

ANNALES DE L'I. H. P., SECTION C

FRANK MERLE

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Annales de l'I. H. P., section C, tome 13, n° 5 (1996), p. 553-565

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Asymptotics for L^2 minimal blow-up solutions of critical nonlinear Schrödinger equation

by

Frank MERLE

Université de Cergy-Pontoise, Centre de Mathématiques
Avenue du Parc 8, Le Campus, 95033 Cergy-Pontoise, France

ABSTRACT. – In this note, we describe the behavior of a sequence $v_n : \mathbb{R}^N \rightarrow \mathbb{C}$ minimal in L^2 such that $\frac{1}{2} \int |\nabla v_n|^2 - \frac{1}{\frac{4}{N}+2} \int |v_n|^{\frac{4}{N}+2} \leq E_0$ and $|v_n|_{H^1} \rightarrow +\infty$.

RÉSUMÉ. – Dans cette note, on explicite de façon optimale le comportement d'une suite $v_n : \mathbb{R}^N \rightarrow \mathbb{C}$ de norme L^2 minimale telle que $\frac{1}{2} \int |\nabla v_n|^2 - \frac{1}{\frac{4}{N}+2} \int |v_n|^{\frac{4}{N}+2} \leq E_0$ et $|v_n|_{H^1} \rightarrow +\infty$.

In the present note, we are interested in the behavior of a sequence $v_n : \mathbb{R}^N \rightarrow \mathbb{C}$ of H^1 functions such that

$$(1) \quad \int |v_n|^2 = \int Q^2,$$

$$(2) \quad E(v_n) = \frac{1}{2} \int |\nabla v_n|^2 - \frac{1}{\frac{4}{N}+2} \int |v_n|^{\frac{4}{N}+2} \leq E_0,$$

$$(3) \quad \int |\nabla v_n|^2 \rightarrow +\infty,$$

where Q is the radial positive symmetric solution of the equation

$$(4) \quad \Delta v + |v|^{\frac{4}{N}} v = v.$$

(See references [1], [4], [7] for existence and uniqueness of Q .)

This problem is related to the asymptotics of minimal blow-up solutions in H^1 of the equation

$$(5) \quad iu_t = -\Delta u - k(x)|u|^{\frac{4}{N}}u \quad \text{and} \quad u(0) = \varphi,$$

where

$$(6) \quad \text{Max}_{x \in \mathbb{R}^N} k(x) = 1.$$

Indeed, for all $\varphi \in H^1$, there is a unique solution in H^1 on $[0, T]$ ([2], [4]) and

$$T = +\infty \text{ or } \lim_{t \rightarrow T} \int |\nabla u(t, x)|^2 = +\infty.$$

In addition, $\forall t$

$$(7) \quad \int |u(t, x)|^2 dx = \int |\varphi(x)|^2 dx$$

$$(8) \quad E_k(u(t)) = E_k(\varphi)$$

where

$$E_k(v) = \frac{1}{2} \int |\nabla v|^2 - \frac{1}{\frac{4}{N} + 2} \int k(x)|v|^{\frac{4}{N}+2}.$$

From [9]

$$(9) \quad \forall v \in H^1, \quad \frac{1}{\frac{4}{N} + 2} \int |v|^{\frac{4}{N}+2} \leq \frac{1}{2} \left(\frac{\int |v|^2}{\int |Q|^2} \right)^{\frac{2}{N}} \int |\nabla v|^2$$

and it follows from (6)-(9) ([6]) that

$$\text{if } |\varphi|_{L^2} < |Q|_{L^2}, \text{ then } T = +\infty.$$

Moreover under some conditions on $k(x)$, for any $\varepsilon > 0$ there are blow-up solutions $u_\varepsilon(t)$ such that

$$|u_\varepsilon(0)|_{L^2}^2 = |\varphi|_{L^2}^2 + \varepsilon \quad ([6]).$$

Thus the questions are about existence of minimal blow-up solution (that is such that $u(t)$ blows up in finite time and $\int |\varphi|^2 = \int Q^2$ and on the

form of these solutions. In the case where $k(x) \equiv 1$, the question has been completely solved (see Merle [5]). The general case is still open. We remark that from (6)-(9), if $u(t)$ is a blow up solution, the sequences $v_n = u(t_n)$ as $t_n \rightarrow T$ satisfies (1)-(3) and we ask about the constrains it imply on v_n .

