



A rigorous derivation of the defocusing cubic nonlinear Schrödinger equation on \mathbb{T}^3 from the dynamics of many-body quantum systems

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Abstract

In this paper, we will obtain a rigorous derivation of the defocusing cubic nonlinear Schrödinger equation on the three-dimensional torus \mathbb{T}^3 from the many-body limit of interacting bosonic systems. This type of result was previously obtained on \mathbb{R}^3 in the work of Erdős, Schlein, and Yau [54–57], and on \mathbb{T}^2 and \mathbb{R}^2 in the work of Kirkpatrick, Schlein, and Staffilani [78]. Our proof relies on an unconditional uniqueness result for the Gross–Pitaevskii hierarchy at the level of regularity $\alpha = 1$, which is proved by using a modification of the techniques from the work of T. Chen, Hainzl, Pavlović and Seiringer [20] to the periodic setting. These techniques are based on the Quantum de Finetti theorem in the formulation of Ammari and Nier [6,7] and Lewin, Nam, and Rougerie [83]. In order to apply this approach in the periodic setting, we need to recall multilinear estimates obtained by Herr, Tataru, and Tzvetkov [74].

Having proved the unconditional uniqueness result at the level of regularity $\alpha = 1$, we will apply it in order to finish the derivation of the defocusing cubic nonlinear Schrödinger equation on \mathbb{T}^3 , which was started in the work of Elgart, Erdős, Schlein, and Yau [50]. In the latter work, the authors obtain all the steps of Spohn’s strategy for the derivation of the NLS [108], except for the final step of uniqueness. Additional arguments are necessary to show that the objects constructed in [50] satisfy the assumptions of the unconditional uniqueness theorem. Once we achieve this, we are able to prove the derivation result. In particular, we show *Propagation of Chaos* for the defocusing Gross–Pitaevskii hierarchy on \mathbb{T}^3 for suitably chosen initial data.

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1. Introduction

1.1. Setup of the problem

In this paper, we will consider the *Gross–Pitaevskii hierarchy* on the three-dimensional torus \mathbb{T}^3 . Given a more general spatial domain $\Lambda = \mathbb{R}^d$ or \mathbb{T}^d , the Gross–Pitaevskii (GP) hierarchy on Λ is a system of infinitely many equations given by:

$$\begin{cases} i \partial_t \gamma^{(k)} + (\Delta_{\vec{x}_k} - \Delta_{\vec{x}'_k}) \gamma^{(k)} = \lambda \cdot \sum_{j=1}^k B_{j,k+1}(\gamma^{(k+1)}) \\ \gamma^{(k)}|_{t=0} = \gamma_0^{(k)}. \end{cases} \tag{1}$$

For $k \in \mathbb{N}$, $\gamma_0^{(k)}$ is a complex-valued function on $\Lambda^k \times \Lambda^k$. Such a function is referred to as a *density matrix of order k* . Each $\gamma^{(k)} = \gamma^{(k)}(t)$ is a time-dependent density matrix of order k . Here, $\Delta_{\vec{x}_k}$ and $\Delta_{\vec{x}'_k}$ denote the Laplace operators in the first and second component of Λ^k . More precisely, $\Delta_{\vec{x}_k} := \sum_{j=1}^k \Delta_{x_j}$ and $\Delta_{\vec{x}'_k} := \sum_{j=1}^k \Delta_{x'_j}$. Furthermore, $B_{j,k+1}$ denotes the *collision operator* given by:

$$B_{j,k+1}(\sigma^{(k+1)}) := \text{Tr}_{k+1}[\delta(x_j - x_{k+1}), \sigma^{(k+1)}],$$

for $\sigma^{(k+1)}$ a density matrix of order $k + 1$. Tr_{k+1} denotes the trace in the variable x_{k+1} . In particular $B_{j,k+1}(\sigma^{(k+1)})$ is a density matrix of order k . A more precise definition is given in (11) below. Finally $\lambda \in \mathbb{R}$ is a non-zero coupling constant. In this paper, we will take $\lambda \in \{1, -1\}$. If $\lambda = 1$, we will call the GP hierarchy *defocusing*, and if $\lambda = -1$, we will call the hierarchy *focusing*.

The Gross–Pitaevskii hierarchy (1) is closely related to the cubic nonlinear Schrödinger equation (NLS) on Λ . We recall that the cubic NLS on Λ is given by:

$$\begin{cases} i \partial_t u + \Delta u = \lambda \cdot |u|^2 u, & \text{on } \mathbb{R}_t \times \Lambda \\ u|_{t=0} = \phi, & \text{on } \Lambda. \end{cases} \tag{2}$$

This problem is called *defocusing* if $\lambda = 1$ and *focusing* if $\lambda = -1$. Given a solution u of (2), we can build a solution to (1) by taking tensor products. In other words,

$$\gamma^{(k)}(t, \vec{x}_k; \vec{x}'_k) := \prod_{j=1}^k u(t, x_j) \overline{u(t, x'_j)} = |u\rangle \langle u|^{\otimes k}(t, \vec{x}_k; \vec{x}'_k) \tag{3}$$

solves (1) for initial data given by $\gamma_0^{(k)} = |\phi\rangle \langle \phi|^{\otimes k}$. Here, we denote by $|\cdot\rangle \langle \cdot|$ the *Dirac bracket*, which is defined as $|f\rangle \langle g|(x, x') := f(x) \overline{g(x')}$. These are called the *factorized solutions*. In this way, we can embed the nonlinear problem (2) into the linear problem (1).

In addition to the structural connection noted above, the study of the GP hierarchy is important in the context of the rigorous derivation of the NLS from the dynamics of many-body quantum systems. In particular, given a potential $V : \Lambda \rightarrow \mathbb{R}$, we can consider the N -body Hamiltonian H_N :

$$H_N := - \sum_{j=1}^N \Delta_j + \frac{1}{N} \sum_{\ell < j}^N V_N(x_\ell - x_j)$$

defined on a dense subspace of $L^2_{sym}(\Lambda^N)$, the space of all permutation-symmetric elements of $L^2(\Lambda^N)$. Here, V_N is a rescaled version of V involving N . Given $\Psi_{N,0} \in L^2_{sym}(\Lambda^N)$, we can study the N -body Schrödinger equation associated to H_N with initial data $\Psi_{N,0}$:

$$\begin{cases} i \partial_t \psi_{N,t} = H_N \psi_{N,t} \\ \psi_{N,t}|_{t=0} = \Psi_{N,0}. \end{cases} \tag{4}$$

The solution $\psi_{N,t}$ belongs to $L^2_{sym}(\Lambda^N)$. In particular, $\|\Psi_{N,t}\|_{L^2(\Lambda^N)} = \|\Psi_{N,0}\|_{L^2(\Lambda^N)}$. We define:

$$\gamma_{N,t}^{(k)} := \text{Tr}_{k+1,\dots,N} |\Psi_{N,t}\rangle\langle\Psi_{N,t}|. \tag{5}$$

Here $\text{Tr}_{k+1,\dots,N}$ denotes the partial trace in x_{k+1}, \dots, x_N . By definition, if $k > N$, we take $\gamma_{N,t}^{(k)} := 0$. This notation is explained in more detail in Subsection 2.3 below.

The sequence $(\gamma_{N,t}^{(k)})_k$ then solves the Bogoliubov–Born–Green–Kirkwood–Yvon (BBGKY) hierarchy:

$$\begin{aligned} & i \partial_t \gamma_{N,t}^{(k)} + (\Delta_{\bar{x}_k} - \Delta_{\bar{x}'_k}) \gamma_{N,t}^{(k)} \\ &= \frac{1}{N} \sum_{\ell < j}^k [V_N(x_\ell - x_j), \gamma_{N,t}^{(k)}] + \frac{N-k}{N} \sum_{j=1}^k \text{Tr}_{k+1} [V_N(x_j - x_{k+1}), \gamma_{N,t}^{(k+1)}]. \end{aligned} \tag{6}$$

We note that, formally, the BBGKY hierarchy converges to the defocusing GP hierarchy as $N \rightarrow \infty$. This heuristic can be made formal in the following sense: given $\phi \in L^2(\Lambda)$ with $\|\phi\|_{L^2(\Lambda)} = 1$, under additional assumptions on the sequence of initial data $(\Psi_N)_N \in \bigoplus_{N \in \mathbb{N}} L^2_{sym}(\Lambda^N)$ in terms of ϕ , one wants to show that there exists a sequence $N_j \rightarrow \infty$, which is independent of $k \in \mathbb{N}$ and $t \in [0, T]$ such that:

$$\text{Tr} |\gamma_{N_j,t}^{(k)} - |S_t(\phi)\rangle\langle S_t(\phi)|^{\otimes k}| \rightarrow 0, \tag{7}$$

as $j \rightarrow \infty$. Here, Tr denotes the trace and S_t denotes the flow map for (2).

We will refer to (7) as a *rigorous derivation of the cubic NLS equation on Λ from the dynamics of many-body quantum systems*. This type of result is also referred to as *Propagation of Chaos* in the sense that the particles become decoupled in the limit. In our paper, we will study this derivation in the context of the defocusing problem.

A strategy for proving (7) was developed by Spohn [108], and it consists of first showing that the sequence $(\gamma_{N,t}^{(k)})_N$ is relatively compact and that the limit points solve the GP hierarchy. Once this is established, one shows the GP hierarchy admits unique solutions in the class of objects obtained in the limit. This is a non-trivial step, due to the fact that the GP hierarchy is an infinite and non-closed system of PDEs.

Spohn [108] applied this approach in order to derive the nonlinear Hartree equation $i \partial_t u + \Delta u = (V * |u|^2) \cdot u$ on \mathbb{R}^3 when $V \in L^\infty(\mathbb{R}^3)$. The case of a Coulomb potential $V(x) = \pm \frac{1}{|x|}$ was later solved by Bardos, Golse, and Mauser [11] and Erdős and Yau [58]. An alternative proof of the latter result was subsequently given by Fröhlich, Knowles, and Schwarz. In a series of works [54–57], Erdős, Schlein, and Yau applied this strategy to obtain a derivation of the defocusing cubic nonlinear Schrödinger equation on \mathbb{R}^3 . The uniqueness step required a use of Feynman graph expansions [54, Sections 9 and 10]. Subsequently, a combinatorial reformulation of the uniqueness proof on \mathbb{R}^3 was given by Klainerman and Machedon [80]. This argument is applicable in a slightly different class of density matrices. The authors prove the uniqueness result in this class by using spacetime estimates reminiscent of their earlier work on null-forms for the nonlinear wave equation [79]. The fact that the objects obtained in the limit in the procedure outlined above satisfy the assumptions needed to apply the boardgame argument was first verified by Kirkpatrick, Schlein, and Staffilani [78] when $\Lambda = \mathbb{R}^2$ and $\Lambda = \mathbb{T}^2$. Related results were subsequently proven by T. Chen and Pavlović [27] and X. Chen and Holmer [36,39] when $\Lambda = \mathbb{R}^3$.

In the author’s joint work with Gressman and Staffilani [68, Proposition 3.3], it was noted that the spacetime estimate for the free evolution operator needed to apply the boardgame argument as in [80] does not hold in the energy space when $\Lambda = \mathbb{T}^3$. This is in sharp contrast to what happens when $\Lambda = \mathbb{R}^3$ or $\Lambda = \mathbb{T}^2$. Heuristically, this is a manifestation of the weaker dispersion on periodic domains, which also becomes weaker as the dimension becomes larger. In particular, it is not possible to apply the spacetime estimate and show uniqueness of solutions in the class from [80] in regularity $\alpha = 1$. A conditional uniqueness result in this class when $\alpha = 1$ is still a relevant open problem, which will require new tools to solve.

Recently, a new approach to studying the uniqueness problem was taken by T. Chen, Hainzl, Pavlović, and Seiringer in [20]. Here, the authors prove an unconditional uniqueness result for the GP hierarchy on \mathbb{R}^3 by means of the *Quantum de Finetti Theorem*. This theorem is a quantum analogue of the theorem of de Finetti on exchangeable sequences of random variables [44,45]. Related results to those of de Finetti were subsequently proven by Dynkin [48], Hewitt and Savage [75], and Diaconis and Freedman [47]. The Quantum de Finetti theorem states that, under certain assumptions, density matrices can be viewed as averages over factorized states. The precise statement is recalled in

Theorem 4.1 below. Similar results were first proven in the C^* algebra context in the work of Hudson and Moody [77], and Størmer [109]. Connections to density matrices were subsequently made in the work of Ammari and Nier [6,7], and Lewin, Nam, and Rougerie [83,84].

We recall from [68, Proposition 5.3] that it is possible to estimate factorized objects at the level of regularity of the energy space, i.e. when the Sobolev exponent is $\alpha = 1$. Thus, it would be reasonable to expect that it is also possible to estimate averages of factorized states, as one obtains from the Quantum de Finetti Theorem. In particular, the first result that we prove is:

Theorem 1 (Unconditional uniqueness for the GP hierarchy on \mathbb{T}^3 when $\alpha = 1$). *Let us fix $T > 0$. Suppose that $(\gamma^{(k)}(t))_k$ is a mild solution to the Gross–Pitaevskii hierarchy in $L^\infty_{t \in [0, T]} \mathfrak{S}^1$ such that, for each $t \in [0, T]$, there exist $\Gamma_{N,t} \in L^2_{sym}(\mathbb{T}^{3N} \times \mathbb{T}^{3N})$ which are non-negative as operators with trace equal to 1, such that:*

$$\text{Tr}_{k+1, \dots, N} \Gamma_{N,t} \xrightarrow{*} \gamma^{(k)}(t)$$

as $N \rightarrow \infty$ in the weak- $*$ topology of the trace class on $L^2_{sym}(\mathbb{T}^{3k})$. Then, the solution $(\gamma^{(k)}(t))_k$ is uniquely determined by the initial data $(\gamma_0^{(k)})_k$.

Theorem 1 is stated as **Theorem 4.6** in Subsection 4.2. Given a Sobolev exponent $\alpha \in \mathbb{R}$, the space \mathfrak{S}^α is defined in **Definition 4.4** below. Furthermore, the concept of a mild solution to the Gross–Pitaevskii hierarchy is given in **Definition 4.5** below. We note that **Theorem 1** applies both in the *defocusing* ($\lambda = 1$) and the *focusing* ($\lambda = -1$) context. Let us remark that **Theorem 1** does not immediately improve on the main result of [68] because the spaces which are used in [68] are different and cannot in general be compared to the ones used in the present paper. In particular, the norms used in [68] are of Hilbert–Schmidt type, whereas the norms used in the present paper are obtained from the trace as in [53–57]. In [68], there is an additional condition involving the collision operator as was the case in the analysis of [80].

Let us note that in [50], Elgart, Erdős, Schlein and Yau complete the first step in the strategy developed by Spohn in the context of the derivation of the defocusing cubic NLS on \mathbb{T}^3 . In other words, they show that the sequence $(\gamma_{N,t}^{(k)})_N$ is relatively compact and that the limit points of this sequence solve the Gross–Pitaevskii hierarchy. Their analysis applies at the level of regularity $\alpha = 1$.

In [50], the authors state the uniqueness step as an open problem. In our paper, we will resolve this problem. As a result, we will obtain a derivation of the defocusing cubic NLS on \mathbb{T}^3 . The result that we prove is:

Theorem 2 (A derivation of the defocusing cubic NLS on \mathbb{T}^3). *Let $\phi \in H^1(\mathbb{T}^3)$ with $\|\phi\|_{L^2(\mathbb{T}^3)} = 1$ be given. Suppose that the sequence of initial data in the N -body Schrödinger equation $(\Psi_N)_N \in \bigoplus_{N \in \mathbb{N}} L^2_{sym}(\mathbb{T}^{3N})$ satisfies the assumptions of:*

- 1) *Bounded energy per particle:* $\sup_{N \in \mathbb{N}} \langle H_N \Psi_N, \Psi_N \rangle_{L^2(\mathbb{T}^{3N})} < \infty$.
- 2) *Asymptotic factorization:* $\text{Tr} |\text{Tr}_{2,3, \dots, N} |\Psi_N\rangle \langle \Psi_N| - |\phi\rangle \langle \phi|| \rightarrow 0$ as $N \rightarrow \infty$.

For fixed $N \in \mathbb{N}$, let $(\gamma_{N,t}^{(k)})_k$ denote the solution of the BBGKY hierarchy evolving from initial data $(\gamma_N^{(k)})_k := \text{Tr}_{k+1, \dots, N} |\Psi_N\rangle \langle \Psi_N|$. Let $T > 0$ be fixed. Then, there exists a sequence $N_j \rightarrow \infty$, such that for all $k \in \mathbb{N}$ and for all $t \in [0, T]$ the convergence (7) holds.

Theorem 2 is given as **Theorem 5.4** in Section 5. We note that this result is stated only in the defocusing context. In the proof of **Theorem 2**, we will use a uniqueness argument based on **Theorem 1** and the analysis of [50]. For details on the application of the uniqueness result in the proof of **Theorem 2**, we refer the reader to (72) and to **Remark 5.9** below. The assumptions of **Theorem 2** are satisfied for purely factorized states $\Psi_N = |\phi\rangle \langle \phi|^{\otimes N}$, and the result holds in this case as is noted in **Corollary 5.5** below.

1.2. Physical interpretation

The Gross–Pitaevskii hierarchy and the nonlinear Schrödinger equation have a related physical interpretation. Both objects occur in the context of Bose–Einstein condensation, a state of matter consisting of dilute bosonic particles which are cooled to a temperature close to absolute zero. At such a temperature, these particles tend to occupy the lowest quantum state, which can be expressed mathematically as the ground state of an energy functional related to the NLS. In this context, the NLS equation is sometimes referred to as the *Gross–Pitaevskii equation*, following the work of Gross [72] and Pitaevskii [98].

The physical phenomenon of Bose–Einstein condensation was theoretically predicted in the pioneering works of Bose [15] and Einstein [49] in 1924–1925. Their prediction was verified by experiments conducted independently by the teams led by Cornell and Wieman [9] and Ketterle [43] in 1995. These two groups were awarded the Physics Nobel Prize in 2001.

1.3. Additional related results

In addition to the references mentioned above, there is a rich literature devoted to the connection between NLS-type equations and hierarchies similar to (1). In addition to the strategy developed by Spohn [108], outlined above, an independent strategy based on Fock space methods was simultaneously developed by Hepp [73] and Ginibre and Velo [66,67]. Spohn’s strategy was subsequently applied in the derivation problem in [1,2,11–14,27,30,33–38,51,59,62,95,111], whereas the Fock space techniques were subsequently applied in [31,32,60,61,69–71]. The question of deriving the nonlinear Hartree equation was revisited in the work of Fröhlich, Tsai, and Yau [63–65]. Once one obtains a derivation of the NLS-type equation, it is natural to ask what is the rate of convergence. This question was first addressed by Rodnianski and Schlein [101]. Subsequent results on this problem have been obtained in [8,14,19,31,32,52,60,69–71,81,82,93,95–97]. The Gross–Pitaevskii hierarchy has been studied at the N -body level by Lieb and Seiringer [88], Lieb, Seiringer and Yngvason [91,92], and Lieb, Seiringer, Yngvason, and Solovej [89]. In these works, the assumption of asymptotic factorization given in the assumption of [Theorem 2](#) was rigorously verified for a sequence of appropriate ground states. For more details on this aspect of the problem, we refer the reader to the expository work [90] and to the references therein. This question was revisited in a recent preprint of Lewin, Nam, and Rougerie [85]. A connection of the above problems with optical lattice models was explored in the work of Aizenman, Lieb, Seiringer, Yngvason, and Solovej [3,4]. An expository survey of many of the above results can be found in [103].

