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# Odd anharmonic oscillators and shape resonances

by

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ABSTRACT. — By adapting a stability argument of Vock and Hunziker it is proved that the Borel sum of the Rayleigh-Schrödinger perturbation expansion of any odd anharmonic oscillator  $p^2 + x^2 + gx^{2k+1}$ , k = 1, 2, ..., is the limit of a sequence of resonances in the standard sense of dilation analyticity. The same method yields, for a suitable class of potentials, an existence proof of shape resonances.

RÉSUMÉ. — En adaptant un argument de stabilité dû à Vock et Hunziker, on montre que la somme de Borel de la série de perturbations de Rayleigh Schrödinger pour un oscillateur anharmonique impair  $p^2 + x^2 + gx^{2k+1}$ ,  $k=1,2,\ldots$ , est la limite d'une suite de résonances au sens usuel de l'analycité par dilatation. La même méthode fournit, pour une classe convenable de potentiels, une preuve d'existence des « résonances de forme ».

#### 1. INTRODUCTION

The purpose of this paper is to complete an earlier investigation [2] on the spectral and perturbation theory of any odd anharmonic oscillator  $p^2 + x^2 + gx^{2k+1}$ ,  $k = 1, 2, \ldots$  The aim is to justify the standard physical picture of these problems (see e. g. Davydov [4] and, for a more recent discussion of this order of ideas, Coleman [3]) given in terms of unstable states generated by the shape of the potential, which should be represented by the Rayleigh-Schrödinger perturbation expansion, in the framework of dilation analyticity and Borel summability.

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Now the differential expression  $H=-d^2/dx^2+x^2+gx^{2^{k+1}}, k\in \mathbb{N}, g>0$  is, when defined on  $C_0^\infty(R)$ , a symmetric operator in  $L^2(R)$  admitting infinitely many distinct self-adjoint extensions, each of them having discrete spectrum. Therefore the differential expression H was first realized as an operator in  $L^2(R)$  for g complex,  $0<\arg(g)<\pi$ , where it represents a holomorphic family of type A with compact resolvents. It turns out that H converges in the generalized sense to  $p^2+x^2$  as  $|g|\to 0, 0<\arg(g)<\pi$ , so that  $\sigma_d(H)\neq\emptyset$  for |g| suitably small, and that the (divergent) Rayleigh-Schrödinger perturbation expansion near any simple eigenvalue of  $p^2+x^2$  is Borel summable to the nearby eigenvalue of H. When analytically continued to  $g\in R$  the eigenvalues of H can be interpreted as resonances of the problem: in particular they are shown to be second sheet poles of a unique generalized resolvent (see e. g. [1] for a definition) of the symmetric operator H.

From this fact one can ask whether it is possible to approximate  $x^2 + gx^{2k+1}$  by dilation analytic potentials so that such generalized resonances become the limit of resonances of self-adjoint operators in the standard sense of dilation analyticity.

This is just what we prove in Sect. 2 of this paper, by means of the following family of approximating potentials

$$V_{\alpha}(x) = (x^2 + gx^{2k+1})(\alpha^2 x^{4k+2} + 1)^{-1/2}, \quad \text{as } \alpha \to 0.$$

Taking into account the result in [2] on the discreteness of the spectrum of  $H(0,\theta)=e^{-\theta}p^2+e^{\theta}x^2+ge^{(2k+1)\theta/2}x^{2k+1}$  for  $\text{Im }\theta>0$  and g>0, the problem reduces to a stability result for the eigenvalues of  $H(0,\theta)$  with respect to the family  $H(\alpha,\theta)=e^{-\theta}p^2+V_{\alpha}(e^{\theta/2}x)$  as  $\alpha\to 0$ . In this connection we apply the stability theory recently developed by Hunziker and Vock [6]: to this end we need to extend the operator class explicitly provided for applications ([6], Example 5) and prove the uniform boundedness of the resolvents which is not given by a control of the numerical ranges.

Since  $V_{\alpha}(x)$  exhibits the typical shape of a barrier, this result examplifies at the same time how in some cases the so-called shape resonances [3] can be proved to exist in the standard sense of dilation analyticity. This last result, obtained by adapting to  $L^2(R_+)$  our arguments valid on  $L^2(R)$ , is briefly described after proving the existence of resonances for  $p^2 + V_{\alpha}(x)$  by the above mentioned stability argument.

## 2. STABILITY OF EIGENVALUES AND APPROXIMATING RESONANCES

The following theorem allows to enlarge the operator class given in [6], Example 5, in which our model is not included.

THEOREM 2.1. — Let  $G \subset L_{1\infty}^2(R^{\nu})$  and  $\{A(V)\}_{V \in G}$  be a class of closable operators in  $L^2(R^{\nu})$  of the form  $A(V) = p^2 + V$ , with  $C_0^{\infty}(R^{\nu})$  as a core and such that for any  $V \in G$  there exists a function  $\gamma_V \colon R^{\nu} \to R$  satisfying the following conditions:

- (1)  $\operatorname{Re}\langle u \cos \gamma_{V}, p^{2}u \rangle + \operatorname{Im}\langle u \sin \gamma_{V}, p^{2}u \rangle$  $\leq \xi \left[ \operatorname{Re}\langle u \cos \gamma_{V}, A(V)u \rangle + \operatorname{Im}\langle u \sin \gamma_{V}, A(V)u \rangle + \eta \langle u, u \rangle \right], \forall u \in C_{0}^{\infty}(\mathbb{R}^{\nu})$
- (2)  $\inf_{x \in \mathbb{R}^{\nu}} \cos \gamma_{\mathbf{V}}(x) \ge C_0 > 0$
- $(3) \quad \gamma_{V}\!\in\! H^{1}_{loc}(R^{\nu}) \quad \text{and} \quad \|\, \nabla\cos\gamma_{V}\,\|_{\infty} \leq C_{1}, \, \|\, \nabla\sin\gamma_{V}\,\|_{\infty} \leq C_{1},$

where  $\xi$ ,  $\eta$ ,  $C_0$  and  $C_1$  are positive constants independent of V. Here  $H^1_{loc}(R^{\nu})$  denotes, as usual, the space of locally square integrable functions which admit first order generalized derivatives belonging to  $L^2_{loc}(R^{\nu})$ ;

$$p^2 = -\sum_{k=1}^{\nu} \partial^2/\partial x_k^2.$$

Then there is b > 0 independent of  $V \in G$  such that  $\forall u \in C_0^{\infty}(\mathbb{R}^{\nu})$ 

$$\|(1+p^2)^{1/2}u\| \le b(\|u\| + \|A(V)u\|)$$
(2.1)

and there exists a sequence of multiplication operators  $\{M_n\}_{n\in\mathbb{N}}$  in  $L^2(\mathbb{R}^{\nu})$  satisfying (i) and (ii) of Hypothesis 3 in [6]. Such Hypothesis is completely fulfilled if, moreover,  $\Delta \neq \emptyset$  ( $\Delta$  as defined in [6]).

Proof. — It suffices to prove the following inequality

$$\langle u, p^2 u \rangle \le \xi_1 [\text{Re} \langle u \cos \gamma_V, A(V)u \rangle + \text{Im} \langle u \sin \gamma_V, A(V)u \rangle + \eta_1 \langle u, u \rangle]$$
 (2.2)

 $\forall u \in C_0^{\infty}(\mathbb{R}^v)$ , with  $\xi_1, \eta_1$  independent of V,  $\xi_1 > 0$ ,  $\eta_1 > 0$ .

In fact (2.1) easily follows from (2.2) by Schwarz' inequality. Without loss of generality we can assume ||u|| = 1. Then

Re 
$$\langle u \cos \gamma_V, p^2 u \rangle + \text{Im} \langle u \sin \gamma_V, p^2 u \rangle$$

$$= \sum_{k=1}^{\nu} \langle (\cos \gamma_{\mathbf{V}}) p_{k} u, p_{k} u \rangle + \sum_{k=1}^{\nu} \operatorname{Re} \langle u p_{k} \cos \gamma_{\mathbf{V}}, p_{k} u \rangle + \sum_{k=1}^{\nu} \operatorname{Im} \langle u p_{k} \sin \gamma_{\mathbf{V}}, p_{k} u \rangle$$

$$\geq C_{0} \langle u, p^{2} u \rangle - \sum_{k=1}^{\nu} |\langle u p_{k} \cos \gamma_{\mathbf{V}}, p_{k} u \rangle| - \sum_{k=1}^{\nu} |\langle u p_{k} \sin \gamma_{\mathbf{V}}, p_{k} u \rangle|$$

$$\geq C_{0} \langle u, p^{2} u \rangle - 2C_{1} \sum_{k=1}^{\nu} \langle u, p_{k}^{2} u \rangle^{1/2}$$

$$\geq C_{0} \langle u, p^{2} u \rangle - C_{1} \lambda^{-1} \sum_{k=1}^{\nu} 2\lambda \langle u, p_{k}^{2} u \rangle^{1/2},$$

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for any  $\lambda > 0$ ,

$$\geq C_0 \langle u, p^2 u \rangle - C_1 \lambda^{-1} \sum_{k=1}^{\nu} [\lambda^2 + \langle u, p_k^2 u \rangle] = (C_0 - C_1 \lambda^{-1}) \langle u, p^2 u \rangle - C_1 \lambda \nu$$

where the second inequality follows from Schwarz' inequality and assumption (3). Here, as usual,  $p_k = -i\partial/\partial x_k$ , k = 1, 2, ..., v. Now we choose  $\lambda > 0$  such that  $C_0 - C_1 \lambda^{-1} > 0$  and (2.2) follows from assumption (1).

In order to show that Hypothesis 3 in [6] is satisfied we proceed as in Lemma 3.2 of [6], choosing  $M_n = 1 - \chi_n$ , where  $\chi_n(x) = \chi(x/n)$ ,  $\chi \in C_0^{\infty}(\mathbb{R}^{\nu})$ ,  $\chi(x) = 1$  for |x| < 1, and the theorem is thus proved.

Consider in L<sup>2</sup>(R) the formal differential expression

$$H = p^2 + x^2 + gx^{2k+1}, \quad k = 1, 2, ...; \quad g \in \mathbb{R} \setminus \{0\}$$

where, as usual,  $p^2 = -d^2/dx^2$ . Since specular arguments apply for g < 0, we can take g > 0 without loss of generality.

By  $\overline{H}$  we mean the closure of the symmetric operator defined by H on  $C_0^\infty(R)$ . Set

$$H(0, \theta) = e^{-\theta}p^2 + e^{\theta}x^2 + ge^{(2k+1)\theta/2}x^{2k+1}$$

For  $\theta \in \mathbb{R}$ ,  $H(0, \theta)$  can be written as  $U(\theta)HU(\theta)^{-1}$ , where  $U(\theta)f(x)=e^{\theta/4}f(e^{\theta/2}x)$  defines the group of unitary dilations  $U(\theta)$  in  $L^2(\mathbb{R})$ .

It is well known ([8], [9]) that  $\overline{H}$  admits infinitely many distinct self-adjoint extensions. The following proposition is proved in [2].

Theorem 2.2. — i) Each self-adjoint extension of  $\overline{H}$  has discrete spectrum.

