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# Probability Theory/Functional Analysis

# The invertibility of adapted perturbations of identity on the Wiener space

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#### Abstract

Let  $(W, H, \mu)$  be the classical Wiener space. Assume that  $U = I_W + \mu$  is an adapted perturbation of identity, i.e.,  $u: W \to H$  is adapted to the canonical filtration of W. We give some sufficient analytic conditions on u which imply the invertibility of the map U. To cite this article: A.S. Üstünel, M. Zakai, C. R. Acad. Sci. Paris, Ser. I 342 (2006). © 2006 Académie des sciences. Published by Elsevier SAS. All rights reserved.

#### Résumé

L'inversibilité des perturbations d'identité adaptées sur l'espace de Wiener. Soit  $(W, H, \mu)$  l'espace de Wiener. Soit  $U = I_W + u$  une perturbation d'identité adaptée, i.e.,  $u: W \to H$  est adaptée à la filtration canonique de W. Nous donnons quelques conditions suffisantes qui impliquent l'inversibilité de l'application U. Pour citer cet article : A.S. Üstünel, M. Zakai, C. R. Acad. Sci. Paris, Ser. I 342 (2006).

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## 1. Preliminaries

Let  $W = C_0([0, 1])$  be the Banach space of continuous functions on [0, 1], with its Borel sigma field denoted by  $\mathcal{F}$ . We denote by H the Cameron–Martin space, namely the space of absolutely continuous functions on [0, 1] with square integrable Lebesgue density:

$$H = \left\{ h \in W: \ h(t) = \int_{0}^{t} \dot{h}(s) \, \mathrm{d}s, \ |h|_{H}^{2} = \int_{0}^{1} \left| \dot{h}(s) \right|^{2} \, \mathrm{d}s < \infty \right\}.$$

 $\mu$  denotes the classical Wiener measure on  $(W, \mathcal{F})$ ,  $(\mathcal{F}_t, t \in [0, 1])$  is the filtration generated by the paths of the Wiener process  $(t, w) \rightarrow W_t(w)$ , where  $W_t(w)$  is defined as w(t) for  $w \in W$  and  $t \in [0, 1]$ . We shall recall briefly

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some well-known functional analytic tools on the Wiener space, we refer the reader to [4,3,5] or to [6] for further details:  $(P_{\tau}, \tau \in \mathbb{R}_+)$  denotes the semi-group of Ornstein–Uhlenbeck on *W*, defined as

$$P_{\tau} f(w) = \int_{W} f(e^{-\tau} w + \sqrt{1 - e^{-2\tau}} y) \mu(dy).$$

Let us recall that  $P_{\tau} = e^{-\tau \mathcal{L}}$ , where  $\mathcal{L}$  is the number operator. We denote by  $\nabla$  the Sobolev derivative which is the extension (with respect to the Wiener measure) of the Fréchet derivative in the Cameron–Martin space direction. The iterates of  $\nabla$  are defined similarly. Note that, if f is real valued, then  $\nabla f$  is a vector and if u is an H-valued map, then  $\nabla u$  is a Hilbert–Schmidt (on H) operator valued map whenever defined. If Z is a separable Hilbert space and if  $p > 1, k \in \mathbb{R}$ , we denote by  $\mathbb{D}_{p,k}(Z)$  the  $\mu$ -equivalence classes of Z-valued measurable mappings  $\xi$ , defined on W such that  $(I + \mathcal{L})^{k/2}\xi$  belongs to  $L^p(\mu, Z)$  and this set, equipped with the norm

$$\|\xi\|_{p,k} = \left\| (I+\mathcal{L})^{k/2} \xi \right\|_{L^p(\mu,Z)} \tag{1}$$

becomes a Banach space. From the Meyer inequalities, we know that the norm defined by

$$\sum_{k=0}^{n} \left\| \nabla^{k} \xi \right\|_{L^{p}(\mu, Z \otimes H^{\otimes k})}, \quad n \in \mathbb{N},$$

is equivalent to the norm  $\|\xi\|_{p,n}$  defined by (1). We denote by  $\delta$  the adjoint of  $\nabla$  under  $\mu$  and recall that, whenever  $u \in \mathbb{D}_{p,0}(H)$  for some p > 1 is adapted, then  $\delta u$  is equal to the Itô integral of the Lebesgue density of u:

$$\delta u = \int_0^1 \dot{u}_s \, \mathrm{d} W_s.$$

#### 2. A sufficient condition for invertibility

Assume that  $u: W \to H$  is adapted, i.e.,  $u(t) = \int_0^t \dot{u}_s \, ds, t \in [0, 1]$  and that  $\dot{u}_s$  is  $\mathcal{F}_s$ -measurable for almost all  $s \in [0, 1]$ . We suppose that  $\rho(-\delta u)$  defined as

$$\rho(-\delta u) = \exp\left[-\delta u - \frac{1}{2}|u|_{H}^{2}\right]$$

is the terminal value of a uniformly integrable martingale. We shall assume that u is in  $\mathbb{D}_{2,0}(H)$ . We have

