



Mathematical Problems in Mechanics

Explicit energy-minimizers of incompressible elastic brittle bars under uniaxial extension

Minimiseurs de l'énergie explicites d'une barre incompressible, élastique, mais fragile, soumise à une extension uniaxiale

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ARTICLE INFO

Article history:

Received 8 March 2010

Accepted after revision 3 September 2010

Presented by John Ball

ABSTRACT

A rectangular bar made of a hyperelastic, but brittle, incompressible homogeneous and isotropic material is subject to uniaxial extension. We prove that the energy minimizers are, depending on the toughness coefficient of the material, either the homogeneous deformation, or the family of deformations for which a horizontal fracture breaks the material in two rectangular pieces, each of which is a rigid motion of the undeformed piece.

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RÉSUMÉ

Une barre rectangulaire faite d'un matériau incompressible, homogène, isotropique, hyperélastique, mais fragile, est soumise à une extension uniaxiale. Nous prouvons que les minimiseurs de l'énergie correspondent, selon le coefficient de ténacité du matériau, soit à une déformation homogène, soit à une famille de déformations pour lesquelles une fracture horizontale casse le matériau en deux parties rectangulaires, chacune étant un mouvement rigide de la pièce non déformée.

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1. Introduction

Let $\Omega := (-R, R) \times (0, L)$ be the reference configuration of a homogeneous, isotropic, incompressible, hyperelastic material. Sivaloganathan and Spector [8] proved that, under some natural conditions on the stored-energy function, the absolute minimizer subject to uniaxial extension is the homogeneous deformation.

In this paper, we allow the material to undergo brittle fracture. The natural space to pose the problem is the space SBV of special bounded variation, and the energy functional of a deformation $u \in SBV(\Omega, \mathbb{R}^2)$ is defined as

$$I[u] := \int_{\Omega} \Phi(|\nabla u(x, y)|) \, dx \, dy + \kappa \mathcal{H}^1(J_u), \quad (1)$$

where $\kappa > 0$ is the toughness coefficient of the material, J_u stands for the set of jump discontinuities of u , and \mathcal{H}^1 denotes the 1-dimensional Hausdorff measure (see [1] for properties of SBV functions and applications to free discontinuity prob-

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lems). Here ∇u denotes the absolutely continuous part of the distributional gradient of u , and the norm $|\cdot|$ of a matrix is taken as the square root of the sum of the squares of its components. The fact that the stored-energy function only depends on the norm of the deformation gradient follows from standard representation theorems, since, as well as the frame-indifference requirement, our material is assumed to be isotropic and incompressible. Hence $\int_{\Omega} \Phi(|\nabla u|)$ represents the elastic energy of the deformation. The function Φ is assumed to be convex and increasing, as usual when dealing with polyconvex functions (see [3]).

We impose the boundary conditions

$$u_2(x, 0) = 0, \quad u_2(x, L) = \lambda L, \quad \text{a.e. } x \in (-R, R), \tag{2}$$

where the equality is taken in the sense of traces, and $u = (u_1, u_2)$. Here $\lambda > 1$ is a given number, and, hence, (2) models uniaxial extension of the material. We impose no boundary condition on the remaining part of the boundary. To avoid the trivial non-uniqueness of minimizers given by horizontal translations, we impose additionally

$$\int_{\Omega} u_1(x, y) \, dx \, dy = 0. \tag{3}$$

We prove that, depending on the parameter κ , the minimizers of I are either the homogeneous deformation, or the family of deformations for which a horizontal fracture breaks the material in two rectangular pieces, each of which is just a translation or a π -rotation of the undeformed piece. To be precise, the homogeneous deformation is given by

$$u^H(x, y) := \left(\frac{1}{\lambda}x, \lambda y \right), \quad (x, y) \in \Omega,$$

whereas, for each $y_0 \in (0, L)$ and $\alpha \in \mathbb{R}$, we have the following family of horizontal-fracture deformations at level y_0 :