The first result in this direction was obtained by Weinstein in [9]. Using the concentration compactness method, he showed that there is a $\theta_n \in \mathbb{R}$, $x_n \in \mathbb{R}^N$ such that

$$(10) \quad v_n = \lambda_n^{\frac{N}{2}} e^{i\theta_n} Q\left(\lambda_n^{\frac{N}{2}}(x - x_n)\right) + \varepsilon_n,$$

where

$$(11) \quad \lambda_n = \frac{|\nabla v_n|_{L^2}}{|\nabla Q|_{L^2}},$$

$$(12) \quad |\varepsilon_n|_{L^2} \xrightarrow{n \rightarrow +\infty} 0 \quad \text{and} \quad \frac{|\nabla \varepsilon_n|_{L^2}}{\lambda_n} \xrightarrow{n \rightarrow +\infty} 0.$$

In [5], Merle then showed that for all $R > 0$, there is a $c > 0$ such that

$$(13) \quad \int_{|x-x_n|>R} |\nabla v_n|^2 \leq c.$$

We now claim the following result

THEOREM. – Let (v_n) be a sequence of H^1 functions satisfying (1)-(3) and $\theta_n(x)$ be such that $v_n = |v_n|e^{i\theta_n}$.

i) *Phase estimates.* There is a $c > 0$ such that

$$\forall n, \quad \int |v_n|^2 |\nabla \theta_n|^2 \leq c.$$

ii) *asymptotics on the modulus.*

There is a $\varepsilon_n(x)$, $x_n \in \mathbb{R}^N$, and $c > 0$ such that

$$\forall x, \quad |v_n(x)| = \lambda_n^{\frac{N}{2}} Q(\lambda_n(x - x_n)) + \varepsilon_n(x)$$

where

$$|\nabla \varepsilon_n|_{L^2} \leq c, \quad |\varepsilon_n|_{L^2} \leq \frac{c}{\lambda_n} \quad \text{and} \quad \lambda_n \left(\frac{|\nabla Q|_{L^2}}{|\nabla v_n|_{L^2}} \right) \xrightarrow{n \rightarrow +\infty} 1.$$

Remark. – This Theorem simplifies some proofs in [5], [6]. The case where v_n is real valued is also related to similar problems for the generalized Korteweg-de Vries equation with critical nonlinearity.

Remark. – The Theorem implies in particular for blow-up solution of equation (5) $u(t, x) = |u(t, x)|e^{i\theta(t, x)}$ and $\int |u(t)|^2 = \int Q^2$ the phase gradient is uniformly bounded at the blow-up: there is a $c > 0$ such that

$$\forall 0 < t < T, \quad \int |u(t, x)|^2 |\nabla \theta(t, x)|^2 dx \leq c.$$

(Of course, we still have $\int |\nabla |u||^2(t, x) \xrightarrow[t \rightarrow T]{} +\infty$.)

Remark. – It is easy to check that the result is optimal. We remark that the residual term in the theorem $\varepsilon_n = O(1)$ (compared to $o(|\nabla v_n|_{L^2})$ in [9]).

Proof of the Theorem. – Let (v_n) a sequence of H^1 function satisfying (1)-(3) and $\theta_n(x)$ such that $v_n = |v_n|e^{i\theta_n}$. We have that

$$(14) \quad \frac{1}{2} \int |\nabla v_n|^2 = \frac{1}{2} \left\{ |\nabla |v_n||^2 + \int |v_n|^2 |\nabla \theta_n|^2 \right\}$$

and

$$(15) \quad E(v_n) = \frac{1}{2} \int |v_n|^2 |\nabla \theta_n|^2 + E(|v_n|).$$

The idea is to apply the variational identity (9) not with v_n but with $|v_n|$. Indeed, since $v_n \in H^1$ we have that $|v_n| \in H^1$. From (9) (applied with $|v_n|$)

$$(16) \quad \frac{1}{\frac{4}{n} + 2} \int |v_n|^{\frac{4}{n} + 2} \leq \frac{1}{2} \left(\frac{\int |v_n|^2}{\int Q^2} \right)^{\frac{2}{n}} \int |\nabla |v_n||^2 \leq \frac{1}{2} \int |\nabla |v_n||^2,$$

or equivalently

$$(17) \quad E(|v_n|) \geq 0.$$

Thus (2), (15), (17) imply that

$$(18) \quad \frac{1}{2} \int |v_n|^2 |\nabla \theta_n|^2 \leq E_0$$

$$(19) \quad E(|v_n|) \leq E_0.$$

Part i). – It is implied by (18).