The Gross–Pitaevskii hierarchy has been studied as a Cauchy problem in its own right in the recent works of T. Chen and Pavlović [22,23,25–27], and in their joint works with Tzirakis [28,29], as well as in the subsequent work of Z. Chen and Liu [41]. The motivation is to study the Gross–Pitaevskii hierarchy as a generalization of the nonlinear Schrödinger equation via the factorized solutions and to prove analogues of the known results for the Cauchy problem for the NLS in the context of the GP hierarchy. By appropriately modifying the collision operator, it is also possible to consider a hierarchy which is related to the quintic NLS [22–26,28,29], as well as the NLS with more general power-type nonlinearities [111]. The Cauchy problem associated to the Hartree equation for infinitely many particles has recently been studied by Lewin and Sabin [86,87]. Randomization techniques similar to those used in the context of the Cauchy problem for nonlinear dispersive equations in the work of Bourgain [17] and Burq and Tzvetkov [18] were used in order to study randomized forms of the Gross–Pitaevskii hierarchy in the author’s joint work with Staffilani [107], as well as in the author’s work [106].

Techniques similar to those used in [20] have been applied in order to show scattering results in the subsequent work of T. Chen, Hainzl, Pavlović and Seiringer [21]. The unconditional uniqueness result of [20] was subsequently extended to lower regularities by Hong, Taliaferro, and Xie in [76]. We note that the methods used in [76] do not directly apply to the periodic setting due to the weaker dispersion. It is an interesting open problem to see if the unconditional uniqueness in the periodic section can be extended to lower regularities. We note that the Quantum de Finetti theorem was recently also applied in the context of the Chern–Simons–Schrödinger hierarchy in the work of X. Chen and Smith [40].

1.4. Main ideas of the proof

The proof of the unconditional uniqueness result in [Theorem 1](#) will be based on the Weak Quantum de Finetti Theorem used in the work of T. Chen, Pavlović, Hainzl, and Seiringer [20] in the non-periodic setting. Several modifications will be necessary in order to apply these methods in the periodic setting. The main point is that one needs to use Strichartz estimates on \mathbb{T}^3 . We will use the trilinear estimate from the work of Herr, Tataru, and Tzvetkov [74]. This result is recalled in [Proposition 2.1](#) below. In particular, using this estimate we are able to prove a product estimate at the level of H^1 regularity in [Proposition 3.1](#). Following the terminology of [20], this result is used in order to prove the *bound on the L^2 level* (41) and the *bound on the H^1 level* (42). As in [20], these estimates allow us to bound the contributions from the factors corresponding to distinguished and regular trees in Subsubsection 4.2.5 below.

The proof of the derivation of the defocusing cubic nonlinear Schrödinger equation on \mathbb{T}^3 , i.e. of [Theorem 2](#), is based on applying [Theorem 1](#) to the solutions of the defocusing Gross–Pitaevskii hierarchy on \mathbb{T}^3 , which were previously constructed in the work of Elgart, Erdős, Schlein, and Yau [50]. In particular, in this paper, the authors construct solutions without a statement about uniqueness. The spaces in which they work do not originally involve the trace norm. In Section 5, we show that the assumptions of [Theorem 1](#) hold for these solutions. As an intermediate step, we will have to use the approximation argument stated in Subsection 5.1 in order to control higher powers of the N -body Hamiltonian as was needed in the assumptions in [50]. We note that such a procedure was used in [54,56,57] in the non-periodic setting, as well as in [78] in the periodic setting.

1.5. Organization of the paper

In Section 2, we will give the relevant notation and we will recall some important facts from Harmonic and Functional analysis. In particular, in Subsection 2.1, we will recall the definition of the function spaces which we will use on \mathbb{T}^3 . An important trilinear estimate from the work of Herr, Tataru, and Tzvetkov [74] is recalled in [Proposition 2.1](#). Subsection 2.2 is devoted to notation and operations concerning density matrices. In Subsection 2.3, we recall some important facts about trace class operators. Section 3 is devoted to the proof of the product estimate at the level of H^1 regularity, which is stated in [Proposition 3.1](#). In Section 4, we prove [Theorem 1](#), which is stated in [Theorem 4.6](#). The variant of the Quantum de Finetti Theorem which we will use is recalled in Subsection 4.1. The unconditional uniqueness result stated in [Theorem 1](#) is proved in Subsection 4.2. Section 5 is devoted to the proof of [Theorem 2](#), which is stated as [Theorem 5.4](#). In Subsection 5.1, we recall the relevant approximation procedure which allows us to control higher powers of H_N applied to the initial data. A comparison of different forms of convergence which appear in the problem is given in Subsection 5.2. In Subsection 5.3, we give a proof of the derivation of the defocusing cubic NLS on \mathbb{T}^3 , stated in [Theorem 2](#).

2. Notation and known facts from harmonic and functional analysis

In our paper, given non-negative quantities A and B , let us denote by $A \lesssim B$ an estimate of the form $A \leq CB$, for some constant $C > 0$. If C depends on a parameter p , we write the inequality as $A \lesssim_p B$ or we emphasize that $C = C(p)$. Throughout our paper, we will take the spatial domain Λ to be the three-dimensional torus $\mathbb{T}^3 = [0, 2\pi]^3$.

2.1. Function spaces

Given $f \in L^2(\Lambda)$, we can define its Fourier transform as:

$$\widehat{f}(n) := \int_{\Lambda} f(x) e^{-i\langle x, n \rangle} dx.$$

Here, $n \in \mathbb{Z}^3$ and $\langle \cdot, \cdot \rangle$ denotes the inner product on \mathbb{R}^3 .

Let us take the following convention for the Japanese bracket $\langle \cdot \rangle$ on Λ :

$$\langle x \rangle := \sqrt{1 + |x|^2}.$$

Given $s \in \mathbb{R}$, the Sobolev space $H^s(\Lambda)$ is the normed space corresponding to:

$$\|f\|_{H^s(\Lambda)} := \left(\sum_{n \in \mathbb{Z}^3} |\widehat{f}(n)|^2 \cdot \langle n \rangle^{2s} \right)^{\frac{1}{2}}.$$

Let us recall the periodic version of the Sobolev embedding estimate:

$$\|\phi\|_{L^6(\Lambda)} \lesssim \|\phi\|_{H^1(\Lambda)}. \tag{8}$$

The estimate (8) can be deduced from the analogous estimate on \mathbb{R}^3 by using an extension argument. More precisely, given $\phi \in H^1(\Lambda)$, we define ϕ^e to be the periodic extension of ϕ to all of \mathbb{R}^3 . We then let $F \in C_0^\infty(\mathbb{R}^3)$ be a function which is identically 1 on Λ . Then, $\|F \cdot \phi^e\|_{L^6(\mathbb{R}^3)} \sim \|\phi\|_{L^6(\Lambda)}$ and $\|F \cdot \phi^e\|_{H^1(\mathbb{R}^3)} \lesssim \|\phi\|_{H^1(\Lambda)}$. By the Sobolev embedding theorem on \mathbb{R}^3 , we know that $\|F \cdot \phi^e\|_{L^6(\mathbb{R}^3)} \lesssim \|F \cdot \phi^e\|_{H^1(\mathbb{R}^3)}$. We can hence deduce (8).

Suppose that $N = 2^j$ is a dyadic integer. We will denote by P_N the Littlewood–Paley projection to frequencies which are of the order N . In order to do this, let us first consider $\Phi \in C_0^\infty(-2, 2)$, a non-negative, even function, such that $\Phi \equiv 1$ on $[-1, 1]$. We define the function Φ_N on \mathbb{Z}^3 by:

$$\begin{cases} \Phi\left(\frac{|\xi|}{N}\right) - \Phi\left(\frac{2|\xi|}{N}\right), & \text{for } N \geq 2. \\ \Phi(|\xi|), & \text{for } N = 1. \end{cases}$$

The operator P_N is defined as:

$$(P_N f)^\wedge(\xi) := \Phi_N(\xi) \cdot \widehat{f}(\xi). \tag{9}$$

In particular,

$$\text{supp}(P_N f)^\wedge \subseteq \{|\xi| \sim N\}.$$

Proposition 2.1 (A consequence of Proposition 3.5 in [74]). *There exists a universal constant $\delta > 0$ such that for any dyadic integers N_1, N_2, N_3 with $N_1 \geq N_2 \geq N_3 \geq 1$, and for any finite interval $I \subseteq \mathbb{R}$, it is the case that:*

$$\begin{aligned} & \|P_{N_1} e^{it\Delta} f_1 \cdot P_{N_2} e^{it\Delta} f_2 \cdot P_{N_3} e^{it\Delta} f_3\|_{L^2(I \times \Lambda)} \\ & \lesssim N_2 N_3 \max\left\{\frac{N_3}{N_1}, \frac{1}{N_2}\right\}^\delta \cdot \|P_{N_1} f_1\|_{L^2(\Lambda)} \cdot \|P_{N_2} f_2\|_{L^2(\Lambda)} \cdot \|P_{N_3} f_3\|_{L^2(\Lambda)}. \end{aligned}$$

Let us note that the implied constant in the above estimate depends on the length $|I|$ of the interval I . It can be taken to be an increasing function of $|I|$. The result of Proposition 2.1 can be deduced by combining [74, Proposition 3.5] together with [74, Proposition 2.10] and the inclusion of spaces given in [74, Proposition 2.8]. We will omit the details. An analogous estimate was shown on the 3D irrational torus in [110, Proposition 4.1].

2.2. Density matrices

Let us fix $k \in \mathbb{N}$. A density matrix of order k on Λ is a function:

$$\gamma^{(k)} : \Lambda^k \times \Lambda^k \rightarrow \mathbb{C}.$$

It is also sometimes called a k -particle density matrix.

Given $\gamma^{(k)} \in L^2(\Lambda^k \times \Lambda^k)$, we say that $\gamma^{(k)} \in L_{\text{sym}}^2(\Lambda^k \times \Lambda^k)$ if

$$\gamma^{(k)}(x_{\sigma(1)}, \dots, x_{\sigma(k)}; x'_{\sigma(1)}, \dots, x'_{\sigma(k)}) = \gamma^{(k)}(x_1, \dots, x_k; x'_1, \dots, x'_k)$$

for all $(\vec{x}_k; \vec{x}'_k) = (x_1, \dots, x_k; x'_1, \dots, x'_k) \in \Lambda^k \times \Lambda^k$ and for all $\sigma \in S^k$. Similarly, if $\Psi \in L^2(\Lambda^k)$, we say that $\Psi \in L_{\text{sym}}^2(\Lambda^k)$ if:

$$\Psi(x_{\sigma(1)}, \dots, x_{\sigma(k)}) = \Psi(x_1, \dots, x_k)$$

for all $\vec{x}_k = (x_1, \dots, x_k) \in \Lambda^k$ and for all $\sigma \in S_k$. Here, we will always assume that it is clear from context whether the object we are considering is a function or a density matrix. In other words, we will be able to distinguish $L^2_{sym}(\Lambda^k \times \Lambda^k)$ from $L^2_{sym}(\Lambda^{2k})$. Let us note that, if $\Psi \in L^2_{sym}(\Lambda^k)$, then $|\Psi\rangle\langle\Psi| \in L^2_{sym}(\Lambda^k \times \Lambda^k)$. We will sometimes consider $\gamma^{(k)}$ as an operator on $L^2(\Lambda^k)$ by means of identifying an integral operator with its kernel. We will use this convention throughout our work.

Given $\gamma^{(k)}$, a density matrix of order k which belongs to $L^2(\Lambda^k \times \Lambda^k)$, its Fourier transform $(\gamma^{(k)})^\wedge$ is defined as follows:

$$(\gamma^{(k)})^\wedge(\vec{n}_k; \vec{n}'_k) := \int_{\Lambda^k \times \Lambda^k} \gamma^{(k)}(\vec{x}_k; \vec{x}'_k) \cdot e^{-i \cdot \sum_{j=1}^k (x_j, n_j) + i \cdot \sum_{j=1}^k (x'_j, n'_j)} d\vec{x}_k d\vec{x}'_k, \tag{10}$$

for $(\vec{n}_k; \vec{n}'_k) = (n_1, \dots, n_k; n'_1, \dots, n'_k) \in \mathbb{Z}^{3k} \times \mathbb{Z}^{3k}$. The convention for (10) is consistent with our definition of factorized solutions (3). We will also write $(\gamma^{(k)})^\wedge$ as $\widehat{\gamma}^{(k)}$ for simplicity of notation.

The differential operator $i\partial_t + (\Delta_{\vec{x}_k} - \Delta_{\vec{x}'_k})$ acts on density matrices of order k . Its associated *free evolution operator* is denoted by $\mathcal{U}^{(k)}(t)$. More precisely, for $\gamma^{(k)}$, a density matrix of order k :

$$\mathcal{U}^{(k)}(t)\gamma^{(k)} := e^{it \sum_{j=1}^k \Delta_{x_j}} \gamma^{(k)} e^{-it \sum_{j=1}^k \Delta_{x'_j}}.$$

In this way, we obtain a solution to:

$$(i\partial_t + (\Delta_{\vec{x}_k} - \Delta_{\vec{x}'_k}))\mathcal{U}^{(k)}(t)\gamma^{(k)} = 0.$$

Given $\alpha \in \mathbb{R}$, we can use the Fourier transform and define the operation $S^{(k,\alpha)}$ of *differentiation of order α* on matrices of order k . Given $\gamma^{(k)}$ a density matrix of order k , we let $S^{(k,\alpha)}\gamma^{(k)}$ be the density matrix of order k whose Fourier transform is given by:

$$\begin{aligned} (S^{(k,\alpha)}\gamma^{(k)})^\wedge(n_1, \dots, n_k; n'_1, \dots, n'_k) \\ := \langle n_1 \rangle^\alpha \cdots \langle n_k \rangle^\alpha \cdot \langle n'_1 \rangle^\alpha \cdots \langle n'_k \rangle^\alpha \cdot \widehat{\gamma}^{(k)}(n_1, \dots, n_k; n'_1, \dots, n'_k). \end{aligned}$$

We now define the *collision operator*. Given $k \in \mathbb{N}$ and $j \in \{1, 2, \dots, k\}$, the collision operator $B_{j,k+1}$ acts linearly on density matrices of order $k+1$ as:

$$\begin{aligned} B_{j,k+1}(\gamma^{(k+1)}) &= \text{Tr}_{k+1}[\delta(x_j - x_{k+1}), \gamma^{(k+1)}] \\ &= \int_{\Lambda} (\delta(x_j - x_{k+1})\gamma^{(k+1)}(\vec{x}_k, x_{k+1}; \vec{x}'_k, x_{k+1}) dx_{k+1} \\ &\quad - \delta(x'_j - x_{k+1})\gamma^{(k+1)}(\vec{x}_k, x_{k+1}; \vec{x}'_k, x_{k+1})) dx_{k+1}. \end{aligned} \tag{11}$$

Here, δ denotes the Dirac delta function. We note that $B_{j,k+1}(\gamma^{(k+1)})$ is a density matrix of order k . We sometimes omit the parenthesis and write $B_{j,k+1}\gamma^{(k+1)}$. The *full collision operator* B_{k+1} is given by:

$$B_{k+1} := \sum_{j=1}^k B_{j,k+1}. \tag{12}$$

2.3. Some facts about trace class operators

Let \mathcal{H} be a separable Hilbert space over \mathbb{C} with inner product (\cdot, \cdot) . Let A be a bounded operator on \mathcal{H} and let $\{\vec{e}_j, j \in \mathbb{N}\}$ be an orthonormal basis of \mathcal{H} . We say that A belongs to the *trace class on \mathcal{H}* if the quantity:

$$\|A\|_1 := \sum_{j=1}^{+\infty} ((A^*A)^{\frac{1}{2}} \vec{e}_j, \vec{e}_j)$$

is finite. If this is the case, the quantity $\|A\|_1$ can be shown to be independent of the choice of basis. The space of all trace-class operators on \mathcal{H} is denoted by $\mathcal{L}^1(\mathcal{H})$.

Given $A \in \mathcal{H}$ and $\{\vec{e}_j, j \in \mathbb{N}\}$ as above, the trace of A :

$$\text{Tr } A := \sum_{j=1}^{+\infty} (A\vec{e}_j, \vec{e}_j)$$

is well-defined and is independent of the choice of basis. This fact is proved in [99, Theorem VI.24]. We will sometimes write $\text{Tr } A$ as $\text{Tr}_{\mathcal{H}} A$ in order to emphasize that the operator A acts on the Hilbert space \mathcal{H} . We note that $\|A\|_1 = \text{Tr}|A| = \text{Tr}(A^*A)^{\frac{1}{2}}$.

Let $\mathcal{K}(\mathcal{H})$ denote the space of all compact operators on \mathcal{H} . It is a normed vector space with respect to the operator norm $\|\cdot\|$. The following duality result then holds:

$$(\mathcal{L}^1(\mathcal{H}), \|\cdot\|_1) = (\mathcal{K}(\mathcal{H}), \|\cdot\|)^* \tag{13}$$

where the duality pairing is given by:

$$(T, K) \in \mathcal{L}^1(\mathcal{H}) \times \mathcal{K}(\mathcal{H}) \mapsto \text{Tr}(TK). \tag{14}$$

For the proof of the above result we refer the reader to [99, Theorem VI.26].

Given $N \in \mathbb{N}$, we let $\mathcal{H}^N := \bigotimes_{j=1}^N \mathcal{H}$. If $A \in \mathcal{L}^1(\mathcal{H}^N)$, and $k \in \{1, 2, \dots, N\}$, we define $\text{Tr}_{k+1, \dots, N} A$ to be the element of $\mathcal{L}^1(\mathcal{H}^k)$ obtained by taking the trace of A in the last $N - k$ factors of \mathcal{H} . If $k > N$, we define $\text{Tr}_{k+1, \dots, N} A$ to be equal to zero. We refer to this procedure as taking the partial trace.