**Theorem 1.** Assume that u satisfies the hypothesis above. For  $\tau \in [0, 1]$ , define  $u_{\tau}$  as to be  $P_{\tau}u$ , where  $P_{\tau}$  is the Ornstein–Uhlenbeck semigroup and assume also that  $E[\rho(-\delta u_{\tau})] = 1$  for  $\tau \in [0, 1]$ . Then the adapted perturbation of identity  $U = I_W + u$  is invertible provided that

$$E\left[\int_{0}^{1} \left| (I_H + \nabla u_\tau)^{-1} \mathcal{L} u_\tau \right|_H \rho(-\delta u_\tau) \,\mathrm{d}\tau \right] < \infty.$$
<sup>(2)</sup>

**Proof.** Note that the map  $u_{\tau}$  is again adapted and  $H - C^1$  (in fact it is even  $H - C^{\infty}$ , cf. [7]). This means that there exists a negligible set  $N \subset W$  (in fact its capacity is null [6]) with  $H + N \subset N$ , such that, for any  $w \in N^c$ , the map  $h \to u_{\tau}(w + h)$  is continuously Fréchet differentiable on H. Consequently  $U_{\tau} = I_W + u_{\tau}$  satisfies the change of variables formula: for any  $f \in C_b(W)$ ,

$$E[f \circ U_{\tau} \rho(-\delta u_{\tau})] = E[f(w)N_{\tau}(w)],$$

where  $N_{\tau}$  is the multiplicity function of  $U_{\tau}$ , namely the cardinality of the set  $U_{\tau}^{-1}(\{w\})$  (cf. [7]). Since  $E[\rho(-\delta u_{\tau})] = 1$ , it follows that  $N_{\tau} = 1 \mu$ -almost surely and this implies the existence of the inverse of  $U_{\tau}$  which is denoted as  $V_{\tau}$ . Note that  $V_{\tau}$  is of the form  $V_{\tau} = I_W + v_{\tau}$ , where  $v_{\tau} : W \to H$  and that the image of  $\mu$  under  $V_{\tau}$ , denoted as  $V_{\tau}\mu$ , is equivalent to  $\mu$  with the Radon–Nikodym density

$$\frac{\mathrm{d}V_{\tau}\mu}{\mathrm{d}\mu} = \rho(-\delta u_{\tau}). \tag{3}$$

$$\frac{\mathrm{d}v_{\tau}}{\mathrm{d}\tau} = -\left((I_H + \nabla u_{\tau})^{-1}\mathcal{L}u_{\tau}\right) \circ V_{\tau}.$$
(4)

Since

$$|v_{\beta}-v_{\alpha}| \leqslant \int\limits_{\alpha}^{\beta} \left| \frac{\mathrm{d}v_{\tau}}{\mathrm{d}\tau} \right|_{H} \mathrm{d}\tau,$$

and since  $L^0(\mu, H)$  is complete, in order to show that  $\lim_{\alpha,\beta\to 0} \mu(\{|v_\alpha - v_\beta| > c\}) = 0$ , for any c > 0, it suffices to show that

$$E\int_{0}^{\kappa}\left|\frac{\mathrm{d}v_{\tau}}{\mathrm{d}\tau}\right|\mathrm{d}\tau<\infty,$$

for some  $\kappa > 0$ . From the relations (3) and (4), we obtain

$$E \int_{\alpha}^{\beta} \left| \frac{\mathrm{d}v_{\tau}}{\mathrm{d}\tau} \right|_{H} \mathrm{d}\tau = E \int_{\alpha}^{\beta} \left| \left( \left( I_{H} + \nabla u_{\tau} \right)^{-1} \mathcal{L}u_{\tau} \right) \circ V_{\tau} \right|_{H} \mathrm{d}\tau \right|_{H} \mathrm{d}\tau$$
$$= E \int_{\alpha}^{\beta} \left| \left( I_{H} + \nabla u_{\tau} \right)^{-1} \mathcal{L}u_{\tau} \right|_{H} \rho(-\delta u_{\tau}) \mathrm{d}\tau$$