- ii) Let  $\theta$  be complex,  $0 < \operatorname{Im} \theta < \min(\pi/4, 2\pi/(2k+3))$ . Then  $H(0, \theta)$  represents a holomorphic family of type A of compact resolvent operators in  $L^2(R)$  if defined on  $D(p^2) \cap D(x^{2k+1}) = H^2(R) \cap L^2_{(2k+1)/2}(R)$  where  $H^2(R)$  and  $L^2_{(2k+1)/2}(R)$  denote the usual Sobolev and weighted  $L^2$  spaces respectively.
  - iii) For any compact  $\Gamma$  contained in the strip

$$0 < \text{Im } \theta < \min (\pi/4, 2\pi/(2k+3)),$$

 $\exists g(\Gamma) > 0$  such that,  $\forall g < g(\Gamma)$ ,  $H(0, \theta)$  admits eigenvalues (independent of  $\theta$ ) for  $\theta \in \Gamma$ .

iv) There is a dense set S of dilation analytic vectors

(for 
$$|\text{Im }\theta| < \min(\pi/4, 2\pi/(2k+3))$$

and a generalized resolvent  $\mathcal{R}(E)$  of the closed symmetric operator  $\overline{H}$  (see e. g. [1] for a definition) such that if  $\psi \in S$  the function

$$f_{\mu}(E) = \langle \mathscr{R}(E)\psi, \psi \rangle$$

which is a priori defined as an analytic function in the upper half-plane  $\{E \mid Im E > 0\}$  has a meromorphic continuation to the lower half-plane  $\{E \mid Im E < 0\}$  across the real axis. The set of singularities  $\{E \mid f_{\psi} \text{ has a pole at E for some } \psi \in S\}$  coincides with

$$\bigcup_{\theta} \{ \sigma(H(0,\theta)) \mid 0 < \text{Im } \theta < \min(\pi/4, 2\pi/(2k+3)) \},$$

and  $\mathcal{R}(E)$  is uniquely determined by the identity

$$\langle \mathcal{R}(\mathbf{E})\phi, \psi \rangle = \langle [\mathbf{H}(0, \theta) - \mathbf{E}]^{-1}\phi(\theta), \psi(\overline{\theta}) \rangle,$$
  
0 < Im \theta < \min (\pi/4, 2\pi/(2k+3)); \phi, \psi \in \mathbf{E}, \psi(\theta)(\theta) = e^{\theta/4}\psi(e^{\theta/2}x); \text{Im } \mathbf{E} > 0.

For  $\alpha > 0$  and  $\theta$  complex let us define

$$H(\alpha, \theta) = e^{-\theta} [p^2 + V_{\alpha}(x, \theta)] = e^{-\theta} \left[ p^2 + \frac{e^{2\theta} x^2 + ge^{(2k+3)\theta/2} x^{2k+1}}{(\alpha^2 e^{(2k+1)\theta} x^{4k+2} + 1)^{1/2}} \right]$$
 (2.3)

on the domain  $D(H(\alpha, \theta)) = H^2(R)$ , as an operator in  $L^2(R)$ . It is easily seen that the complex-valued potential  $V_{\alpha}(x, \theta)$  is an analytic function of  $\theta$  in the strip  $|\operatorname{Im} \theta| < \pi/(4k + 2)$ . Since  $V_{\alpha}(x, \theta)$  is bounded, by standard arguments (see e. g. Kato [7])  $H(\alpha, \theta)$  is a holomorphic family of type A for  $|\operatorname{Im} \theta| < \pi/(4k + 2)$ .

Theorem 2.3. — Let 
$$|\operatorname{Im} \theta| < \pi/(4k + 2)$$
. Then

$$\sigma_{ess}(\mathbf{H}(\alpha, \theta)) = \{ z = -g\alpha^{-1} + \lambda e^{-\theta} \mid \lambda \ge 0 \} \cup \{ z = g\alpha^{-1} + \lambda e^{-\theta} \mid \lambda \ge 0 \}.$$

Proof. — Set  $D' = \{u \in H^2(R) \mid \text{ there is a compact subset } K(u) \text{ of } R \text{ such that supp } u \subset K(u) \}$  and  $H'(\alpha, \theta) = H(\alpha, \theta) \upharpoonright D'$ . Now we restrict D' by imposing the additional conditions u(0) = u'(0) = 0; let D'' denote the manifold obtained by this restriction and  $H''(\alpha, \theta) = H'(\alpha, \theta) \upharpoonright D''$ . Then  $H''(\alpha, \theta) = H'_1(\alpha, \theta) \oplus H'_2(\alpha, \theta)$  (for the « decomposition » method see [5] or [8]) where  $H'_1(\alpha, \theta)$  and  $H'_2(\alpha, \theta)$  are generated in  $L^2(R_-)$  and  $L^2(R_+)$  respectively by the differential expression  $e^{-\theta}[p^2 + V_{\alpha}(x, \theta)]$  in the same way as  $H'(\alpha, \theta)$  was. Then  $\sigma_{ess}(\overline{H}''(\alpha, \theta)) = \sigma_{ess}(\overline{H}'_1(\alpha, \theta)) \cup \sigma_{ess}(\overline{H}'_2(\alpha, \theta))$ , where  $\overline{H}''(\alpha, \theta)$  denotes any finite dimensional closed extension of  $H''(\alpha, \theta)$  and similarly for  $\overline{H}'_1$  and  $\overline{H}'_2$  (again see [5]): in particular  $\overline{H}''(\alpha, \theta)$  can be  $H(\alpha, \theta)$ . Now,  $(e^{-\theta}V_{\alpha}(x, \theta) - g\alpha^{-1})$  is relatively compact with respect to  $(e^{-\theta}p^2 + g\alpha^{-1})$  in  $L^2(R_+)$ , since  $(e^{-\theta}V_{\alpha}(x, \theta) - g\alpha^{-1}) \to 0$  as  $x \to +\infty$  (see [7]), so that

$$\sigma_{ess}(\overline{\mathbf{H}}_2'(\alpha,\theta)) = \{ z = g\alpha^{-1} + \lambda e^{-\theta} \mid \lambda \ge 0 \}.$$

An analogous argument works for  $\overline{H}_1'(\alpha,\theta)$  and the theorem is thus proved. From now on,  $\theta$  will be fixed in the strip  $0<\operatorname{Im}\theta<\pi/(4k+2)$ ; for any  $\alpha\geq 0$  we shall use the simplified notation  $H(\alpha)$  to denote the differential operator defined by  $H(\alpha)=e^{\theta}H(\alpha,\theta)=p^2+V_{\alpha},\ D(H(0))=H^2(R)\cap L^2_{(2k+1)/2}(R)$  and  $D(H(\alpha))=H^2(R)$   $\forall \alpha>0$ .