Throughout the paper, we will primarily consider the case when $\mathcal{H} = L^2(\Lambda^k)$ with respect to the Lebesgue measure for some $k \in \mathbb{N}$. In this case, we will write $\mathcal{L}^1(\mathcal{H})$ as \mathcal{L}_k^1 and $\mathcal{K}(\mathcal{H})$ as \mathcal{K}_k . These are subspaces of the space of density matrices of order k on Λ . We observe that if $A \in \mathcal{L}_k^1$ is an integral operator with kernel $\gamma^{(k)} \in L^2(\Lambda^k \times \Lambda^k)$, then:

$$\text{Tr } A = \int_{\Lambda^k \times \Lambda^k} \gamma^{(k)}(\vec{x}_k; \vec{x}_k) d\vec{x}_k.$$

For a more detailed discussion on trace class operators, we refer the reader to [42, Sections 18 and 19] and [99, Section VI].

3. Multilinear estimates

We consider the spatial domain $\Lambda := \mathbb{T}^3$. The main result of this section is:

Proposition 3.1. *Suppose that $s \in [0, 1]$ and suppose that I is bounded time interval. Then, the following estimate holds:*

$$\begin{aligned} & \left\| |\nabla_x|^s (e^{it\Delta} f_1 \cdot \overline{e^{it\Delta} f_2} \cdot e^{it\Delta} f_3) \right\|_{L^2(I \times \Lambda)} \\ & \lesssim \min \left\{ \|f_1\|_{H^s} \cdot \|f_2\|_{H^1} \cdot \|f_3\|_{H^1}, \|f_1\|_{H^1} \cdot \|f_2\|_{H^s} \cdot \|f_3\|_{H^1}, \|f_1\|_{H^1} \cdot \|f_2\|_{H^1} \cdot \|f_3\|_{H^s} \right\}, \end{aligned}$$

whenever f_1, f_2, f_3 are functions on Λ for which the right-hand side is well-defined. Here, the implied constant depends on s and on the length of I .

Proof. Suppose that f_1, f_2, f_3 are as in the assumptions of the proposition. We will dyadically localize the factors according to (9). Namely, for fixed dyadic integers N_1, N_2, N_3 , we want to estimate the expression:

$$\left\| |\nabla_x|^s (P_{N_1} e^{it\Delta} f_1 \cdot P_{N_2} \overline{e^{it\Delta} f_2} \cdot P_{N_3} e^{it\Delta} f_3) \right\|_{L^2(I \times \Lambda)}$$

which is:

$$\begin{aligned} & \lesssim_s \max\{N_1, N_2, N_3\}^s \cdot \left\| P_{N_1} e^{it\Delta} f_1 \cdot P_{N_2} \overline{e^{it\Delta} f_2} \cdot P_{N_3} e^{it\Delta} f_3 \right\|_{L^2(I \times \Lambda)} \\ & = \max\{N_1, N_2, N_3\}^s \cdot \left\| P_{N_1} e^{it\Delta} f_1 \cdot P_{N_2} e^{it\Delta} f_2 \cdot P_{N_3} e^{it\Delta} f_3 \right\|_{L^2(I \times \Lambda)}. \end{aligned} \tag{15}$$

Hence, we do not have to keep track of complex conjugates throughout the proof. Let us define:

$$\psi_j := e^{it\Delta} f_j \quad \text{for } j = 1, 2, 3. \tag{16}$$

Since $\|P_{N_1}\psi_1 \cdot P_{N_2}\psi_2 \cdot P_{N_3}\psi_3\|_{L^2(I \times \Lambda)}$ is symmetric under the permutation of the functions f_j , we may assume without loss of generality that $N_1 \geq N_2 \geq N_3 \geq 1$. In this case, the expression in (15) equals:

$$N_1^s \cdot \|P_{N_1}\psi_1 \cdot P_{N_2}\psi_2 \cdot P_{N_3}\psi_3\|_{L^2(I \times \Lambda)} =: K_{N_1, N_2, N_3}.$$

We will now estimate:

$$K := \sum_{\substack{N_1, N_2, N_3 \\ N_1 \geq N_2 \geq N_3}} K_{N_1, N_2, N_3}.$$

In order to do this, we will estimate each term K_{N_1, N_2, N_3} by considering the different cases determined by the relative sizes of N_1 and N_2 .

Case 1: $N_1 \gg N_2$

In this case, we will estimate K_{N_1, N_2, N_3} by using duality. Namely, let us take $u \in L^2(I \times \Lambda)$ such that $\|u\|_{L^2(I \times \Lambda)} = 1$ and let us note that:

$$\int_I \int_\Lambda u \cdot (P_{N_1}\psi_1 \cdot P_{N_2}\psi_2 \cdot P_{N_3}\psi_3) dx dt = \int_I \int_\Lambda P_{N_0}u \cdot (P_{N_1}\psi_1 \cdot P_{N_2}\psi_2 \cdot P_{N_3}\psi_3) dx dt,$$

for some $N_0 \sim N_1$. More precisely, N_0 here denotes an appropriate dilation of the dyadic integer N_1 and we extend the definition of P_{N_0} accordingly.

Consequently, we need to estimate the following sum:

$$\begin{aligned} K^1 &:= \sum_{\substack{N_0, N_1, N_2, N_3 \\ N_0 \sim N_1, N_1 \gg N_2 \geq N_3}} N_1^s \cdot \left| \int_I \int_\Lambda P_{N_0}u \cdot (P_{N_1}\psi_1 \cdot P_{N_2}\psi_2 \cdot P_{N_3}\psi_3) dx dt \right| \\ &\leq \sum_{\substack{N_0, N_1, N_2, N_3 \\ N_0 \sim N_1, N_1 \gg N_2 \geq N_3}} N_1^s \cdot \|P_{N_0}u\|_{L^2(I \times \Lambda)} \cdot \|P_{N_1}\psi_1 \cdot P_{N_2}\psi_2 \cdot P_{N_3}\psi_3\|_{L^2(I \times \Lambda)} \\ &\lesssim_{|I|} \sum_{\substack{N_0, N_1, N_2, N_3 \\ N_0 \sim N_1, N_1 \gg N_2 \geq N_3}} \|P_{N_0}u\|_{L^2(I \times \Lambda)} \cdot N_1^s \cdot N_2 \cdot N_3 \cdot \max\left\{\frac{N_3}{N_1}, \frac{1}{N_2}\right\}^\delta \\ &\quad \cdot \|P_{N_1}f_1\|_{L^2(\Lambda)} \cdot \|P_{N_2}f_2\|_{L^2(\Lambda)} \cdot \|P_{N_3}f_3\|_{L^2(\Lambda)}. \end{aligned} \tag{17}$$

Here, we used the Cauchy–Schwarz inequality on $I \times \Lambda$, (16) and Proposition 2.1. By dyadic localization, the expression in (17) is:

$$\begin{aligned} &\lesssim \sum_{\substack{N_0, N_1, N_2, N_3 \\ N_0 \sim N_1, N_1 \gg N_2 \geq N_3}} \|P_{N_0}u\|_{L^2(I \times \Lambda)} \cdot \max\left\{\frac{N_3}{N_1}, \frac{1}{N_2}\right\}^\delta \cdot \|P_{N_1}f_1\|_{H^s(\Lambda)} \cdot \|P_{N_2}f_2\|_{H^1(\Lambda)} \cdot \|P_{N_3}f_3\|_{H^1(\Lambda)} \\ &\lesssim \sum_{\substack{N_0, N_1, N_2, N_3 \\ N_0 \sim N_1, N_1 \gg N_2 \geq N_3}} \|P_{N_0}u\|_{L^2(I \times \Lambda)} \cdot \left\{ \left(\frac{N_3}{N_1}\right)^{\frac{\delta}{2}} \cdot \left(\frac{N_2}{N_1}\right)^{\frac{\delta}{2}} + \frac{1}{N_2^{\frac{\delta}{2}} \cdot N_3^{\frac{\delta}{2}}} \right\} \\ &\quad \cdot \|P_{N_1}f_1\|_{H^s(I)} \cdot \|P_{N_2}f_2\|_{H^1(\Lambda)} \cdot \|P_{N_3}f_3\|_{H^1(\Lambda)}. \end{aligned}$$

The latter estimate follows from the fact that $N_1 \geq N_2 \geq N_3$. We use the Cauchy–Schwarz inequality in N_2 and in N_3 to deduce that this quantity is:

$$\lesssim \sum_{\substack{N_0, N_1 \\ N_0 \sim N_1}} \|P_{N_0} u\|_{L^2(I \times \Lambda)} \cdot \|P_{N_1} f_1\|_{H^s(\Lambda)} \cdot \|f_2\|_{H^1(\Lambda)} \cdot \|f_3\|_{H^1(\Lambda)}$$

which by using the Cauchy–Schwarz inequality in $N_0 \sim N_1$ is:

$$\lesssim \|u\|_{L^2(I \times \Lambda)} \cdot \|f_1\|_{H^s(\Lambda)} \cdot \|f_2\|_{H^1(\Lambda)} \cdot \|f_3\|_{H^1(\Lambda)} = \|f_1\|_{H^s(\Lambda)} \cdot \|f_2\|_{H^1(\Lambda)} \cdot \|f_3\|_{H^1(\Lambda)},$$

since $\|u\|_{L^2(I \times \Lambda)} = 1$ by assumption. In particular:

$$K^1 \lesssim \|f_1\|_{H^s(\Lambda)} \cdot \|f_2\|_{H^1(\Lambda)} \cdot \|f_3\|_{H^1(\Lambda)}. \tag{18}$$

Since $s \in [0, 1]$, we know that $N_1^s \cdot N_2 \cdot N_3 \leq N_1 \cdot N_2^s \cdot N_3$, and hence we can replace the factor of $N_1^s \cdot N_2 \cdot N_3$ in (17) by $N_1 \cdot N_2^s \cdot N_3$. By the above argument, it follows that:

$$K^1 \lesssim \|f_1\|_{H^1(\Lambda)} \cdot \|f_2\|_{H^s(\Lambda)} \cdot \|f_3\|_{H^1(\Lambda)}. \tag{19}$$

Similarly, since $N_1^s \cdot N_2 \cdot N_3 \leq N_1 \cdot N_2 \cdot N_3^s$, it follows that:

$$K^1 \lesssim \|f_1\|_{H^1(\Lambda)} \cdot \|f_2\|_{H^1(\Lambda)} \cdot \|f_3\|_{H^s(\Lambda)}. \tag{20}$$

We use duality and (18), (19), and (20) in order to deduce that:

$$\begin{aligned} & \sum_{\substack{N_1, N_2, N_3 \\ N_1 \gg N_2 \geq N_3}} K_{N_1, N_2, N_3} \\ & \lesssim \min\{\|f_1\|_{H^s} \cdot \|f_2\|_{H^1} \cdot \|f_3\|_{H^1}, \|f_1\|_{H^1} \cdot \|f_2\|_{H^s} \cdot \|f_3\|_{H^1}, \|f_1\|_{H^1} \cdot \|f_2\|_{H^1} \cdot \|f_3\|_{H^s}\}. \end{aligned} \tag{21}$$

The implied constant depends on s and the length of I by construction.

Case 2: $N_1 \sim N_2$

In this case, we will estimate:

$$K^2 := \sum_{\substack{N_1, N_2, N_3 \\ N_1 \geq N_2 \geq N_3; N_1 \sim N_2}} K_{N_1, N_2, N_3} = \sum_{\substack{N_1, N_2, N_3 \\ N_1 \geq N_2 \geq N_3; N_1 \sim N_2}} N_1^s \cdot \|P_{N_1} \psi_1 \cdot P_{N_2} \psi_2 \cdot P_{N_3} \psi_3\|_{L^2(I \times \Lambda)}.$$

By (16) and Proposition 2.1, this expression is:

$$\begin{aligned} & \lesssim_{|I|} \sum_{\substack{N_1, N_2, N_3 \\ N_1 \geq N_2 \geq N_3; N_1 \sim N_2}} N_1^s \cdot N_2 \cdot N_3 \cdot \max\left\{\frac{N_3}{N_1}, \frac{1}{N_2}\right\}^\delta \\ & \cdot \|P_{N_1} f_1\|_{L^2(\Lambda)} \cdot \|P_{N_2} f_2\|_{L^2(\Lambda)} \cdot \|P_{N_3} f_3\|_{L^2(\Lambda)}. \end{aligned} \tag{22}$$

Since $N_1 \sim N_2$ and $N_3 \geq 1$, this is:

$$\begin{aligned} & \lesssim \sum_{\substack{N_1, N_2, N_3 \\ N_1 \geq N_2 \geq N_3; N_1 \sim N_2}} N_1^s \cdot N_2 \cdot N_3 \cdot \left(\frac{N_3}{N_1}\right)^\delta \cdot \|P_{N_1} f_1\|_{L^2(\Lambda)} \cdot \|P_{N_2} f_2\|_{L^2(\Lambda)} \cdot \|P_{N_3} f_3\|_{L^2(\Lambda)} \\ & \lesssim \sum_{\substack{N_1, N_2, N_3 \\ N_1 \geq N_2 \geq N_3; N_1 \sim N_2}} \left(\frac{N_3}{N_1}\right)^\delta \cdot \|P_{N_1} f_1\|_{H^s(\Lambda)} \cdot \|P_{N_2} f_2\|_{H^1(\Lambda)} \cdot \|P_{N_3} f_3\|_{H^1(\Lambda)}. \end{aligned}$$

By using the Cauchy–Schwarz inequality in N_3 , we can bound this sum by:

$$\lesssim \sum_{\substack{N_1, N_2 \\ N_1 \sim N_2}} \|P_{N_1} f_1\|_{H^s(\Lambda)} \cdot \|P_{N_2} f_2\|_{H^1(\Lambda)} \cdot \|f_3\|_{H^1(\Lambda)}.$$

By using the Cauchy–Schwarz inequality in $N_1 \sim N_2$, this sum is:

$$\lesssim \|f_1\|_{H^s(\Lambda)} \cdot \|f_2\|_{H^1(\Lambda)} \cdot \|f_3\|_{H^1(\Lambda)}.$$

As in Case 1, we can replace $N_1^s \cdot N_2 \cdot N_3$ by $N_1 \cdot N_2^s \cdot N_3$ and $N_1 \cdot N_2 \cdot N_3^s$ in (22) and we can hence deduce that:

$$\begin{aligned} K^2 &= \sum_{\substack{N_1, N_2, N_3 \\ N_1 \geq N_2 \geq N_3; N_1 \sim N_2}} K_{N_1, N_2, N_3} \\ &\lesssim \min\{\|f_1\|_{H^s} \cdot \|f_2\|_{H^1} \cdot \|f_3\|_{H^1}, \|f_1\|_{H^1} \cdot \|f_2\|_{H^s} \cdot \|f_3\|_{H^1}, \|f_1\|_{H^1} \cdot \|f_2\|_{H^1} \cdot \|f_3\|_{H^s}\}. \end{aligned} \tag{23}$$

The implied constant depends on s and the length of I by construction.

The proposition now follows from the estimates (21) and (23). \square

4. An unconditional uniqueness result

In this section, we will use a version of the Quantum de Finetti Theorem in order to obtain an unconditional uniqueness result for the Gross–Pitaevskii hierarchy on \mathbb{T}^3 . The unconditional uniqueness result will hold in a class of density matrices with regularity $\alpha = 1$, and hence for higher regularities. In particular, we will prove an analogue of the non-periodic unconditional uniqueness result proved by T. Chen, Hainzl, Pavlović, and Seiringer [20].

4.1. The Weak Quantum de Finetti Theorem

Let \mathcal{H} be a separable Hilbert space. We denote by $\mathcal{H}_{sym}^k := \bigotimes_{sym}^k \mathcal{H}$ the associated bosonic k -particle space. More precisely, this is the space obtained from the quotient under the action of the symmetric group S_k on $\underbrace{\mathcal{H} \times \mathcal{H} \times \dots \times \mathcal{H}}_{k \text{ times}}$

by $\sigma \cdot (h_1, h_2, \dots, h_k) := (h_{\sigma(1)}, h_{\sigma(2)}, \dots, h_{\sigma(k)})$.

We will use the following variant of the Quantum de Finetti Theorem, which is called the *Weak Quantum de Finetti Theorem*:

Theorem 4.1 (Weak Quantum de Finetti [6,7,83]). *Suppose that the sequence $(\Gamma_N)_N$ satisfies the following assumptions for all $N \in \mathbb{N}$:*

- i) Γ_N is a bounded self-adjoint operator on \mathcal{H}_{sym}^N .
- ii) $\Gamma_N \geq 0$ as an operator on \mathcal{H}_{sym}^N .
- iii) $\text{Tr}_{\mathcal{H}_{sym}^N} \Gamma_N = 1$.

Moreover, suppose that for all $k \in \mathbb{N}$, the corresponding sequence of k -particle marginals $\gamma_N^{(k)} := \text{Tr}_{k+1, \dots, N} \Gamma_N$ converges to $\gamma^{(k)} \in \mathcal{H}_{sym}^k$ as $N \rightarrow \infty$ in the weak- $*$ topology of the trace class, i.e. $\gamma_N^{(k)} \rightharpoonup^* \gamma^{(k)}$ in the trace class on \mathcal{H}_{sym}^k .

Under these assumptions, there exists a unique Borel probability measure supported on the unit ball \mathcal{B} of \mathcal{H} , which is invariant under multiplication by complex numbers of modulus one such that, for all $k \in \mathbb{N}$:

$$\gamma^{(k)} = \int_{\mathcal{B}} (|\phi\rangle\langle\phi|^{\otimes k}) d\mu(\phi). \tag{24}$$

Theorem 4.1 is proved in [83, Theorem 2.2] and is based on equivalent results proved in [6] and [7], where μ is called a *Wigner measure*. The connection between Wigner measures and de Finetti measures is explained in [5, Section 6.4].

Remark 4.2. By the invariance of μ under multiplication by complex numbers of modulus one, we mean that for all $z \in \mathbb{C}$ with $|z| = 1$ and for all $\phi \in \mathcal{H}$, it is the case that $(z \cdot \mu)(\phi) := \mu(z \cdot \phi) = \mu(\phi)$.

Remark 4.3. There is also a Strong Quantum de Finetti Theorem. This is the original result obtained in the work of Hudson and Moody [77], and Størmer [109] in the context of C^* algebras. In the context of density matrices results, analogues of this result were also obtained in the work of Ammari and Nier [6,7], and Lewin, Nam, and Rougerie [83]. In the strong version, the density matrices $\gamma^{(k)} \in \mathcal{H}^k$ are not assumed to be weak- $*$ limits. Instead, they are assumed to have the property that $\text{Tr}_{\mathcal{H}^k} \gamma^{(k)} = 1$ and that they satisfy the admissibility property $\gamma^{(k)} = \text{Tr}_{k+1} \gamma^{(k+1)}$. In this case, it is possible to prove (24) with the measure μ supported on the unit sphere \mathcal{S} of \mathcal{H} . This strong version can be used to prove an unconditional uniqueness result with the corresponding assumptions, as was done in [20,76]. For our applications in Section 5, we will use the unconditional uniqueness result obtained from the Weak Quantum de Finetti Theorem in Theorem 4.6 below.