Hence the hypothesis (2) implies the existence of the limit  $\lim_{\tau \to 0} v_{\tau}$  in  $L^{1}(\mu, H)$  which we shall denote by v. Since  $v_{\tau} = -u_{\tau} \circ V_{\tau}$  and since  $(\rho(-\delta u_{\tau}), \tau \in [0, 1])$  is uniformly integrable,  $V\mu$  is absolutely continuous with respect to  $\mu$  and we have also the identity  $v = -u \circ V$ , where  $V = I_{W} + v$ . Now it is easy to see that  $U \circ V = V \circ U = I_{W}$   $\mu$ -almost surely.  $\Box$ 

Combining Theorem 1 with the inequality of T. Carleman (cf. [1] or [2], Corollary XI.6.28) which says:

$$\|\det_2(I_H + A)(I_H + A)^{-1}\| \leq \exp \frac{1}{2} (\|A\|_2^2 + 1),$$

for any Hilbert–Schmidt operator A, where the left hand side is the operator norm,  $det_2(I_H + A)$  denotes the modified Carleman–Fredholm determinant and  $\|\cdot\|_2$  denotes the Hilbert–Schmidt norm, we get

**Theorem 2.** Assume that  $u \in \mathbb{D}_{2,1}(H)$  such that  $E[\rho(-\delta u_{\tau})] = 1$  and that

$$E\left[e^{\frac{1}{2}\|\nabla u\|_{2}^{2}}\int_{0}^{1}P_{\tau}\left(\rho(-\delta u_{\tau})|\mathcal{L}u_{\tau}|_{H}\right)\mathrm{d}\tau\right]<\infty.$$

Then U satisfies the conclusions of Theorem 1.

**Proof.** The integrand in the relation (2) can be upperbounded as follows:

$$\left| (I_H + \nabla u_\tau)^{-1} \mathcal{L} u_\tau \right|_H \leq \exp \frac{1}{2} \left( \| \nabla u_\tau \|_2^2 + 1 \right) |\mathcal{L} u_\tau|_H$$
$$\leq |\mathcal{L} u_\tau|_H P_\tau \left( \exp \frac{1}{2} \left( \| \nabla u \|_2^2 + 1 \right) \right)$$

where the second line follows from the Jensen inequality. Here there is no term with det<sub>2</sub> since,  $\nabla u_{\tau}$  being quasinilpotent, its Carleman–Fredholm determinant is always equal to one. We then use the symmetry of  $P_{\tau}$  with respect to  $\mu$ .  $\Box$  **Corollary 1.** Suppose that *u* is adapted,  $E[\rho(-\delta u_{\tau})] = 1$  for all  $\tau \in [0, 1]$ . Let  $\varepsilon > 0$  be given and assume further that  $u \in \mathbb{D}_{\underline{\varepsilon+1},2}(H)$  and that the following relation holds:

$$E\left[\left(1+\mathrm{e}^{-\mathrm{e}(1+\varepsilon)\delta u}\right)\exp\left(\frac{1+\varepsilon}{2}\|\nabla u\|_{2}^{2}\right)\right]<\infty.$$
(5)

*Then*  $U = I_W + u$  *is*  $\mu$ *-almost surely invertible.* 

**Proof.** Let  $C_{\varepsilon}$  represent the left-hand side of the relation (5), then using the Hölder inequality we get

$$E\left[\int_{0}^{1} \left| (I_{H} + \nabla u_{\tau})^{-1} \mathcal{L} u_{\tau} \right|_{H} \rho(-\delta u_{\tau}) \, \mathrm{d}\tau \right] \leqslant C_{\varepsilon}^{\frac{1}{1+\varepsilon}} \|u\|_{\frac{1+\varepsilon}{\varepsilon},2}.$$

Hence the conclusion follows.  $\Box$ 

**Remark.** If we take  $\varepsilon = 1$  in Corollary 1, then it is easy to see, using the Wiener chaos expansion for  $E[|\mathcal{L}P_{\tau}u|_{H}^{2}]$  that

$$E\int_0^1 |\mathcal{L}P_{\tau}u|_H^2 \,\mathrm{d}\tau \leqslant \|u\|_{2,1}^2.$$

**Remark.** In the case where *u* is not adapted, the condition (5) with  $\varepsilon = 1$  is sufficient for the measure theoretic degree of the map *U* to be one as it is proven in Theorem 9.3.2 of [7].

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