$$\begin{aligned} u^{F_{1,y_0,\alpha}}(z) &:= \begin{cases} z + (\alpha(L - y_0), 0) & \text{if } z \in \Omega_{y_0}^-, \\ z + (-\alpha y_0, (\lambda - 1)L) & \text{if } z \in \Omega_{y_0}^+, \end{cases} & u^{F_{2,y_0,\alpha}}(z) &:= \begin{cases} -z + (\alpha(L - y_0), 0) & \text{if } z \in \Omega_{y_0}^-, \\ z + (-\alpha y_0, (\lambda - 1)L) & \text{if } z \in \Omega_{y_0}^+, \end{cases} \\ u^{F_{3,y_0,\alpha}}(z) &:= \begin{cases} z + (\alpha(L - y_0), 0) & \text{if } z \in \Omega_{y_0}^-, \\ -z + (-\alpha y_0, (\lambda + 1)L) & \text{if } z \in \Omega_{y_0}^+, \end{cases} & u^{F_{4,y_0,\alpha}}(z) &:= \begin{cases} -z + (\alpha(L - y_0), 0) & \text{if } z \in \Omega_{y_0}^-, \\ -z + (-\alpha y_0, (\lambda + 1)L) & \text{if } z \in \Omega_{y_0}^+. \end{cases} \end{aligned}$$

We have denoted $\Omega_{y_0}^- := (-R, R) \times (0, y_0)$ and $\Omega_{y_0}^+ := (-R, R) \times (y_0, L)$. All maps $u^{F_{i,y_0,\alpha}}$ represent realistic deformations, in the sense that all are one-to-one. However, the boundary conditions (2) model uniaxial extension only for $u^{F_{1,y_0,\alpha}}$.

The proof, which of course uses many ideas of Sivaloganathan and Spector [8], also depends on the rigidity result for SBV deformations due to Chambolle, Giacomini and Ponsiglione [4], and on the structure theorem on the boundary of sets of finite perimeter by Federer [7].

2. Main theorem

The following result follows at once from polar decomposition. Recall that $SO(2)$ denotes the set of orthogonal 2×2 matrices with positive determinant.

Lemma 1. *Let $A \in \mathbb{R}^{2 \times 2}$ satisfy $\det A = 1$. Then $|A| \geq \sqrt{2}$, with equality if and only if $A \in SO(2)$.*

Theorem 2. *Let $\Phi : [\sqrt{2}, \infty) \rightarrow \mathbb{R}$ be a convex and strictly increasing function. Let*

$$R > 0, \quad L > 0, \quad \lambda > 1, \quad \kappa > 0, \quad \kappa_0 := L(\Phi(\sqrt{\lambda^2 + \lambda^{-2}}) - \Phi(\sqrt{2})).$$

Let $\Omega := (-R, R) \times (0, L)$. Let \mathcal{A} be the set of $u \in SBV(\Omega, \mathbb{R}^2)$ such that $\det \nabla u = 1$ a.e., and conditions (2) and (3) hold. Define I as in (1). Then the following assertions hold:

- i) *If $\kappa < \kappa_0$ then the set of minimizers of I in \mathcal{A} is $\{u^{F_{i,y_0,\alpha}} : i \in \{1, 2, 3, 4\}, y_0 \in (0, L), \alpha \in \mathbb{R}\}$.*
- ii) *If $\kappa = \kappa_0$ then the set of minimizers of I in \mathcal{A} is $\{u^{F_{i,y_0,\alpha}} : i \in \{1, 2, 3, 4\}, y_0 \in (0, L), \alpha \in \mathbb{R}\} \cup \{u^H\}$.*
- iii) *If $\kappa > \kappa_0$ then the set of minimizers of I in \mathcal{A} is $\{u^H\}$.*

Proof. An immediate computation shows that, for any $i \in \{1, 2, 3, 4\}$, any $y_0 \in (0, L)$ and any $\alpha \in \mathbb{R}$,

$$I[u^{F_{i,y_0,\alpha}}] = 2RL\Phi(\sqrt{2}) + 2R\kappa, \quad I[u^H] = 2RL\Phi(\sqrt{\lambda^2 + \lambda^{-2}}). \tag{4}$$