4.2. An unconditional uniqueness result

We henceforth consider the Gross–Pitaevskii hierarchy on $\Lambda = \mathbb{T}^3$:

$$\begin{cases} i \partial_t \gamma^{(k)} + (\Delta_{\vec{x}_k} - \Delta_{\vec{x}'_k}) \gamma^{(k)} = \lambda \cdot \sum_{j=1}^k B_{j,k+1}(\gamma^{(k+1)}) \\ \gamma^{(k)}|_{t=0} = \gamma_0^{(k)}. \end{cases} \tag{25}$$

We will assume that $\lambda \in \{1, -1\}$. The results of this section will apply to both the defocusing case $\lambda = 1$ and the focusing case $\lambda = -1$. In Section 5, we will restrict our attention to the defocusing case.

Following [20], we will use the spaces \mathfrak{H}^α :

Definition 4.4. Given $\alpha \geq 0$, \mathfrak{H}^α denotes the set of all sequences $(\gamma^{(k)})_k$ such that, for each k :

- i) $\gamma^{(k)} \in L^2_{\text{sym}}(\Lambda^k \times \Lambda^k)$ and $\gamma(\vec{x}_k, \vec{x}'_k) = \overline{\gamma^{(k)}(\vec{x}'_k; \vec{x}_k)}$ for all (\vec{x}_k, \vec{x}'_k) in $\Lambda^k \times \Lambda^k$.
- ii) $S^{(k,\alpha)} \gamma^{(k)} \in \mathcal{L}^1_k$.
- iii) There exists $M > 0$, independent of k , for which $\text{Tr}(|S^{(k,\alpha)} \gamma^{(k)}|) \leq M^{2k}$.

Variants of such spaces were used previously in the work of Erdős, Schlein, and Yau [53–57] and in related works. A different class of spaces, based on a Hilbert–Schmidt norm, was used in the work of Klainerman and Machedon [80].

We adopt the terminology from [20]:

Definition 4.5. Given $T > 0$, we say that $(\gamma^{(k)})_k = (\gamma^{(k)}(t))_k$ is a mild solution to the Gross–Pitaevskii hierarchy (25) in $L^\infty_{t \in [0, T]} \mathfrak{H}^\alpha$ if, for all $k \in \mathbb{N}$:

- i) $(\gamma^{(k)}(0))_k = (\gamma_0^{(k)})_k \in \mathfrak{H}^\alpha$.
- ii) For all $t \in [0, T]$, the following integral equation is valid:

$$\gamma^{(k)}(t) = \mathcal{U}^{(k)}(t) - i\lambda \int_0^t \mathcal{U}^{(k)}(t-s) B_{k+1} \gamma^{(k+1)}(s) ds.$$

iii)

$$\sup_{t \in [0, T]} \text{Tr}(|S^{(k,\alpha)} \gamma^{(k)}(t)|) \leq M^{2k}$$

for some $M > 0$ which is independent of k and t .

For point ii), we recall the definition of the full collision operator B_{k+1} in (12).

In what follows, let S_t denote the flow map of the cubic nonlinear Schrödinger equation:

$$\begin{cases} iu_t + \Delta u = \lambda \cdot |u|^2 u \\ u|_{t=0} = \phi. \end{cases} \tag{26}$$

More precisely, we let $S_t(\phi)(x) := u(x, t)$. We note that the Cauchy theory of this problem is well-understood [16] and hence the map S_t is well-defined.

With these definitions in mind, we will prove:

Theorem 4.6 (An unconditional uniqueness result). *Let $(\gamma_0^{(k)})_k \in \mathfrak{S}^1$. Suppose that $(\gamma^{(k)}(t))_k$ is a mild solution to the Gross–Pitaevskii hierarchy (25) in $L^\infty_{t \in [0, T]} \mathfrak{S}^1$ with initial data $(\gamma^{(k)}(0))_k = (\gamma_0^{(k)})_k$ such that, for each $t \in [0, T]$, $\gamma^{(k)}(t)$ satisfies the assumptions of Theorem 4.1 with $\mathcal{H} = L^2(\Lambda)$. More precisely, we assume that there exist $\Gamma_{N,t} \in L^2_{\text{sym}}(\Lambda^N \times \Lambda^N)$ as in Theorem 4.1 such that $\gamma_{N,t}^{(k)} := \text{Tr}_{k+1, \dots, N} \Gamma_{N,t}$ satisfies $\gamma_{N,t}^{(k)} \xrightarrow{*} \gamma^{(k)}(t)$ as $N \rightarrow \infty$ in \mathcal{L}^1_k . Then, the conclusion is that $(\gamma^{(k)}(t))_k$ is uniquely determined by the initial data $(\gamma_0^{(k)})_k$.*

Furthermore, suppose that $(\gamma_0^{(k)})_k$ satisfies, for all $k \in \mathbb{N}$:

$$\gamma_0^{(k)} = \int_{\mathcal{B}} (|\phi\rangle\langle\phi|^{\otimes k}) d\mu(\phi) \tag{27}$$

for μ a Borel probability measure which is invariant under multiplication by complex numbers of modulus one and which is supported on the unit ball \mathcal{B} of $L^2(\Lambda)$. Then, for all $k \in \mathbb{N}$ and for all $t \in [0, T]$, it is the case that:

$$\gamma^{(k)}(t) = \int_{\mathcal{B}} (|S_t(\phi)\rangle\langle S_t(\phi)|^{\otimes k}) d\mu(\phi). \tag{28}$$

The remainder of this section is devoted to the proof of Theorem 4.6. Our method will follow the combinatorial graph expansion from [20]. The main difference from the non-periodic setting is that we will estimate the terms coming from the distinguished vertices by using Proposition 3.1. More precisely, we will have to give a different proof of the L^2 estimate (41) and of the H^1 estimate (42) below. Namely, the proof of these estimates in the non-periodic setting [20, Lemma 7.1] relies on the use of Strichartz estimates on \mathbb{R}^3 , which are not available on \mathbb{T}^3 . We overcome this difficulty by applying Proposition 3.1. In addition, in bounding the trilinear term $\| |\phi|^2 \phi \|_{L^2_x}$, we need to recall the Sobolev embedding on \mathbb{T}^3 in (8). For completeness of the exposition, we will recall the main ideas of the graph expansion and how it is used to prove unconditional uniqueness results. For the full details of the construction and the methods, we refer the interested reader to [20].

Let us suppose that $(\gamma_1^{(k)}(t))_k$ and $(\gamma_2^{(k)}(t))_k$ are solutions to (25) with initial data $(\gamma_0^{(k)})_k$ which satisfy the assumptions of the theorem. We need to show that $\gamma^{(k)} := \gamma_1^{(k)} - \gamma_2^{(k)}$ is identically zero on $[0, T]$ for all $k \in \mathbb{N}$. In particular, it suffices to show that:

$$\text{Tr}(|\gamma^{(k)}(t)|) = 0 \tag{29}$$

for all $k \in \mathbb{N}$ and for all $t \in [0, T]$.

By Theorem 4.1, there exist unique time-dependent Borel probability measures $\mu_t^{(1)}$ and $\mu_t^{(2)}$, which are invariant under multiplication by complex numbers of modulus one and which are supported on the unit ball \mathcal{B} of $L^2(\Lambda)$ such that for all $k \in \mathbb{N}$ and $t \in [0, T]$:

$$\gamma_j^{(k)}(t) = \int_{\mathcal{B}} (|\phi\rangle\langle\phi|^{\otimes k}) d\mu_t^{(j)}(\phi)$$

for $j = 1, 2$. Hence, we can write:

$$\gamma^{(k)}(t) = \int_{\mathcal{B}} (|\phi\rangle\langle\phi|^{\otimes k}) d\mu_t(\phi), \tag{30}$$

where $\mu_t := \mu_t^1 - \mu_t^2$ is a signed Borel measure on \mathcal{B} . The assumption that $(\gamma_1^{(k)}(t))_k$ and $(\gamma_2^{(k)}(t))_k \in L^\infty_{t \in [0, T]} \mathfrak{H}^1$ implies that for $j = 1, 2$, there exists $M_j > 0$, such that for all $k \in \mathbb{N}$ and for all $t \in [0, T]$:

$$\text{Tr}(|\gamma_j^{(k)}(t)|) = \int_{\mathcal{B}} \|\phi\|_H^{2k} d\mu_t^{(j)}(\phi) \leq M_j^k. \tag{31}$$

Let us note that, by the Duhamel expansion:

$$\gamma^{(k)}(t) = \int_0^t \int_0^{t_1} \cdots \int_0^{t_{r-1}} (-i\lambda)^r \cdot J^k(\underline{t}_r) d\underline{t}_r, \tag{32}$$

where $\underline{t}_r := (t_1, t_2, \dots, t_r)$ and:

$$J^k(\underline{t}_r) := \mathcal{U}^{(k)}(t - t_1) B_{k+1} \mathcal{U}^{(k+1)}(t_1 - t_2) B_{k+2} \cdots \mathcal{U}^{(k+r-1)}(t_{r-1} - t_r) B_{k+r} \gamma^{(k+r)}(t_r). \tag{33}$$

By using the same reductions as in [20, Section 4.1], based on the Erdős–Schlein–Yau combinatorial method in the boardgame form [80], it is possible to write:

$$\gamma^{(k)}(t) = (-i\lambda)^r \cdot \sum_{\sigma \in \mathcal{N}_{k,r} D(\sigma,t)} \int_{\mathcal{B}} J^k(\phi; \sigma; t, t_1, \dots, t_r) d\mu_t(\phi) d\underline{t}_r. \tag{34}$$

Here $\mathcal{N}_{k,r}$ denotes the set of all monotonically increasing mappings $\sigma : \{k + 1, k + 2, \dots, k + r\} \rightarrow \{1, 2, \dots, k + r - 1\}$ such that $\sigma(j) < j$ for all $j \in \{k + 1, k_2, \dots, k + r\}$ and $D(\sigma, t)$ denotes a union of simplices contained in $[0, t]^r$, whose interiors are mutually disjoint. Moreover, for $\sigma \in \mathcal{N}_{k,r}$ and $\phi \in \mathcal{B}$, we write:

$$J^k(\phi; \sigma; t, t_1, \dots, t_r; \vec{x}_k; \vec{x}'_k) := \mathcal{U}^{(k)}(t - t_1) B_{\sigma(k+1), k+1} \mathcal{U}^{(k+1)}(t_1 - t_2) \cdots \mathcal{U}^{(k+r-1)}(t_{r-1} - t_r) B_{\sigma(k+r), k+r} (|\phi\rangle \langle \phi|^{\otimes(k+r)}) (\vec{x}_k; \vec{x}'_k). \tag{35}$$

In (35), we use the convention as in [20] that the time variable t_ℓ is attached to the interaction operator $B_{\sigma(k+\ell), k+\ell}$.

4.2.1. A reformulation of the Duhamel expansion following [20]

In [20, Eq. (4.23)], it is noted that $J^k(\phi; \sigma; t, t_1, \dots, t_r; \vec{x}_k; \vec{x}'_k)$ has the following product structure:

$$J^k(\phi; \sigma; t, t_1, \dots, t_r; \vec{x}_k; \vec{x}'_k) = \prod_{j=1}^k J_j^1(\phi; \sigma_j; t, t_{\ell_{j,1}}, \dots, t_{\ell_{j,m_j}}; x_j; x'_j). \tag{36}$$

The interaction operators which appear in J_j^1 are labelled “internally” with the maps σ_j , $j = 1, \dots, k$. Relative to J_j^1 , σ_j is in upper echelon form in the sense of [80], i.e. it is monotonically increasing. We note that the construction of [20] carries over to the periodic setting without any changes. In particular:

$$\text{Tr}(|\gamma^{(k)}(t)|) \leq C^{k+r} \cdot \sum_{\ell=1}^2 \sup_{\sigma \in \mathcal{N}_{k,r} [0,t]^r} \int_{\mathcal{B}} \prod_{j=1}^k \text{Tr}(|J_j^1(\phi; \sigma_j; t, t_{\ell_{j,1}}, \dots, t_{\ell_{j,m_j}})|) d\mu_t^{(\ell)}(\phi) d\underline{t}_r. \tag{37}$$

4.2.2. The precise definition of the graph expansion (after [20])

Let us now recall the precise definition of the graph structure following [20, Section 5]. One assigns to the expression (35) and the product structure (36) a union of k disjoint binary tree graphs. As is noted in [20], these graphs have already appeared as substructures associated to the graphs in the more involved analysis of [57]. More precisely, one assigns:

- 1) An *internal vertex* v_ℓ associated to each collision operator $B_{\sigma(k+\ell), k+\ell}$, for $\ell = 1, 2, \dots, r$. Hence, the time variable t_ℓ can be thought of as being attached to the vertex v_ℓ .
- 2) A *root vertex* w_j to each factor J_j^1 , for $j = 1, 2, \dots, k$.

3) A leaf vertex u_i to the factor $(|\phi\rangle\langle\phi|)(x_i; x'_i)$ occurring in $|\phi\rangle\langle\phi|^{\otimes(k+r)}(\vec{x}_{k+r}, \vec{x}'_{k+r})$. Here $i = 1, 2, \dots, k + r$.

The equivalence relation “ \sim ” of connectivity between vertices is given by:

- 1) Suppose that $j \in \{1, 2, \dots, k\}$ is given. If ℓ is the *smallest element* of $\{1, 2, \dots, r\}$ such that $\sigma(k + \ell) = j$, then the root vertex w_j is connected to the internal vertex v_ℓ , i.e. these two vertices are equivalent. We say that w_j is the *parent vertex* of v_ℓ and that v_ℓ is the *child vertex* of w_j .
- 2) If there exists no internal vertex connected to w_j as above, then the root vertex w_j is connected to the leaf vertex u_j . We say that w_j is the *parent vertex* of u_j and u_j is the *child vertex* of w_j .
- 3) Suppose that $\ell \in \{1, 2, \dots, r\}$. If there exists $\ell' \in \{1, 2, \dots, r\}$ with $\ell' > \ell$ such that $k + \ell = \sigma(k + \ell')$ or $\sigma(k + \ell) = \sigma(k + \ell')$, then we say that $v_\ell \sim v_{\ell'}$. In this case, it is said that v_ℓ is a *parent vertex* of $v_{\ell'}$ and that $v_{\ell'}$ is a *child vertex* of v_ℓ . The two child vertices of v_ℓ will be denoted by $v_{\kappa_-(\ell)}$ and $v_{\kappa_+(\ell)}$, where by convention one takes $\kappa_-(\ell) < \kappa_+(\ell)$.
- 4) If there exists no internal vertex $v_{\ell'}$ with $\ell' > \ell$ and $k + \ell = \sigma(k + \ell')$ as above, then v_ℓ is connected to the leaf vertex u_ℓ . In this case, we say that v_ℓ is the *parent vertex* of u_ℓ and that u_ℓ is a *child vertex* of v_ℓ . Furthermore, if there exists no internal vertex $v_{\ell'}$ such that $\sigma(k + \ell) = \sigma(k + \ell')$, then v_ℓ is connected to the leaf vertex $u_{\sigma(k+\ell)}$. In this case, we say that v_ℓ is the *parent vertex* of $u_{\sigma(k+\ell)}$ and that $u_{\sigma(k+\ell)}$ is a *child vertex* of v_ℓ .

In this way, one obtains a graph which is a disjoint union of k binary trees. These binary trees are denoted by τ_j , for $j = 1, 2, \dots, k$. The root of τ_j is given by w_j . We note that the σ_j in (36) can be viewed as $\sigma|_{\tau_j}$. In other words, this is the internal labeling of the collision operators that respects the ordering of the vertices of τ_j . Moreover, τ_j has m_j internal vertices corresponding in this way to time labels $t_{\ell_{j,1}}, t_{\ell_{j,2}}, \dots, t_{\ell_{j,m_j}}$.

The internal vertex v_r is called the *distinguished vertex*. The other internal vertices are called *regular*. The two children vertices corresponding to v_r are called the *distinguished leaves*. All the other leaves are called *regular*. The tree τ_j is called *distinguished* if $v_r \in \tau_j$ and otherwise it is called *regular*. For more details on this construction, we refer the reader to [20, Section 5].

4.2.3. A detailed analysis of the distinguished tree graph

Let us now recall the analysis of the distinguished tree from [20, Section 6]. As above, the analysis transfers to the periodic setting. For completeness, we will summarize the notation and the main ideas.

Let τ_j be a distinguished tree. This tree contains m_j internal vertices $(v_{\ell_{j,a}})_{a=1}^{m_j}$ and $m_j + 1$ leaf vertices $(u_{j,s})_{s=1}^{m_j+1}$. The internal vertices are enumerated by $a \in \{1, 2, \dots, m_j\}$, and a corresponds to the label in the collision operator $B_{\sigma_j(a+1),a+1}$. It is advantageous to continue this labeling and to denote the leaf vertices by $a \in \{m_j + 1, \dots, 2m_j + 1\}$.

By the semigroup property of $\mathcal{U}^{(k)}(t)$, it is possible to show that the term in (36) corresponding to the distinguished tree τ_j equals:

$$\begin{aligned} & J_j^1(\phi; \sigma_j; t, t_{\ell_{j,1}}, \dots, t_{\ell_{j,m_j}}; x_j; x'_j) \\ &= \mathcal{U}^{(1)}(t - t_{\ell_{j,1}}) B_{1,2} \mathcal{U}^{(2)}(t_{\ell_{j,1}} - t_{\ell_{j,2}}) B_{\sigma_j(3),3} \cdots B_{\sigma_j(a),a} \mathcal{U}^{(a)}(t_{\ell_{j,a-1}} - t_{\ell_{j,a}}) B_{\sigma_j(a+1),a+1} \cdots \\ & \quad \mathcal{U}^{(m_j)}(t_{\ell_{j,m_j-1}} - t_{\ell_{j,m_j}}) B_{\sigma_j(m_j+1),m_j+1} (|\phi\rangle\langle\phi|)^{\otimes(m_j+1)}. \end{aligned}$$

We now want to bound the integral:

$$\int_{[0,T]^{m_j-1}} \text{Tr}(|J_j^1(\phi; \sigma_j; t, t_{\ell_{j,1}}, \dots, t_{\ell_{j,m_j}})|) dt_{\ell_{j,m_j-1}} dt_{\ell_{j,m_j-1}} \cdots dt_{\ell_{j,1}}. \tag{38}$$

In order to estimate the expression in (38), one associates to the vertex labeled by a a kernel θ_a . It can be seen that, at a regular leaf vertex $\theta_a(x; x') := \phi(x) \cdot \overline{\phi(x')}$. Moreover, at the distinguished vertex, $\theta_{m_j}(x; x') := \tilde{\psi}(x) \cdot \overline{\phi(x')} - \phi(x) \cdot \overline{\tilde{\psi}(x')}$, for $\tilde{\psi} := |\phi|^2 \phi$. Using an inductive formula, it is shown in [20, Lemma 6.1] that for every $a \in \{1, 2, \dots, m_j\}$, the kernel θ_a can be written as:

$$\theta_a(x; x') = \sum_{\beta_a} c_{\beta_a}^a \chi_{\beta_a}^a(x) \overline{\psi_{\beta_a}^a(x')} \tag{39}$$

for at most $2^{m_j - \alpha + 1}$ coefficients $c_{\beta_a}^a \in \{-1, 1\}$. By using additional product identities as in [20, Section 6.4], it is possible to inductively obtain the factors $\chi_{\beta_a}^a$ and $\psi_{\beta_a}^a$. We will omit the details. Finally, without loss of generality, it is possible to assume that $\psi_{\beta_1}^1$ and its iterates, i.e. the terms of the form $\psi_{\beta_{\kappa_+^q(1)}}^{\kappa_+^q(1)}$ are functions of $\tilde{\psi} = |\phi|^2 \phi$. We say that these factors are *distinguished*. Here, $\kappa_+^q(1)$ denotes the q -th iterate of κ_+ applied to the vertex with label 1.