LEMMA 2.4. — The operator family  $\{H(\alpha)\}_{\alpha\geq 0}$  satisfies assumptions (1), (2), (3) of Theorem 2.1.

*Proof.* — For simplicity we assume  $\theta$  of the form  $\theta = i\theta_0$ , with

$$0 < \theta_0 = \text{Im } \theta < \pi/(4k + 2).$$

Let 
$$\alpha \ge 0$$
 and set  $f_{\alpha}(x) = (\alpha^2 e^{(2k+1)i\theta_0} x^{4k+2} + 1)^{1/2}$ . Then

$$H(\alpha) = p^2 + (e^{2i\theta_0}x^2 + ge^{(2k+3)i\theta_0/2}x^{2k+1})\overline{f_{\alpha}}(x) |f_{\alpha}(x)|^{-2}.$$

Now we proceed to calculate the terms of the right hand side of the inequality to be proved.

$$\begin{aligned} \operatorname{Re}\langle u, \cos \gamma_{\alpha} \operatorname{H}(\alpha) u \rangle &= \operatorname{Re}\langle u, \cos \gamma_{\alpha} p^{2} u \rangle \\ &+ \langle u, \cos \gamma_{\alpha} (x^{2} \cos 2\theta_{0} + g x^{2k+1} \cos (2k+3)\theta_{0}/2) \mid f_{\alpha} \mid^{-2} \operatorname{Re} f_{\alpha} u \rangle \\ &+ \langle u, \cos \gamma_{\alpha} (x^{2} \sin 2\theta_{0} + g x^{2k+1} \sin (2k+3)\theta_{0}/2) \mid f_{\alpha} \mid^{-2} \operatorname{Im} f_{\alpha} u \rangle. \end{aligned}$$

Similarly

$$\operatorname{Im}\langle u, \sin \gamma_{\alpha} H(\alpha) u \rangle = \operatorname{Im}\langle u, \sin \gamma_{\alpha} p^{2} u \rangle$$

$$+ \langle u, \sin \gamma_{\alpha} (x^{2} \cos 2\theta_{0} + g x^{2k+1} \cos (2k+3)\theta_{0}/2) | f_{\alpha}|^{-2} \operatorname{Im} f_{\alpha} u \rangle$$

$$- \langle u, \sin \gamma_{\alpha} (x^{2} \sin 2\theta_{0} + g x^{2k+1} \sin (2k+3)\theta_{0}/2) | f_{\alpha}|^{-2} \operatorname{Re} f_{\alpha} u \rangle.$$

Then

Re 
$$\langle u, \cos \gamma_{\alpha} H(\alpha) u \rangle + \text{Im} \langle u, \sin \gamma_{\alpha} H(\alpha) u \rangle$$
  
= Re  $\langle u, \cos \gamma_{\alpha} p^{2} u \rangle + \text{Im} \langle u, \sin \gamma_{\alpha} p^{2} u \rangle$   
+  $\langle u, x^{2} | f_{\alpha} |^{-2} [\text{Re} f_{\alpha} \cos (2\theta_{0} + \gamma_{\alpha}) + \text{Im} f_{\alpha} \sin (2\theta_{0} + \gamma_{\alpha})] u \rangle$   
+  $\langle u, g x^{2k+1} | f_{\alpha} |^{-2} [\text{Re} f_{\alpha} \cos (\gamma_{\alpha} + (2k+3)\theta_{0}/2) + \text{Im} f_{\alpha} \sin (\gamma_{\alpha} + (2k+3)\theta_{0}/2)] u \rangle$ 

Let us construct  $\gamma_{\alpha}(x)$  so that the second term of the right hand side of the last inequality is positive and the third one vanishes.

If  $\gamma_{\alpha}(x) \neq -(2k+3)\theta_0/2$  then  $\sin \left[\gamma_{\alpha}(x) + (2k+3)\theta_0/2\right] \neq 0$  and the condition Re  $f_{\alpha}[\cos \gamma_{\alpha} + (2k+3)\theta_0/2] + \operatorname{Im} f_{\alpha} \sin \left[\gamma_{\alpha} + (2k+3)\theta_0/2\right] = 0$  is equivalent to the following

- 
$$\tan \arg (f_{\alpha}) = \cot \left[ \gamma_{\alpha} + (2k+3)\theta_{0}/2 \right] = \tan \left[ \pi/2 - \gamma_{\alpha} - (2k+3)\theta_{0}/2 \right].$$

One solution to this equation is given by

$$\gamma_{\alpha}(x) = \pi/2 + \arg(f_{\alpha}(x)) - (2k+3)\theta_0/2$$
 (2.4)

Note that for any  $\alpha > 0$  inf  $\arg(f_{\alpha}(x)) = 0$  and  $\sup_{x} \arg(f_{\alpha}(x)) = (2k+1)\theta_{0}/2$ . On the other hand,  $f_{0}(x) = 1$ , for all x; hence

$$\inf_{x} \cos \gamma_{\alpha}(x) \ge \cos (\pi/2 - \theta_0) = \sin \theta_0 = C_0 > 0,$$

for any  $\alpha \ge 0$ , and for this choice of  $\gamma_{\alpha}(x)$  assumption (2) of Theorem 2.1 is satisfied.