Part ii). – We claim that it is as a consequence of (18)-(19). We prove it in three steps:

- from Weinstein's results, we first obtain rough estimates on $|v_n|$,
- we then choose appropriate approximation parameters,
- we conclude the proof using a convexity property in certain directions of E near Q (and use in a crucial way that $|v_n|$ is a real-valued function).

Step 1: First asymptotics. – Since

$$(20) \quad \int |\nabla v_n|^2 = \int |\nabla |v_n||^2 + \int |v_n|^2 |\nabla \theta_n|^2 \xrightarrow{n \rightarrow +\infty} +\infty,$$

and

$$\forall n, \quad \int |v_n|^2 |\nabla \theta_n|^2 \leq c,$$

we have

$$(21) \quad \int |\nabla |v_n||^2 \xrightarrow{n \rightarrow +\infty} +\infty.$$

Moreover,

$$(22) \quad \int |v_n|^2 = \int Q^2 \text{ and } E(|v_n|) \leq E_0.$$

We conclude from Weinstein's result on the existence of $\hat{x}_n, \hat{\varepsilon}_n$ such that

$$(23) \quad |v_n|(x) = \hat{\lambda}_n^{N/2} Q(\hat{\lambda}_n x - \hat{x}_n) + \hat{\varepsilon}_n(x)$$

where

$$(24) \quad \hat{\lambda}_n = \frac{|\nabla |v_n||_{L^2}}{|\nabla Q|_{L^2}},$$

$$(25) \quad |\nabla \hat{\varepsilon}_n|_{L^2} = o(\hat{\lambda}_n), \quad |\hat{\varepsilon}_n|_{L^2} = o(1).$$

In order to obtain better estimates on the rest (that is $|\nabla \varepsilon_n|_{L^2} \leq c$), we have to choose appropriate parameters λ_n, x_n and use the structure of the functional $E(\cdot)$ near Q .

Step 2: Choice of the parameters of approximation. – Let us first renormalize the problem. We consider

$$(26) \quad w_{n, \lambda_1, x_1}(x) = \left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^{\frac{N}{2}} |v_n| \left((\lambda_1 x + \hat{x}_n + x_1) \frac{1}{\hat{\lambda}_n} \right)$$

$$(27) \quad \tilde{\varepsilon}_{n,\lambda_1,x_1}(x) = \left(\frac{\lambda_1}{\hat{\lambda}_n}\right)^{\frac{N}{2}} \hat{\varepsilon}_n \left((\lambda_1 x + \hat{x}_n + x_1) \frac{1}{\hat{\lambda}_n} \right).$$

We have from (23),

$$(28) \quad w_{n,\lambda_1,x_1}(x) = \lambda_1^{N/2} Q(\lambda_1 x + x_1) + \tilde{\varepsilon}_{n,\lambda_1,x_1}(x)$$

where

$$(29) \quad \frac{|\nabla \tilde{\varepsilon}_n|_{L^2}}{\lambda_1} + |\tilde{\varepsilon}_n|_{L^2} \xrightarrow{n \rightarrow +\infty} 0.$$

We write (28) as follows

$$w_{n,\lambda_1,x_1}(x) = Q(x) + \varepsilon_{n,\lambda_1,x_1}(x)$$

where

$$(30) \quad \varepsilon_{n,\lambda_1,x_1}(x) = \left[\lambda_1^{N/2} Q(\lambda_1 x + x_1) - Q(x) \right] + \tilde{\varepsilon}_{n,\lambda_1,x_1}(x).$$

From the implicit function Theorem, we derive easily for $|\tilde{\varepsilon}_n|_{H^1}$ small enough the existence of $\lambda_{1,n}, x_{1,n}$ such that

$$(31) \quad \forall i = 1, \dots, N, \quad \int \varepsilon_{n,\lambda_{1,n},x_{1,n}} x_i Q = 0$$

$$(32) \quad \int \varepsilon_{n,\lambda_{1,n},x_{1,n}} |x|^2 Q = 0.$$

Moreover, from (29)

$$(33) \quad (\lambda_{1,n}, x_{1,n}) \xrightarrow{n \rightarrow +\infty} (1, 0).$$