4.2.4. Recursive L^2 and H^1 bounds for the distinguished tree

We now summarize the main ideas from [20, Section 7] in which one inductively uses L^2 and H^1 bounds in order to estimate the expression in (38). As we will see, at this step, a modification of the argument will be required to deal with the periodic case.

By the definition of the kernel θ_1 and by (39), the expression in (38) is:

$$\leq \sum_{\beta_1} \int_{[0, T]^{m_j - 1}} \|\psi_{\beta_1}^1\|_{L_x^2} \cdot \|\chi_{\beta_1}^1\|_{L_x^2} dt_{\ell_{j, m_j - 1}} dt_{\ell_{j, m_j - 2}} \cdots dt_{\ell_{j, 1}}. \tag{40}$$

In what follows, one denotes by $\tau_{j, a}$ the subtree of τ_j which is rooted at the vertex a . Given $\tau_{j, a}$, d_a denotes the number of non-leaf (i.e. internal and root) vertices in $\tau_{j, a}$. Finally, one defines:

$$\int_{[0, T]^{d_a}} (\cdots) \prod_{a' \in \tau_{j, a}} dt_{\ell_{j, a'}}$$

to denote the integral with respect to the time variables which correspond to the non-leaf vertices in $\tau_{j, a}$. If $\tau_{j, a} = \emptyset$, we just evaluate the integrand. If we emphasize that $a' \neq m_j$, then we will not be integrating with respect to $t_{\ell_{j, m_j}}$. The following estimates hold for all $a \in \{1, 2, \dots, m_j - 1\}$:

$$\begin{aligned} & \int_{[0, T]^{d_a}} \|\psi_{\beta_a}^a\|_{L_x^2} \cdot \|\chi_{\beta_a}^a\|_{H_x^1} \prod_{a' \in \tau_{j, a}} dt_{\ell_{j, a'}} \\ & \leq CT^{\frac{1}{2}} \cdot \left(\int_{[0, T]^{d_{\kappa_-(a)}}} \|\psi_{\beta_{\kappa_-(a)}}^{\kappa_-(a)}\|_{H_x^1} \cdot \|\chi_{\beta_{\kappa_-(a)}}^{\kappa_-(a)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa_-(a)}} dt_{\ell_{j, a'}} \right) \\ & \quad \cdot \left(\int_{[0, T]^{d_{\kappa_+(a)}}} \|\psi_{\beta_{\kappa_+(a)}}^{\kappa_+(a)}\|_{L_x^2} \cdot \|\chi_{\beta_{\kappa_+(a)}}^{\kappa_+(a)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa_+(a)}} dt_{\ell_{j, a'}} \right) \end{aligned} \tag{41}$$

and

$$\begin{aligned} & \int_{[0, T]^{d_a}} \|\psi_{\beta_a}^a\|_{H_x^1} \cdot \|\chi_{\beta_a}^a\|_{H_x^1} \prod_{a' \in \tau_{j, a}} dt_{\ell_{j, a'}} \\ & \leq CT^{\frac{1}{2}} \cdot \left(\int_{[0, T]^{d_{\kappa_-(a)}}} \|\psi_{\beta_{\kappa_-(a)}}^{\kappa_-(a)}\|_{H_x^1} \cdot \|\chi_{\beta_{\kappa_-(a)}}^{\kappa_-(a)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa_-(a)}} dt_{\ell_{j, a'}} \right) \\ & \quad \cdot \left(\int_{[0, T]^{d_{\kappa_+(a)}}} \|\psi_{\beta_{\kappa_+(a)}}^{\kappa_+(a)}\|_{H_x^1} \cdot \|\chi_{\beta_{\kappa_+(a)}}^{\kappa_+(a)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa_+(a)}} dt_{\ell_{j, a'}} \right). \end{aligned} \tag{42}$$

The estimate in (41) is referred to as the *bound on the L^2 level*, and the estimate in (42) is referred to as the *bound on the H^1 level*. The L^2 estimate (41) is proved by using Proposition 3.1 when $s = 0$. In particular, we use:

$$\|(e^{it\Delta} f_1)(x) \cdot \overline{(e^{it\Delta} f_2)(x)} \cdot (e^{it\Delta} f_3(x))\|_{L_{[0, T]}^2 L_x^2} \leq C \|f_1\|_{H_x^1} \cdot \|f_2\|_{H_x^1} \cdot \|f_3\|_{L_x^2}. \tag{43}$$

Here, we recall that, by assumption, $T \in [0, 1]$. The estimate (43) was also used in the non-periodic setting [20, Section 7]. In this case, the trilinear estimate was deduced from Strichartz estimates on \mathbb{R}^3 . The rest of the proof

of (41) proceeds as in [20, Section 7]. We will omit the details. The estimate (42) is proved by using Proposition 3.1 when $s = 1$.

4.2.5. Estimating the factors corresponding to distinguished and regular trees

We will now record the bounds for the contributions from the distinguished tree and from the regular trees in (37).

For the distinguished tree, it is necessary to estimate the expression in (38). By applying (41), it is possible to bound the right-hand side of (40) by:

$$\sum_{\beta_{\kappa-(1)}, \beta_{\kappa+(1)}} CT^{\frac{1}{2}} \cdot \left(\int_{[0, T]^{d_{\kappa-(1)}}} \|\psi_{\beta_{\kappa-(1)}}^{\kappa-(1)}\|_{H_x^1} \cdot \|\chi_{\beta_{\kappa-(1)}}^{\kappa-(1)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa-(1)}} dt_{a'} \right) \cdot \left(\int_{[0, T]^{d_{\kappa+(1)}-1}} \|\psi_{\beta_{\kappa+(1)}}^{\kappa+(1)}\|_{L_x^2} \cdot \|\chi_{\beta_{\kappa+(1)}}^{\kappa+(1)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa+(1)}, a' \neq m_j} dt_{a'} \right). \tag{44}$$

The key point is that the distinguished factor $\psi_{\beta_{\kappa+(1)}}^{\kappa+(1)}$ is estimated in the L^2 norm. This is possible to do because in Proposition 3.1, when $s = 0$, we can choose which factor we want to estimate in L^2 . By iterating (42) as in [20, Section 8], it follows that:

$$\int_{[0, T]^{d_{\kappa-(1)}}} \|\psi_{\beta_{\kappa-(1)}}^{\kappa-(1)}\|_{H_x^1} \cdot \|\chi_{\beta_{\kappa-(1)}}^{\kappa-(1)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa-(1)}} dt_{a'} \leq C_1^{d_{\kappa-(1)}} \cdot T^{\frac{d_{\kappa-(1)}}{2}} \cdot \|\phi\|_{H_x^1}^{2b_{\kappa+(1)}}. \tag{45}$$

Here, b_a denotes the number of regular, i.e. non-distinguished leaf vertices of the subtree $\tau_{j, a}$, which is rooted at a . $C_1 > 0$ denotes the constant in (42).

Moreover, by combining (41) and (42) as in [20, Section 8], it follows that, for some constant $C_1 > 0$:

$$\int_{[0, T]^{d_{\kappa+(1)}-1}} \|\psi_{\beta_{\kappa+(1)}}^{\kappa+(1)}\|_{L_x^2} \cdot \|\chi_{\beta_{\kappa+(1)}}^{\kappa+(1)}\|_{H_x^1} \prod_{a' \in \tau_{j, \kappa+(1)}, a' \neq m_j} dt_{a'} \leq C_1^{d_{\kappa+(1)}-1} \cdot T^{\frac{d_{\kappa+(1)}-1}{2}} \cdot \|\phi\|_{H_x^1}^{2b_{\kappa+(1)}+1} \cdot \|\phi\|_{L_x^2}^2. \tag{46}$$

By applying the Sobolev embedding result in (8), it follows that the right-hand side of (46) is:

$$\leq C_2^{d_{\kappa+(1)}-1} \cdot T^{\frac{d_{\kappa+(1)}-1}{2}} \cdot \|\phi\|_{H_x^1}^{2b_{\kappa+(1)}+4}, \tag{47}$$

for some constant $C_2 > 0$.

Using the expansion (39), the estimates (45), (47), and counting the number of terms, it follows that, for some constant $C_3 > 0$:

$$\int_{[0, T]^{m_j-1}} \text{Tr}(|J_j^1(\phi; \sigma_j; t, t_{\ell_{j,1}}, \dots, t_{\ell_{j,m_j}})|) dt_{\ell_{j,m_j-1}} \cdots dt_{\ell_{j,1}} \leq C_3^{m_j} \cdot T^{\frac{m_j-1}{2}} \cdot \|\phi\|_{H_x^1}^{2m_j+2}. \tag{48}$$

This is the bound that one uses for the distinguished tree τ_j .

Suppose now that τ_j is a regular tree. In this case, we can iteratively use the estimate (42) and obtain the bound as in [20, Proposition 8.2]:

$$\int_{[0, T]^{m_j}} \text{Tr}(|J_j^1(\phi; \sigma_j; t, t_{\ell_{j,1}}, \dots, t_{\ell_{j,m_j}})|) dt_{\ell_{j,m_j}} dt_{\ell_{j,m_j-1}} \cdots dt_{\ell_{j,1}} \leq C_4^{m_j} \cdot T^{\frac{m_j}{2}} \cdot \|\phi\|_{H_x^1}^{2m_j+2} \tag{49}$$

for some constant $C_4 > 0$. This is the bound that we use for the regular trees.

4.2.6. *The proof of Theorem 4.6*

By using the previous results, we will now give the proof of Theorem 4.6:

Proof. Combining (37), (48), and (49), the result of Theorem 4.6 is now proved in the same way as in [20, Section 8]. □

5. The derivation of the defocusing cubic NLS on \mathbb{T}^3

We will now apply the above unconditional uniqueness theorem in order to obtain a derivation of the defocusing cubic NLS on \mathbb{T}^3 , following the program of Elgart, Erdős, Schlein, and Yau [50] and Erdős, Schlein, and Yau [53–57]. Let us note that the results of Section 4 apply both in the focusing and in the defocusing regime. The results of this section will apply only in the *defocusing* regime. More precisely, in what follows, we will set $\lambda = 1$ in (25) and (26). As before, S_t will denote the corresponding flow map. The main point that we want to address in the analysis of this section is that, for appropriately chosen initial data, the solution to the Gross–Pitaevskii hierarchy obtained as a limit in [50] satisfies the assumptions of Theorem 4.6. In this context, it is possible to apply the unconditional uniqueness result from Theorem 4.6 and uniquely characterize the limit point $\Gamma_{\infty,t}$. This will be an important idea in the proof of the main result, which is stated as Theorem 5.4 below.

For completeness, let us now briefly summarize the main results of [50]. Suppose that $V : \mathbb{R}^3 \rightarrow \mathbb{R}$ is a smooth, non-negative function with compact support such that $\int_{\mathbb{R}^3} V(x) dx = 1$. The bosons in \mathbb{T}^3 interact via a two-body potential which is given by:

$$V_N(x) := N^{3\beta} V(N^\beta x),$$

for $\beta > 0$ a parameter. The N -body Hamiltonian for the N weakly coupled bosons is given by:

$$H_N := - \sum_{j=1}^N \Delta_j + \frac{1}{N} \sum_{\ell < j}^N V_N(x_\ell - x_j). \tag{50}$$

In [50], it is assumed that:

$$\beta \in \left(0, \frac{3}{5}\right). \tag{51}$$

In what follows, we will make the same assumption on β . H_N acts on a dense subset of $L^2_{sym}(\Lambda^N \times \Lambda^N)$. We remark that V_N is originally a function on \mathbb{R}^3 . Nevertheless, for N large, V_N is supported in a small neighborhood of 0. Hence, it can be extended to a periodic function on \mathbb{R}^3 . By this procedure, it can be thought of as a function on \mathbb{T}^3 . We will use this convention in our paper.

Let us fix $\Psi_{N,0} \in L^2_{sym}(\Lambda^N)$. Let $\Psi_{N,t}$ be the solution of the N -body Schrödinger equation associated to H_N with initial data $\Psi_{N,0}$ given in (4). We recall from (5) that if we take $\gamma_{N,t}^{(k)} := \text{Tr}_{k+1, \dots, N} |\Psi_{N,t}\rangle \langle \Psi_{N,t}|$, the sequence $(\gamma_{N,t}^{(k)})_k$ solves the BBGKY hierarchy (6).

In [50,53], the authors define for $\Gamma \in \bigoplus_{k \in \mathbb{N}} L^2_{sym}(\Lambda^k \times \Lambda^k)$ the following quantities:

$$\|\Gamma\|_{H_-} := \sum_{k=1}^{+\infty} 2^{-k} \cdot \|\gamma^{(k)}\|_{L^2(\Lambda^k \times \Lambda^k)}^2 \tag{52}$$

$$\|\Gamma\|_{H_+} := \sup_{k \geq 1} 2^k \cdot \|\gamma^{(k)}\|_{L^2(\Lambda^k \times \Lambda^k)}^2. \tag{53}$$

One then defines H_- to be the set of $\Gamma \in \bigoplus_{k \in \mathbb{N}} L^2_{sym}(\Lambda^k \times \Lambda^k)$ for which $\|\Gamma\|_{H_-}$ is finite. The space H_+ is defined as $\{\Gamma \in \bigoplus_{k \in \mathbb{N}} L^2(\Lambda^k \times \Lambda^k), \lim_{k \rightarrow \infty} 2^k \cdot \|\gamma^{(k)}\|_{L^2(\Lambda^k \times \Lambda^k)}^2 = 0\}$, with the norm given by (53). H_- and H_+ are then Banach spaces. Furthermore:

$$(H_+)^* = H_-$$

with respect to the pairing:

$$(\Gamma_1, \Gamma_2) \mapsto \sum_{k=1}^{+\infty} \langle \gamma_1^{(k)}, \gamma_2^{(k)} \rangle_{L^2(\Lambda^k \times \Lambda^k)},$$

whenever $\Gamma_1 = (\gamma_1^{(k)})_k$ and $\Gamma_2 = (\gamma_2^{(k)})_k$ in $\bigoplus_{k \in \mathbb{N}} L^2_{sym}(\Lambda^k \times \Lambda^k)$.

Given $T > 0$, $C([0, T], H_-)$ denotes the space of functions of $t \in [0, T]$ which take values in H_- , and which are continuous with respect to the weak- $*$ topology on H_- . The space H_+ is separable, and hence it is possible to find a countable dense subset in the unit ball of H_+ , which is denoted by $\{\mathcal{J}_{i,\ell}, \ell \in \mathbb{N}\}$. The metric ϱ on H_- is defined as:

$$\begin{aligned} \varrho(\Gamma_1, \Gamma_2) &:= \sum_{\ell=1}^{+\infty} 2^{-\ell} \cdot \left| \sum_{k=1}^{+\infty} \int_{\Lambda^k \times \Lambda^k} \overline{\mathcal{J}_{i,\ell}^{(k)}(\vec{x}_k; \vec{x}'_k)} \cdot [\gamma_1^{(k)}(\vec{x}_k; \vec{x}'_k) - \gamma_2^{(k)}(\vec{x}_k; \vec{x}'_k)] d\vec{x}_k d\vec{x}'_k \right|, \end{aligned} \tag{54}$$

whenever $\Gamma_1 = (\gamma_1^{(k)})_k, \Gamma_2 = (\gamma_2^{(k)})_k$ in H_- .

The metric $\widehat{\varrho}$ on $C([0, T], H_-)$ is defined as:

$$\widehat{\varrho}(\Gamma_1, \Gamma_2) := \sup_{t \in [0, T]} \varrho(\Gamma_1(t), \Gamma_2(t)),$$

whenever $\Gamma_1, \Gamma_2 \in C([0, T], H_-)$. For more details on these spaces, we refer the reader to [50,53].

The main result of [50] is the following:

Theorem 1 in [50]. *Given $N \in \mathbb{N}$, let H_N be as in (50). Suppose that $\gamma_{N,0} \in L^2_{sym}(\Lambda^N \times \Lambda^N)$ is such that $\gamma_{N,0}$ is self-adjoint, $\gamma_{N,0} \geq 0$ as an operator on $L^2(\Lambda^N)$ and $\text{Tr } \gamma_{N,0} = 1$. Furthermore, suppose that there exists a constant $C_1 > 0$ such that for all $k, N \in \mathbb{N}$:*

$$\text{Tr } H_N^k \gamma_{N,0} \leq C_1^k N^k. \tag{55}$$

Let us fix $T > 0$ and let $\Gamma_{N,t} = (\gamma_{N,t}^{(k)})_k$ be the solution to the BBGKY hierarchy (6) with initial data $\Gamma_{N,0} = (\gamma_{N,0}^{(k)})_k$. Then:

- i) The sequence $(\Gamma_{N,t})_N$ is relatively compact in $C([0, T], H_-)$ with respect to the metric $\widehat{\varrho}$.
- ii) If $\Gamma_{\infty,t} = (\gamma_{\infty,t}^{(k)})_k$ is a limit point of $(\Gamma_{N,t})_N$ with respect to the metric $\widehat{\varrho}$, then there exists a constant $C_2 > 0$ such that:

$$\text{Tr} |S^{(k,1)} \gamma_{\infty,t}^{(k)}| \leq C_2^k,$$

for all $k \in \mathbb{N}$ and for all $t \in [0, T]$.

- iii) $\text{Tr } \gamma_{\infty,t}^{(k)} = 1$ for all $k \in \mathbb{N}$ and for all $t \in [0, T]$.
- iv) $\Gamma_{\infty,t}$ solves the Gross–Pitaevskii hierarchy (25) with $\lambda = 1$.

Remark 5.1. In the assumption (55), we note that the operator H_N acts only in the first N components.

Remark 5.2. $\Gamma_{\infty,t}$ constructed in [50, Theorem 1] is a mild solution of the Gross–Pitaevskii hierarchy in the sense of Definition 4.5. The precise argument showing that the solution of type as is constructed in [50] is a mild solution is given in [54, Eq. (8.31)].

Remark 5.3. The fact that we can take $\lambda = 1$ in (25) follows from the assumption that $\int_{\mathbb{R}^3} V(x) dx = 1$.