Now we need to show that if  $\gamma_{\alpha}(x)$  is given by (2.4) then

Re 
$$f_{\alpha}(x)$$
 cos  $(2\theta_0 + \gamma_{\alpha}(x)) + \text{Im } f_{\alpha}(x)$  sin  $(2\theta_0 + \gamma_{\alpha}(x)) \ge 0$ ,  $\forall x$  (2.5)

We have  $2\theta_0 + \gamma_{\alpha}(x) = \pi/2 + \arg(f_{\alpha}(x)) - (2k-1)\theta_0/2$ . Thus for  $\alpha = 0$ , (2.5) is trivially verified. For  $\alpha > 0$ ,

$$\sup_{x} (\gamma_{\alpha}(x) + 2\theta_{0}) = \theta_{0} + \pi/2 \quad \text{and} \quad \inf_{x} (\gamma_{\alpha}(x) + 2\theta_{0}) = \pi/2 - (2k-1)\theta_{0}/2;$$

therefore  $\sin(\gamma_{\alpha}(x)+2\theta_0) > 0$  for all x and (2.5) is now equivalent to  $\cot(2\theta_0 + \gamma_{\alpha}(x)) \ge -\tan\arg(f_{\alpha}(x))$ , i. e.

$$\tan\left((2k-1)\theta_0/2 - \arg\left(f_{\alpha}(x)\right)\right) \ge \tan\left(-\arg\left(f_{\alpha}(x)\right)\right)$$

and this last inequality is certainly satisfied  $\forall x$ .

Now the inequality in assumption (1) of Theorem 2.1 holds with  $\xi = 1$  and  $\eta = 0$  for any  $\alpha \ge 0$ . In order to prove part (3) we need a more explicit expression for  $\arg(f_{\alpha}(x))$ ,  $\alpha > 0$ :  $\operatorname{Im} f_{\alpha}^{2}(x) = a^{2}x^{4k+2} \sin(2k+1)\theta_{0}$  and  $\operatorname{Re} f_{\alpha}^{2}(x) = \alpha^{2}x^{4k+2} \cos(2k+1)\theta_{0} + 1$ . Thus,

$$\forall x \neq 0$$
,  $\arg(f_{\alpha}(x)) = (1/2) \arccos[\cot(2k+1)\theta_0 + (\alpha^2 x^{4k+2} \sin(2k+1)\theta_0)^{-1}].$ 

An easy calculation yields

$$\frac{d}{dx}(\cos \gamma_{\alpha})(x) = -(2k+1)\sin \gamma_{\alpha}(x) \left[\alpha^{2} x^{4k+1} \sin (2k+1)\theta_{0}\right] [h(x)]^{-1}$$

where  $h(x) = (1+a^2)\alpha^4 x^{8k+4} \sin^2(2k+1)\theta_0 + 2a\alpha^2 x^{4k+2} \sin(2k+1)\theta_0 + 1$ , with  $a = \cot(2k+1)\theta_0$ . If we assume  $\theta_0 < \pi/(8k+4)$  then the rational term in the last equation is easily seen to be bounded by 1, uniformly in  $\alpha$ . Since  $\frac{d}{dx}(\cos \gamma_0)(x) = 0$ ,  $\forall x$ , we conclude that

$$\left\| \frac{d}{dx} \cos \gamma_{\alpha} \right\|_{\infty} \le (2k+1) = C_1 < +\infty, \quad \forall \alpha \ge 0.$$

Similarly one can prove  $\left\| \frac{d}{dx} \sin \gamma_{\alpha} \right\|_{\infty} \le C_1$  and this completes the proof of Lemma 2.4.

Now set  $\Delta = \{ z \in \mathbb{C} \mid \text{ there is } \overline{\alpha}(z) > 0 \text{ s. t. } z \notin \sigma(H(\alpha)) \text{ and } \| (H(\alpha) - z)^{-1} \| \text{ is uniformly bounded for } 0 \le \alpha < \overline{\alpha}(z) \}.$ 

It is well known (see [7]) that  $\Delta$  is open and, given any compact  $\Gamma \subset \Delta$  there is  $\overline{\alpha}(\Gamma) > 0$  such that  $\| (H(\alpha) - z)^{-1} \|$  is uniformly bounded for  $z \in \Gamma$  and  $0 \le \alpha < \overline{\alpha}(\Gamma)$ .

Unlike the situation presented in [6], Example 5, we are dealing with operators whose numerical ranges invade the whole complex plane as  $\alpha \to 0$ . For this reason we need the following.

LEMMA 2.5. — The operator family  $\{H(\alpha)\}_{\alpha \geq 0}$  satisfies  $\Delta \neq \emptyset$ .

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*Proof.* — By Theorem 2.2  $\sigma_{ess}(H(0)) = \emptyset$  and by Theorem 2.3 for any  $z \in C$  there exists  $\overline{\alpha}(z) > 0$  such that  $z \notin \sigma_{ess}(H(\alpha))$  for  $0 \le \alpha < \overline{\alpha}(z)$ . Thus, Lemma 5.1 of [6] applies and if  $z \notin \sigma_d(H(0))$  then  $z \in \Delta$ , unless there exist two sequences  $\{\alpha_n\} \subset \mathbb{R}_+$  and  $\{u_n\} \subset \mathbb{L}^2(\mathbb{R})$  such that  $\alpha_n \to 0$ ,  $u_n \in D(H(\alpha_n))$ ,  $\|u_n\| \to 0$ ,  $u_n \to 0$  and  $\|(z - H(\alpha_n))u_n\| \to 0$ . In order to exclude the second alternative we notice that  $u_n$  can be chosen in  $C_0^\infty(\mathbb{R})$ , since it is a core for  $H(\alpha_n)$ . Now fix  $\chi \in C_0^\infty(\mathbb{R})$  and let  $M_n = 1 - \chi_n$  be the sequence of multiplication operators specified in the proof of Theorem 2.1. If we define

