(T, Z)) - h(\mathcal{L}_X T, Z) - h(T, \mathcal{L}_X Z),
\]

in which the last two terms are classically defined as locally integrable functions, but the first one is defined in the distributional sense, only.

**Definition 3.1.** A weakly regular \(T^2\)-symmetric Riemannian manifold is a compact \(C^\infty\)-manifold \(\Sigma\) endowed with a tensor field \(h\) which enjoys the following regularity and symmetry properties:

(i) **Riemannian structure.** The field \(h\) is a Riemannian metric with components in \(L^\infty(\Sigma)\), and the associated volume form is continuous.
(ii) **Symmetry property.** On the Riemannian manifold \((\Sigma, h)\), there exists an action of the Lie group \(T^2\) generated by two (linearly independent, commuting) Killing fields \(X,Y\) of class \(C^\infty\) (with closed orbits), satisfying in particular \(\mathcal{L}_X h = 0\) and \(\mathcal{L}_Y h = 0\) understood in the distributional sense (1).

(iii) **Regularity of the metric on the orbits.** The functions \(h(X,X), h(X,Y), \text{ and } h(Y,Y)\) belong to the Sobolev space \(H^1(\Sigma; h)\).

(iv) **Regularity of the area function.** The orbit area function

\[
R^2 := h(X,X) h(Y,Y) - h(X,Y)^2
\]

(which then lies in \(W^{1,1}(\Sigma; h)\)) is required to belong to \(C^1(\Sigma)\).

(v) **Regularity of the metric on the orthogonal of the orbits.** There exists a smooth vector field \(\Theta\) defined on \(\Sigma\) such that \((\Theta, X, Y)\) form a basis of commuting vector fields and the vector field

\[
Z := \Theta + a X + b Y, \quad Z \in \{X,Y\}^\perp, \quad \mathcal{L}_X Z = \mathcal{L}_Y Z = 0
\]

implies the regularity \(h(Z,Z) \in H^1(\Sigma; h)\).

We refer to such a triple \((X,Y,Z)\) as an adapted frame on \(\Sigma\). We emphasize that no regularity is required on the derivatives of the “cross-terms” \(h(X,\Theta)\) and \(h(Y,\Theta)\). We also need to consider the second fundamental form of the initial slice.

**Definition 3.2.** A weakly regular \(T^2\)-symmetric triple \((\Sigma, h, K)\) is a weakly regular \(T^2\)-symmetric Riemannian manifold \((\Sigma, h)\) with adapted frame \((X,Y,Z)\), satisfying the following conditions:

(i) **Integrability.** The field \(K\) is a symmetric 2-tensor field on \(\Sigma\) such that its components in the basis \((X,Y,Z)\) lie in \(L^2(\Sigma)\).

(ii) **Symmetry property.** Moreover, \(K\) is invariant under the action of the Lie group \(T^2\) generated by \((X,Y)\), that is, \(\mathcal{L}_X K = \mathcal{L}_Y K = 0\) understood in the sense of distributions (1).

(iii) **Additional regularity** (associated with the time derivative of the area function):

\[
\text{Tr}^{(2)}(K) := h(X,X) K(X,X) + 2h(X,Y) K(X,Y) + h(Y,Y) K(Y,Y) \in C(\Sigma).
\]

Our notion of initial data set follows the above two definitions, as we define the curvature of the manifold and impose the constraint equations in the distribution sense. Such a statement requires to take into account our symmetry assumption in order for the terms to make sense as distributions.

**Acknowledgements**

The first author (P.L.F.) was partially supported by the Agence Nationale de la Recherche (ANR) through the grant 06-2-134423 entitled “Mathematical Methods in General Relativity” (Math-GR). The second author (J.S.) thanks the Albert Einstein Institute for financial support. This work was done when the second author (J.S.) was visiting the Laboratoire Jacques-Louis Lions with the support of the ANR.

**References**