In the discussion that follows, let us fix $\phi \in H^1(\Lambda)$ with $\|\phi\|_{L^2(\Lambda)} = 1$. By properly choosing the initial data, we will show that Theorem 4.6 implies that $\Gamma_{\infty,t}$ is uniquely determined and that $\gamma_{\infty,t}^{(k)} = |S_t(\phi)\rangle\langle S_t(\phi)|^{\otimes k}$. Let us observe that, by the assumption (55) of [50, Theorem 1], it is not possible to directly take factorized initial

data $\gamma_{N,0} = |\phi^{\otimes N}\rangle\langle\phi^{\otimes N}| = |\phi\rangle\langle\phi|^{\otimes N}$. Namely, then (55) only holds when $k = 1$. In order to obtain the full range of k , it is necessary to perform an additional approximation argument which is outlined in Subsection 5.1 below.

Let us consider a sequence $(\Psi_N)_N \in \bigoplus_{N \in \mathbb{N}} L^2_{sym}(\Lambda^N)$ which satisfies the following assumptions:

1) *Bounded energy per particle:*

$$\sup_N \frac{1}{N} \langle H_N \Psi_N, \Psi_N \rangle_{L^2(\Lambda^N)} < +\infty. \tag{56}$$

Here, H_N is the N -body Hamiltonian defined in (50).

2) *Asymptotic factorization:*

$$\text{Tr} |\gamma_N^{(1)} - |\phi\rangle\langle\phi|| \rightarrow 0 \tag{57}$$

as $N \rightarrow \infty$. Here, $\gamma_N^{(1)} := \text{Tr}_{2,3,\dots,N} |\Psi_N\rangle\langle\Psi_N|$.

The main result of this section is the following:

Theorem 5.4 (A derivation of the defocusing cubic NLS on \mathbb{T}^3). *Let $\phi \in H^1(\Lambda)$ with $\|\phi\|_{L^2(\Lambda)} = 1$ be given. Suppose that the sequence $(\Psi_N)_N \in \bigoplus_{N \in \mathbb{N}} L^2(\Lambda^N)$ satisfies the assumptions (56) and (57). For fixed $N \in \mathbb{N}$, we define: $\gamma_N^{(k)} := \text{Tr}_{k+1,\dots,N} |\Psi_N\rangle\langle\Psi_N|$. Let $(\gamma_{N,t}^{(k)})_k$ denote the solution of the BBGKY hierarchy evolving from initial data $(\gamma_N^{(k)})_k$.*

Let us fix $T > 0$. There exists a sequence $N_j \rightarrow \infty$ as $j \rightarrow \infty$ such that, for all $k \in \mathbb{N}$ and for all $t \in [0, T]$:

$$\text{Tr} |\gamma_{N_j,t}^{(k)} - |S(t)\phi\rangle\langle S(t)\phi|^{\otimes k}| \rightarrow 0 \tag{58}$$

as $j \rightarrow \infty$.

In particular, we obtain:

Corollary 5.5 (Evolution of purely factorized states). *In the assumptions (56) and (57), it is possible to take purely factorized states $\Psi_N = \phi^{\otimes N}$ for $\phi \in H^1(\Lambda)$ with $\|\phi\|_{L^2(\Lambda)} = 1$. Thus, (58) holds if the initial data is purely factorized.*

Remark 5.6. The derivation method that we will present applies only in the defocusing setting. In order to study the focusing problem, one would need to take potentials V which would be negative on some set. We need to assume the non-negativity of the potential V in order to apply [50, Theorem 1]. The necessity of this assumption can be seen in the analysis of [50, Section 3].

5.1. Approximation of the initial data

We would like to use [50, Theorem 1] in order to prove Theorem 5.4. In order to do this, we will first apply an approximation procedure to the sequence $(\Psi_N)_N$ satisfying (56) and (57) so that the initial data satisfies the assumption (55) of [50, Theorem 1]. Let us choose $\zeta \in C^\infty_0(\mathbb{R})$ with $0 \leq \zeta \leq 1$, $\zeta \equiv 1$ on $[0, 1]$ and $\zeta \equiv 0$ on $(2, +\infty)$. Finally, let us fix $\kappa > 0$ to be a small parameter.

As in [54, Proposition 5.1], [56, Proposition 9.1], and [57, Proposition 8.1], we define the following approximation to Ψ_N :

$$\tilde{\Psi}_N^\kappa := \frac{\zeta(\kappa \cdot \frac{H_N}{N}) \Psi_N}{\|\zeta(\kappa \cdot \frac{H_N}{N}) \Psi_N\|_{L^2(\Lambda^N)}}.$$

In other words, we smoothly cut off the high frequencies in Ψ_N , and we normalize the result in $L^2(\Lambda^N)$.

The function $\tilde{\Psi}_N^\kappa \in L^2_{sym}(\Lambda^N)$ then satisfies the following properties:

- 1) $\langle H_N^k \tilde{\Psi}_N^\kappa, \tilde{\Psi}_N^\kappa \rangle = \text{Tr } H_N^k |\tilde{\Psi}_N^\kappa\rangle \langle \tilde{\Psi}_N^\kappa| \leq \frac{2^k N^k}{\kappa^k}$.
- 2) For some universal constant $C > 0$:

$$\sup_N \|\Psi_N - \tilde{\Psi}_N^\kappa\|_{L^2(\Lambda^N)} \leq C\kappa^{\frac{1}{2}}. \tag{59}$$

- 3) We let $\tilde{\gamma}_{\kappa;N}^{(k)} := \text{Tr}_{k+1,\dots,N} |\tilde{\Psi}_N^\kappa\rangle \langle \tilde{\Psi}_N^\kappa|$. For $\kappa > 0$ small enough, it is then the case that:

$$\tilde{\gamma}_{\kappa;N}^{(k)} \rightharpoonup^* |\phi\rangle \langle \phi|^{\otimes k} \tag{60}$$

as $N \rightarrow \infty$ in the weak-* topology on \mathcal{L}_k^1 .

This approximation procedure was first used in the non-periodic setting [54,56,57]. The proofs of the analogues of points 1) and 2) in the non-periodic setting [54, Proposition 5.1 i), ii)] and [57, Proposition 8.1 i), ii)] carry over to the periodic setting. We note that the proof of the analogue of 1), given in [54], relies on a Sobolev type inequality whose analogue on \mathbb{T}^3 was proven in Appendix A of [50]. The proof of the non-periodic analogue of point 3), given in [54, Proposition 5.1 iii)], which builds on the previous work of Michelangeli [94], also carries over to the periodic setting by using an appropriate extension argument. In [54], it is assumed that $\beta < \frac{2}{3}$, which holds since we are considering $\beta \in (0, \frac{3}{5})$ by (51). A proof of the non-periodic variant of 3) is also given in [57, Proposition 8.1 iii)]. We will omit the details. Let us note that the same approximation argument was also used in the $2D$ periodic setting in [78].

We observe that, by 1), $|\tilde{\Psi}_N^\kappa\rangle \langle \tilde{\Psi}_N^\kappa|$ satisfies the assumption (55) with implied constant $\frac{2}{\kappa}$. Let us fix $T > 0$ and let $\tilde{\Gamma}_{\kappa;N,t} = (\tilde{\gamma}_{\kappa;N,t}^{(k)})_k$ be the solution to the BBGKY hierarchy (6) with initial data $\tilde{\Gamma}_{\kappa;N,0} = (\tilde{\gamma}_{\kappa;N}^{(k)})_k$ on the time interval $[0, T]$. In particular, for all $k \in \mathbb{N}$ and for all $t \in [0, T]$:

$$\tilde{\gamma}_{\kappa;N,t}^{(k)} = \text{Tr}_{k+1,\dots,N} |\tilde{\Psi}_{N,t}^\kappa\rangle \langle \tilde{\Psi}_{N,t}^\kappa|, \tag{61}$$

where $\tilde{\Psi}_{N,t}^\kappa$ is the solution of:

$$\begin{cases} i \partial_t \tilde{\Psi}_{N,t}^\kappa = H_N \tilde{\Psi}_{N,t}^\kappa \\ \tilde{\Psi}_{N,t}^\kappa|_{t=0} = \tilde{\Psi}_N^\kappa. \end{cases}$$

By [50, Theorem 1], there exists a limit point $\tilde{\Gamma}_{\kappa;\infty,t} = (\tilde{\gamma}_{\kappa;\infty,t}^{(k)})_k$ of $(\tilde{\Gamma}_{\kappa;N,t})_N$ with respect to the metric \hat{q} which solves the GP hierarchy. In other words, there exists $N_j \rightarrow \infty$ such that:

$$\tilde{\Gamma}_{\kappa;N_j,t} \rightarrow \tilde{\Gamma}_{\kappa;\infty,t} \tag{62}$$

as $j \rightarrow \infty$ with respect to the metric \hat{q} .

Let us note that, for fixed $k \in \mathbb{N}$ and $t \in [0, T]$:

$$\text{Tr } \tilde{\gamma}_{\kappa;N,t}^{(k)} = \text{Tr} |\tilde{\Psi}_{N,t}^\kappa\rangle \langle \tilde{\Psi}_{N,t}^\kappa| = \text{Tr} |\tilde{\Psi}_N^\kappa\rangle \langle \tilde{\Psi}_N^\kappa| = \|\tilde{\Psi}_N^\kappa\|_{L^2(\Lambda^N)}^2 = 1. \tag{63}$$

Here, we used the fact that $\text{Tr} |\tilde{\Psi}_{N,t}^\kappa\rangle \langle \tilde{\Psi}_{N,t}^\kappa| = \|\tilde{\Psi}_{N,t}^\kappa\|_{L^2(\Lambda^N)}^2$ is conserved in time. Furthermore, we note that, for all $(\vec{x}_k; \vec{x}'_k) \in \Lambda^k \times \Lambda^k$:

$$\tilde{\gamma}_{\kappa;N,t}^{(k)}(\vec{x}_k; \vec{x}'_k) = \int_{\Lambda^{N-k}} \tilde{\Psi}_{N,t}^\kappa(\vec{x}_k, \vec{y}_{N-k}) \cdot \overline{\tilde{\Psi}_{N,t}^\kappa(\vec{x}'_k, \vec{y}_{N-k})} d\vec{y}_{N-k}.$$

Hence, for all $\vec{x}_k \in \Lambda^k$:

$$|\tilde{\gamma}_{\kappa;N,t}^{(k)}(\vec{x}_k; \vec{x}_k)| = \tilde{\gamma}_{\kappa;N,t}^{(k)}(\vec{x}_k; \vec{x}_k) = \int_{\Lambda^{N-k}} |\tilde{\Psi}_{N,t}^\kappa(\vec{x}_k, \vec{y}_{N-k})|^2 d\vec{y}_{N-k}.$$

In particular:

$$\text{Tr} |\tilde{\gamma}_{\kappa;N,t}^{(k)}| = \text{Tr } \tilde{\gamma}_{\kappa;N,t}^{(k)} = 1. \tag{64}$$

5.2. A comparison of different forms of convergence

Let us henceforth fix κ to be sufficiently small so that (60) holds. We would like to use Theorem 4.6 and deduce that $\tilde{\gamma}_{\kappa; \infty, t}^{(k)} = |S_t(\phi)\rangle\langle S_t(\phi)|^{\otimes k}$ for all $k \in \mathbb{N}$ and for all $t \in [0, T]$. We can not directly apply Theorem 4.6 since the convergence in (62) is given in terms of the metric $\widehat{\varrho}$, and in the assumptions of the theorem, the convergence of each component is given with respect to the weak- $*$ topology on \mathcal{L}_k^1 .

The arguments in [50, Section 4.1] imply that the sequence $(\tilde{\Gamma}_{\kappa; N, t})_N$ satisfies the following equicontinuity result: For all $k \geq 1$ and for all $\mathcal{Z}^{(k)} \in W^{1, \infty}(\Lambda^k \times \Lambda^k)$, it is the case that, for every $\epsilon > 0$, there exists $\delta > 0$ such that:

$$\left| \langle \mathcal{Z}^{(k)}, \tilde{\gamma}_{\kappa; N, t}^{(k)} - \tilde{\gamma}_{\kappa; N, s}^{(k)} \rangle_{L^2(\Lambda^k \times \Lambda^k)} \right| \leq \epsilon, \tag{65}$$

for all $s, t \in [0, T]$ with $|s - t| \leq \delta$. The quantity δ depends on k and $\mathcal{Z}^{(k)}$, but is independent of N .

We would like to prove an equicontinuity result of the type (65) in the weak- $*$ topology on \mathcal{L}_k^1 . In order to do this, we first prove the following:

Lemma 5.7. *Let us fix $k \in \mathbb{N}$. Then, integral operators on $L^2(\Lambda^k)$ with kernels in $W^{1, \infty}(\Lambda^k \times \Lambda^k)$ are dense in \mathcal{K}_k with respect to the operator norm topology.*

Proof. Suppose that $\mathcal{J}^{(k)} \in \mathcal{K}_k$. Let $\epsilon > 0$ be given. We can then find a finite-rank operator T_1 such that $\|\mathcal{J}^{(k)} - T_1\| < \frac{\epsilon}{2}$. The operator T_1 will have a kernel of the form: $K_1(\vec{x}_k; \vec{x}'_k) = \sum_{j=1}^M \alpha_j \cdot f_j(\vec{x}_k) \cdot g_j(\vec{x}'_k)$ for some $M \in \mathbb{N}, \alpha_1, \alpha_2, \dots, \alpha_M \in \mathbb{C}$, and for some $f_1, \dots, f_M, g_1, \dots, g_M \in L^2(\Lambda^k)$. In particular, $K_1 \in L^2(\Lambda^k \times \Lambda^k)$. Since $W^{1, \infty}(\Lambda^k \times \Lambda^k)$ is dense in $L^2(\Lambda^k \times \Lambda^k)$ with respect to $\|\cdot\|_{L^2(\Lambda^k \times \Lambda^k)}$ (we can see this by truncating Fourier series), it follows that there exists $\mathcal{Z}^{(k)} \in W^{1, \infty}(\Lambda^k \times \Lambda^k)$ such that $\|\mathcal{Z}^{(k)} - K_1\|_{L^2(\Lambda^k \times \Lambda^k)} < \frac{\epsilon}{2}$. We denote the integral operator associated with $\mathcal{Z}^{(k)}$ by T . It follows that $\|T - T_1\| = \|\mathcal{Z}^{(k)} - K_1\|_{L^2(\Lambda^k \times \Lambda^k)} < \frac{\epsilon}{2}$ and hence $\|\mathcal{J}^{(k)} - T\| \leq \|\mathcal{J}^{(k)} - T_1\| + \|T - T_1\| < \epsilon$. The density result now follows. \square

Let us return to the proof of the equicontinuity result in the weak- $*$ topology on \mathcal{L}_k^1 . Suppose that $\mathcal{J}^{(k)} \in \mathcal{L}_k^1$. Let $\epsilon > 0$ be given. Then, by Lemma 5.7, we can find $\mathcal{J}_1^{(k)}$, an integral operator with kernel $\mathcal{Z}^{(k)} \in W^{1, \infty}(\Lambda^k \times \Lambda^k)$ such that $\|\mathcal{J}^{(k)} - \mathcal{J}_1^{(k)}\| \leq \frac{\epsilon}{4}$. In particular, for all $s, t \in [0, T]$, it is the case that:

$$\begin{aligned} |\text{Tr}(\mathcal{J}^{(k)} - \mathcal{J}_1^{(k)})(\tilde{\gamma}_{\kappa; N, t}^{(k)} - \tilde{\gamma}_{\kappa; N, s}^{(k)})| &\leq |\text{Tr}(\mathcal{J}^{(k)} - \mathcal{J}_1^{(k)})(\tilde{\gamma}_{\kappa; N, t}^{(k)})| + |\text{Tr}(\mathcal{J}^{(k)} - \mathcal{J}_1^{(k)})(\tilde{\gamma}_{\kappa; N, s}^{(k)})| \\ &\leq \|\mathcal{J}^{(k)} - \mathcal{J}_1^{(k)}\| \cdot \text{Tr}|\tilde{\gamma}_{\kappa; N, t}^{(k)}| + \|\mathcal{J}^{(k)} - \mathcal{J}_1^{(k)}\| \cdot \text{Tr}|\tilde{\gamma}_{\kappa; N, s}^{(k)}| \leq \frac{\epsilon}{4} + \frac{\epsilon}{4} = \frac{\epsilon}{2}. \end{aligned}$$

In the last line, we used the duality pairing (14) as well as (64).

Furthermore, we notice that:

$$\text{Tr} \mathcal{J}_1^{(k)}(\tilde{\gamma}_{\kappa; N, t}^{(k)} - \tilde{\gamma}_{\kappa; N, s}^{(k)}) = \langle \mathcal{Z}^{(k)}, (\tilde{\gamma}_{\kappa; N, t}^{(k)} - \tilde{\gamma}_{\kappa; N, s}^{(k)}) \rangle_{L^2(\Lambda^k \times \Lambda^k)}.$$

Hence, from (65), it follows that there exists $\delta > 0$ (which depends on k and $\mathcal{J}_1^{(k)}$) such that for all $s, t \in [0, T]$ with $|s - t| \leq \delta$:

$$|\text{Tr} \mathcal{J}_1^{(k)}(\tilde{\gamma}_{\kappa; N, t}^{(k)} - \tilde{\gamma}_{\kappa; N, s}^{(k)})| \leq \frac{\epsilon}{2}.$$

In particular, by the triangle inequality, it follows that for all $s, t \in [0, T]$ with $|s - t| \leq \delta$:

$$|\text{Tr} \mathcal{J}^{(k)}(\tilde{\gamma}_{\kappa; N, t}^{(k)} - \tilde{\gamma}_{\kappa; N, s}^{(k)})| \leq \epsilon. \tag{66}$$

Remark 5.8. Let us note that similar density arguments were used in order to obtain equicontinuity results in the periodic setting in [53, Lemma 9.2] the non-periodic setting in [54, Lemma 7.2], [56, Lemma 6.2], [57, Lemma 6.2].

Let us recall from (13) that $\mathcal{L}_k^1 = \mathcal{K}_k^*$. Since \mathcal{K}_k is separable, there exists a countable dense subset of the unit ball of \mathcal{K}_k given by $\{\mathcal{J}_{c,\ell}^{(k)}, \ell \in \mathbb{N}\}$. As in [54,56,57], using the operators $\mathcal{J}_{c,\ell}^{(k)}$, it is possible to define a metric on \mathcal{L}_k^1 as:

$$\eta_k(\gamma_1^{(k)}, \gamma_2^{(k)}) := \sum_{\ell=1}^{+\infty} 2^{-\ell} \cdot |\text{Tr } \mathcal{J}_{c,\ell}^{(k)}(\gamma_1^{(k)} - \gamma_2^{(k)})|. \tag{67}$$

The topology induced by the metric η_k and the weak-* topology are equivalent on the unit ball of \mathcal{L}_k^1 , c.f. [102, Theorem 3.16].