$$M_n^+(x) = M^+(x/n) = \begin{cases} 0, & \text{if } x < 0 \\ 1 - \chi(x/n), & \text{if } x \ge 0 \end{cases}$$

and

$$\mathbf{M}_{n}^{-}(x) = \mathbf{M}^{-}(x/n) = \begin{cases} 1 - \chi(x/n), & \text{if } x \le 0 \\ 0, & \text{if } x > 0, \end{cases}$$

then  $M_n = M_n^+ + M_n^-$  and by (ii) of Hypothesis 3 in [6], which follows from Theorem 2.1 and Lemma 2.4, there exists a > 0 such that, for each n,

$$\lim_{m \to +\infty} \| M_n u_m \|^2 = \lim_{m \to +\infty} \sup (\| M_n^+ u_m \|^2 + \| M_n^- u_m \|^2) \ge a^2.$$

Thus, for each n

either 
$$\lim_{m} \sup \|\mathbf{M}_{n}^{+} u_{m}\| \ge a/2 \tag{2.6a}$$

or 
$$\lim_{m} \sup || \mathbf{M}_{n}^{-} \mathbf{u}_{m} || \ge a/2$$
 (2.6b)

Then the second alternative can be specified either with a sequence  $v_n^+ = M_n^+ u_{m(n)}$ , such that  $v_n^+(x) = 0$  for x < n, or with a sequence  $v_n^- = M_n^- u_{m(n)}$ , such that  $v_n^-(x) = 0$  for x > -n, by suitably choosing m = m(n).

In fact, let (2.6a) hold. We have

$$||(z - H(\alpha_m))M_n^+ u_m|| \le ||(z - H(\alpha_m))u_m|| + ||[M_n^+, H(\alpha_m)]u_m|| \quad (2.7)$$

and, since

$$[\mathbf{M}_{n}^{+}, \mathbf{H}(\alpha)] = [\mathbf{M}_{n}^{+}, p^{2}] = 2\mathrm{in}^{-1} \frac{d\mathbf{M}^{+}}{dx} (x/n) - n^{-2} \frac{d^{2}\mathbf{M}^{+}}{dx^{2}} (x/n),$$

it follows from (2.1) that

$$\|[\mathbf{M}_{n}^{+}, \mathbf{H}(\alpha)]u\| \le \operatorname{cn}^{-1}(\|\mathbf{H}(\alpha)u\| + \|u\|)$$
 (2.8)

where c > 0 is a constant independent of  $\alpha$  and n. Now, combining (2.6a), (2.7) and (2.8), a suitable choice of m = m(n) and normalization yield

$$\lim_n \| \left(z - H(\alpha_{\mathit{m(n)}}) v_n^+ \| = 0, \qquad v_n^+ \underset{\overrightarrow{W}}{\rightarrow} 0 \quad \text{and} \quad \| \, v_n^+ \| = 1 \, .$$

The analogous properties can be specified for the sequence  $\{v_n^-\}$  if (2.6b) holds. Now the contradiction follows from

$$|| [H(\alpha_{m(n)}) - z]v_n^+|| \ge \operatorname{dist}(z, \operatorname{E}_n^+(\alpha_{m(n)}))$$

where  $E_n^+(\alpha) = \{ \langle H(\alpha)v, v \rangle \mid v \in C_0^\infty(\mathbb{R}), ||v|| = 1, v(x) = 0 \text{ for } x < n \}$ . In fact we have

$$\operatorname{Re} \langle (p^2 + V_{\alpha})v, v \rangle \ge \langle (\operatorname{Re} V_{\alpha})v, v \rangle$$
 and  $\lim_{\substack{x \to +\infty \\ a \to 0}} \operatorname{Re} V_{\alpha}(x) = +\infty$ 

i. e. there are not two sequences  $\{x_k\} \to +\infty$  and  $\{\alpha_k\} \to 0$  such that Re  $V_{\alpha_k}(x_k)$  is bounded from above. Thus, since  $\alpha_{m(n)} \to 0$ , for any  $\delta > 0$  there exists  $n_0 \in \mathbb{N}$  such that  $\| [H(\alpha_{m(n)}) - z]v_n^+ \| \ge \delta > 0$ ,  $\forall n \ge n_0$ . Similarly we obtain a contradiction if (2.6b) holds, since  $\lim_{\substack{x \to -\infty \\ \alpha \to 0}} \operatorname{Im} V_{\alpha}(x) = -\infty$ , and this completes the proof of Lemma 2.5.

COROLLARY 2.6. — Every eigenvalue of H(0) is stable (in the sense of Kato [7]) with respect to the family  $\{H(\alpha)\}_{\alpha\geq 0}$ .

Proof. — By Lemma 2.4 and 2.5 the family  $\{H(\alpha)\}_{\alpha\geq 0}$  satisfies Hypothesis 3 of [6] with the operators  $M_n$  specified in the proof of Theorem 2.1. Moreover, by Theorems 2.2 and 2.3, for any fixed  $\lambda\in C$  and  $\delta>0$  there exists  $\alpha_0>0$  such that  $\mathrm{dist}\,(\lambda,\,\sigma_{ess}(H(\alpha)))\geq\delta$  for  $0\leq\alpha\leq\alpha_0$ . Finally we can decompose  $M_n=M_n^++M_n^-$ , as specified in the proof of Lemma 2.5, where  $M_n^+$  and  $M_n^-$  also satisfy (iii) of Hypothesis 3 in [6]. Since  $\lim_{x\to+\infty} \mathrm{Re}\,V_\alpha(x)=+\infty$ , proceeding as in the last part of the preceding lemma,  $\lim_{x\to+\infty} \mathrm{Re}\,V_\alpha(x)=+\infty$ , proceeding as in the last part of the preceding lemma,