Let $u \in \mathcal{A}$. We first show that $I[u]$ is bounded from below by a convex combination of $I[u^{F_{i,y_0,\alpha}}]$ and $I[u^H]$. Define F_u as the set of $x \in (-R, R)$ such that $J_{u(x,\cdot)} \neq \emptyset$. By slicing properties of BV functions (see, e.g., [1, Sect. 3.11]), $u(x, \cdot) \in$

$SBV((0, L), \mathbb{R}^2)$ for a.e. $x \in (-R, R)$, and, in particular, $u(x, \cdot) \in W^{1,1}((0, L), \mathbb{R}^2)$ for a.e. $x \in (-R, R) \setminus F_u$. Moreover, by [1, Rk. 3.109],

$$J_{u(x,\cdot)} = \{y \in (0, L) : (x, y) \in J_u\} \quad \text{a.e. } x \in (-R, R). \tag{5}$$

Reasoning as in [8], we obtain that

$$\int_{(-R,R) \setminus F_u} \int_0^L \Phi(|\nabla u(x, y)|) \, dy \, dx \geq (2R - \mathcal{L}^1(F_u))L\Phi(\sqrt{\lambda^2 + \lambda^{-2}}). \tag{6}$$

Here \mathcal{L}^1 denotes the Lebesgue measure over \mathbb{R} . The proof of [8] assumed that the deformation is of class C^1 , but standard slicing techniques show that the result also holds for Sobolev maps. Let $\pi : \mathbb{R}^2 \rightarrow \mathbb{R}$ be the projection onto the first coordinate. Using that projections have Lipschitz constant 1, and Eq. (5), we obtain that

$$\mathcal{H}^1(J_u) \geq \mathcal{L}^1(\pi J_u) = \mathcal{L}^1(\{x \in (-R, R) : (\{x\} \times (0, L)) \cap J_u \neq \emptyset\}) = \mathcal{L}^1(F_u). \tag{7}$$

Therefore, thanks to Lemma 1,

$$\int_{F_u \times (0,L)} \Phi(|\nabla u|) \, dx \, dy + \kappa \mathcal{H}^1(J_u) \geq \mathcal{L}^1(F_u)L\Phi(\sqrt{2}) + \kappa \mathcal{L}^1(F_u). \tag{8}$$

Putting (6), (8) and (4) together, we conclude that

$$I[u] \geq \mathcal{L}^1(F_u)(\kappa - \kappa_0) + 2RL\Phi(\sqrt{\lambda^2 + \lambda^{-2}}) = \frac{\mathcal{L}^1(F_u)}{2R}I[u^{F_i, y_0, \alpha}] + \left(1 - \frac{\mathcal{L}^1(F_u)}{2R}\right)I[u^H], \tag{9}$$

for any $i \in \{1, 2, 3, 4\}$, any $y_0 \in (0, L)$ and any $\alpha \in \mathbb{R}$. This shows that $u^{F_i, y_0, \alpha}$ is a minimizer when $\kappa \leq \kappa_0$, and that u^H is a minimizer when $\kappa \geq \kappa_0$. We prove now that no more minimizers can exist.

Let $\kappa < \kappa_0$. Then $I[u^{F_i, y_0, \alpha}] < I[u^H]$ for any $i \in \{1, 2, 3, 4\}$, any $y_0 \in (0, L)$ and any $\alpha \in \mathbb{R}$. Let $u \in \mathcal{A}$ satisfy $I[u] = I[u^{F_i, y_0, \alpha}]$. By (9), (7), (8) and Lemma 1, we obtain that $\mathcal{H}^1(J_u) = \mathcal{L}^1(F_u) = 2R$ and $\nabla u \in SO(2)$ a.e. In order to apply the area formula (see, e.g., [1, Th. 2.91]), we notice that for \mathcal{H}^1 -a.e. $z \in J_u$, a unit vector $e(z) \in \mathbb{R}^2$ spans the approximate tangent space $T_z J_u$ of J_u at z . Therefore, the tangential differential of π along J_u at z is given by $\pi e(z)$. By the area formula and (5), we obtain that

$$2R = \mathcal{H}^1(J_u) \geq \int_{J_u} |\pi e(z)| \, d\mathcal{H}^1(z) = \int_{-R}^R \mathcal{H}^0(\{y \in (0, L) : (x, y) \in J_u\}) \, dx \geq 2R, \tag{10}$$

and, hence, we can take $e(z) = (1, 0)$ for \mathcal{H}^1 -a.e. $z \in J_u$.