Indeed, let us note

$$\begin{aligned} \text{for } i = 1, \dots, N, \quad \rho_i(\lambda_1, x_1) &= \int \varepsilon_{n,\lambda_1,x_1} x_i Q, \\ \rho_{N+1}(\lambda_1, x_1) &= \int \varepsilon_{n,\lambda_1,x_1} |x|^2 Q. \end{aligned}$$

From (30), we have

$$\begin{aligned} \frac{\partial \varepsilon_{n,1,0}}{\partial x_{1,i}} &= \partial_i Q + \partial_i \tilde{\varepsilon}_{n,1,0} \\ \frac{\partial \varepsilon_{n,1,0}}{\partial \lambda_1} &= \frac{N}{2} Q + x \cdot \nabla Q + \left(\frac{N}{2} \tilde{\varepsilon}_{n,1,0} + x \cdot \nabla \tilde{\varepsilon}_{n,1,0} \right), \end{aligned}$$

where $x_1 = (x_{1,1}, \dots, x_{1,N})$.

Therefore, from (29) and integration by parts,

– for $i = 1, \dots, N$, and $j = 1, \dots, N$,

$$\begin{aligned} \frac{\partial \rho_i}{\partial x_{1,j}}(1, 0) &= \int \partial_j Q x_i Q + o(1) = -2\delta_{i,j} \int Q^2 + o(1), \\ \frac{\partial \rho_i}{\partial \lambda_1}(1, 0) &= \int \left(\frac{N}{2} Q + x \cdot \nabla Q \right) x_i Q + o(1) = o(1), \\ \frac{\partial \rho_{N+1}}{\partial x_{1,j}}(1, 0) &= \int \partial_j Q |x|^2 Q + o(1) = o(1), \\ \frac{\partial \rho_{N+1}}{\partial \lambda_1}(1, 0) &= \int \left(\frac{N}{2} Q + x \cdot \nabla Q \right) |x|^2 Q + o(1) \\ &= \frac{N}{2} \int |x|^2 Q^2 - \frac{N}{2} \int |x|^2 Q^2 - \frac{1}{2} \int x \cdot x Q^2 + o(1) \\ &= -\frac{1}{2} \int |x|^2 Q^2 + o(1). \end{aligned}$$

Therefore, the implicit function theorem implies the existence of $(\lambda_{1,n}, x_{1,n})$ such that (31)-(33) hold.

In conclusion, we have proved the following. There exist $(\lambda_{1,n}, x_{1,n}) \xrightarrow[n \rightarrow +\infty]{} (1, 0)$ such that

$$(34) \quad w_{n, \lambda_{1,n}, x_{1,n}}(x) = Q(x) + \varepsilon_{n, \lambda_{1,n}, x_{1,n}}(x)$$

where

$$(35) \quad \forall i = 1, \dots, n, \quad \int \varepsilon_{n, \lambda_{1,n}, x_{1,n}} x_i Q = 0,$$

$$(36) \quad \int \varepsilon_{n, \lambda_{1,n}, x_{1,n}} |x|^2 Q = 0,$$

$$(37) \quad \|\varepsilon_{n, \lambda_{1,n}, x_{1,n}}\|_{H^1} \xrightarrow[n \rightarrow +\infty]{} 0.$$

We now note

$$\begin{aligned} w_n &= w_{n, \lambda_{1,n}, x_{1,n}}, \\ \varepsilon_n &= \varepsilon_{n, \lambda_{1,n}, x_{1,n}}. \end{aligned}$$

Step 3 : Conclusion of the proof. – Geometry of energy functions at Q .

We now use convexity properties of a functional (related to E) and the fact $\int w_n^2 = \int Q^2$ to conclude the proof. Let

$$H(v) = \frac{1}{2} \int |\nabla v|^2 - \frac{1}{\frac{4}{N} + 2} \int |v|^{\frac{4}{N} + 2} + \frac{1}{2} \int v^2 = E(v) + \frac{1}{2} \int v^2,$$

and

$$H_2(v) = \frac{1}{2} \int |\nabla v|^2 - \frac{4}{N} \int Q^{\frac{4}{N}} v^2 + \frac{1}{2} \int v^2.$$

We know that Q is a critical point of H , and H_2 is the quadratic part of H near Q (where v is real-valued). Moreover, it is classical that for $|\varepsilon|_{H^1} \leq 1$,

$$(38) \quad H(Q + \varepsilon) - H(Q) = H_2(\varepsilon) + \tilde{H}_2(\varepsilon)$$

where $|\tilde{H}_2(\varepsilon)| = o(|\varepsilon|_{H^1}^2)$.