we find  $\alpha_1 > 0$  and  $n_0 \in \mathbb{N}$  such that

$$d_n^+(\lambda,\alpha) = \inf_{u \in \mathrm{D}(\mathrm{H}(\alpha)), \, ||\mathrm{M}_n^+ u|| \, = \, 1} ||\, (\lambda \, - \, \mathrm{H}(\alpha)) \mathrm{M}_n^+ u \, || \, \geq \, \delta \, > \, 0$$

for  $\alpha < \alpha_1$  and  $n > n_0$ . Similarly, since  $\lim_{\substack{x \to -\infty \\ \alpha \to 0}} \text{Im } V_{\alpha}(x) = -\infty, \ d_n^-(\lambda, \alpha) \ge \delta$ 

for  $\alpha < \alpha_1$  and  $n > n_0$ . Therefore we can apply Theorem 5.8 of [6] and the stability of eigenvalues is proved.

As a consequence, we can now prove the result announced in the introduction. For convenience, we return to the more explicit notation of Theorem 2.2, by replacing  $H(\alpha)$  with  $e^{\theta}H(\alpha, \theta)$ .

Theorem 2.7. — Let  $H(\alpha,\theta)$  be the holomorphic family of type A in  $L^2(R)$  defined for  $|\operatorname{Im} \theta| < \pi/(4k+2)$  and  $\alpha > 0$  by (2.3), self-adjoint for  $\theta \in R$ . Then there exists  $g_0 > 0$  such that for  $0 < g < g_0$  there is  $\alpha_g > 0$  with the property that  $\forall \alpha \in (0,\alpha_g)$ ,  $H(\alpha,0)$  admits second sheet poles of the resolvent in the following sense: if  $S \subset L^2(R)$  is the dense set of all dilation analytic vectors for  $|\operatorname{Im} \theta| < \pi/(4k+2)$  and  $\psi \in S$ , the scalar product

$$g_{\psi}(\mathbf{E}) = \langle [\mathbf{H}(\alpha, 0) - \mathbf{E}]^{-1} \psi, \psi \rangle,$$

which is a priori analytic in  $C\setminus \sigma(H(\alpha, 0))$ , admits a meromorphic continua-

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tion onto the second sheet across the cut  $[-g\alpha^{-1}, +\infty) = \sigma_{ess}(H(\alpha, 0))$  (from the upper half-plane of the first sheet) to the region

$$\{ E \mid -\tan (\operatorname{Im} \theta)(\operatorname{Re} E + g\alpha^{-1}) \leq \operatorname{Im} E \leq 0 \} \setminus \sigma_{ess}(H(\alpha, \theta)),$$

for any  $\theta$ ,  $0 < \operatorname{Im} \theta < \pi/(4k + 2)$ . For  $i \in \mathbb{N}$ , let  $\alpha(i) > 0$  be such that for  $\alpha < \alpha(i)$  the eigenvalue  $E_i(\alpha)$  of  $H(\alpha, \theta)$  exists by stability with respect to the *i*-th eigenvalue of  $H(0, \theta)$ . Then for fixed  $\bar{i} \in \mathbb{N}$  and  $\alpha \le \min{(\alpha(1), \ldots, \alpha(\bar{i}))}$  the set  $\{E(\alpha) \mid E(\alpha) \text{ is a pole of } g_{\psi}(E) \text{ for some } \psi \in S \}$  contains at least the finite set of eigenvalues  $\{E_i(\alpha) \mid i=1,2,\ldots,\bar{i}\}$ . For each  $i \in \mathbb{N}$ ,  $E_i(\alpha)$  tends to the corresponding second sheet pole of the generalized resolvent  $\mathscr{R}(E)$  (specified in Theorem 2.2) of the closure of the symmetric operator  $p^2 + x^2 + gx^{2k+1}$ .

*Proof.* — It is a consequence of Corollary 2.6 since, by the usual analyticity arguments (see e. g. [10])

$$g_{\psi}(\mathbf{E}) = \langle [\mathbf{H}(\alpha, \theta) - \mathbf{E}]^{-1} \psi(\theta), \psi(\overline{\theta}) \rangle$$

identically, for  $0 < \text{Im } \theta < \pi/(4k+2)$ , Im E > 0, where  $\psi(\theta)(x) = e^{\theta/4}\psi(e^{\theta/2}x)$  for  $\psi \in S$ .

An analogous stability argument can be used to obtain an existence proof of shape resonances for a suitable class of Schrödinger operators in  $L^2(\mathbb{R}_+)$ . In this case the results of Vock and Hunziker [6] can be directly applied, having care to set the necessary Dirichlet condition at the origin. Explicitly, if we set

$$V_0(z) = z^2 - gz^{2k+1}, \quad k \in \mathbb{N}, \quad g > 0, \quad z \in \mathbb{C},$$
 (2.9)

then for the operator family defined in  $L^2(R_+)$  by the differential expression  $A(0, \theta) = e^{-\theta}p^2 + V_0(e^{\theta/2}x) = e^{-\theta}p^2 + e^{2\theta}x^2 - ge^{(2k+3)\theta/2}x^{2k+1}$ ,  $(0 < \text{Im } \theta < \min(\pi/4, 2\pi/(2k+3)))$  we can prove results analogous to Theorem 2.2. In particular  $A(0, \theta)$  is a holomorphic family of type A in  $L^2(R_+)$  for  $0 < \text{Im } \theta < \min(\pi/4, 2\pi/(2k+3))$  when defined on

$$D = D(A(0, \theta)) = \{ u \mid u \in H^2(R_+) \cap L^2_{(2k+1)/2}(R_+), u(0) = 0 \}$$
 (2.10)

and we can show the existence of eigenvalues  $E_i = E_i(g)$  (for  $0 < g < g_0$ , for some  $g_0 > 0$ ) with Im  $E_i < 0$ ,  $i \in N$ , which tend to the odd eigenvalues of the harmonic oscillator as  $g \to 0$ .