By the rigidity result of Chambolle, Giacomini and Ponsiglione [4], there exist $I \subset \mathbb{N}$ and a Caccioppoli partition (i.e., a partition made of sets of finite perimeter) $\{E_i\}_{i \in I}$ of Ω such that

$$J_u = \bigcup_{i \in I} \partial^* E_i \cap \Omega, \quad \mathcal{H}^1\text{-a.e.} \tag{11}$$

and

$$u(z) = \sum_{i \in I} (R_i z + b_i) \chi_{E_i}(z), \quad \text{a.e. } z \in \Omega, \tag{12}$$

for some $\{R_i\}_{i \in I} \subset SO(2)$ and $\{b_i\}_{i \in I} \subset \mathbb{R}^2$. Here, ∂^* denotes the measure-theoretic boundary, which, for sets of finite perimeter, coincides \mathcal{H}^1 -a.e. with Federer’s boundary and with the reduced boundary (see, e.g., [1]). Moreover, thanks to the structure theorem of the boundary of planar sets of finite perimeter by Federer [7, 4.2.25] (see also [2, Cor. 1]), for each $i \in I$ there exist a $J_i \subset \mathbb{N}$ and a disjoint family $\{C_{ij}\}_{j \in J_i}$ of Jordan curves such that

$$\partial^* E_i = \bigcup_{j \in J_i} C_{ij}, \quad \mathcal{H}^1\text{-a.e.} \tag{13}$$

Consequently, $T_z C_{ij}$ is spanned by $(1, 0)$ for \mathcal{H}^1 -a.e. $z \in C_{ij} \cap \Omega$, every $i \in I$, and every $j \in J_i$. Fix $i \in I$ and $j \in J_i$. There exists a one-to-one Lipschitz arc-length parametrization $\gamma : [0, \mathcal{H}^1(C_{ij})] \rightarrow C_{ij}$ of C_{ij} (see, e.g., [2, Lemma 3]), and $T_z C_{ij}$ is spanned by $\gamma'(\gamma^{-1}(z))$ for \mathcal{H}^1 -a.e. $z \in C_{ij}$ (see, e.g., [1, Prop. 2.88]). Hence $\gamma'(t)$ is proportional to $(1, 0)$ for a.e. $t \in [0, \mathcal{H}^1(C_{ij})]$ such that $\gamma(t) \in \Omega$. This implies that C_{ij} describes a rectangle, with one edge contained in $\{-R\} \times [0, L]$, and another in $\{R\} \times [0, L]$. From the equality $\mathcal{H}^1(J_u) = 2R$ we obtain that there exists $y_0 \in (0, L)$ such that $J_u = (-R, R) \times \{y_0\}$, \mathcal{H}^1 -a.e.

Applying the fact that open connected sets with finite perimeter only admit the trivial Caccioppoli partition formed by themselves (see [2, Prop. 2]) to $(-R, R) \times (0, y_0)$ and to $(-R, R) \times (y_0, L)$, we obtain that Ω admits only one Caccioppoli partition satisfying (11), namely,

$$\{(-R, R) \times (0, y_0), (-R, R) \times (y_0, L)\}.$$

We then conclude from (12) that there exist $R_1, R_2 \in SO(2)$ and $b_1, b_2 \in \mathbb{R}^2$ such that

$$u(z) = (R_1 z + b_1)\chi_{(-R,R) \times (0,y_0)}(z) + (R_2 z + b_2)\chi_{(-R,R) \times (y_0,L)}(z), \quad \text{a.e. } z \in \Omega.$$

Conditions (2) and (3) imply, after some calculations, that

$$R_1 \in \{-I, I\}, \quad b_1 = (\alpha(L - y_0), 0), \quad (R_2, b_2) \in \{(I, (-\alpha y_0, (\lambda - 1)L)), (-I, (-\alpha y_0, (\lambda + 1)L))\},$$

for some $\alpha \in \mathbb{R}$.