From a result of Weinstein [8] (see also Kwong [4]), we have the following convexity property of H_2 at Q .

PROPOSITION 1. See [8]. – (Directions of convexity of H at Q in the set of real-valued functions.)

There is a constant $c_1 > 0$ such that $\forall \varepsilon \in H^1$

If

$$(i) \quad \forall i = 1, \dots, N, \quad \int \varepsilon x_i Q = 0$$

$$(ii) \quad \int \varepsilon |x|^2 Q = 0,$$

$$(iii) \quad \int \varepsilon Q = 0,$$

then

$$H_2(\varepsilon) \geq c_1 \left(\int |\nabla \varepsilon|^2 + \varepsilon^2 \right) = c_1 |\varepsilon|_{H^1}^2.$$

Remark. – We have here a strict convexity property (up to the invariance of the equation) except in the direction Q which is not true for the quadratic part of H for complex valued functions (see [8]).

Remark. – This proposition is optimal. Other functions can be chosen also.

Using now crucially estimates on the L^2 norm, we obtain the following

PROPOSITION 2 (Control of the Q direction by the L^2 norm).

Assume

$$(i) \quad \forall i = 1, \dots, N, \quad \int \varepsilon x_i Q = 0,$$

$$(ii) \quad \int \varepsilon |x|^2 Q = 0,$$

$$(iii) \quad \int (Q + \varepsilon)^2 = \int Q^2,$$

then there are $c_1 > 0$ and $c_2 > 0$ such that

$$|\nabla \varepsilon|_{L^2} + |\varepsilon|_{L^2} \leq c_2 \text{ implies } H_2(\varepsilon) \geq c_1 (|\nabla \varepsilon|_{L^2}^2 + |\varepsilon|_{L^2}^2).$$

Remark. – We need control on 3 directions to obtain estimates on $|\varepsilon|_{H^1}$ with $H_2(Q + \varepsilon)$. Two directions can be controlled using the invariance of the equation. The last one is controlled by the condition of minimality on the L^2 norm (among sequence satisfying (2)).

Proof of Proposition 2. – Let us note

$$\tilde{H}_2(v_1, v_2) = \frac{1}{2} \int \nabla v_1 \nabla v_2 - \frac{4}{N} \int Q^{\frac{4}{N}} v_1 v_2 + \frac{1}{2} \int v_1 v_2.$$

We can write

$$\varepsilon = z + aQ + b|x|^2Q$$

with

$$\int zQ = \int zx_iQ = \int z|x|^2Q = 0 \quad \text{for } i = 1, \dots, N.$$

Indeed a and b have to satisfy

$$\begin{aligned} \int \varepsilon Q &= a \int Q^2 + b \int |x|^2 Q^2 \\ 0 &= a \int |x|^2 Q^2 + b \int |x|^4 Q^2 \end{aligned}$$

or equivalently

$$b = -a \left(\frac{\int |x|^2 Q^2}{\int |x|^4 Q^2} \right)$$

$$a \left(\frac{\int Q^2 \int |x|^4 Q^2 - (\int |x|^2 Q^2)^2}{\int |x|^4 Q^2} \right) = \int \varepsilon Q$$

(which has always a solution since from the Schwarz inequality and the fact $|x|^2 Q \neq Q$, $\int |x|^2 Q^2 < (\int Q^2 \int |x|^4 Q^2)^{1/2}$).