Let $C([0, T], \mathcal{L}_k^1)$ denote the space of all functions from $[0, T]$ which take values in \mathcal{L}_k^1 and which are continuous with respect to the topology given by the metric η_k given in (67). On $C([0, T], \mathcal{L}_k^1)$, one defines the metric:

$$\widehat{\eta}_k(\gamma_1^{(k)}(\cdot), \gamma_2^{(k)}(\cdot)) := \sup_{t \in [0, T]} \eta_k(\gamma_1^{(k)}(t), \gamma_2^{(k)}(t)). \tag{68}$$

Finally, one defines τ_{prod} to be the product metric on $\bigoplus_{k \in \mathbb{N}} C([0, T], \mathcal{L}_k^1)$ obtained from the metric $\widehat{\eta}_k$ defined in (68).

5.3. Proof of Theorem 5.4

We will now use the above facts and prove Theorem 5.4:

Proof. It follows from (66) that the sequence $(\widetilde{\gamma}_{\kappa; N, t}^{(k)})_N$ is equicontinuous with respect to η_k . More precisely, given $\epsilon > 0$, we find $M > 1$ such that $\sum_{\ell > M} 2^{-\ell} < \frac{\epsilon}{4}$. We note that

$$|\text{Tr } \mathcal{J}_{c,\ell}^{(k)}(\widetilde{\gamma}_{\kappa; N, t}^{(k)} - \widetilde{\gamma}_{\kappa; N, s}^{(k)})| \leq \|\mathcal{J}_{c,\ell}^{(k)}\| \cdot \text{Tr}|\widetilde{\gamma}_{\kappa; N, t}^{(k)} - \widetilde{\gamma}_{\kappa; N, s}^{(k)}| \leq 2.$$

Here, we used the triangle inequality and (64). Hence:

$$\eta_k(\widetilde{\gamma}_{\kappa; N, t}^{(k)}, \widetilde{\gamma}_{\kappa; N, s}^{(k)}) \leq \sum_{\ell=1}^M 2^{-\ell} \cdot |\text{Tr } \mathcal{J}_{c,\ell}^{(k)}(\widetilde{\gamma}_{\kappa; N, t}^{(k)} - \widetilde{\gamma}_{\kappa; N, s}^{(k)})| + \frac{\epsilon}{2},$$

which can be made $\leq \epsilon$ by using (66) if we choose $|t - s|$ to be sufficiently small. By the Arzelà–Ascoli theorem, it follows that the sequence $(\widetilde{\gamma}_{\kappa; N, t}^{(k)})_N$ is relatively compact with respect to $\widehat{\eta}_k$. By using an additional diagonal argument, it follows that $(\widetilde{\Gamma}_{\kappa; N, t}^{(k)})_N$ is relatively compact with respect to τ_{prod} . In particular, it follows that there exists $\overline{\Gamma}_{\kappa; \infty, t} = (\overline{\gamma}_{\kappa; \infty, t}^{(k)})_k \in \bigoplus_{k \in \mathbb{N}} C([0, T], \mathcal{L}_k^1)$ and $N_j \rightarrow \infty$ such that:

$$\widetilde{\Gamma}_{\kappa; N_j, t}^{(k)} \rightarrow \overline{\Gamma}_{\kappa; \infty, t} \tag{69}$$

as $j \rightarrow \infty$ with respect to τ_{prod} . Let us note that “ \cdot ” in the above notation does not mean a complex conjugate. By taking subsequences, we can arrange for the N_j in (62) and (69) to be the same. Let us note that the sequence $(N_j)_j$ depends on κ .

We now want to show that, for all $t \in [0, T]$:

$$\widetilde{\Gamma}_{\kappa; \infty, t} = \overline{\Gamma}_{\kappa; \infty, t} \tag{70}$$

as elements of $L^2(\Lambda^k \times \Lambda^k)$. Let us fix $k \in \mathbb{N}$ and $t \in [0, T]$. Suppose that $F \in L^2(\Lambda^k \times \Lambda^k)$. The integral operator T with kernel F is Hilbert–Schmidt and hence is compact, i.e. it belongs to \mathcal{K}_k . Since $\widetilde{\gamma}_{\kappa; N_j, t}^{(k)} \xrightarrow{*} \overline{\gamma}_{\kappa; \infty, t}^{(k)}$ in the weak-* topology on \mathcal{L}_k^1 as $j \rightarrow \infty$, it follows from (14) that:

$$\text{Tr}(T \widetilde{\gamma}_{\kappa; N_j, t}^{(k)}) \rightarrow \text{Tr}(T \overline{\gamma}_{\kappa; \infty, t}^{(k)})$$

as $j \rightarrow \infty$, which can be rewritten as:

$$\begin{aligned} & \int_{\Lambda^k \times \Lambda^k} F(\vec{x}_k; \vec{x}'_k) \cdot \widetilde{\gamma}_{\kappa; N_j, t}^{(k)}(\vec{x}_k; \vec{x}'_k) d\vec{x}_k d\vec{x}'_k \\ & \rightarrow \int_{\Lambda^k \times \Lambda^k} F(\vec{x}_k; \vec{x}'_k) \cdot \overline{\gamma}_{\kappa; \infty, t}^{(k)}(\vec{x}_k; \vec{x}'_k) d\vec{x}_k d\vec{x}'_k. \end{aligned} \tag{71}$$

In particular:

$$\tilde{\gamma}_{\kappa; N_j, t}^{(k)} \rightharpoonup \tilde{\gamma}_{\kappa; \infty, t}^{(k)}$$

weakly in $L^2(\Lambda^k \times \Lambda^k)$ as $j \rightarrow \infty$.

Let us observe that (62) implies that:

$$\tilde{\gamma}_{\kappa; N, t}^{(k)} \rightarrow \tilde{\gamma}_{\kappa; \infty, t}^{(k)}$$

strongly in $L^2(\Lambda^k \times \Lambda^k)$ as $j \rightarrow \infty$. In particular, $(\tilde{\gamma}_{\kappa; N_j, t}^{(k)})_k$ also converges to $\tilde{\gamma}_{\kappa; \infty, t}^{(k)}$ weakly in $L^2(\Lambda^k \times \Lambda^k)$.

By the uniqueness of weak limits, it follows that $\tilde{\gamma}_{\kappa; \infty, t}^{(k)} = \bar{\gamma}_{\kappa; \infty, t}^{(k)}$ as elements of $L^2(\Lambda^k \times \Lambda^k)$. In particular, we can deduce (70). We will henceforth write the limit as $\tilde{\Gamma}_{\kappa; \infty, t}^{(k)}$.

To summarize, we know from [50, Theorem 1] that $\tilde{\Gamma}_{\kappa; \infty, t}^{(k)}$ is a mild solution of the Gross–Pitaevskii hierarchy which belongs to $L_{t \in [0, T]}^\infty \mathfrak{H}^1$. Let us note that, by construction:

$$\tilde{\gamma}_{\kappa; N, t}^{(k)} = \text{Tr}_{k+1, \dots, N} |\tilde{\Psi}_{N, t}^\kappa\rangle \langle \tilde{\Psi}_{N, t}^\kappa| \geq 0$$

when considered as an operator on $L^2(\Lambda^k)$. Furthermore, $\tilde{\gamma}_{\kappa; N, t}^{(k)}$ is bounded and self-adjoint on $L^2(\Lambda^k)$. We recall from (63) that $\text{Tr} \tilde{\gamma}_{\kappa; N, t}^{(k)} = 1$. Finally, we note that for all $k \in \mathbb{N}$ and for all $t \in [0, T]$, it is the case that $\tilde{\gamma}_{\kappa; N_j, t}^{(k)} \rightharpoonup^* \tilde{\gamma}_{\kappa; \infty, t}^{(k)}$ as $j \rightarrow \infty$ in the weak-* topology on \mathcal{L}_k^1 . By (60), we can deduce that, for all $k \in \mathbb{N}$, it is the case that $\tilde{\gamma}_{\kappa; \infty, 0}^{(k)} = |\phi\rangle \langle \phi|^{\otimes k}$.

In particular, it follows that $\tilde{\Gamma}_{\kappa; \infty, t}^{(k)}$ is a mild solution of the Gross–Pitaevskii hierarchy with initial data given by $\tilde{\Gamma}_{\kappa; \infty, 0}^{(k)} = (|\phi\rangle \langle \phi|^{\otimes k})_k$, which satisfies the assumptions of Theorem 4.6. We can now apply Theorem 4.6 and deduce that:

$$\tilde{\Gamma}_{\kappa; \infty, t}^{(k)} = (|S_t(\phi)\rangle \langle S_t(\phi)|^{\otimes k})_k. \tag{72}$$

Consequently:

$$\tilde{\gamma}_{\kappa; N_j, t}^{(k)} \rightharpoonup^* |S(t)\phi\rangle \langle S(t)\phi|^{\otimes k} \tag{73}$$

as $j \rightarrow \infty$ in the weak-* topology on \mathcal{L}_k^1 .

Let us note that by (64) and the fact that $\|\phi\|_{L^2} = 1$:

$$\text{Tr} |\tilde{\gamma}_{\kappa; N_j, t}^{(k)}| = \text{Tr} ||S(t)\phi\rangle \langle S(t)\phi|^{\otimes k}| = 1. \tag{74}$$

We observe that (73) and (74) imply that:

$$\text{Tr} |\tilde{\gamma}_{\kappa; N_j, t}^{(k)} - |S(t)\phi\rangle \langle S(t)\phi|^{\otimes k}| \rightarrow 0 \tag{75}$$

as $j \rightarrow \infty$. In other words, weak-* convergence in \mathcal{L}_k^1 and convergence of the norms implies strong convergence in \mathcal{L}_k^1 . This fact was already noted in [56, Proposition 9.1 iii)]. It follows from a more general fact about convergence in the trace class proved by Arazy [10]. For related results, we refer the reader to [104], [105, Addendum H], and [83, Corollary 2.4], as well as [46] and [100]. Let us note that, in (75), the sequence $(N_j)_j$ can be taken to be independent of $k \in \mathbb{N}$ and $t \in [0, T]$. We recall that $(N_j)_j$ does depend on κ .

Let us now prove an analogue of (75) which is obtained by starting from the original functions Ψ_N . In other words, we start from $(\Psi_N)_N \in \bigoplus_{N \in \mathbb{N}} L^2(\Lambda^N)$ which satisfies (56) and (57), and we consider:

$$\gamma_{N, t}^{(k)} := \text{Tr}_{k+1, \dots, N} |\Psi_{N, t}\rangle \langle \Psi_{N, t}|$$

as in (5).

We want to estimate $\text{Tr} |\gamma_{N,t}^{(k)} - \tilde{\gamma}_{\kappa;N,t}^{(k)}|$ for $\tilde{\gamma}_{\kappa;N,t}^{(k)}$ as in (61). Let us note that:

$$\begin{aligned} \text{Tr} |\gamma_{N,t}^{(k)} - \tilde{\gamma}_{\kappa;N,t}^{(k)}| &= \text{Tr}_{1,2,\dots,k} |\text{Tr}_{k+1,\dots,N} |\Psi_{N,t}\rangle \langle \Psi_{N,t}| - \text{Tr}_{k+1,\dots,N} |\tilde{\Psi}_{\kappa;N,t}\rangle \langle \tilde{\Psi}_{\kappa;N,t}|| \\ &\leq \text{Tr} |\Psi_{N,t}\rangle \langle \Psi_{N,t}| - |\tilde{\Psi}_{\kappa;N,t}\rangle \langle \tilde{\Psi}_{\kappa;N,t}|| \\ &\leq \text{Tr} |\Psi_{N,t}\rangle \langle \Psi_{N,t} - \tilde{\Psi}_{\kappa;N,t}|| + \text{Tr} |(\Psi_{N,t} - \tilde{\Psi}_{\kappa;N,t})\rangle \langle \Psi_{N,t}|| \\ &\quad + \text{Tr} |(\Psi_{N,t} - \tilde{\Psi}_{\kappa;N,t})\rangle \langle (\Psi_{N,t} - \tilde{\Psi}_{\kappa;N,t})|| \\ &\leq 2 \|\Psi_{N,t}\|_{L^2(\Lambda^N)} \cdot \|\Psi_{N,t} - \tilde{\Psi}_{\kappa;N,t}\|_{L^2(\Lambda^N)} + \|\Psi_{N,t} - \tilde{\Psi}_{\kappa;N,t}\|_{L^2(\Lambda^N)}^2 \\ &= 2 \|\Psi_N\|_{L^2(\Lambda^N)} \cdot \|\Psi_N - \tilde{\Psi}_{\kappa;N}\|_{L^2(\Lambda^N)} + \|\Psi_N - \tilde{\Psi}_{\kappa;N}\|_{L^2(\Lambda^N)}^2. \end{aligned} \tag{76}$$

Above, we used the triangle inequality, the Cauchy–Schwarz inequality and the fact that $\|\Psi_{N,t}\|_{L^2(\Lambda^N)}$ and $\|\Psi_{N,t} - \tilde{\Psi}_{\kappa;N,t}\|_{L^2(\Lambda^N)}$ are conserved in time. We note that (57) implies that $\|\Psi_N\|_{L^2(\Lambda^N)} \rightarrow 1$ as $N \rightarrow \infty$. In particular, $(\|\Psi_N\|_{L^2(\Lambda^N)})_N$ is bounded. By using (59) and the fact that $\|\Psi_N\|_{L^2(\Lambda^N)} \lesssim 1$, it follows that the expression in (76) is:

$$\leq C' \cdot \kappa^{\frac{1}{2}} \tag{77}$$

for some $C' > 0$, which is independent of k, N and κ .

Let us recall that, the sequence $(N_j)_j$ for which (75) holds, depends on κ . We will now emphasize this by writing $N_j = N_j(\kappa)$. Consequently, by (77), we note that:

$$\begin{aligned} \text{Tr} |\gamma_{N_j(\kappa),t}^{(k)} - |S(t)\phi\rangle \langle S(t)\phi|^{\otimes k}| &\leq \text{Tr} |\tilde{\gamma}_{\kappa;N_j(\kappa),t}^{(k)} - |S(t)\phi\rangle \langle S(t)\phi|^{\otimes k}| + \text{Tr} |\gamma_{N_j(\kappa),t}^{(k)} - \tilde{\gamma}_{\kappa;N_j(\kappa),t}^{(k)}| \\ &\leq \text{Tr} |\tilde{\gamma}_{\kappa;N_j(\kappa),t}^{(k)} - |S(t)\phi\rangle \langle S(t)\phi|^{\otimes k}| + C' \cdot \kappa^{\frac{1}{2}}. \end{aligned}$$

We take a sequence $\kappa_n := \frac{1}{n}$, and we apply (75) and a diagonal argument in order to deduce that there exists a new sequence of positive integers tending to infinity, which is again denoted by $(N_j)_j$, which is independent of k and t , and has the property that:

$$\text{Tr} |\gamma_{N_j,t}^{(k)} - |S(t)\phi\rangle \langle S(t)\phi|^{\otimes k}| \rightarrow 0 \tag{78}$$

as $j \rightarrow \infty$. The theorem now follows from (78). \square

Remark 5.9. We note that the place where we applied the uniqueness result from Theorem 4.6 is in (72) above in order to deduce that $\tilde{\Gamma}_{\kappa;\infty,t}$ is uniquely determined. In general, we can deduce that $\Gamma_{\infty,t}$ in [50, Theorem 1] is uniquely determined provided that it satisfies the assumptions of Theorem 4.6. As we saw above, this was true in the intermediate step (72) of deriving the defocusing cubic nonlinear Schrödinger equation on \mathbb{T}^3 .

Conflict of interest statement

The author certifies no connections or involvement with any organization or entity which has financial or non-financial interest in the subject matter discussed in the manuscript.