THEOREM 2.8. — Let  $V_{\alpha}(x)$ ,  $\alpha > 0$ ,  $x \in \mathbb{R}_+$  be a family of real-valued functions enjoying the following properties:

- i) there is  $\theta_0 \in (0, \min{(\pi/4, 2\pi/(2k+2))})$  such that, for any  $\alpha > 0$ ,  $V_{\alpha}(x)$  is the restriction to  $z \in \mathbb{R}_+$  of a function  $V_{\alpha}(z)$  holomorphic at least in the sector  $\{z \in \mathbb{C} \mid |\arg z| < \theta_0\}$ , and bounded near z = 0;
- ii) for  $\alpha > 0$ ,  $|\operatorname{Im} \theta| < \theta_0$ ,  $V_{\alpha}(e^{\theta/2}x) \to c(\alpha)$  as  $x \to +\infty$ , for some  $c(\alpha) \in \mathbb{R}$ , and  $\lim_{\alpha \to 0} c(\alpha) = -\infty$ ;

- iii) for any fixed  $\theta$  with  $0 < \text{Im } \theta < \theta_0$ ,  $V_{\alpha}(e^{\theta/2}x) \to V_0(e^{\theta/2}x)$  as  $\alpha \to 0$  uniformly on the compact subsets of  $R_+(V_0(z))$  as defined in (2.9);
- iv) the family  $\{e^{\theta}V_{\alpha}(e^{\theta/2}x)\}_{\alpha\geq 0}$  satisfies the hypotheses of Theorem 6.1 in [6], with  $C_0^{\infty}(R)$  replaced by D, as defined in (2.10). Then:
- (1) for any  $\alpha > 0$  the differential expression  $e^{-\theta}p^2 + V_{\alpha}(e^{\theta/2}x)$  defines a holomorphic family of type A in  $L^2(R_+)$ , denoted by  $A(\alpha, \theta)$ , for  $|\operatorname{Im} \theta| < \theta_0$ , if defined on  $D(A(\alpha, \theta)) = \{ u \mid u \in H^2(R_+), u(0) = 0 \}$ , with D as a core. Furthermore  $\sigma_{ess}(A(\alpha, \theta)) = \{ z = c(\alpha) + \lambda e^{-\theta} \mid \lambda \ge 0 \}$ .

Let  $E_0$  be an eigenvalue of  $A(0,\theta)$  (whence  $\operatorname{Im} E_0 < 0$ ); suppose that for some  $\varepsilon > 0$ ,  $n_0 \in \mathbb{N}$  and  $\alpha_0 > 0$ , dist  $(E_0, F_n(\alpha)) \ge \varepsilon > 0$ ,  $\forall n \ge n_0$ ,  $0 < \alpha < \alpha_0$ , where  $F_n(\alpha) = \{ \langle A(\alpha, \theta)u, u \rangle \mid ||u|| = 1, u \in D(A(\alpha, \theta)), \text{ supp } u \subset [n, +\infty) \}$ . Then:

- (2)  $E_0$  is a stable eigenvalue with respect to the family  $\{A(\alpha, \theta)\}_{\alpha \geq 0}$ ; in particular there is  $\overline{\alpha} > 0$  such that for all  $\alpha \in (0, \overline{\alpha})$  there exists  $E(\alpha) \in \sigma_d(A(\alpha, \theta))$  with Im  $E(\alpha) < 0$  and  $E(\alpha) \to E_0$  as  $\alpha \to 0$ ;
- (3) if S is the dense set of dilation analytic vectors for  $|\operatorname{Im} \theta| < \theta_0$ , then for all  $\alpha \in (0, \overline{\alpha})$  and  $\psi \in S$  the function  $h_{\psi}(E) = \langle [A(\alpha, 0) E]^{-1}\psi, \psi \rangle$ , which is *a priori* analytic in  $C \setminus \sigma(A(\alpha, 0))$ , admits a meromorphic continuation onto the second sheet to an open connected region of the lower half plane (from the upper half plane of the first sheet) across the cut

$$\sigma_{ess}(A(\alpha, 0)) = [c(\alpha), + \infty).$$

Moreover there is  $\psi \in S$  such that the continuation of  $h_{\psi}$  has a simple pole at  $E(\alpha)$ .

Natural examples are self-adjoint operators of the form

$$A(\alpha,0) = -d^2/dx^2 + (x^2 + \alpha P^{(2k)}(x) - gx^{2k+1})(\alpha x^{\nu\rho} + 1)^{1/\nu}.$$

Here k, v and  $\rho$  are fixed  $(k \in \mathbb{N}, v > 0 \text{ and } \rho = 2k + 1) \alpha > 0$ ,  $P^{(2k)}(x)$  is an arbitrary polynomial of order not larger than 2k,

$$D(A(\alpha,0)) = \left\{ \, u \mid u \in H^2(\mathbb{R}_+), \ u(0) = 0 \, \right\}.$$

Since in this case  $\lim_{\substack{\alpha \to 0 \\ x \to +\infty}} \operatorname{Im} \left[ e^{\theta} V_{\alpha}(e^{\theta/2}x) \right] = -\infty$ , every eigenvalue  $E_i$ 

of  $A(0, \theta)$  satisfies the required inequality dist  $(E_i, F_n(\alpha)) \ge \varepsilon_i > 0$  for  $n \ge n_i$ ,  $0 < \alpha < \alpha_i$ . So, by the previous theorem,  $A(\alpha, 0)$  provides an example of shape resonances in the standard sense of dilation analyticity.

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