On the other hand, we have from $\int (Q + \varepsilon)^2 = \int Q^2$

$$2 \int Q \varepsilon = - \int \varepsilon^2$$

$$2 \left(a \int Q^2 + b \int |x|^2 Q^2 \right) = - \int \varepsilon^2$$

$$2a \left(\frac{\int Q^2 \int |x|^4 Q^2 - (\int |x|^2 Q^2)^2}{\int |x|^4 Q^2} \right) = - \int z^2 + O(a^2 + b^2)$$

or equivalently,

$$ac_0 = - \int z^2 + O(a^2) \quad \text{where } c_0 \neq 0$$

which implies that

$$a = O \left(\int z^2 \right) \quad \text{and} \quad b = O \left(\int z^2 \right)$$

and for $|\varepsilon|_{H^1}$ small enough

$$|\varepsilon|_{H^1}^2 \geq |z|_{H^1}^2 \geq \frac{1}{2} |\varepsilon|_{H^1}^2.$$

On the other hand, by bilinearity and Proposition 1, we have for $|\varepsilon|_{H^1}$ small enough

$$\begin{aligned} H_2(\varepsilon) &= H_2(z) + 2a\tilde{H}_2(z, Q) + 2b\tilde{H}_2(z, |x|^2 Q) + 2ab\tilde{H}_2(Q, |x|^2 Q) \\ &\quad + a^2 H_2(Q) + b^2 H_2(|x|^2 Q) \\ &\geq H_2(z) - c(|z|_{H^1}(|a| + |b|) + a^2 + b^2) \\ &\geq H_2(z) - c(|z|_{H^1}^3 + |z|_{H^1}^4) \\ &\geq c_1(|z|_{H^1}^2) - c(|z|_{H^1} + |z|_{H^1}^4) \\ &\geq \frac{c_1}{2} |z|_{H^1}^2 \\ &\geq \frac{c_1}{4} |\varepsilon|_{H^1}^2. \end{aligned}$$

This concludes the proof of Proposition 2.

As a corollary of Proposition 2 and (38), we have

COROLLARY. – *There are $c_1 > 0$ and $c_2 > 0$ such that if*

$$(i) \quad \forall i = 1, \dots, N, \quad \int \varepsilon x_i Q = 0$$

$$(ii) \quad \int \varepsilon |x|^2 Q = 0$$

$$(iii) \quad \int (Q + \varepsilon)^2 = \int Q^2$$

$$(iv) \quad |\nabla \varepsilon|_{L^2} + |\varepsilon|_{L^2} \leq c_2$$

then

$$H(Q + \varepsilon) - H(Q) \geq c_1 \left(\int \nabla \varepsilon^2 + \int \varepsilon^2 \right).$$

We now apply the corollary. If $w_n = Q + \varepsilon_n$, we have
 - $|\varepsilon_n|_{H^1} \xrightarrow[n \rightarrow +\infty]{} 0$, and in particular there is n_0 such that

$$\forall n \geq n_0, \quad |\varepsilon_n|_{H^1} \leq c_2,$$

$$- \forall i = 1, \dots, N, \quad \int \varepsilon_n x_i Q = 0.$$

In addition,

$$\begin{aligned} H(Q) &= \frac{1}{2} \int |\nabla Q|^2 - \frac{1}{\frac{4}{N} + 2} \int |Q|^{\frac{4}{N} + 2} + \frac{1}{2} \int Q^2 \\ &= E(Q) + \frac{1}{2} \int Q^2 \\ &= \frac{1}{2} \int Q^2 \end{aligned}$$

(since the Pohozaev identity for equation (4) yields $E(Q) = 0$), and

$$\begin{aligned} H(Q + \varepsilon_n) &= H(w_n) = H \left(\left(\frac{\lambda_1}{\hat{\lambda}_n} \right)^{\frac{N}{2}} |v_n| \left(\frac{x \lambda_1}{\hat{\lambda}_n} + \hat{x}_n + x_1 \right) \right) \\ &= E \left(\left(\frac{\lambda_1}{\hat{\lambda}_n} \right)^{\frac{N}{2}} |v_n| \left(\frac{x \lambda_1}{\hat{\lambda}_n} \right) \right) + \frac{1}{2} \int |v_n|^2 \\ &= \left(\frac{\lambda_1}{\hat{\lambda}_n} \right)^2 E(|v_n|) + \frac{1}{2} \int Q^2. \end{aligned}$$