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References

- [1] R. Adami, C. Bardos, F. Golse, A. Teta, Towards a rigorous derivation of the cubic nonlinear Schrödinger equation in dimension one, *Asymptot. Anal.* 40 (2) (2004) 93–108.
- [2] R. Adami, F. Golse, A. Teta, Rigorous derivation of the cubic NLS in dimension one, *J. Stat. Phys.* 127 (6) (2007) 1193–1220.
- [3] M. Aizenman, E.H. Lieb, R. Seiringer, J.P. Solovej, J. Yngvason, Bose–Einstein quantum phase transition in an optical lattice model, *Phys. Rev. A* 70 (2004) 023612.
- [4] M. Aizenman, E.H. Lieb, R. Seiringer, J.P. Solovej, J. Yngvason, Bose–Einstein condensation as a quantum phase transition in an optical lattice, in: *Mathematical Physics of Quantum Mechanics*, in: *Lecture Notes in Physics*, vol. 690, 2006, pp. 199–215.
- [5] Z. Ammari, *Systèmes hamiltoniens en théorie quantique des champs: dynamique asymptotique et limite classique*, Habilitation à Diriger des Recherches, University of Rennes I, February 2013.
- [6] Z. Ammari, F. Nier, Mean field limit for bosons and infinite dimensional phase-space analysis, *Ann. Henri Poincaré* 9 (2008) 1503–1574.
- [7] Z. Ammari, F. Nier, Mean field propagation of Wigner measures and BBGKY hierarchies for general bosonic states, *J. Math. Pures Appl.* 95 (2011) 585–626.
- [8] I. Anapolitanos, Rate of Convergence towards the Hartree-von Neumann limit in the mean-field regime, *Lett. Math. Phys.* 98 (1) (Oct. 2011) 1–31.
- [9] M.H. Anderson, J.R. Ensher, M.R. Matthews, C.E. Wieman, E.A. Cornell, Observations of Bose–Einstein condensation in a dilute atomic vapor, *Science* 269 (1995) 198–201.
- [10] J. Arazy, More on convergence in unitary metric spaces, *Proc. Am. Math. Soc.* 83 (1981) 44–48.
- [11] C. Bardos, F. Golse, N. Mauser, Weak coupling limit of the N -particle Schrödinger equation, *Methods Appl. Anal.* 7 (2000) 275–293.
- [12] N. Benedikter, M. Porta, B. Schlein, Mean-field evolution of Fermionic systems, preprint, <http://arxiv.org/abs/1305.2768>, 2013.
- [13] N. Benedikter, M. Porta, B. Schlein, Mean-field dynamics of Fermions with relativistic dispersion, preprint, <http://arxiv.org/abs/1311.6270>, 2013.
- [14] N. Benedikter, G. de Oliveira, B. Schlein, Quantitative derivation of the Gross–Pitaevskii equation, preprint, <http://arxiv.org/abs/1208.0373>, 2012.
- [15] S.N. Bose, Plancks Gesetz und Lichtquantenhypothese, *Z. Phys.* 26 (1924) 178.
- [16] J. Bourgain, Fourier transform restriction phenomena for certain lattice subsets and applications to nonlinear evolution equations, I: Schrödinger equations, *Geom. Funct. Anal.* 3 (1993) 107–156.
- [17] J. Bourgain, Periodic nonlinear Schrödinger equation and invariant measures, *Commun. Math. Phys.* 166 (1) (1994) 1–26.
- [18] N. Burq, N. Tzvetkov, Random data Cauchy theory for supercritical wave equations. I. Local theory, *Invent. Math.* 173 (3) (2008) 449–475.
- [19] L. Chen, J.O. Lee, B. Schlein, Rate of convergence towards Hartree dynamics, *J. Stat. Phys.* 144 (4) (2011) 872–903.
- [20] T. Chen, C. Hainzl, N. Pavlović, R. Seiringer, Unconditional uniqueness for the cubic Gross–Pitaevskii hierarchy via quantum de Finetti, preprint, <http://arxiv.org/abs/1307.3168>, 2013.
- [21] T. Chen, C. Hainzl, N. Pavlović, R. Seiringer, On the well-posedness and scattering for the Gross–Pitaevskii hierarchy via quantum de Finetti, *Lett. Math. Phys.* 104 (7) (2014) 871–891.
- [22] T. Chen, N. Pavlović, On the Cauchy problem for focusing and defocusing Gross–Pitaevskii hierarchies, *Discrete Contin. Dyn. Syst.* 27 (2) (2010) 715–739.
- [23] T. Chen, N. Pavlović, Recent results on the Cauchy problem for focusing and defocusing Gross–Pitaevskii hierarchies, *Math. Model. Nat. Phenom.* 5 (4) (2010) 54–72.
- [24] T. Chen, N. Pavlović, The quintic NLS as the mean field limit of a Boson gas with three-body interactions, *J. Funct. Anal.* 260 (4) (2011) 959–997.
- [25] T. Chen, N. Pavlović, A new proof of existence of solutions for focusing and defocusing Gross–Pitaevskii hierarchies, *Proc. Am. Math. Soc.* 141 (1) (2013) 279–293.
- [26] T. Chen, N. Pavlović, Higher order energy conservation and global well-posedness for Gross–Pitaevskii hierarchies, *Commun. Partial Differ. Equ.* 39 (9) (2014) 1597–1634.
- [27] T. Chen, N. Pavlović, Derivation of the cubic NLS and Gross–Pitaevskii hierarchy from many-body dynamics in $d = 2, 3$ based on spacetime norms, *Ann. Henri Poincaré* 15 (3) (2014) 543–588.
- [28] T. Chen, N. Pavlović, N. Tzirakis, Energy conservation and blowup of solutions for focusing and defocusing Gross–Pitaevskii hierarchies, *Ann. Inst. Henri Poincaré (C), Anal. Non Linéaire* 27 (5) (2010) 1271–1290.
- [29] T. Chen, N. Pavlović, N. Tzirakis, Multilinear Morawetz identities for the Gross–Pitaevskii hierarchy, in: *Recent Advances in Harmonic Analysis and Partial Differential Equations*, in: *Contemp. Math.*, vol. 581, Amer. Math. Soc., Providence, RI, 2012, pp. 39–62.
- [30] T. Chen, K. Taliaferro, Positive semidefiniteness and global well-posedness of solutions to the Gross–Pitaevskii hierarchy, preprint, <http://arxiv.org/abs/1305.1404>, 2013.
- [31] X. Chen, The Grillakis–Machedon–Margetis second order corrections to mean field evolution for weakly interacting bosons in the case of 3-body interactions, preprint, <http://arxiv.org/abs/0911.4153>, 2009.
- [32] X. Chen, Second order corrections to mean field evolution for weakly interacting bosons in the case of three-body interactions, *Arch. Ration. Mech. Anal.* 203 (2012) 455–497.
- [33] X. Chen, Collapsing estimates and the Rigorous derivation of the 2d cubic nonlinear Schrödinger equation with anisotropic switchable quadratic traps, *J. Math. Pures Appl.* (9) 98 (4) (2012) 450–478.
- [34] X. Chen, On the rigorous derivation of the 3D cubic nonlinear Schrödinger equation with a quadratic trap, *Arch. Ration. Mech. Anal.* 210 (2) (2013) 365–408.

- [35] X. Chen, J. Holmer, On the rigorous derivation of the 2D cubic nonlinear Schrödinger equation from 3D quantum many-body dynamics, *Arch. Ration. Mech. Anal.* 210 (3) (2013) 909–954.
- [36] X. Chen, J. Holmer, On the Klainerman–Machedon conjecture of the quantum BBGKY hierarchy with self-interaction, *J. Eur. Math. Soc.* (2013), in press.
- [37] X. Chen, J. Holmer, Focusing quantum many-body dynamics: the Rigorous derivation of the 1D focusing cubic nonlinear Schrödinger equation, preprint, <http://arxiv.org/abs/1308.3895>, 2013.
- [38] X. Chen, J. Holmer, Focusing quantum many-body dynamics II: the Rigorous derivation of the 1D focusing cubic nonlinear Schrödinger equation from 3D, preprint, <http://arxiv.org/abs/1407.8457>, 2014.
- [39] X. Chen, J. Holmer, Correlation structures, many-body scattering processes and the derivation of the Gross–Pitaevskii hierarchy, preprint, <http://arxiv.org/abs/1409.1425>, 2014.
- [40] X. Chen, P. Smith, On the unconditional uniqueness of solutions to the infinite radial Chern–Simons–Schrödinger hierarchy, preprint, <http://arxiv.org/abs/1406.2649>, 2014.
- [41] Z. Chen, C. Liu, On the Cauchy problem for Gross–Pitaevskii hierarchies, *J. Math. Phys.* 52 (3) (2011) 032103.
- [42] J. Conway, *A Course in Operator Theory*, Graduate Studies in Mathematics, vol. 21, American Mathematical Society, Providence, RI, 1991.
- [43] K.B. Davis, M.O. Mewes, M.R. Andrews, N.J. van Druten, D.S. Durfee, D.M. Kurn, W. Ketterle, Bose–Einstein condensation in a gas of sodium atoms, *Phys. Rev. Lett.* 75 (22) (1995) 3969–3973.
- [44] B. de Finetti, Funzione caratteristica di un fenomeno aleatorio, *Mem. R. Accad. Lincei* 4, 86–133.
- [45] B. de Finetti, La pr’évision: ses lois logiques, ses sources subjectives, *Ann. Inst. Henri Poincaré* 7 (1) (1937).
- [46] G. Dell’ Antonio, On the limits of sequences of normal states, *Commun. Pure Appl. Math.* 20 (1967) 413–429.
- [47] P. Diaconis, D. Freedman, Finite exchangeable sequences, *Ann. Probab.* 8 (1980) 745–764.
- [48] E.B. Dynkin, Classes of equivalent random quantities, *Usp. Mat. Nauk* 8 (1953) 125–130.
- [49] A. Einstein, Quantentheorie des einatomigen idealen Gases, *Sitzber. Preuss. Akad. Wiss. (Berl.)*, Phys.-math. Kl. 1 (3) (1925) 18–25.
- [50] A. Elgart, L. Erdős, B. Schlein, H.-T. Yau, Gross–Pitaevskii equation as the mean field limit of weakly coupled bosons, *Arch. Ration. Mech. Anal.* 179 (2006) 265–283.
- [51] A. Elgart, B. Schlein, Mean field dynamics of boson stars, *Commun. Pure Appl. Math.* 60 (4) (2007) 500–545.
- [52] L. Erdős, B. Schlein, Quantum dynamics with mean field interactions: a new approach, *J. Stat. Phys.* 134 (5) (2009) 859–870.
- [53] L. Erdős, B. Schlein, H.-T. Yau, Derivation of the Gross–Pitaevskii hierarchy for the dynamics of Bose–Einstein condensate, *Commun. Pure Appl. Math.* 59 (12) (2006) 1659–1741.
- [54] L. Erdős, B. Schlein, H.-T. Yau, Derivation of the cubic non-linear Schrödinger equation from quantum dynamics of many-body systems, *Invent. Math.* 167 (3) (2007) 515–614.
- [55] L. Erdős, B. Schlein, H.-T. Yau, Rigorous derivation of the Gross–Pitaevskii equation, *Phys. Rev. Lett.* 98 (4) (2007) 040404.
- [56] L. Erdős, B. Schlein, H.-T. Yau, Rigorous derivation of the Gross–Pitaevskii equation with a large interaction potential, *J. Am. Math. Soc.* 22 (4) (2009) 1099–1156.
- [57] L. Erdős, B. Schlein, H.-T. Yau, Derivation of the Gross–Pitaevskii equation for the dynamics of Bose–Einstein condensate, *Ann. Math. (2)* 172 (1) (2010) 291–370.
- [58] L. Erdős, H.-T. Yau, Derivation of the nonlinear Schrödinger equation from a many body Coulomb system, *Adv. Theor. Math. Phys.* 5 (6) (2001) 1169–1205.
- [59] J. Fröhlich, S. Graffi, S. Schwarz, Mean-field and classical limit of many body Schrödinger dynamics for bosons, *Commun. Math. Phys.* 271 (3) (2007) 681–697.
- [60] J. Fröhlich, A. Knowles, A. Pizzo, Atomism and quantization, *J. Phys. A* 40 (12) (2007) 3033–3045.
- [61] J. Fröhlich, A. Knowles, S. Schwarz, On the mean-field limit of bosons with Coulomb two-body interaction, *Commun. Math. Phys.* 288 (3) (2009) 1023–1059.
- [62] J. Fröhlich, E. Lenzmann, Mean-field limit of quantum Bose gases and nonlinear Hartree equation, in: *Sem. É.D.P. 2003–2004, Exp. No. XIX*, in: *Sémin. Équ. Dériv. Partielles, École Polytech, Palaiseau, 2004*, 26 pp.
- [63] J. Fröhlich, T.-P. Tsai, H.-T. Yau, On a classical limit of quantum theory and the non-linear Hartree equation, in: *Conférence Moshé Flato 1999, Vol. I (Dijon)*, in: *Math. Phys. Stud.*, vol. 21, Kluwer Acad. Publ., Dordrecht, 2000, pp. 189–207.
- [64] J. Fröhlich, T.-P. Tsai, H.-T. Yau, On a classical limit of quantum theory and the non-linear Hartree equation, in: *Gafa 2000, Tel Aviv, 1999*, in: *Geom. Funct. Anal.*, special volume, part I, 2000, pp. 57–78.
- [65] J. Fröhlich, T.-P. Tsai, H.-T. Yau, On the point-particle (Newtonian) limit of the non-linear Hartree equation, *Commun. Math. Phys.* 225 (2) (2002) 223–274.
- [66] J. Ginibre, G. Velo, The classical field limit of scattering theory for nonrelativistic many-boson systems I, *Commun. Math. Phys.* 66 (1) (1979) 37–76.
- [67] J. Ginibre, G. Velo, The classical field limit of scattering theory for nonrelativistic many-boson systems II, *Commun. Math. Phys.* 68 (1) (1979) 45–68.
- [68] P. Gressman, V. Sohinger, G. Staffilani, On the uniqueness of solutions to the periodic 3D Gross–Pitaevskii hierarchy, *J. Funct. Anal.* 266 (7) (2014) 4705–4764.
- [69] M. Grillakis, M. Machedon, Pair excitations and the mean field approximation of interacting bosons, I, *Commun. Math. Phys.* 324 (2) (2013) 601–636.
- [70] M. Grillakis, M. Machedon, D. Margetis, Second order corrections to mean field evolution of weakly interacting bosons, I, *Commun. Math. Phys.* 294 (2010) 273–301.
- [71] M. Grillakis, M. Machedon, D. Margetis, Second order corrections to mean field evolution of weakly interacting bosons, II, *Adv. Math.* 228 (3) (2011) 1788–1815.

- [72] E. Gross, Structure of a quantized vortex in boson systems, *Nuovo Cimento* 20 (1961) 454–466.
- [73] K. Hepp, The classical limit for quantum mechanical correlation functions, *Commun. Math. Phys.* 35 (1974) 265–277.
- [74] S. Herr, D. Tataru, N. Tzvetkov, Global well-posedness of the energy-critical nonlinear Schrödinger equation with small initial data in $H^1(\mathbb{T}^3)$, *Duke Math. J.* 159 (2) (2011) 329–349.
- [75] E. Hewitt, L.J. Savage, Symmetric measures on Cartesian products, *Trans. Am. Math. Soc.* 80 (1955) 470–501.
- [76] Y. Hong, K. Taliaferro, Z. Xie, Unconditional uniqueness of the cubic Gross–Pitaevskii hierarchy with low regularity, preprint, <http://arxiv.org/abs/1402.5347>, 2014.
- [77] R.L. Hudson, G.R. Moody, Locally normal symmetric states and an analogue of de Finetti’s theorem, *Z. Wahrscheinlichkeitstheor. Verw. Geb.* 33 (1975/1976) 343–351.
- [78] K. Kirkpatrick, B. Schlein, G. Staffilani, Derivation of the cubic nonlinear Schrödinger equation from quantum dynamics of many-body systems: the periodic case, *Am. J. Math.* 133 (1) (2011) 91–130.
- [79] S. Klainerman, M. Machedon, Space-time estimates for null forms and the local existence theorem, *Commun. Pure Appl. Math.* 46 (1993) 169–177.
- [80] S. Klainerman, M. Machedon, On the uniqueness of solutions to the Gross–Pitaevskii hierarchy, *Commun. Math. Phys.* 279 (1) (2008) 169–185.
- [81] A. Knowles, P. Pickl, Mean-field dynamics: singular potentials and rate of convergence, *Commun. Math. Phys.* 298 (1) (2010) 101–139.
- [82] J.O. Lee, Rate of convergence towards semi-relativistic Hartree dynamics, *Ann. Henri Poincaré* 14 (2) (2013) 313–346.
- [83] M. Lewin, P.T. Nam, N. Rougerie, Derivation of Hartree’s theory for generic mean-field Bose systems, *Adv. Math.* 254 (2014) 570–621.
- [84] M. Lewin, P.T. Nam, N. Rougerie, Remarks on the quantum de Finetti theorem for bosonic systems, preprint, <http://arxiv.org/abs/1310.2200>, 2013.
- [85] M. Lewin, P.T. Nam, N. Rougerie, The mean-field approximation and the non-linear Schrödinger functional for trapped Bose gases, preprint, <http://arxiv.org/abs/1405.3220>, 2014.
- [86] M. Lewin, J. Sabin, The Hartree equation for infinitely many particles. I. Well-posedness theory, *Commun. Math. Phys.* (2013), in press.
- [87] M. Lewin, J. Sabin, The Hartree equation for infinitely many particles. II. Dispersion and scattering in 2D, preprint, <http://arxiv.org/abs/1310.0604>, 2013.
- [88] E. Lieb, R. Seiringer, Proof of Bose–Einstein condensation for dilute trapped gases, *Phys. Rev. Lett.* 88 (2002) 170409–1–170409–4.
- [89] E. Lieb, R. Seiringer, J.P. Solovej, The quantum-mechanical many-body problem: the Bose gas, in: *Perspectives in Analysis*, in: *Math. Phys. Stud.*, vol. 27, Springer, Berlin, 2005, pp. 97–183.
- [90] E. Lieb, R. Seiringer, J.P. Solovej, J. Yngvason, *The Mathematics of the Bose Gas and Its Condensation*, Oberwolfach Seminars, vol. 34, Birkhäuser Verlag, Basel, 2005.
- [91] E. Lieb, R. Seiringer, J. Yngvason, Bosons in a trap: a rigorous derivation of the Gross–Pitaevskii energy functional, *Phys. Rev. A* 61 (2000) 043602.
- [92] E. Lieb, R. Seiringer, J. Yngvason, A rigorous derivation of the Gross–Pitaevskii energy functional for a two-dimensional Bose gas. Dedicated to Joel L. Lebowitz, *Commun. Math. Phys.* 224 (1) (2001) 17–31.
- [93] J. Lührmann, Mean-field quantum dynamics with magnetic fields, *J. Math. Phys.* 53 (2) (2012) 022105, 19 pp.
- [94] A. Michelangeli, Equivalent definitions of asymptotic 100% BEC, *Il Nuovo Cimento B* 123 (2) (2008) 181–192.
- [95] A. Michelangeli, B. Schlein, Dynamical collapse of boson stars, *Commun. Math. Phys.* 311 (3) (2012) 645–687.
- [96] P. Pickl, Derivation of the time dependent Gross–Pitaevskii equation with external fields, *J. Stat. Phys.* 140 (1) (2010) 76–89.
- [97] P. Pickl, A simple derivation of mean field limits for quantum systems, *Lett. Math. Phys.* 97 (2) (2011) 151–164.
- [98] L. Pitaevskii, Vortex lines in an imperfect Bose gas, *Sov. Phys. JETP* 13 (1961) 451–454.
- [99] M. Reed, B. Simon, *Methods of Modern Mathematical Physics, Vol. 1: Functional Analysis*, revised and enlarged edition, Academic Press, San Diego, CA, 1975.
- [100] D.W. Robinson, Normal and locally normal states, *Commun. Math. Phys.* 19 (1970) 219–234.
- [101] I. Rodnianski, B. Schlein, Quantum fluctuations and rate of convergence towards mean field dynamics, *Commun. Math. Phys.* 291 (1) (2009) 31–61.
- [102] W. Rudin, *Functional Analysis*, second edition, McGraw–Hill Series in Higher Mathematics, McGraw–Hill Book Co., New York, NY, 1991.
- [103] B. Schlein, Derivation of effective evolution equations from microscopic quantum dynamics, in: D. Ellwood, I. Rodnianski, G. Staffilani, J. Wunsch (Eds.), *Evolution Equations*, in: *Clay Mathematics Proceedings*, vol. 17, 2013, 572 pp.
- [104] B. Simon, Convergence in trace ideals, *Proc. Am. Math. Soc.* 83 (1981) 39–43.
- [105] B. Simon, *Trace Ideals and Their Applications*, second edition, *Mathematical Surveys and Monographs*, vol. 120, American Mathematical Society, Providence, RI, 2005.
- [106] V. Sohinger, Local existence of solutions to randomized Gross–Pitaevskii hierarchies, *Trans. Am. Math. Soc.* (2014), in press.
- [107] V. Sohinger, G. Staffilani, Randomization and the Gross–Pitaevskii hierarchy, preprint, <http://arxiv.org/abs/1308.3714>, 2013.
- [108] H. Spohn, Kinetic equations from Hamiltonian dynamics, *Rev. Mod. Phys.* 52 (3) (1980) 569–615.
- [109] E. Størmer, Symmetric states of infinite tensor products of C^* algebras, *J. Funct. Anal.* 3 (1969) 48–68.
- [110] N. Strunk, Strichartz estimates for Schrödinger equations on irrational tori in two and three dimensions, preprint, <http://arxiv.org/abs/1401.6080>, 2014.
- [111] Z. Xie, Derivation of a nonlinear Schrödinger equation with a general power-type nonlinearity, preprint, <http://arxiv.org/abs/1305.7240>, 2013.