Let $\kappa > \kappa_0$, and let $u \in \mathcal{A}$ satisfy $I[u] = I[u^H]$. Then, by (9), $\mathcal{L}^1(F_u) = 0$, so by (6), $\mathcal{H}^1(J_u) = 0$. Hence $u \in W^{1,1}(\Omega, \mathbb{R}^2)$, and we conclude, thanks to [8], that $u = u^H$.

Finally, let $\kappa = \kappa_0$, and let $u \in \mathcal{A}$ satisfy $I[u] = I[u^H]$. If $\mathcal{L}^1(F_u) = 0$, we proceed as in the case $\kappa > \kappa_0$ to conclude that $u = u^H$. If $\mathcal{L}^1(F_u) = 2R$, we proceed as in the case $\kappa < \kappa_0$ to conclude that $u = u^{F_i, y_0, \alpha}$ for some $i \in \{1, 2, 3, 4\}$, some $y_0 \in (0, L)$ and some $\alpha \in \mathbb{R}$. Suppose, for contradiction, that $0 < \mathcal{L}^1(F_u) < 2R$. From (6) we obtain that $|\nabla u| = \sqrt{\lambda^2 + \lambda^{-2}}$ a.e. in $((-R, R) \setminus F_u) \times (0, L)$. As $u(x, \cdot) \in W^{1,1}((0, L), \mathbb{R}^2)$ for a.e. $x \in (-R, R) \setminus F_u$, we conclude, reasoning as in [8], that

$$\nabla u(z) = \text{diag}\left(\frac{1}{\lambda}, \lambda\right) \quad \text{a.e. } z \in ((-R, R) \setminus F_u) \times (0, L), \tag{14}$$

whereas from (8) and Lemma 1 we obtain that

$$\nabla u(z) \in SO(2) \quad \text{a.e. } z \in F_u \times (0, L) \tag{15}$$

and $\mathcal{H}^1(J_u) = \mathcal{L}^1(F_u)$. Call

$$K := SO(2) \cup \text{diag}\left(\frac{1}{\lambda}, \lambda\right).$$

From (14) and (15) we obtain that $\nabla u \in K$ a.e. The result of Chaudhuri and Müller [5] provides, in particular, a rigidity estimate for Sobolev deformations with gradient in K . They proved an L^2 estimate, but a similar proof shows an L^p estimate for $p \in (1, \infty)$ (see also [6], where an alternative proof of the L^p rigidity estimate is given). Therefore, we can apply the SBV rigidity result of [4, Th. 1.2] to obtain that there exist $I \subset \mathbb{N}$ and a Caccioppoli partition $\{E_i\}_{i \in I}$ of Ω such that equalities (11) and (12) hold for some $\{R_i\}_{i \in I} \subset K$ and $\{b_i\}_{i \in I} \subset \mathbb{R}^2$. As above, for each $i \in I$ there exist a $J_i \subset \mathbb{N}$ and a disjoint family $\{C_{ij}\}_{j \in J_i}$ of Jordan curves such that (13) holds. A similar reasoning as in (10) concludes that $T_z J_u$ is spanned by $(1, 0)$ for \mathcal{H}^1 -a.e. $z \in J_u$. As above, each C_{ij} describes a rectangle, with one edge contained in $\{-R\} \times [0, L]$, and another in $\{R\} \times [0, L]$. This is a contradiction with the fact that $0 < \mathcal{H}^1(J_u) < 2R$. \square

If Φ is not assumed to be strictly increasing, but only monotonically increasing, then we still obtain that u^H is a minimizer of I in \mathcal{A} for $\kappa \geq \kappa_0$, and each $u^{F_i, y_0, \alpha}$ is a minimizer of I in \mathcal{A} for $\kappa \leq \kappa_0$, but some other minimizers may exist too.

Acknowledgements

Supported by Project MTM2009-07662 of the Spanish Ministry of Science and Innovation, and by Project PI2010-01 of the Basque Department of Education, Universities and Research.

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