Therefore $\forall n \geq n_0$

$$(40) \quad \left(\frac{\lambda_1}{\hat{\lambda}_n} \right)^2 E(|v_n|) > c_3 \left(|\varepsilon_n|_{H^1}^2 \right)$$

or equivalently from (19), (24) and the fact that $\lambda_1 \rightarrow 1$,

$$(41) \quad |\varepsilon_n|_{H^1}^2 \leq \frac{c}{2} \frac{1}{\int |\nabla |v_n||^2} \leq c \frac{1}{\int |\nabla v_n|^2},$$

where c is independent of n . Thus,

$$w_n = Q + \varepsilon_n,$$

with

$$(42) \quad |\varepsilon_n|_{H^1}^2 \leq \frac{c}{\int |\nabla v_n|^2}.$$

Therefore from (26), there is x_n such that

$$(43) \quad |v_n|(x) = \left(\frac{\hat{\lambda}_n}{\lambda_1} \right)^{\frac{N}{2}} Q \left(x \left(\frac{\hat{\lambda}_n}{\lambda_1} \right) + x_n \right) + \left(\frac{\hat{\lambda}_n}{\lambda_1} \right)^{\frac{N}{2}} \varepsilon_n \left(\frac{\hat{\lambda}_n}{\lambda_1} x + x_n \right).$$

We remark that from (19), (42), the fact that $\lambda_1 \rightarrow 1$,

$$\begin{aligned} \frac{\hat{\lambda}_n}{\lambda_1} \frac{1}{\left(\int \nabla v_n^2 \right)^{\frac{1}{2}}} &= \frac{1}{\lambda_1} \left(\frac{\int |\nabla |v_n||^2}{\int |\nabla v_n|^2 \int |\nabla Q|^2} \right)^{\frac{1}{2}} \xrightarrow{n \rightarrow +\infty} 1, \\ \left| \left(\frac{\hat{\lambda}_n}{\lambda_1} \right)^{\frac{N}{2}} \varepsilon_n \left(\frac{\hat{\lambda}_n}{\lambda_1} x + x_n \right) \right|_{L^2}^2 &= |\varepsilon_n|_{L^2}^2 \leq \frac{c}{\int |\nabla v_n|^2}, \\ \left| \nabla \left(\frac{\hat{\lambda}_n}{\lambda_1} \right)^{\frac{N}{2}} \varepsilon_n \left(\frac{\hat{\lambda}_n}{\lambda_1} x + x_n \right) \right|_{L^2}^2 &\leq c \left(\frac{\hat{\lambda}_n^2}{\int |\nabla v_n|^2} \right) \leq c \end{aligned}$$

conclude the proof of the Theorem.

REFERENCES

- [1] H. BERESTYCKI and P. L. LIONS, Nonlinear scalar field equations I. Existence of a ground state; II. Existence of infinitely many solutions, *Arch. Rational Mech. Anal.*, Vol. **82**, 1983, pp. 313-375.
- [2] J. GINIBRE and G. VELO, On a class of nonlinear Schrödinger equations I, II. The Cauchy problem, general case, *J. Funct. Anal.*, Vol. **32**, 1979, pp. 1-71.
- [3] T. KATO, On nonlinear Schrödinger equations, *Ann. Inst. Henri Poincaré, Physique Théorique*, Vol. **49**, 1987, pp. 113-129.
- [4] M. K. KWONG, Uniqueness of positive solutions of $\Delta u - u + u^p = 0$ in \mathbb{R}^N , *Arch. Rational Mech. Anal.*, Vol. **105**, 1989, pp. 243-266.
- [5] F. MERLE, Determination of blow-up solutions with minimal mass for Schrödinger equation with critical power, *Duke J.*, Vol. **69**, 1993, pp. 427-454.
- [6] F. MERLE, Nonexistence of minimal blow-up solutions of equations $iu_t = -\Delta u - k(x)|u|^{4/N}u$ in \mathbb{R}^N , preprint.
- [7] W. A. STRAUSS, Existence of solitary waves in higher dimensions, *Commun. Math. Phys.*, Vol. **55**, 1977, pp. 149-162.
- [8] M. I. WEINSTEIN, Modulational stability of ground states of the nonlinear Schrödinger equations, *SIAM J. Math. Anal.*, Vol. **16**, 1985, pp. 472-491.
- [9] M. I. WEINSTEIN, Nonlinear Schrödinger equations and sharp interpolation estimates, *Commun. Math. Phys.*, Vol. **87**, 1983, pp. 567-576.

*(Manuscript received October 25, 1994;
revised version received January 26, 1995.)*